Status of stellar and astrophysics constraints on new (feebly-interacting) physics

Maurizio Giannotti, Barry University

> FIP 2020, August 31- September 4, 2020

Probing the low energy frontier with stars.

PHYSICAL REVIEW

Study e.m.

the sun

properties of

neutrinos from

VOLUME 132, NUMBER 3

1 NOVEMBER 1963

Electromagnetic Properties of the Neutrino

JEREMY BERNSTEIN* AND MALVIN RUDERMAN[†] Department of Physics, New York University, New York, New York

AND

GERALD FEINBERG[‡] Department of Physics, Columbia University, New York, New York (Received 11 June 1963)

In this note we make a detailed survey of the experimental information on the neutrino charge, charge radius, and magnetic moment. Both weak-interaction data and astrophysical results can be used to give precise limits to these quantities, independent of the supposition that the weak interactions are charge conserving.

For decades stars have been successfully used to probe the physics of neutrinos, axions, dark photons, and other Feebly Interacting Particles.

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For decades stars have been successfully used to probe the physics of neutrinos, axions, dark photons, and other Feebly Interacting Particles.

Why are stars such good FIP labs?



Feeble is good!

Light particles, $m \leq T$, can be efficiently thermally produced in stellar core





Feeble is good!

Example: the SM process

 $e^+ e^- \rightarrow \bar{\nu} \nu$

has a branching ratio 10^{-19} times smaller than the process



$$e^+ e^- \rightarrow \gamma \gamma$$

D. D. Clayton, "Principles of Stellar Evolution and Nucleosynthesis", Chicago (1984)

Yet, it controls the evolution of a massive star during the late evolutionary stages.

Another very rare process, $\gamma \rightarrow \overline{\nu}\nu$, controls the stellar evolution during the RGB.

Stars as Laboratory for FIPs



 $\leftarrow T_{\rm eff}$

Stars as Laboratory for FIPs



Figures from G. Raffelt, Stars as Laboratories (1996).

8



Most of the stellar life is spent burning H into He in the core







After the H in the core is exhausted, a light star star moves in the RGB. The surface luminosity keeps increasing, till the Heflash. That is the tip of the RGB



G. Raffelt, Stars as Laboratories (1996).



G. Raffelt, Stars as Laboratories (1996).

reported in the past few years.

... after the He-flash, He is ignited and the star moves to the HB.





The number ratio of HB and RGB is the R-parameter



$$R = \frac{N_{\rm HB}}{N_{\rm RGB}}$$

The number ratio of HB and RGB is the R-parameter



$$R = \frac{N_{\rm HB}}{N_{\rm RGB}}$$

The current analysis shows a slight discrepancy between the predicted and observed R-parameter.



Ayala et al., PRL 113 (2014)

Lighter stars go to the AGB at the end of the central He, and end up as CO White Dwarfs.





Stellar Evolution: WD

WD Variables (WDV)

Measures of the period change rate in WD variables offer a way to test the cooling of WDs

 $\dot{P}/P \propto \dot{T}/T$

Star	<i>P</i> (s)	$\dot{P}_{\rm obs}({\rm s/s})$	$\dot{P}_{\rm th}({\rm s}/{\rm s})$
G117 - B15A	215	$(4.2\pm0.7) imes10^{-15}$	$(1.25 \pm 0.09) imes 10^{-15}$
R548	213	$(3.3 \pm 1.1) imes 10^{-15}$	$(1.1 \pm 0.09) imes 10^{-15}$
PG 1351+489	489	$(2.0\pm0.9) imes10^{-13}$	$(0.81\pm0.5) imes10^{-13}$
L 19-2 (113)	113	$(3.0\pm0.6) imes10^{-15}$	$(1.42\pm0.85) imes10^{-15}$
L 19-2 (192)	192	$(3.0\pm0.6) imes10^{-15}$	$(2.41 \pm 1.45) imes 10^{-15}$

L. Di Luzio, M.G., E. Nardi, L. Visinelli, Phys.Rept. 870 (2020)

Observations over the past ~30 yr showed consistently $\dot{P}_{\rm obs} > \dot{P}_{\rm th}$, which seems to imply an overly efficient cooling.

Many works starting form Isern, Hernanz, Garcia-Berro (1992)

More massive stars exhibit <u>blue loops</u> during the He-burning phase.



G. Raffelt, Stars as Laboratories (1996).

More massive stars exhibit <u>blue loops</u> during the He-burning phase.

log L/L⊙

Current observations show:

1) a small <u>red-shift</u> of the bluest point of the blue loop in the high luminosity region of the CMD and

2) too many blue stars (B/R problem).

Additional cooling in the Heburinig stage could alleviate this problem

- R. C. Dohm-Palmer, E. D. Skillman (2002)
- K. B. W. McQuinn et al. (2011)
- Friedland, M.G., Wise (2013)
- Carosi et al. (2013)



Stellar Cooling Anomalies: New Physics?

Overly efficient energy loss in stars led to new physics speculation • Long standing WD hint. *Isern, Hernanz, Garcia-Berro (1992), Córsico, Althaus, Miller Bertolami, Kepler (2019), and many other works*

• Less solid indication of anomalous cooling in RGB O. Straniero et al., to appear.

• Very recent claim of a possible hint form Red Clumps Stars Mori, Kusakabe, Balantekin, Kajino, Famiano (2020)

• Indication from HB stars (R-parameter) Ayala, Dominguez, M.G., Mirizzi, Straniero (2014); Straniero et al. (2015)

• Massive stars (blue loop) A. Friedland, M.G., M. Wise (2013), Carosi et al. (2014)

ALPs offer the best explanation

M.G., Irastorza, Redondo, Ringwald (2016)

Looking for FIPs in Stars



During the early evolutionary phases, which are very relevant to the study of FIPs,

$$T_c \sim (1 - 10) \,\mathrm{keV}$$

and

 $\rho_c \sim (10^2 - 10^6) \,\mathrm{g \, cm^{-3}}$

Other stages, when neutrino cooling dominates, are very fast

Axions and Axion-Like Particles (ALPs)

Axions are among the most well motivated FIPs:

Theoretically motivated by the strong CP problem and as DM candidates.

$$L_{int} = -i\sum g_{ai} a \bar{\psi}_i \gamma_5 \psi_i - \frac{1}{4}g_{a\gamma} aF\tilde{F}$$

<u>OCD axions</u> parameter space restricted from relations among couplings and mass.

<u>ALPs</u>: no restrictions in the parameter space



Di Luzio, M.G., Nardi, Visinelli, Phys.Rept. 870 (2020)



Axion-photon coupling ($g_{a\gamma}$)



Most relevant process in stars: Primakoff





Axion-photon coupling $(g_{a\gamma})$





Axion-electron coupling (g_{ae}) : Hints from stars?

Long standing hint from WDV.

WDLF also shows excess
energy loss, interpretable as:
ALPs good fit;

• HP good fit

• μ_{ν} very bad fit



M.G., Irastorza, Redondo, Ringwald, Saikawa, (2017)

					Preliminary
Less solid		photometric	distance	<i>g</i> ₁₃	<i>g</i> ₁₃
evidence for a		sample	scale	best valu	e bound
hint from RGBT	→ M5	VI+JK	ZAHB	0	2.30
	47 Tuc	VI+ JK	parallax	0.45	1.87
<u>No hint</u> from	NGC 362	VI+JK	parallax	0	1.37
Capozzi, Ranen (2020) [arXiv:2007.03694]	22 GCs	JK	ZAHB	0.60	1.38
	16 GCs	JK	KINEMATIC	0.35	1.05

Straniero et al., to appear

Still insufficient experimental sensitivity for direct detection of solar ALPs



 $g_{ae} \times 10^{13}$

Xe1T,

LUX,

Xenon 1T did not observe solar ALPs...



 $g_{ae} \times 10^{13}$

Xe1T hint

Xenon 1T did not observe solar ALPs...

... However, it might have observed DM ALPs, with $g_{ae} \sim (0.5 - 0.7) \times 10^{-13}$ (RGB hint), mass a few keV, and weakly coupled to photons

Takahashi, Yamada, Yin (2020) [arXiv:2006.10035]



 $g_{ae} \times 10^{13}$

Stellar hints on axion couplings ($g_{ae}, g_{a\gamma}$)

A global analysis of RGB, WDLF, WDV and R-parameter gives a preference to some energy loss unaccounted in the SM and explainable by axions coupled to photons and electrons.

- M.G., Irastorza, Redondo, Ringwald, Saikawa, (2017)
- Di Luzio, Fedele, M.G., Mescia, Nardi (in preparation)



Interlude: Neutrino magnetic moment (μ_{ν})



Interlude: Neutrino magnetic moment (μ_{ν})

A $\mu_{\nu} \neq 0$ would help the neutrino detection (ν scattering on electrons and nuclei)



$$L = -\frac{1}{4}F_{\mu\nu}^2 - \frac{1}{4}X_{\mu\nu}^2 - \frac{\epsilon}{2}F_{\mu\nu}F^{\mu\nu} + \frac{m_X^2}{2}X_{\mu}X^{\mu} + eJ_{\rm em}^{\mu}A_{\mu}$$

M. Fabbrichesi, E. Gabrielli, G. Lanfranchi (2020) [arXiv:2005.01515]

Stellar bounds

An, Pospelov, Pradler, Phys.Lett.B 725 (2013); An, Pospelov, Pradler, Ritz, Phys.Lett.B 747 (2015)

Sun
$$\rightarrow \omega_{\rm pl}(r=0) \sim 300 \,\mathrm{eV}$$

HB $\rightarrow \omega_{\rm pl}(r=0) \sim 2.6 \,\mathrm{keV}$
RGB $\rightarrow \omega_{\rm pl}(r=0) \sim 25 \,\mathrm{keV}$

$$L = -\frac{1}{4}F_{\mu\nu}^2 - \frac{1}{4}X_{\mu\nu}^2 - \frac{\epsilon}{2}F_{\mu\nu}F^{\mu\nu} + \frac{m_X^2}{2}X_{\mu}X^{\mu} + eJ_{\rm em}^{\mu}A_{\mu}$$

Solar bound Vinyoles et al., JCAP 10 (2015)

At lower masses the bound weakens. Solar L-mode dominates:

 $\epsilon \cdot m_X \le 1.8 \times 10^{-12} \,\mathrm{eV}, \ m_X \lesssim 1 \,\mathrm{eV}$



$$L = -\frac{1}{4}F_{\mu\nu}^2 - \frac{1}{4}X_{\mu\nu}^2 - \frac{\epsilon}{2}F_{\mu\nu}F^{\mu\nu} + \frac{m_X^2}{2}X_{\mu}X^{\mu} + eJ_{\rm em}^{\mu}A_{\mu}$$

HP can fit well the WDLF

Stellar hints

However, conflict btw WDV and other stellar bounds



M.G., I. Irastorza, J. Redondo, A. Ringwald, JCAP 1605 (2016)

$$L = -\frac{1}{4}F_{\mu\nu}^2 - \frac{1}{4}X_{\mu\nu}^2 - \frac{\epsilon}{2}F_{\mu\nu}F^{\mu\nu} + \frac{m_X^2}{2}X_{\mu}X^{\mu} + eJ_{\rm em}^{\mu}A_{\mu}$$

Xenon 1T did not see solar HP. Their spectrum is too soft...

An, Pospelov, Pradler, Ritz (2020) [arXiv:2006.13929]

$$L = -\frac{1}{4}F_{\mu\nu}^2 - \frac{1}{4}X_{\mu\nu}^2 - \frac{\epsilon}{2}F_{\mu\nu}F^{\mu\nu} + \frac{m_X^2}{2}X_{\mu}X^{\mu} + eJ_{\rm em}^{\mu}A_{\mu}$$



m_X [keV]

$$L = -\frac{1}{4}F_{\mu\nu}^2 - \frac{1}{4}X_{\mu\nu}^2 - \frac{\epsilon}{2}F_{\mu\nu}F^{\mu\nu} + \frac{m_X^2}{2}X_{\mu}X^{\mu} + eJ_{\rm em}^{\mu}A_{\mu}$$

Xenon 1T did not see solar HP. Or did it?

non-minimal scenarios may suppress production in HB and RGB stars but not in the sun. Environmental effects are induced through dynamical kinetic mixing, emerging through interaction with a scalar field.

Chakraborty, Jung, Loladze, Okui, Tobioka (2020) [arXiv:2008.10610]



Exploring higher masses: Supernovae

Supernovae have an internal temperature of ~30 MeV and density of ~ 10^{14} g cm⁻³.

In this conditions more massive FIPs can be created.

However, a reliable description of the FIP production in SN is difficult

Several recent improvements and revisitations

- C. Hanhart, D. R. Phillips, S. Reddy (2001)
- Chang, Essig, McDermott (2018);
- Chang, Essig, McDermott (2019);
- P. Carenza et al. (2019);
- Ertas and Kahlhoefer (2020)
- G. Lucente et al. (2020)
- ...

Supernova 1987A

SN1987A: SNe cannot cool too fast (ν -signal). Roughly, $L_a \lesssim L_{\nu}$

Very old bounds however:

- Emission rate is hard to calculate
- Very few data

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- Ertas and Kahlhoefer (2020)
- G. Lucente et al. (2020)

• ...

Axion-like particles (g_{aN})



- Carenza, Fischer, M.G., Guo, Martinez-Pinedo, Mirizzi (2019);
- Di Luzio, M.G., Nardi, Visinelli Phys.Rept. 870 (2020)

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- Ertas and Kahlhoefer (2020)
- G. Lucente et al. (2020)

• ...

Axion-like particles $(g_{a\gamma})$



- Jaeckel, Malta, Redondo (2017);
- Ertas and Kahlhoefer (2020);
- G. Lucente et al. (2020);

Supernova 1987A

SN1987A: SNe cannot cool too fast (ν -signal). Roughly, $L_a \lesssim L_{\nu}$

Very old bounds however:

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- ...

Dark Photons



Chang, Essig, McDermott (2019);

Stars may produce FIPs copiously.

- Solar ALPs and HP are searched by terrestrial experiments.
- SNe can produce enormous quantities of FIPs ($\sim 10^{52}$ erg/s).
- Very strong limits from SN 1987A [Payez et al. (2015), De Rocco et al. (2020)]
- and from diffuse gamma ray from all past SNe [Calore et al. (2020), De Rocco et al. (2020)]



- Calore, Carenza, M.G., Jaeckel, Mirizzi (2020)
- DeRocco, Graham, Kasen, Marques-Tavares, Rajendran (2020)

Surprisingly strong bounds on ALPs from other stars too, particularly form supergiants. ALPs oscillate into X-ray photons. Bounds from NuSTAR

- *M. Xiao et al. Nustar bounds from Betelgeuse, (in preparation);*
- Dessert, Foster, Safdy, X-ray Searches for Axions from Super Star Clusters (2020)

A few NS, observed by XMM- Newton and Chandra, exhibit an unexplainable excess. Is it due to $a \rightarrow \gamma$?

Buschmann, Co, Dessert, Safdy, X-ray Search for Axions from Nearby Isolated Neutron Stars (2019)



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- Solar ALPs and HP are searched by terrestrial experiments.
- SNe can produce enormous quantities of FIPs ($\sim 10^{52}$ erg/s).

Direct detection of SN ALPs? Helioscopes as SN-scopes



Ge, Hamaguchi, Ichimura, Ishidoshiro, Kanazawa (2020) [arXiv:2008.03924]

Keeps coherence up to higher masses because of much higher energy

Stellar basins for FIPs?

FIPs trapped in stellar gravitational field. Does it matter?

Perhaps so. A FIP basin would build up during the stellar lifetime, assuming the depletion time is slow enough





Stellar basins for FIPs?

FIPs trapped in stellar gravitational field. Does it matter?

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10-10 Solar Dark Photon Basin Limit 10-1 Solar cooling 10-12 XENON1T (Solar) kinetic mixing ϵ 10-13 10-14 $\tau = 10^7 v$ basin excess $(\tau = 2.3 \times 10^9 \text{y})$ 10⁻¹⁵ $= 2.3 \times 10^9 \text{v}$ HB cooling RG cooling $--- \tau = 4.6 \times 10^9 \text{y}$ $\tau = 4.6 \times 10^9 \text{y}$, future 10-16 10⁻¹ 10⁰ 10¹ 10^{2} 10³ 10⁴ 10⁵ dark photon mass m [eV]

Ken Van Tilburg (2020) [arXiv:2006.12431]

Did Xenon 1T observe Dark Photons in solar basin?

Robert Lasenby, Ken Van Tilburg (2020) [arXiv:2008.08594]

The Roaring Twenties

An exciting decade for FIPs astrophysics!

Great improvements in stellar census, photometry, astrometry,...

GAIA, since 2014, Data Release 3 expected soon.

JWST, launch scheduled for Oct. 31, 2021

Vera Rubin Telescope (LSST), beginning of operation schedule for 2021.

In addition: XRISM, soft X-ray 0.4-13 keV, early 2020s Perhaps, less improvements in the 0(10-100) MeV γ -ray detectors, although possibly better resolution. Gamma 400, starts in 2026. Other proposals: eAstrogam, Amego (not yet approved) GAIA DR3: https://www.cosmos.esa.int/ web/gaia/earlydr3

https://www.jwst.nasa.gov

LSST DM group (Alex Drlica-Wagner et al.) (2019) [arXiv:1902.01055]

Fantin, Cote, McConnachie (2020) [arXiv:2007.01312]

https://heasarc.gsfc.nasa.gov/ docs/xrism/

Conclusions

Stars are (excellent) FIPs labs + FIPs factories.

Stellar anomalies have stimulated research in stellar evolution and FIPs models.

Several new bounds, proposals and ideas in stellar FIPs this year. In part motivated by the Xenon 1T anomaly.

A lot of progress expected in the near future. Stay tuned!

Backup Slides

SNe and other monsters

Supernovae, Neutron Stars, and even Black Holes offers unique ways to study FIPs.

Application of these monster stars to the physics of FIPs is relatively recent.

Several recent revisitations for ALPs and HP.

- C. Hanhart, D. R. Phillips, S. Reddy (2001)
- Chang, Essig, McDermott (2018);
- Chang, Essig, McDermott (2019);
- P. Carenza et al. (2019);
- G. Lucente et al. (2020)
- ...

BH superradiance

Tests coupling to gravity. No assumption that the boson is initially present, i.e. there is no requirement for the boson to be the DM.

A. Arvanitaki, S. Dubovsky, PhysRevD.83.044026 (2011);

A. Arvanitaki, M. Baryakhtar, X. Huang, PhysRevD.91.084011 (2015)

V. Cardoso et al. JCAP 1803 (03) (2018)





Figures form V. Cardoso et al. JCAP 1803 (03) (2018)

SN Factories

Calore, Carenza, M.G., Jaeckel, Mirizzi (2020)

Marques-Tavares, Rajendran (2020)

- Diffuse gamma spectrum from ALPs and HP
- Bounds from gamma ray from SN 1987A
- Excess of galactic positrons from HP

Payez, et al., JCAP 02 (2015)

DeRocco, Graham, Kasen,

DeRocco, Graham, Kasen, Marques-Tavares, Rajendran (2020)



SN-scope

Table 1: List of SN progenitor candidates with having a mass $\gtrsim 10 M_{\odot}$ and within 250 pc from the Earth. We basically use the values listed in the Hipparcos catalogue [78]; otherwise, we show the reference for the source.

HIP	Common Name	Distance (pc)	Mass (M_{\odot})	RA (J2000)	Dec (J2000)
65474	Spica/ α Virginis	77(4)	11.43 ± 1.15 [79]	13:25:11.58	-11:09:40.8
81377	ζ Ophiuchi	112(3)	20.0 [80]	16:37:09.54	-10:34:01.5
71860	lpha Lupi	142(3)	10.1 ± 1.0 [81]	14:41:55.76	$-47{:}23{:}17.5$
80763	Antares/ α Scorpii	170(30)	11 - 14.3 [82]	16:29:24.46	-26:25:55.2
107315	Enif/ ϵ Pegasi	211(8)	11.7(8) [81]	21:44:11.16	+09:52:30.0
27989	Betelgeuse/ α Orionis	$222^{+48}_{-34} \ [83]$	$11.6^{+5.0}_{-3.9} \ [84]$	05:55:10.31	$+07{:}24{:}25.4$



Figure 1: The position of the SN progenitors in Table 1 on the Mollweide projection of the celestial sphere, where the red and blue dots correspond to the spectral types of K/M and O/B, respectively. We also show by the gray dots the progenitors with d > 250 pc and $M \gtrsim 10 M_{\odot}$ listed in Table A1 in Ref. [87].

From S. F. Ge et al., IAXO as SN-scope (2020) [arXiv:2008.03924]

Scalars

Bounds on scalar particles produced in stellar plasma have been recently revised



10⁴

*m*_φ/eV

10⁵

10⁶

RG (resonant)

1000

1.×10⁻³⁴

100

Surprisingly strong bounds on ALPs from other stars too, particularly form supergiants. ALPs oscillate into X-ray photons. Bounds from NuSTAR

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White Dwarfs in the Era of the LSST and its Synergies with Space-Based Missions

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Submitted to ApJ

Can stellar bounds be evaded?

see, e.g., De rocco et al. (2020) [arXiv:2006.15112]. Chakraborty, et al (2020) [arXiv:2008.10610]

If so, it may open the way to the interpretation of more astrophysical anomalies in terms of FIPs.

see, e.g., Pallathadka et al. (2020) [arXiv:2008.08100].

Mori, Balantekin, Kajino, Famiano (2020)

	photometric	distance	<i>g</i> ₁₃	<i>8</i> 13
	sample	scale	best value	bound
M5	VI+JK	ZAHB	0	2.30
47 Tuc	VI+ JK	parallax	0.45	1.87
NGC 362	VI+JK	parallax	0	1.37
22 GCs	JK	ZAHB	0.60	1.38
16 GCs	JK	KINEMATIC	0.35	1.05

Capozzi Raffelt: $\mu_{12} < 0.77\,(1.50)$ at $68\,\%\,(95\%)~{\rm CL}$

Straniero et al. (2020)

$$g_{ae} = 0.60^{+0.33}_{-0.53}$$

Axion-photon coupling $(g_{a\gamma})$

