

M-Theory as a Dynamical System

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Research Themes, References & Collaborators

AdS/CFT Correspondence & Integrability

- Work on **Holographic Defects** in collaboration with M. de Leeuw and C. Kristjansen (my INPP talk last month):
 - [1]. M. de Leeuw, C. Kristjansen, and G. Linardopoulos. "Scalar One-point Functions and Matrix Product States of AdS/dCFT". In: *Phys.Lett.* **B781** (2018), p. 238. arXiv: 1802.01598 [hep-th].
 - [2]. M. de Leeuw, C. Kristjansen, and G. Linardopoulos. "One-point Functions of Non-protected Operators in the $SO(5)$ Symmetric D3-D7 dCFT". In: *J. Phys.* **A50** (2017), p. 254001. arXiv: 1612.06236 [hep-th].
- Work on the **Spectral Problem of AdS/CFT** with M. Axenides and E. Floratos:
 - [3]. M. Axenides, E. Floratos, and G. Linardopoulos. "The Omega-Infinity Limit of Single Spikes". In: *Nucl. Phys.* **B907** (2016), p. 323. arXiv: 1511.03587 [hep-th],in collaboration with E. Floratos:
 - [4]. E. Floratos and G. Linardopoulos. "Large-Spin and Large-Winding Expansions of Giant Magnons and Single Spikes". In: *Nucl.Phys.* **B897** (2015), p. 229. arXiv: 1406.0796 [hep-th],in collaboration with E. Floratos and G. Georgiou:
 - [5]. E. Floratos, G. Georgiou, and G. Linardopoulos. "Large-Spin Expansions of GKP Strings". In: *JHEP* **03** (2014), p. 018. arXiv: 1311.5800 [hep-th],and single-authored work:
 - [6]. G. Linardopoulos. "Classical Strings and Membranes in the AdS/CFT Correspondence". <http://www.didaktorika.gr/eadd/handle/10442/35838?locale=en>. PhD thesis. National and Kapodistrian University of Athens, 2015.
 - [7]. G. Linardopoulos. "Large-Spin Expansions of Giant Magnons". In: *PoS (CORFU2014)*, p. 154. arXiv: 1502.01630 [hep-th].

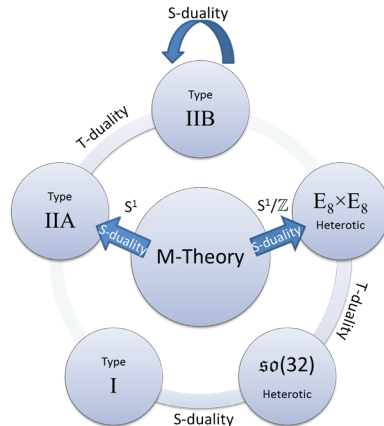
M-Theory and Relativistic Membranes

- Work on the **BMN Matrix Model** in collaboration with M. Axenides and E. Floratos (the present talk):
 - [8]. M. Axenides, E. Floratos, and G. Linardopoulos. "Multipole stability of spinning M2 branes in the classical limit of the BMN matrix model". In: *Phys. Rev.* **D97** (2018), p. 126019. arXiv: 1712.06544 [hep-th].
 - [9]. M. Axenides, E. Floratos, and G. Linardopoulos. "M2-brane Dynamics in the Classical Limit of the BMN Matrix Model". In: *Phys. Lett.* **B773** (2017), p. 265. arXiv: 1707.02878 [hep-th].
- Work on **Relativistic Membranes** with M. Axenides and E. Floratos (see talks on 2012 & 2013 annual meetings of INPP):
 - [10]. M. Axenides, E. Floratos, and G. Linardopoulos. "Stringy Membranes in AdS/CFT". In: *JHEP* **08** (2013), p. 089. arXiv: 1306.0220 [hep-th].
- Work on **Minimal Hypersurfaces** with J. Hoppe and T. Turgut:
 - [11]. J. Hoppe, G. Linardopoulos, and T.O. Turgut. "New Minimal Hypersurfaces in $\mathbb{R}^{(k+1)(2k+1)}$ and S^{2k^2+3k} ". In: *Mathematische Nachrichten* (2016). arXiv: 1602.09101 [math.DG].

What is M-theory?

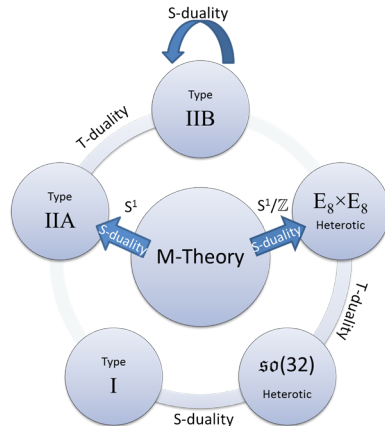
Why do we need M-theory?

- By the end of the first superstring revolution (1984-1994), five different 10-dimensional superstring theories had emerged:
Types I, II (IIA, IIB), Heterotic ($\mathfrak{so}(32)$, $E_8 \times E_8$).



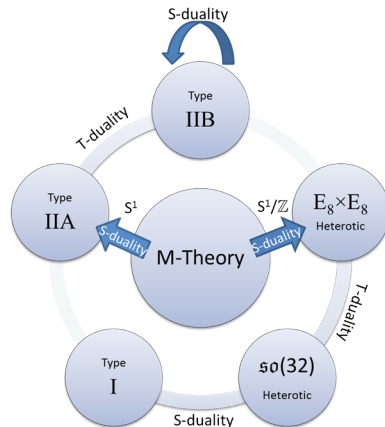
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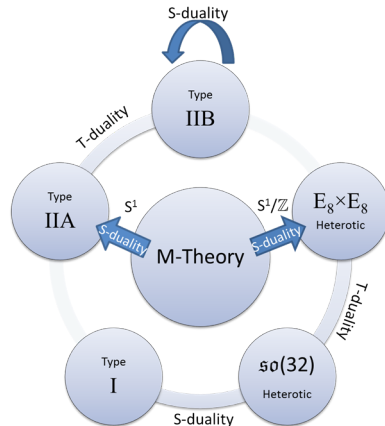
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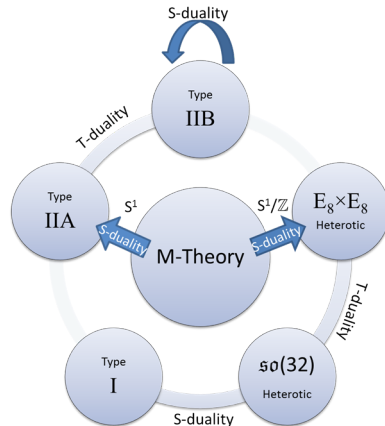
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- The letter "M" stands for "magic, mystery and matrix" according to one of its founders, Edward Witten...

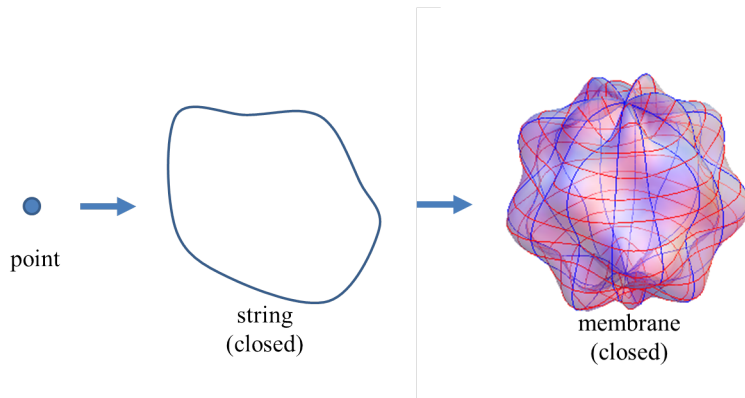


Membranes

- The idea behind membrane theory is simple: replace 1-dimensional strings with 2-dimensional membranes...

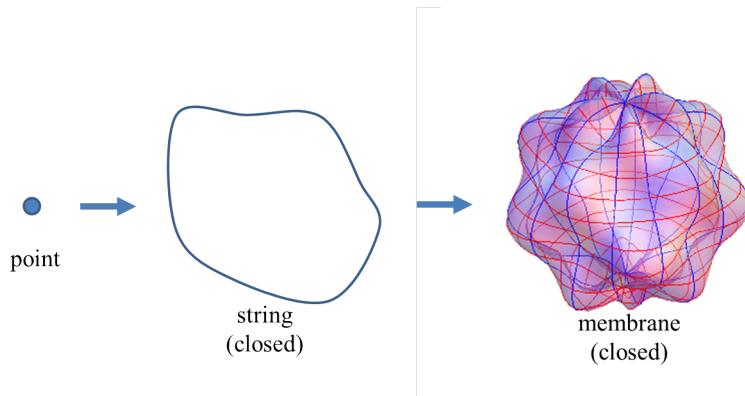
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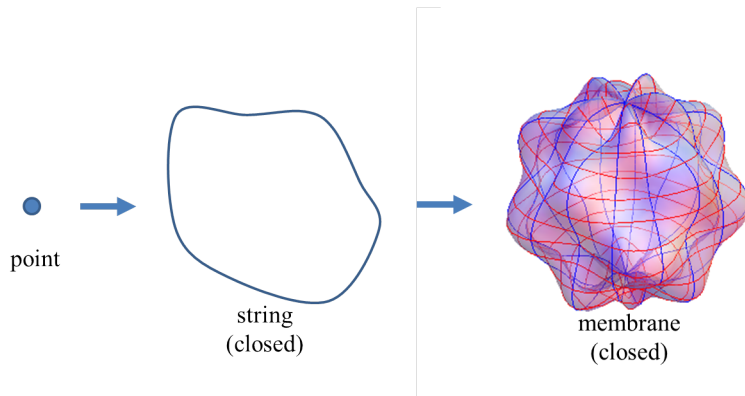
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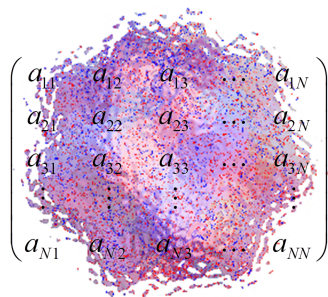
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- Like strings, membranes can be supersymmetrized... we thus obtain supermembranes...
- There are reasons to believe that membranes (or "M2-branes") are the fundamental objects of 11-dimensional M-theory, just like strings are the fundamental objects of 10-dimensional string theory...

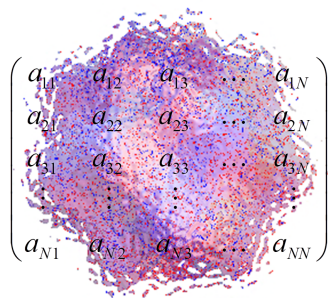
M(atrix) theory

- According to the matrix theory conjecture of Banks, Fischler, Shenker and Susskind (BFSS, 1996), a theory of matrix-quantized supermembranes provides a realization of M-theory in flat space-time...


$$\begin{pmatrix} a_{11} & a_{12} & a_{13} & \cdots & a_{1N} \\ a_{21} & a_{22} & a_{23} & \cdots & a_{2N} \\ a_{31} & a_{32} & a_{33} & \cdots & a_{3N} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ a_{N1} & a_{N2} & a_{N3} & \cdots & a_{NN} \end{pmatrix}$$

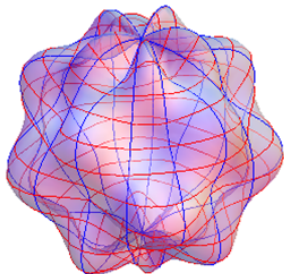
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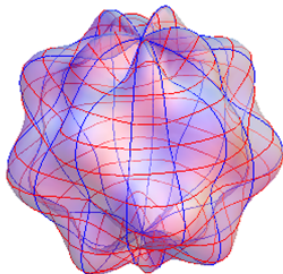
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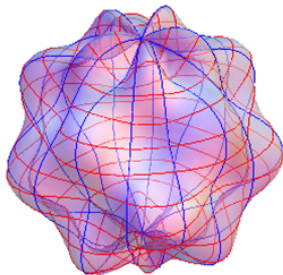
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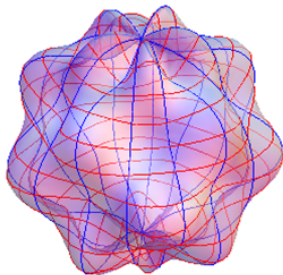
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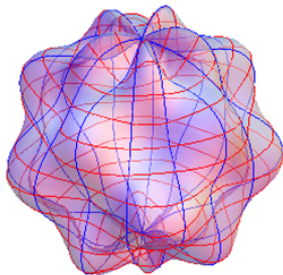
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- Very recently, M(atrrix) theory has been applied to the study of chaotic phenomena that take place on the horizons of black holes...



The black hole membrane paradigm

- Black holes (BHs) are regions of spacetime where the force of gravity is so strong that nothing (not even light) can escape...



Gargantua black hole - *Interstellar* movie.

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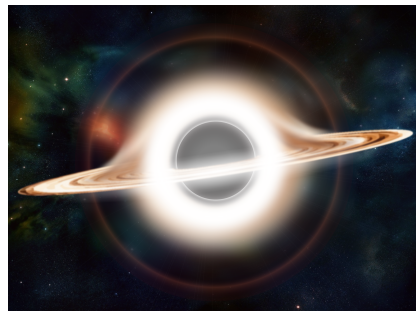
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- One such mechanism is *fast scrambling* or ultra-fast thermalization... In general it is believed that chaotic phenomena are a dominant feature of BH horizons...



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- Because it is inherently nonlocal, M(atrix) theory turns out to be a valuable tool in the study of BH information processing...



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Membranes as dynamical systems

Membranes in plane-wave backgrounds

General setup

- Consider the 11-dimensional maximally supersymmetric plane-wave background:

$$ds^2 = -2dx^+ dx^- - \left[\frac{\mu^2}{9} \sum_{i=1}^3 x^i x^i + \frac{\mu^2}{36} \sum_{j=1}^6 y^j y^j \right] dx^+ dx^+ + dx^i dx^i + dy^j dy^j, \quad F_{123+} = \mu.$$

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- The Hamiltonian of a bosonic relativistic membrane in the above background reads

$$H = \frac{T}{2} \int d^2\sigma \left[p_x^2 + p_y^2 + \frac{1}{2} \{x^i, x^j\}^2 + \frac{1}{2} \{y^i, y^j\}^2 + \{x^i, y^j\}^2 + \frac{\mu^2 x^2}{9} + \frac{\mu^2 y^2}{36} - \frac{\mu}{3} \epsilon_{ijk} \{x^i, x^j\} x^k \right],$$

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- The corresponding equations of motion and the Gauß-law constraint are given by:

$$\ddot{x}_i = \{ \{x_i, x_j\}, x_j \} + \{ \{x_i, y_j\}, y_j \} - \frac{\mu^2}{9} x_i + \frac{\mu}{2} \epsilon_{ijk} \{x_j, x_k\}, \quad \sum_{i=1}^3 \{ \dot{x}^i, x^i \} + \sum_{j=1}^6 \{ \dot{y}^j, y^j \} = 0$$

$$\ddot{y}_i = \{ \{y_i, y_j\}, y_j \} + \{ \{y_i, x_j\}, x_j \} - \frac{\mu^2}{36} y_i.$$

The ansatz

- The following $\mathfrak{so}(3)$ -invariant ansatz automatically satisfies the Gauß-law constraint:

$$x_i = \tilde{u}_i(\tau) e_i, \quad y_j = \tilde{v}_j(\tau) e_j, \quad y_{j+3} = \tilde{w}_j(\tau) e_j, \quad i, j = 1, 2, 3.$$

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The coordinates e_i are given by

$$(e_1, e_2, e_3) \equiv (\cos \phi \sin \theta, \sin \phi \sin \theta, \cos \theta), \quad \phi \in [0, 2\pi), \quad \theta \in [0, \pi],$$

satisfy the $\mathfrak{so}(3)$ Poisson algebra and are orthonormal:

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- As we will see, the above ansatz leads to an interesting dynamical system with stable and unstable solutions that describe rotating and pulsating membranes of spherical topology.
- Similar work in flat space has previously been carried out by (Axenides-Floratos, 2007).

The Hamiltonian system

- Here's the Hamiltonian of the membrane:

$$\begin{aligned}
 H = \frac{2\pi T}{3} & \left[\tilde{p}_u^2 + \tilde{p}_v^2 + \tilde{p}_w^2 + \tilde{u}_1^2 \tilde{u}_2^2 + \tilde{u}_2^2 \tilde{u}_3^2 + \tilde{u}_3^2 \tilde{u}_1^2 + \tilde{r}_1^2 \tilde{r}_2^2 + \tilde{r}_2^2 \tilde{r}_3^2 + \tilde{r}_3^2 \tilde{r}_1^2 + \tilde{u}_1^2 (\tilde{r}_2^2 + \tilde{r}_3^2) + \tilde{u}_2^2 (\tilde{r}_3^2 + \tilde{r}_1^2) + \right. \\
 & \left. + \tilde{u}_3^2 (\tilde{r}_1^2 + \tilde{r}_2^2) + \frac{\mu^2}{9} (\tilde{u}_1^2 + \tilde{u}_2^2 + \tilde{u}_3^2) + \frac{\mu^2}{36} (\tilde{r}_1^2 + \tilde{r}_2^2 + \tilde{r}_3^2) - 2\mu \tilde{u}_1 \tilde{u}_2 \tilde{u}_3 \right], \quad \tilde{r}_j^2 \equiv \tilde{v}_j^2 + \tilde{w}_j^2, \quad j = 1, 2, 3.
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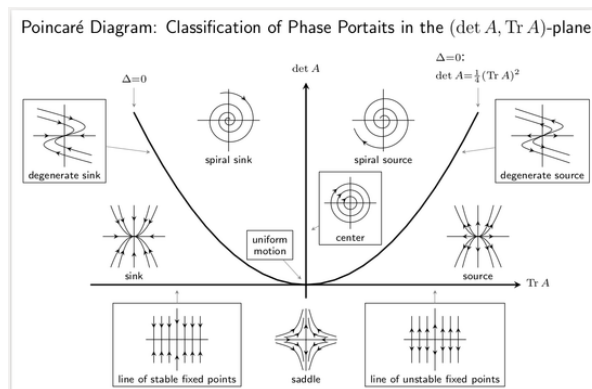
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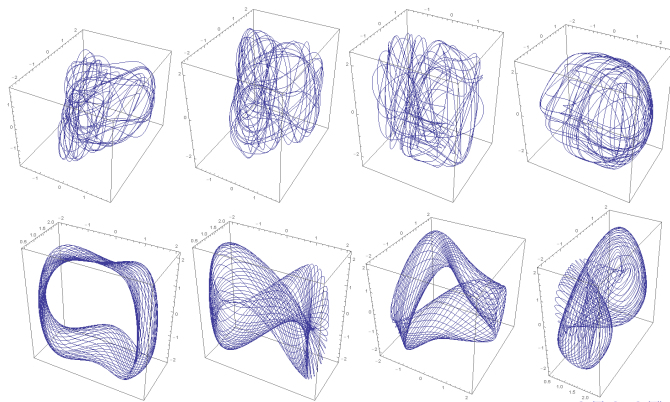
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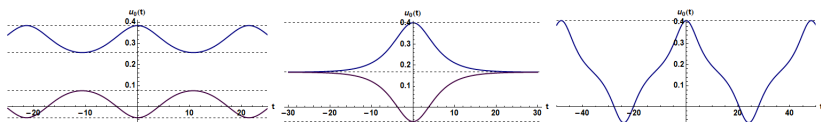
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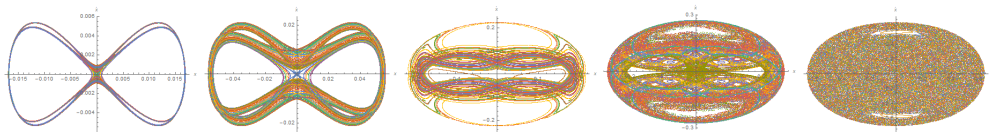


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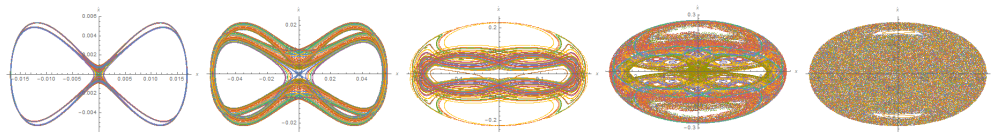
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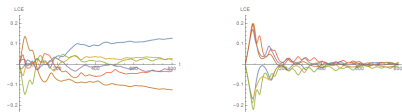


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What is its spectrum of Lyapunov characteristic exponents?



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