



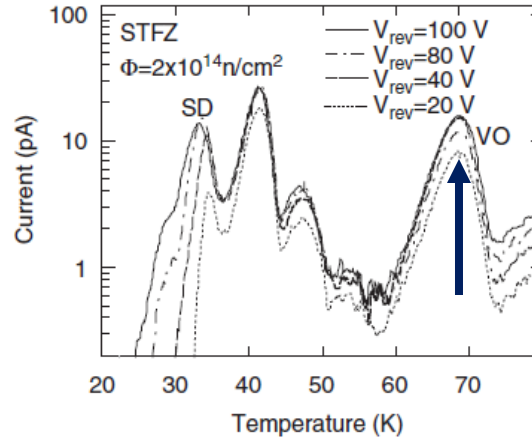
Optimization of the priming procedure for thermally stimulated currents with heavily irradiated silicon detectors

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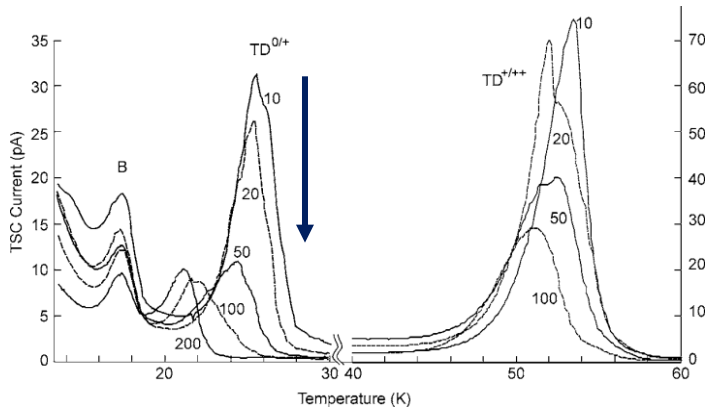
2010

Normally the peak height increases with increasing bias voltage

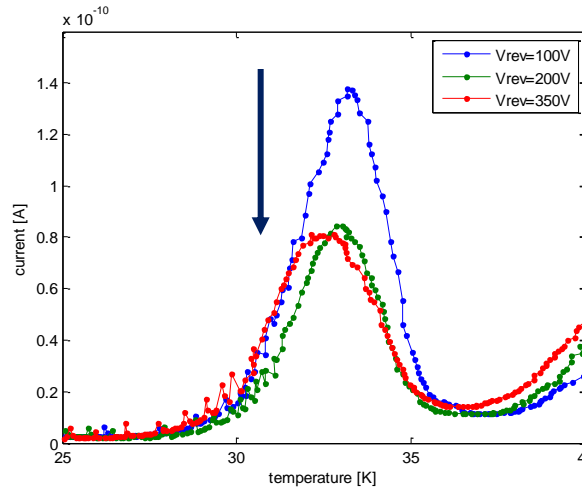


Scaringella et al. *Nucl. Instrum Meth A* **570** (2007) 322–329

In some case it decreases



Bruzzi et al. *J. Appl. Phys.* **99** (2006) 093706



TSC peak saturation does not correspond to full volume emission



Materials and methods



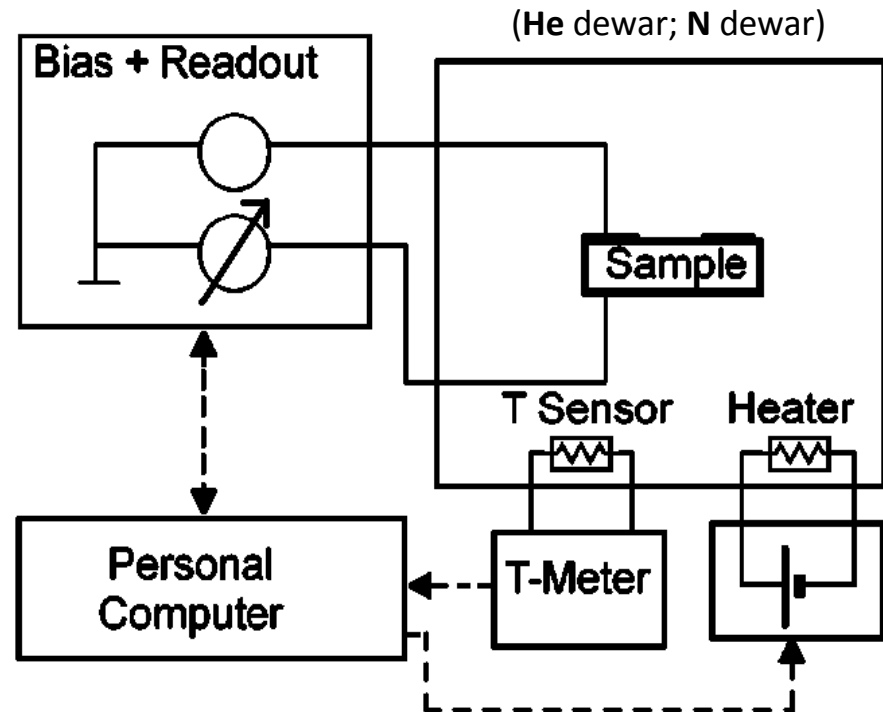
• Samples

- Material: n-type MCz Si produced by Okmetic (Finland) with 900 Ωcm resistivity, $\langle 100 \rangle$ orientation and 280 μm thickness.
- Devices: p-on-n planar diodes 0.5x0.5 cm^2
- Procurement: WODEAN, Thanks to E. Fretwurst, G. Lindstroem
- Irradiation: reactor neutrons at the Jozef Stefan Institute, Ljubljana
- Fluence: 10^{13} - 10^{14} - 10^{15} $n_{\text{eq}}/\text{cm}^2$
- Annealing: 1 year RT

- Material: p-type MCz Si produced by Okmetic (Finland) with 2k Ωcm resistivity, $\langle 100 \rangle$ orientation and 280 μm thickness
- devices: n-on-p square diodes 0.5x0.5 cm^2
- Procurement: SMART, Thanks to D. Creanza, N. Pacifico
- Irradiation: reactor neutrons at the Jozef Stefan Institute, Ljubljana
- Fluence: 10^{14} - 10^{16} $1/\text{cm}^2$
- Annealing: few days RT (+ 80min 80° C)

• Set-up

- He dewar: $T=[5K, RT]$
- N dewar: $T=[90K, RT]$





Materials and methods



- 1. Thermally Stimulated Currents with cryogenic equipments: measurements performed on liquid He vapors, to ensure stable temperatures down to 4.2K minimize thermal inertia and mismatch.**
- 2. Zero Bias Thermally Stimulated Currents measured at high fluence in the high temperature range (100-200K), to avoid background current subtraction, and increase resolution in deep levels analysis. Used in all T range to study residual electric field and charge of defects.**

Issues:

- (a) Optimize the priming procedure to evidence radiation induced defects and correctly evaluate their concentration;**
- (b) Get information about electric field distribution;**
- (c) Correlate defects with transport properties.**



Materials and methods



• TSC

– Priming:

- Forward current priming:

- Cooling from T_{p0} down to T_{fill} under polarization V_{cool}
- Filling at T_{fill} with polarization V_{fill} : I_{fill}

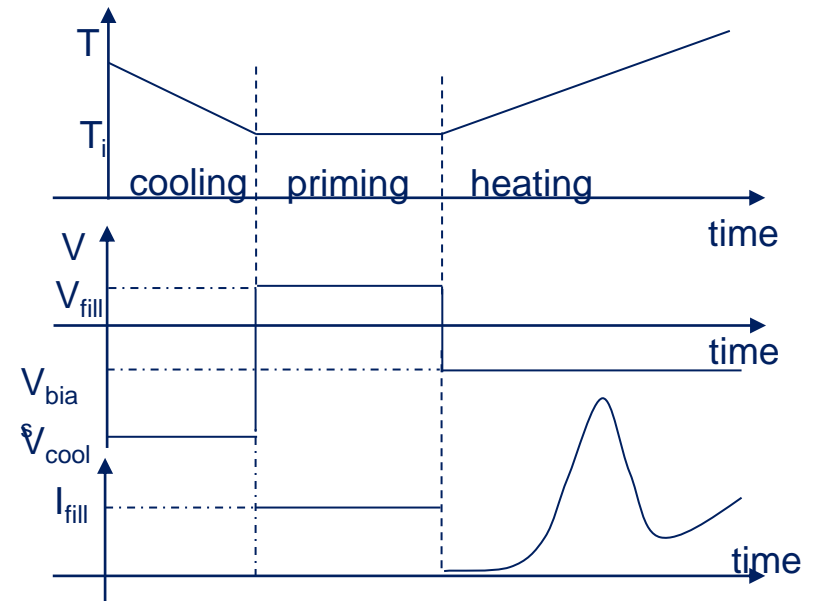
- Optical priming:

- Cooling down to T_{fill}
- Illumination at T_{fill} with IR LED ($\lambda=850\text{nm}$, full penetration) with polarization V_{fill}

– Stimulation

- Thermal scan

- Apply bias voltage V_{bias} (TSC)
- or
- Apply $0V$ (ZBTSC)
- Heating to T_{max}





Experimental results: varying priming



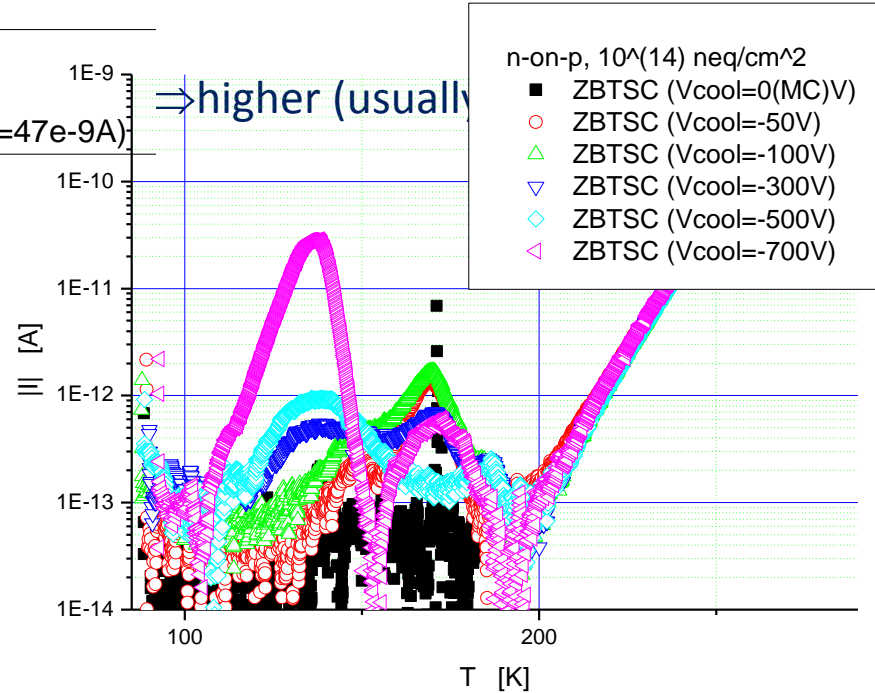
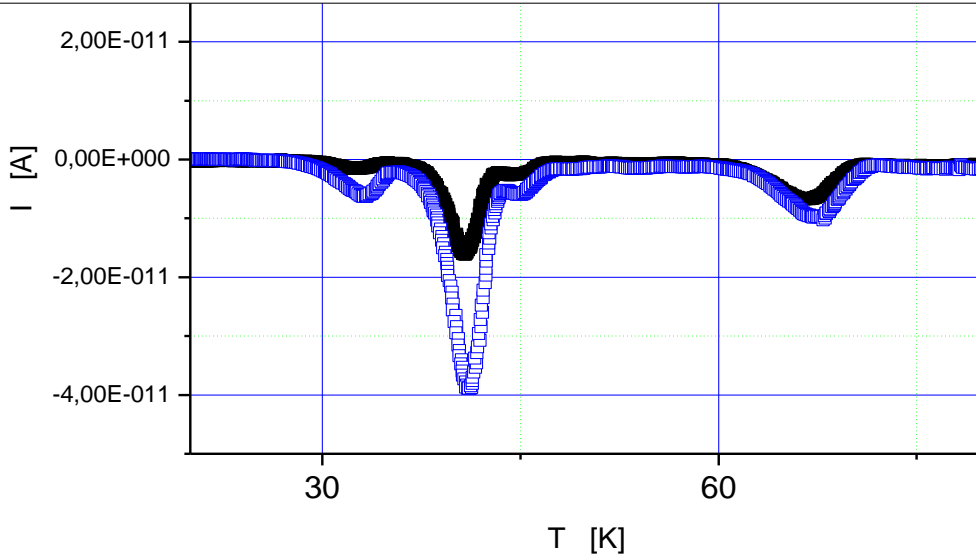
– Forward current priming: varying I_{fill} , V_{cool}

1 Lower filling current \Rightarrow lower TSC

n-on-p, 10^{15} neq/cm² annealed 80min, 80°C

■ $V_{bias}=-100V$ ($V_{cool}=0V$, $T_{p0}=220K$, $V_{fill}=700V$: $I_{fill}=30e-9A$)

□ $V_{bias}=-100V$ ($V_{cool}=-100V$, $T_{p0}=220K$, $V_{fill}=[700,690]V$: $I_{fill}=47e-9A$)



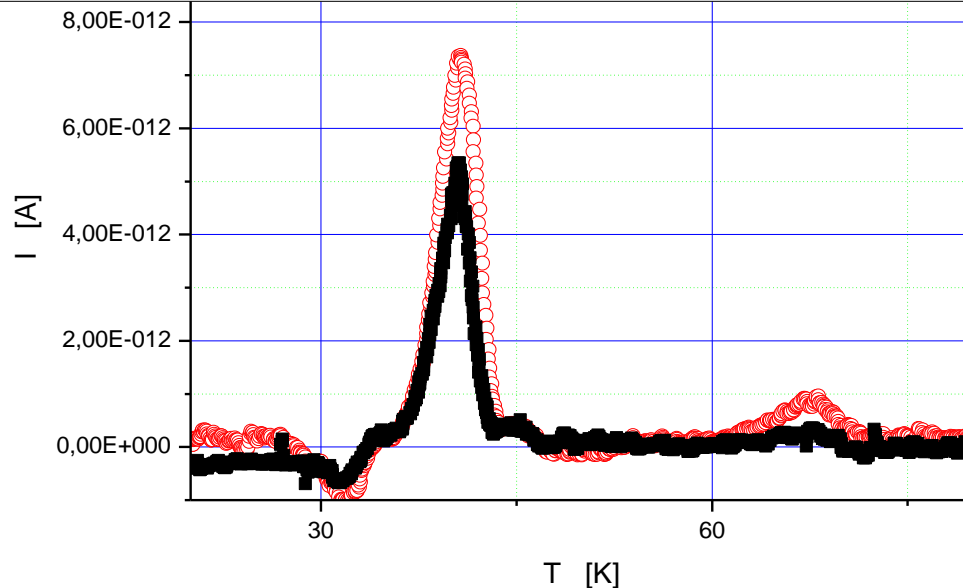


Experimental results: varying priming



– Forward current priming: varying I_{fill}
(independent of V_{cool})

– Filling (n-on-p, 10^{15} neq/cm²) dual field)
○ ZBTSC ($T_{p0}=20K$, $V_{cool}=0V$ (MC) ($T_{amb},0V$ (MC)), $V_{fill}=[+700,+680]V$: $I_{fill}=52e-9A$)
■ ZBTSC ($T_{p0}=20K$, $V_{cool}=-100V$ ($T_{amb},-100V$), $V_{fill}=[+700,+680]V$: $I_{fill}=24e-9A$)

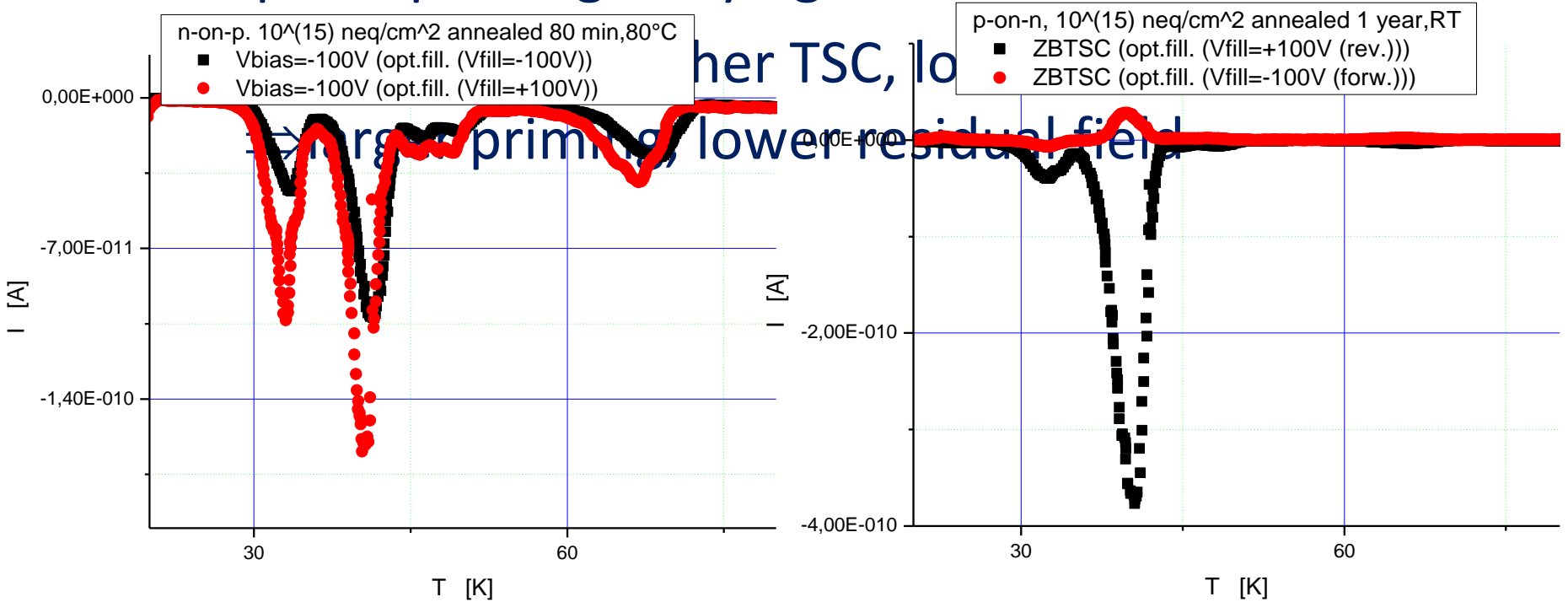




Experimental results: varying priming



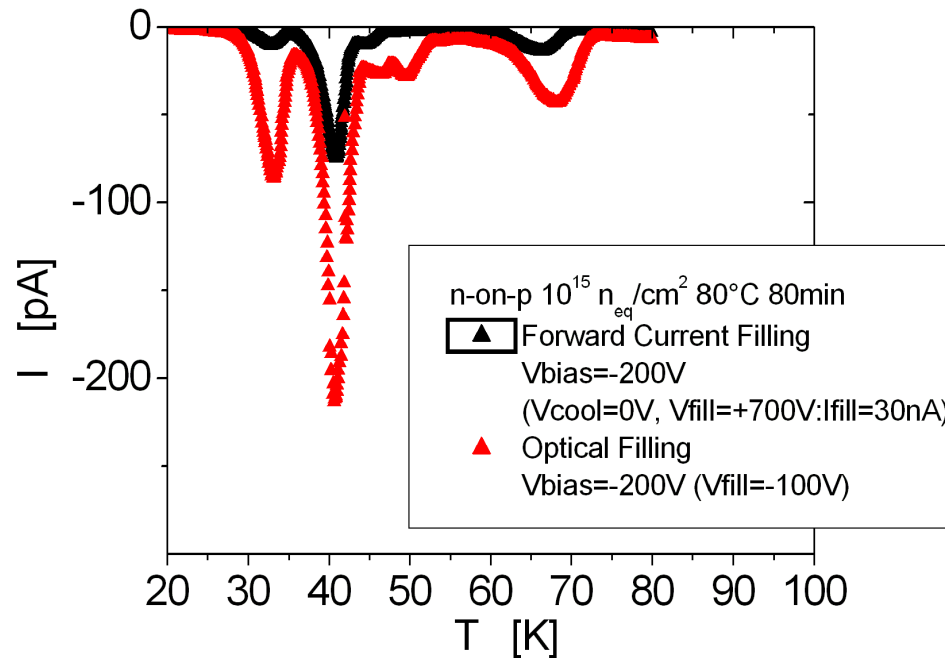
– Optical priming: varying Vfill





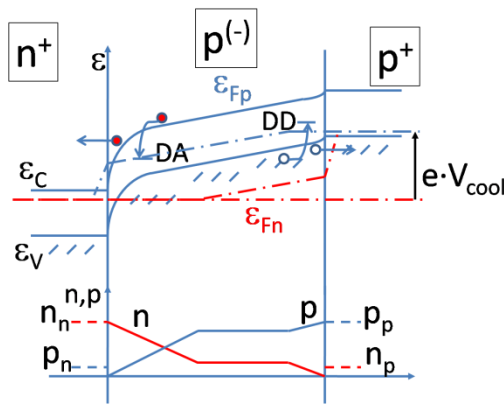
Experimental results: varying priming

- Varying procedure: forward current- vs. optical-priming
- Higher ϵ

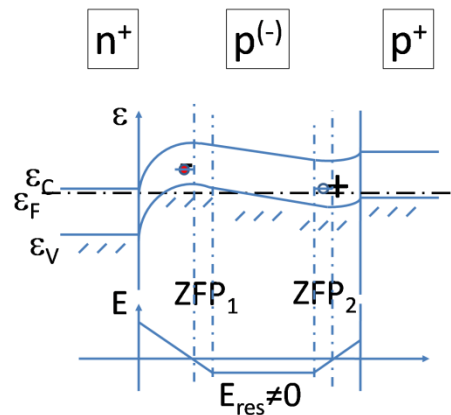


- Priming mechanisms
 - Forward current priming

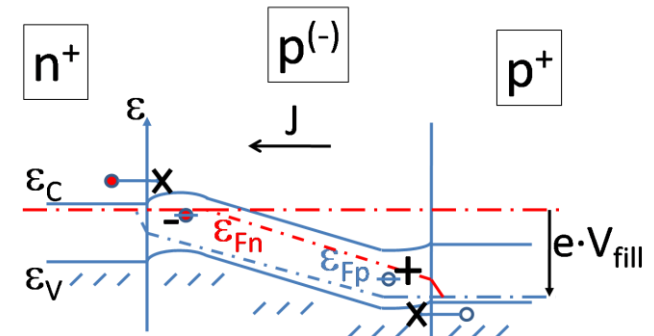
Reverse cooling



Metastable condition



Forward filling



- Space charge enlargement during cooling* $\frac{dJ}{dx} = -\frac{d\rho}{dt}$
- Asymmetrical filling leading to DJ**

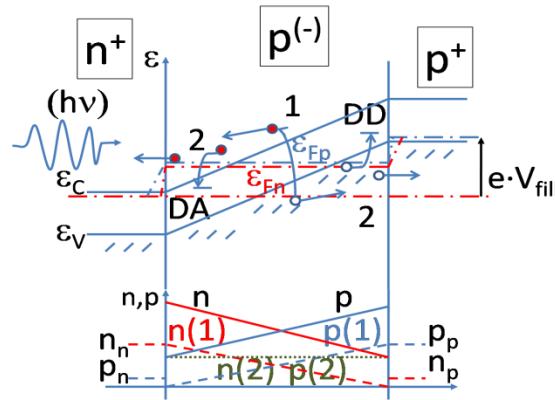
*Simmons and Taylor, JAP, 1972

**Eremin et al., NIMA, 2002

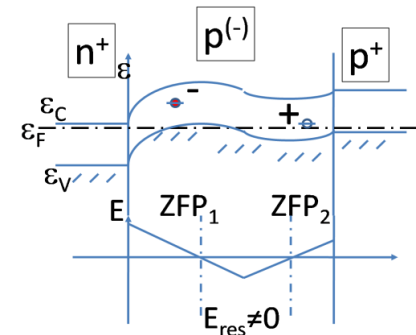
– Priming mechanisms

- Optical priming

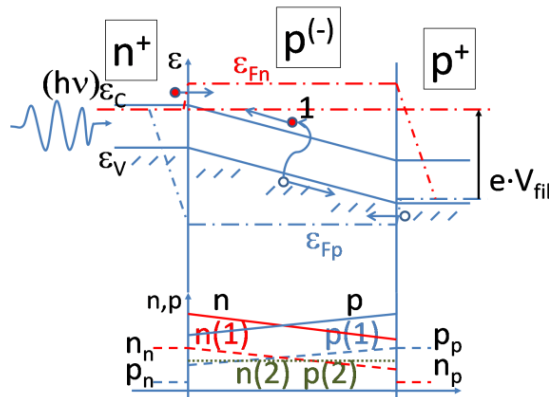
Reverse filling



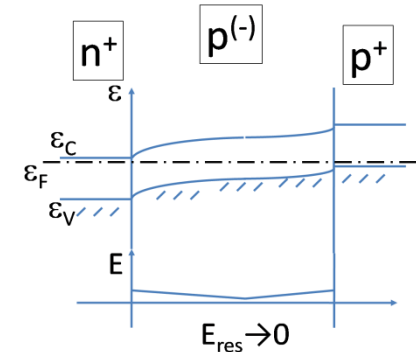
Metastable condition



Forward filling

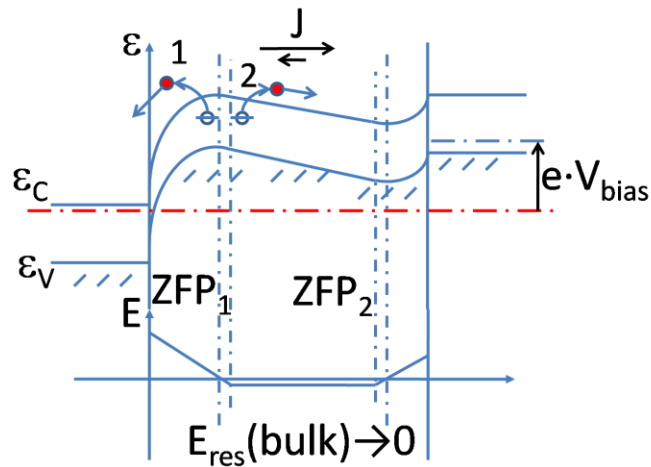


Metastable condition

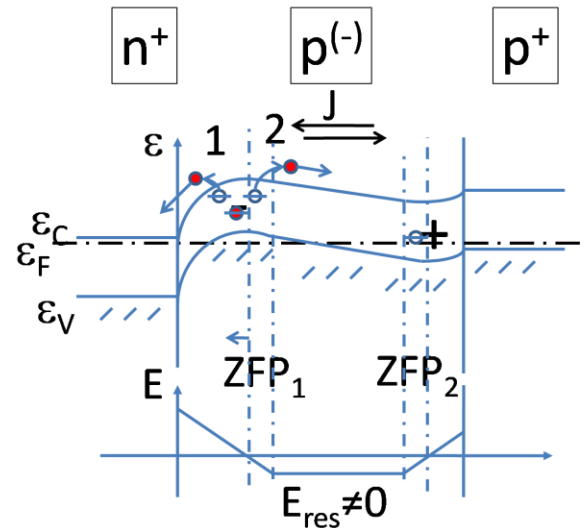


– Stimulation

- TSC n^+ $p^{(-)}$ p^+



- ZBTSC



– On the priming procedure:

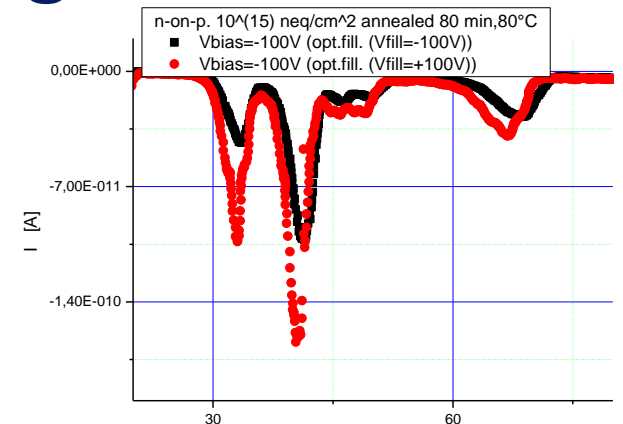
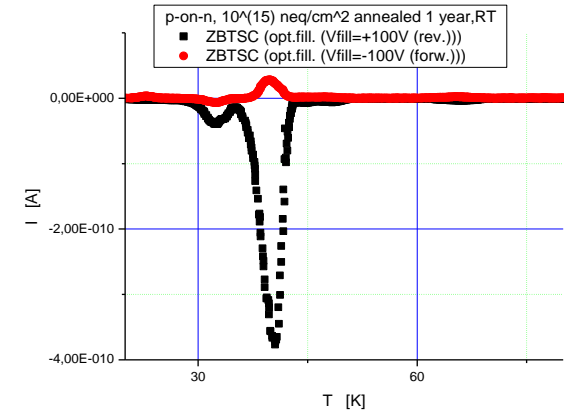
Negligible ZBTSC means:

- a) Null priming
- b) Uniform priming

If TSC after same priming is not negligible:

– There is priming

and it is uniform!!!

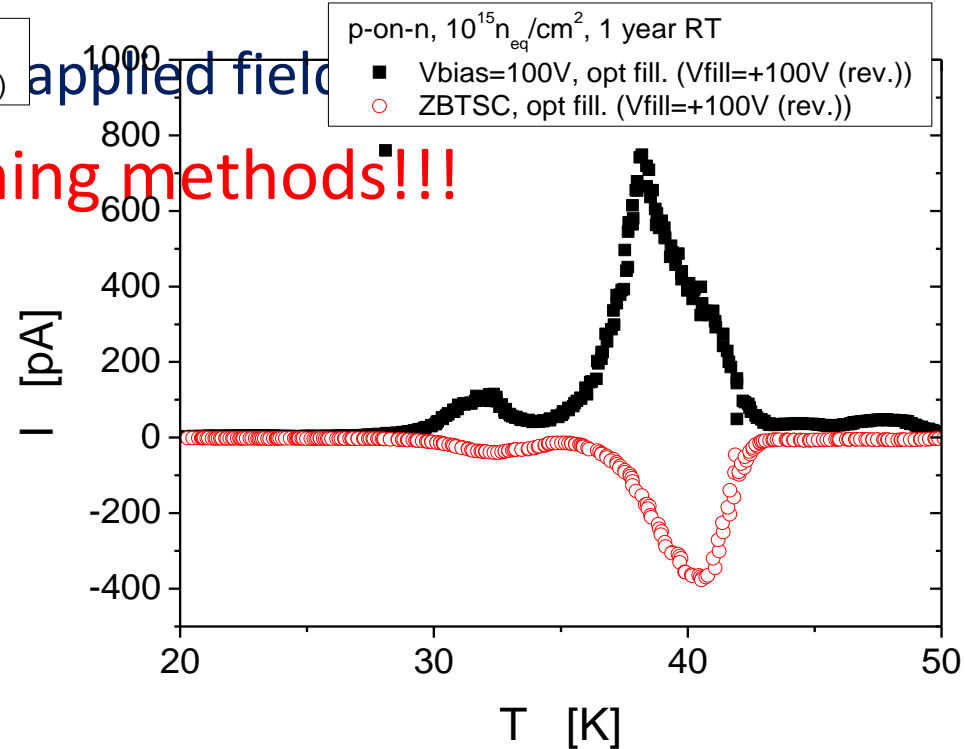
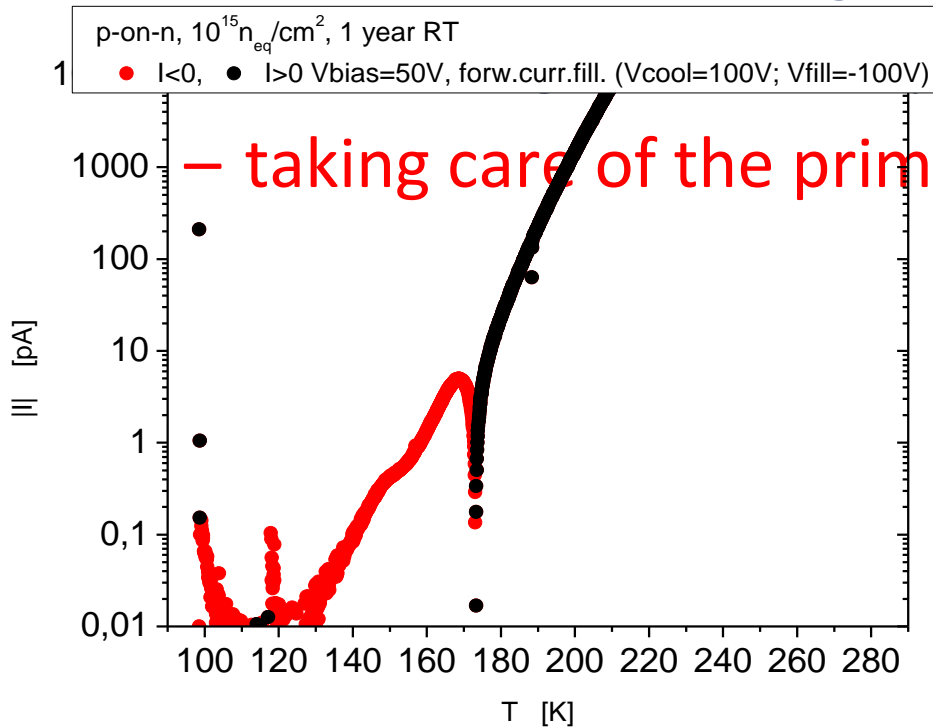




Experimental results: particular phenomena from priming conditions



– Residual field strength:

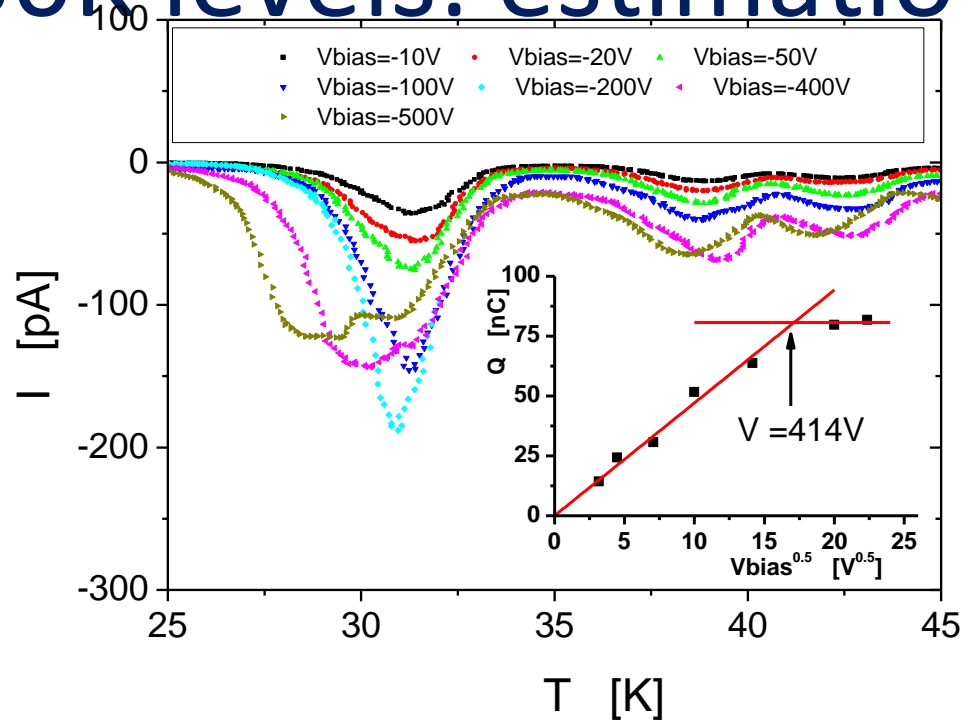




Experimental results: defect studies



-SD30K levels: estimations



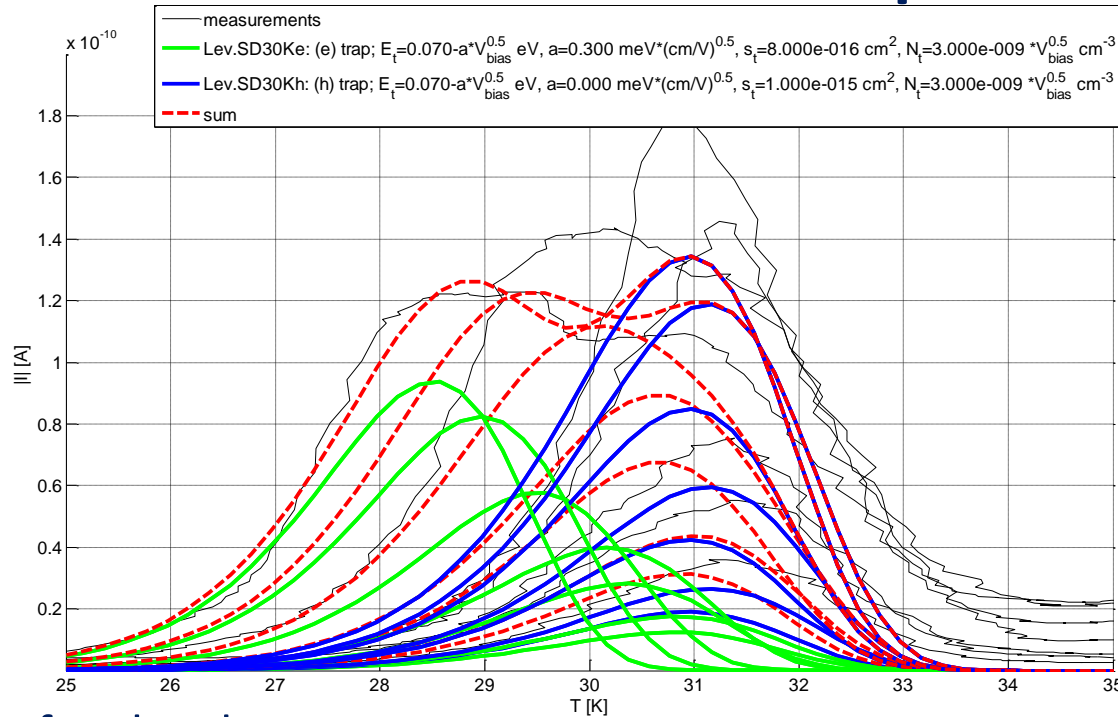
We studied the 30K shallow defect in irradiated MCZ Si using the correct priming procedure. Now emitted charge increases with square root of bias up to saturation for $V \sim V_{dep}$.



Experimental results: defect studies



-SD30K: two main components



- Defect levels:

One (green) shows Poole-Frenkel effect thus it is charged at RT, contributes to Neff

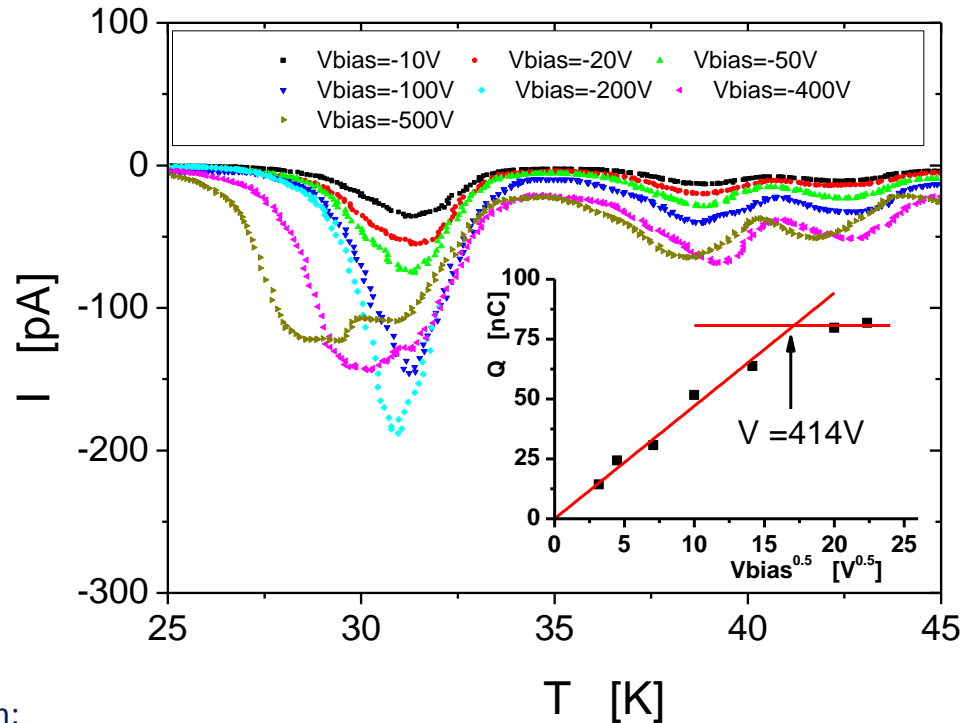
The other (blue) shows no P-F, does not contribute to the effective doping



Experimental results: defect studies



– SD30K levels: estimations



- N_t estimation:
 - From charge at saturation: [SD30K levels] $\cong Q_{sat}/Vol. \cong 2.388 \cdot 10^{14} cm^{-3}$ →reliable
 - From fits: [SD30K charged] $\sim 2 \cdot 10^{13} cm^{-3}$; [SD30K not charged] $\sim 3 \cdot 10^{13} cm^{-3}$ →confirm
- V_{dep} estimation, from Q_{sat} : $V_{dep} \sim 414V$ ($V_{dep}^{(nominal)} = 378V$) →reliable
- V_{dep} evaluation from [SD30K charged] $\cong N_{eff}$: $V_{dep} \sim 5kV!!!$ →?!? → shallow level compensates other negatively charged defects usually responsible of type inversion!!!



Experimental results: TSC and ZBTSC



A comparison between TSC and ZBTSC allows one to evaluate the kind of carrier released by the trap:

- Electron emission:
 - negative SCR (p-type) \Rightarrow field reduction \Rightarrow ZBTSC peak delayed during T scan with respect to TSC peak
 - Positive SCR (n-type) \Rightarrow field increase \Rightarrow peak anticipated
- Hole emission:
 - negative SCR (p-type) \Rightarrow field reduction \Rightarrow peak delayed
 - Positive SCR (n-type) \Rightarrow field increase \Rightarrow peak anticipated

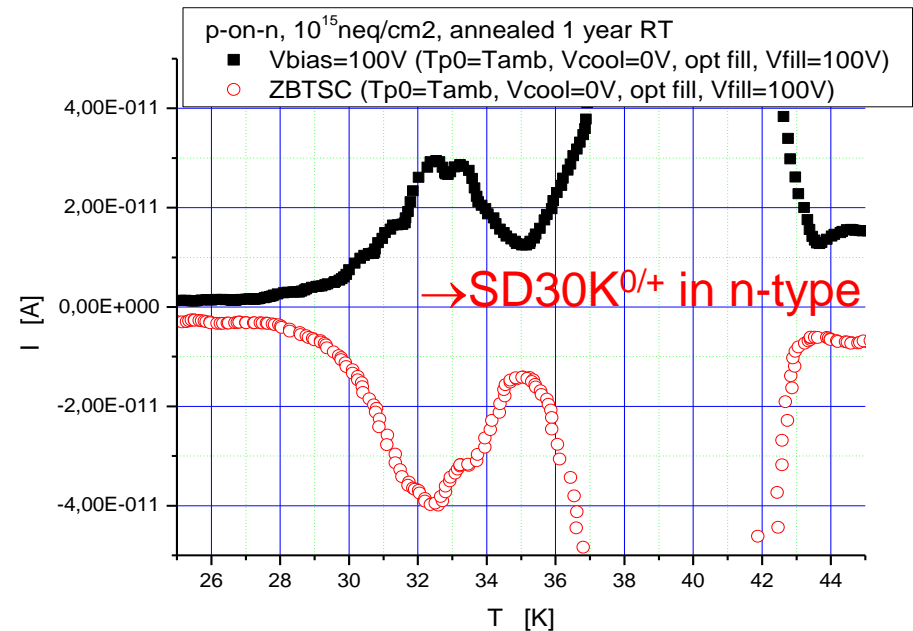
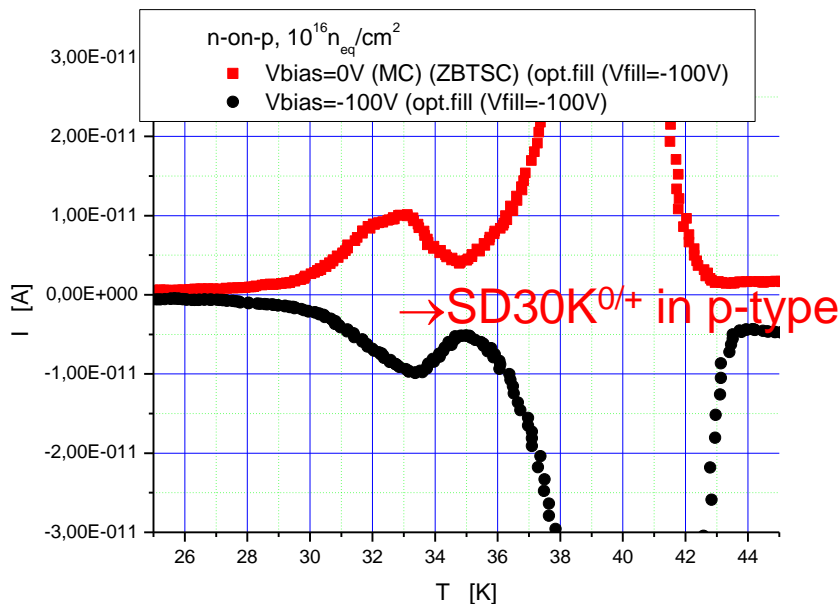


Experimental results: TSC and ZBTSC



For the shallow level at 30K we find electron emission:

- negative SCR (p-type) \Rightarrow field reduction \Rightarrow peak delayed
- Positive SCR (n-type) \Rightarrow field increase \Rightarrow peak anticipated





Work in progress

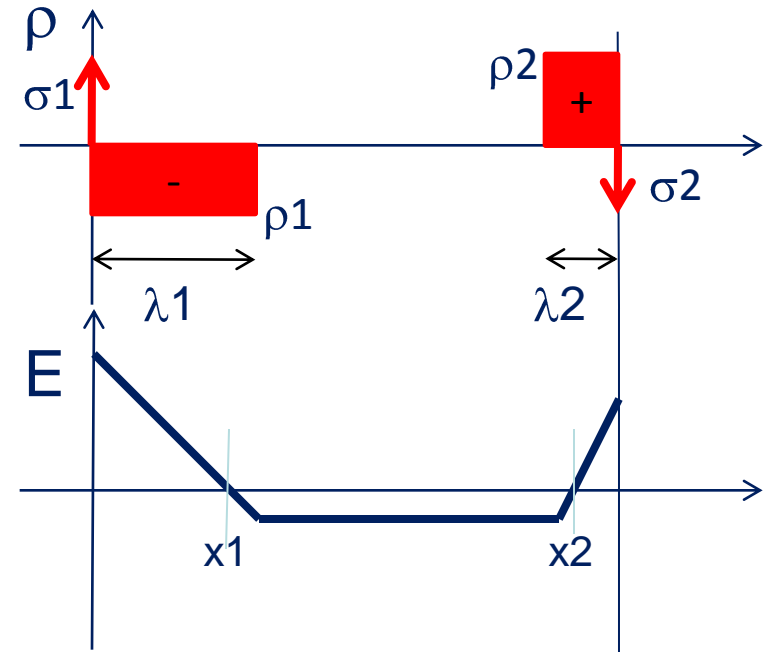


We are working on a model taking into account of fundamental relations:

- Poisson equation
- Continuity equation
- Current definition
- Boundary conditions:
 - » $E(0^-)=0$
 - » $E(d^+)=0$
 - » $q(\phi(d)-\phi(0))=\varepsilon_g$

⇒ Relation between differentials

- Zero-field plane definition



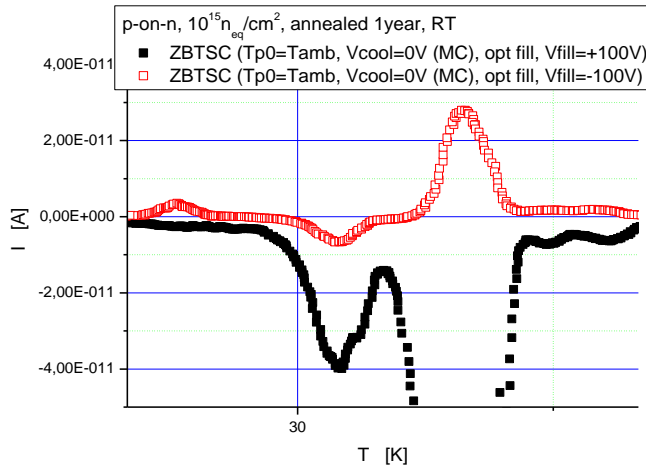
$$d\sigma_1 - (\lambda_1 d\rho_1 + \rho_1 d\lambda_1) + (\lambda_2 d\rho_2 + \rho_2 d\lambda_2) - d\sigma_2 = 0$$

$$\sigma_1 - (\lambda_1 d\rho_1 + \rho_1 d\lambda_1) - \left(\frac{\lambda_1^2}{2d} d\rho_1 + \frac{\rho_1 \lambda_1}{d} d\lambda_1 \right) - \left(\frac{\lambda_2^2}{2d} d\rho_2 + \frac{\rho_2 \lambda_2}{d} d\lambda_2 \right) = 0$$

$$dx_1 = \left(d\lambda_1 - \frac{\lambda_1}{d} d\lambda_1 \right) - \left(\lambda_2 - \frac{\lambda_2^2}{2d} \right) \left(\frac{\rho_1 d\rho_2 - \rho_2 d\rho_1}{\rho_1^2} \right) - \left(\frac{\rho_2}{\rho_1} \right) \left(d\lambda_1 - \frac{\lambda_1}{d} d\lambda_1 \right)$$

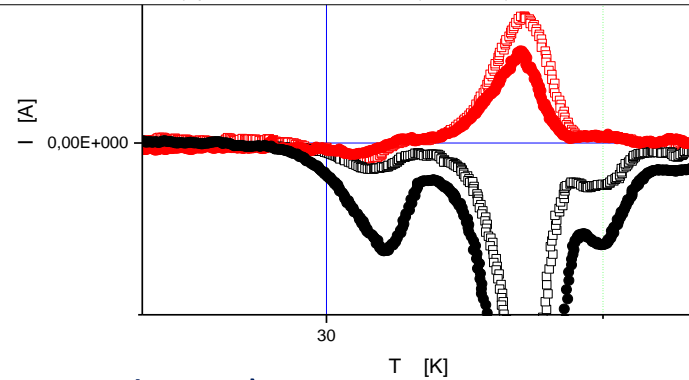
$$\Rightarrow J \approx \rho_1 \frac{dx_1}{dt}$$

1. Measurements suggest:



n-on-p, $10^{15} n_{eq}/cm^2$, annealed 80min, $80^\circ C$

- ZBTSC ($T_{p0}=20K$, $V_{cool}=0V$ (MC) ($T_{amb}, 0V$ (MC)), $V_{fill}=[+700,+680]V$: $I_{fill}=52e-9A$)
- $V_{bias}=-100V$ ($T_{p0}=220K$, $V_{cool}=0V$ ($T_{amb}, 0V$), $V_{fill}=+700V$: $I_{fill}=30e-9A$)
- ZBTSC ($T_{p0}=20K$, $V_{cool}=-100V$ ($T_{amb}, -100V$), $V_{fill}=[+700,+680]V$: $I_{fill}=24e-9A$)
- $V_{bias}=-100V$ ($T_{p0}=220K$, $V_{cool}=-100V$ ($T_{amb}, 0V$), $V_{fill}=[+700,+690]V$: $I_{fill}=47e-9A$)



- Negative (independently on priming and type)

2. Model infer:

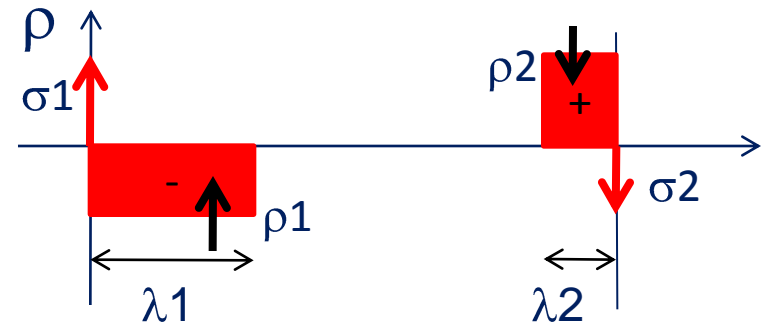
- Hp.s:-SD30K emission, recombination in the bulk

$$\begin{cases} d\rho_2 < 0 \\ d\rho_1 < 0 \end{cases}$$

$$J < 0 \Rightarrow \rho_1 \lambda_1 > \rho_2 \lambda_2$$

- And relation with trap type!

⇒ WORK IN PROGRESS





Conclusions

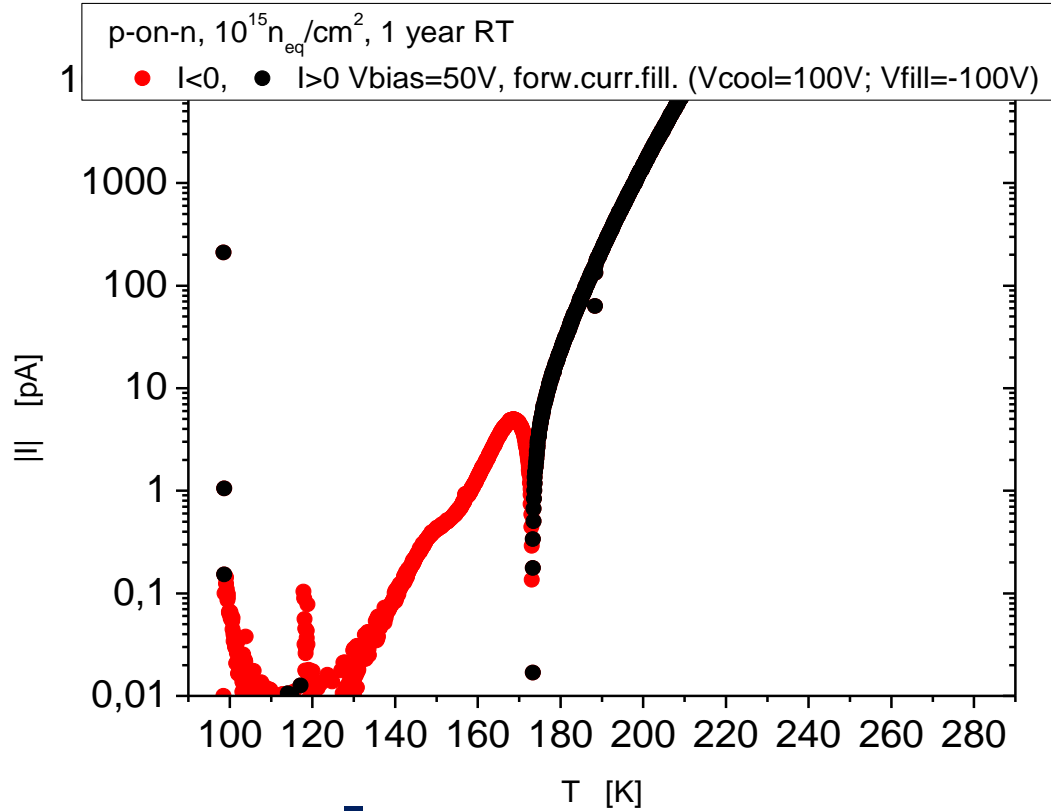


We studied both p- and n-type MCz Si pad detectors after irradiation with reactor neutrons up to 10^{15}cm^{-2} .

-Priming procedures were investigated to optimize the visualization of traps by TSC and to evidence the electric field distribution within the irradiated device. Best priming process found is optical priming in forward polarization (to be checked without polarization).

-The various features observed experimentally by changing the operative parameters are qualitatively described by band diagrams. Best description accounts for the presence of a residual electric field (polarization of the irradiated Si bulk) due to charges frozen at traps in bulk and barriers close to the electrodes. Residual electric field in the bulk can be so high, in highly irradiated device, to dominate the current emission both in TSC and ZB-TSC.

-The method has been applied to the shallow level at 30K. A reliable trap concentration was found ($2 \times 10^{13}\text{cm}^{-3}$), electron-trap nature of defect was determined by comparison of TSC and ZBTSC emissions. Results indicate that this defect is main responsible of compensation in heavily irradiated MCZ Si detectors.



It remember.



Spares



- Useful references
- References on methods and on mechanisms
- References on results
- Addictional results
- Discussion on model



Spares: introduction

- What they (and you) have said more about
 - On the role of contacts:
 - “It should be noted that the space charge is always concentrated near the electrodes and consequently the role of contact phenomena in the behaviour of TSD currents is **sufficient**.”*
 - “the fast retrapping case, when the influence of the contacts on the TSDC curve is **more apparent**”**
 - On the interpretation of measurements:
 - “**starting polarity** of the discharging current is different for various irradiation times, depending on the degree of filling of the traps.”***
 - “the integral of the external current measures only a **small fraction** of the charge lost. This arises because the space charge which is driven deeper into the material contributes to a displacement current at $x=0$ which is a significant fraction of the conduction current at that point and opposed it in direction.”****
 - On how should be the priming for:
 - “the observation of an external current requires not only a time changing density of trapped charge but also a **change in the relative distribution** of trapped charge to give a time-changing zero-field point.”*****
 - “The calculation illustrate the important fact that the largest fraction of the total charge is observed externally when the **initial charge is located within one half** of the thickness of the sample.”*****
 - On the target of the investigations:
 - “The external measurements of current and charge transferred between the electrodes is found to depend strongly on the **ratio of transport-to-trapping times**. [...] For large values of this parameter, a significant fraction of the internal charge can be observed externally, and the externally measured quantities become relatively easy to interpret in terms of the release and retrapping of carriers within the insulator.”*****

*Kostsova and Kostsov, JPD, 2008

**Kostsova and Salman, PSS, 1995

***Petre et al., JAP, 1994

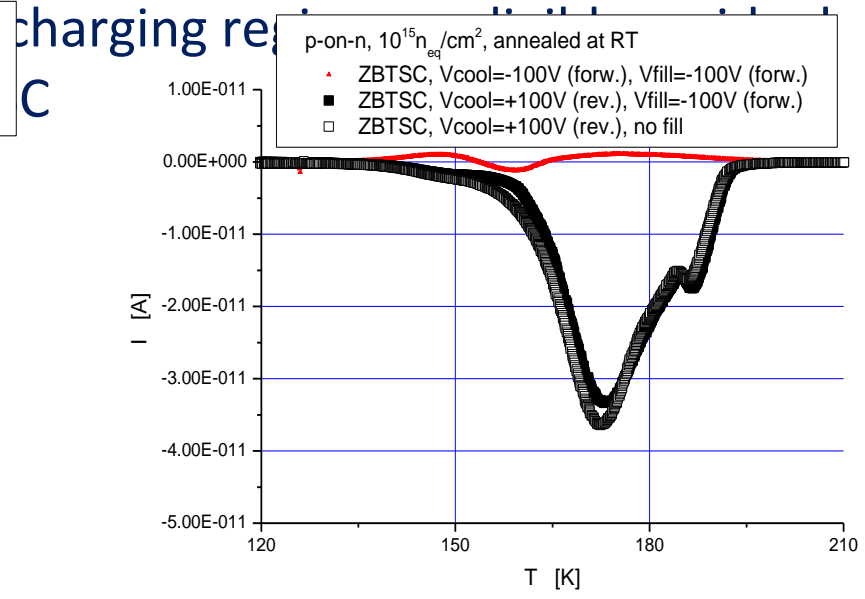
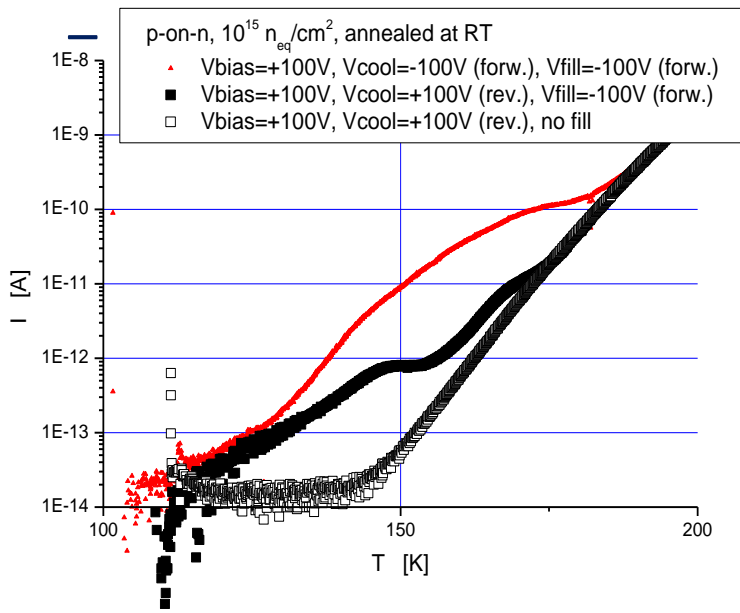
****Wintle, JAP, 1971

*****Monteith and Hauser, JAP, 1967

Spares: experimental results: varying priming



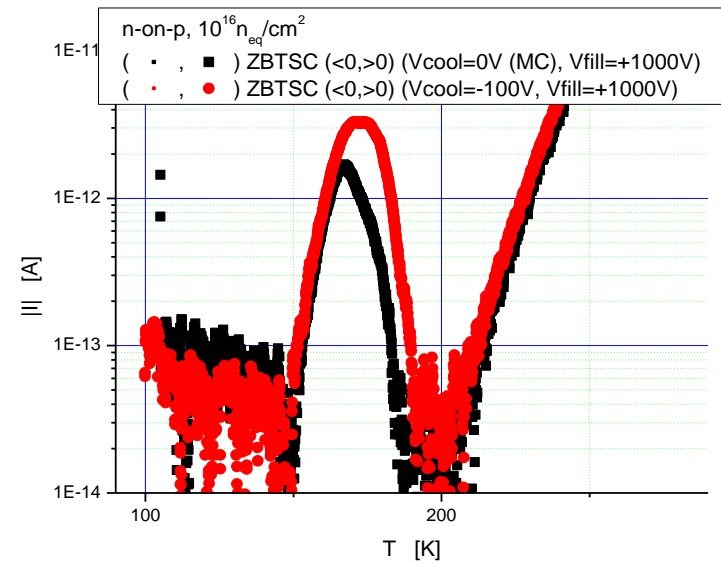
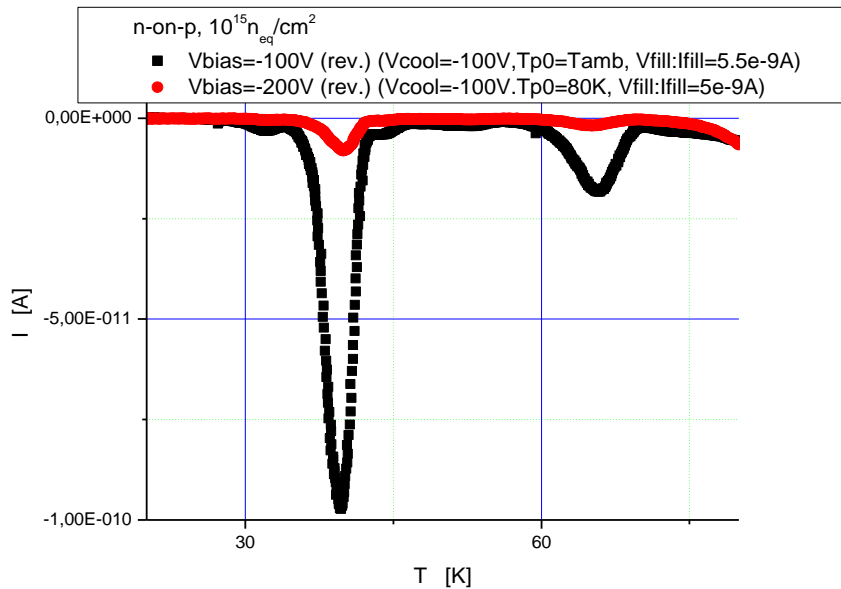
– Forward current priming: varying Vcool



Spares: experimental results: varying priming



- Forward current priming: varying methods (Tp0 and Tfill)



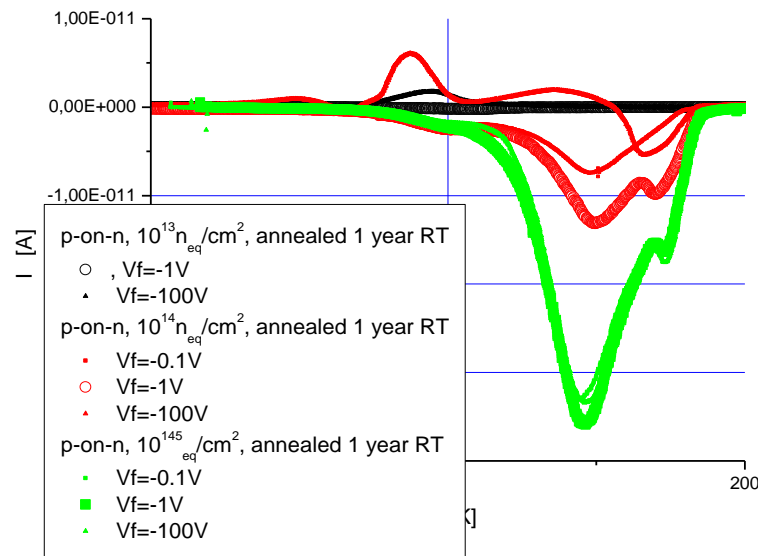
- Higher temperature of reverse polarization \Rightarrow higher residual field
- Lower temperature of forward filling \Rightarrow lower current \Rightarrow higher effects of residual field
- Effects are highlighted in ZBTSC \Rightarrow residual field is the protagonist

Spares: experimental results: varying priming



– Varying irradiation fluence

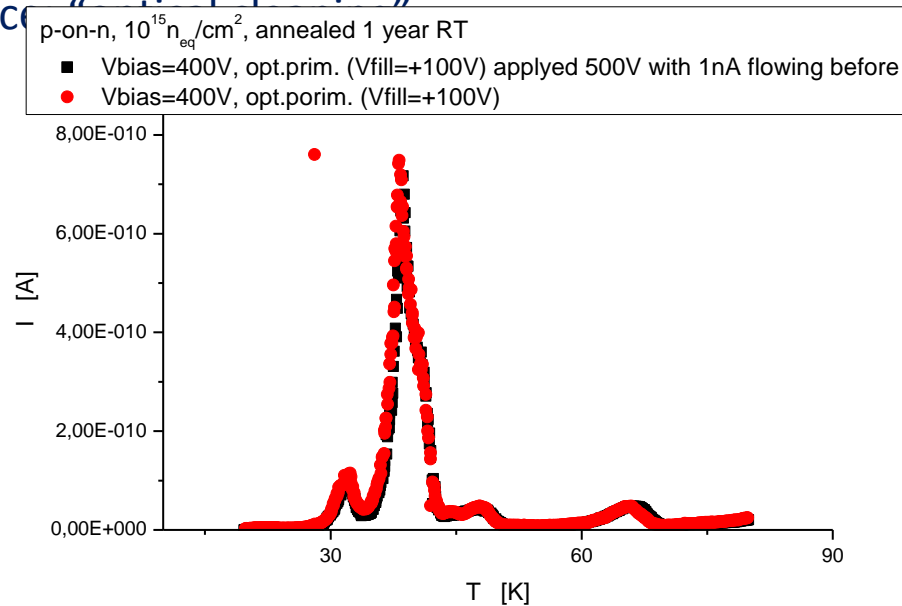
- Effects increase, increasing fluence \Rightarrow higher space charge, higher residual field



Spares: experimental results: varying priming



- Optical priming: varying conditions before illumination
- No dependence “...”

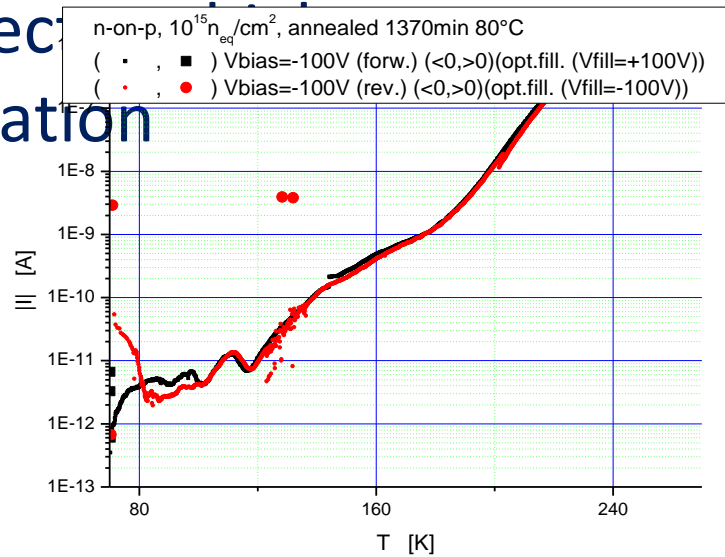


Spares: experimental results: varying priming



– At different temperatures (same T_{fill})

– Lower effective residual field relaxation times \Rightarrow residual field relaxation

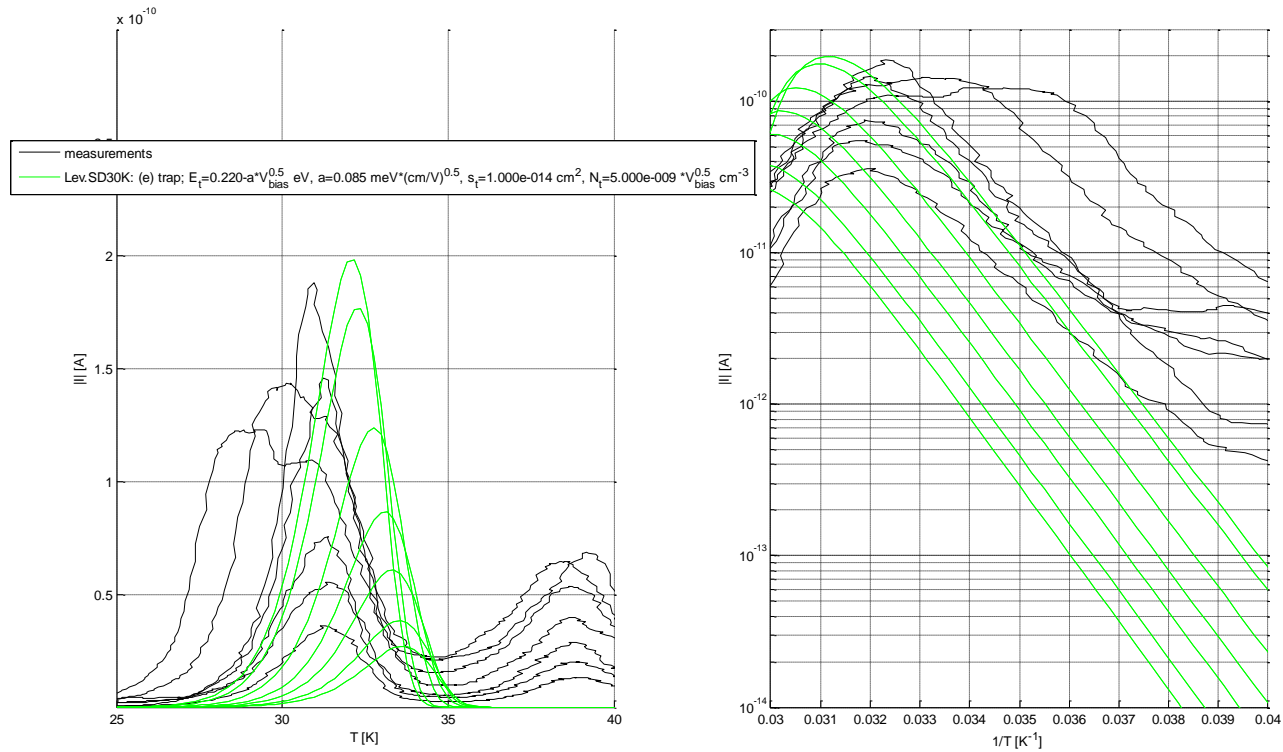


Spares: experimental results: defect studies



– SD30K: fits

- Scaringella* on MCz, STFZ, proton irradiated

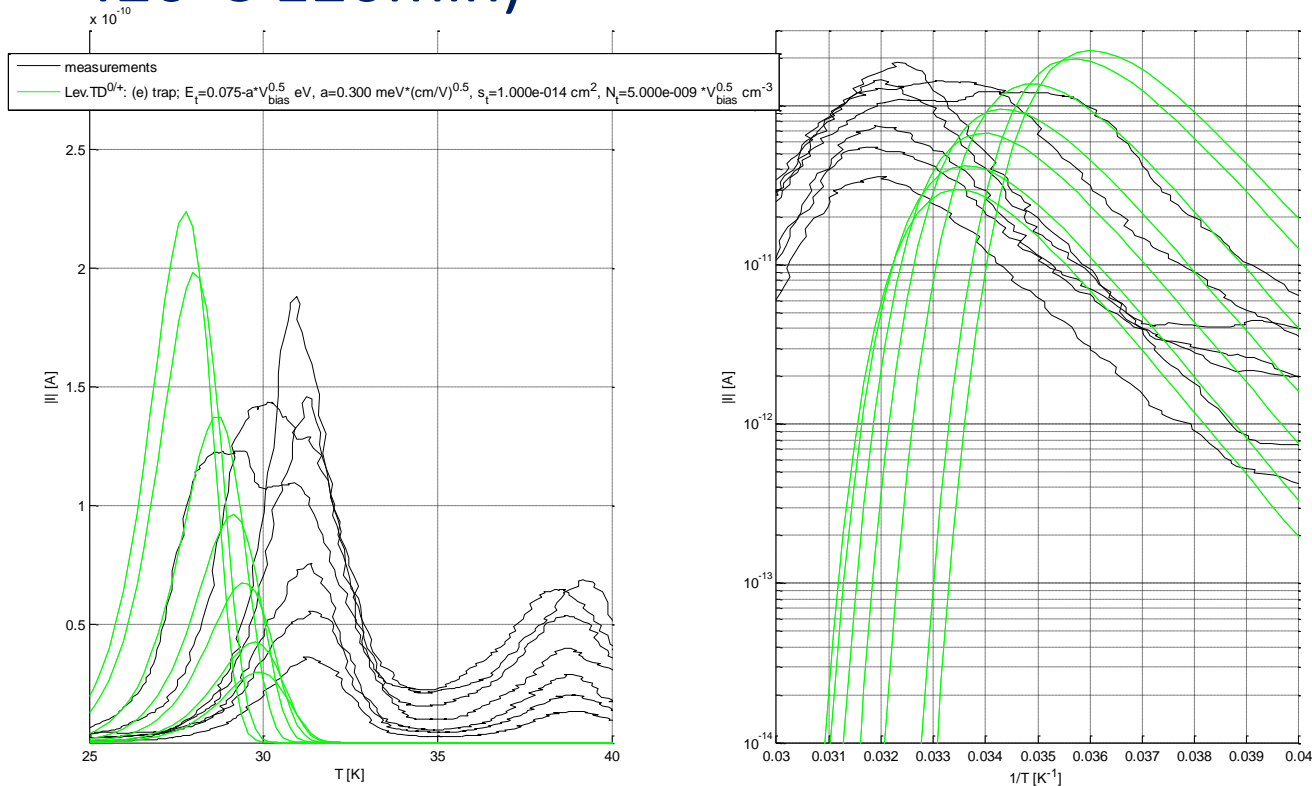


Spares: experimental results: defect studies



– SD30K: fits

- Bruzzi* on p-type MCz, (before annealing at 420°C 120min)

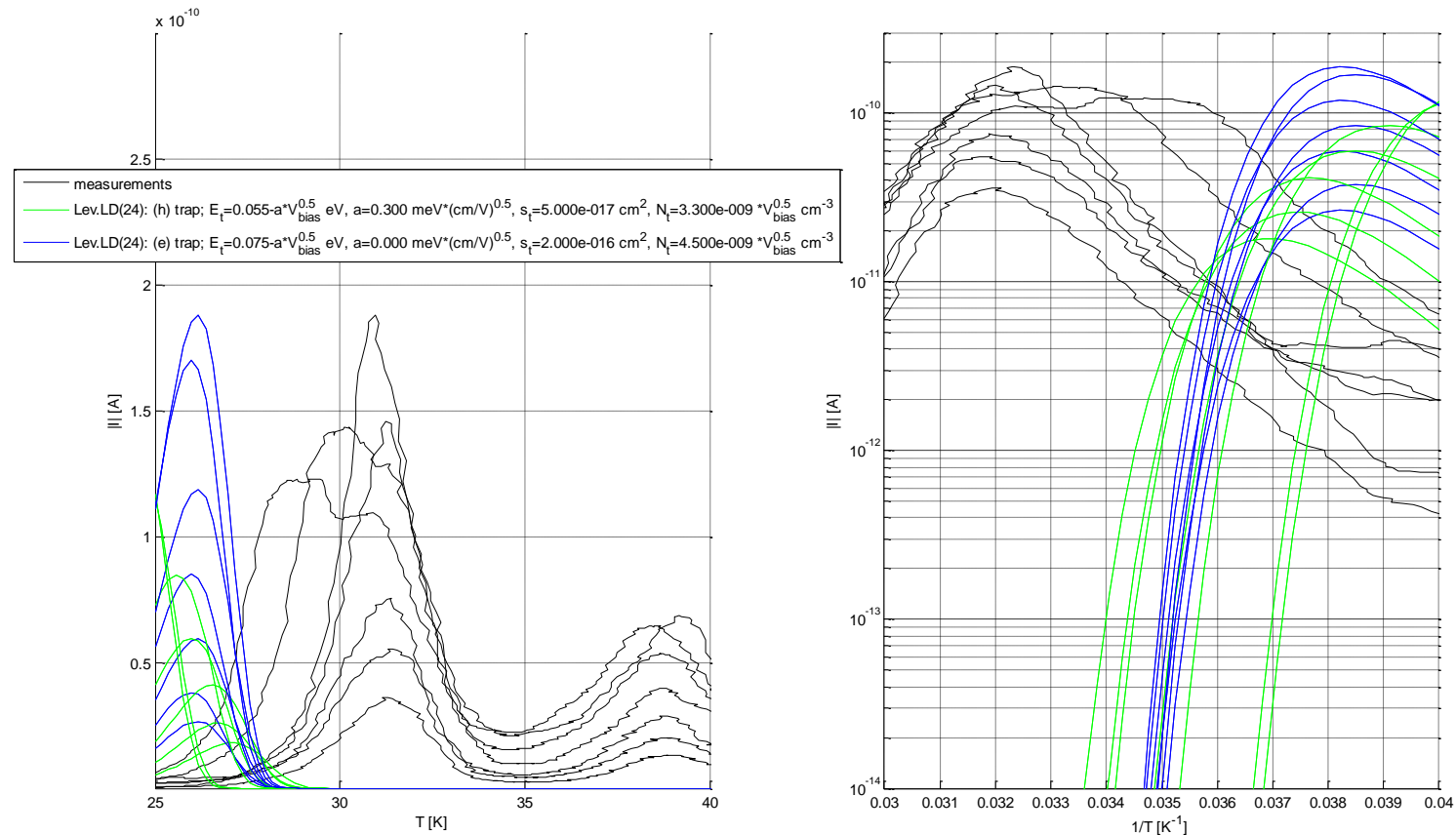


Spares: experimental results: defect studies



– SD30K: fits

- Menichelli* on DOFZ, STFZ, even unirradiated





Spares: discussion



– Stimulation mechanisms: modeling ZBTSC

– Lindmayer * on insulators: $x^* : \int_0^d dx \int_{x^*}^{x'} n_t(x') dx'$

$$dQ_{ext} = qn_t dx^* \Rightarrow I_{ext} = \frac{dQ_{ext}}{dt} = \frac{qn_t dx^*}{dt}$$

– Monteith** on insulators: $Q_{ext} = f\left(\frac{C_n N_t}{\mu}\right)$

$$\frac{C_n N_t}{\mu} \gg \frac{\varepsilon_t}{k_B T} \Rightarrow I_{ext} \propto \frac{dn_t}{dt}$$

– Gross*** on dielectrics:

$$\varepsilon E(0, t) = \varepsilon E(x, t) - \int_0^x \rho(x', t) dx' ; \quad x = x^* \Rightarrow \varepsilon E(0, t) = - \int_0^{x^*} \rho(x', t) dx'$$

$$j_D(0, t) = \frac{\varepsilon \partial E(0, t)}{\partial t} = - \int_0^{x^*} \frac{\partial \rho(x', t)}{\partial t} dx' + \rho(x^*, t) \frac{dx^*}{dt} = i(x^*, t) - i(0, t) + \rho(x^*, t) \frac{dx^*}{dt}$$

$$J = \frac{1}{d} \int_0^d i(x', t) dx' = j_D(0, t) + i(0, t) = i(x^*, t) + \rho(x^*, t) \frac{dx^*}{dt}$$

*Lindmayer, JAP, 1965

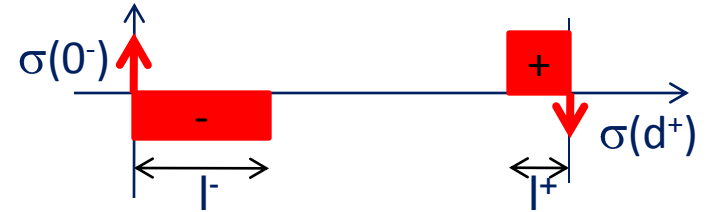
**Monteith and Hauser, JAP, 1967

***Gross and Perlman, JAP, 1972



Spares: discussion

– Interpreting measurements:



• Fast and slow relaxing*:

– Fast relaxation:

$$\sigma(0^-) = \frac{\varepsilon V}{d} \left[\rho^- l^- - \frac{\rho^- l^{-2} + \rho^+ l^{+2}}{2d} \right] \xrightarrow{v=0} - \left[\rho^- l^- - \frac{\rho^- l^{-2} + \rho^+ l^{+2}}{2d} \right]$$

– Slow relaxation:

$$\sigma(0^-) = - \left[\rho^- l^- - \frac{\rho^- l^{-2} + \rho^+ l^{+2}}{2d} \right] \xrightarrow{T \uparrow} - \left[\rho_{eq}^- l_{eq}^- - \frac{\rho_{eq}^- l_{eq}^{-2} + \rho_{eq}^+ l_{eq}^{+2}}{2d} \right]$$

• Trap type distinguish comparing with TSC**:

– (e^-) emission:

$$\left| E_{space \text{ ch arg } e} \right| : \begin{cases} \uparrow, & \rho > 0 \\ \downarrow, & \rho < 0 \end{cases} \Rightarrow j_{PK} : T_{PK} : \begin{cases} \downarrow, & \rho > 0 \\ \uparrow, & \rho < 0 \end{cases}$$

*Calderwood, and Scaife, PT Royal Society of London, 1970

**Pintilie, NIMA, 2000



Spares: discussion



– On model:

- Wintle on dielectrics*: $J(t) \propto [E(d,t)^2 - E(0,t)^2]$

– On measurements:

- Residual current value**: $J(t) \propto E_{bulk} > E_{applied}$

- Peak dependence on priming voltage***:

$$\left\{ \begin{array}{l} \text{Asymmetry} : \text{MAX at } V_{dep} / 2 \\ Q_t \propto \sqrt{V} \\ E_{anode} \propto V; E_{cathode} \propto \sqrt{V} \end{array} \right.$$

$$\Rightarrow I_{redistribution} \propto \left\{ \begin{array}{l} Q \propto \sqrt{V} \\ \text{Asymmetry} \end{array} \right. \propto \left\{ \begin{array}{l} \sqrt{V}, \quad V < V_{dep} \\ \rightarrow 0, \quad V > V_{dep} \end{array} \right. ; I_{injection}^{(anode)} \propto E_{anode} \propto V ; I_{injection}^{(cathode)} \propto E_{cathode} \propto \sqrt{V}$$

*Wintle, JAP, 1971

**Simmons and Taylor, JAP, 1972

***Kostsova and Kostsov, JPD, 2008

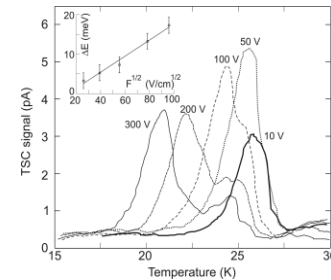


Spares: useful references

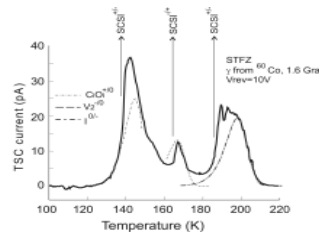


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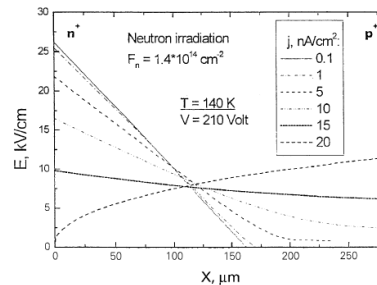
- Pinitilie [Ekstein 2009] group show: $V_{bias}=100V \Rightarrow I_{pk}(SD30K) >$ respect to V_{bias} 400V on EPIDO
- Bruzzi [Bruzzi 2006]: peak decrease with bias in n-on-p
- Scaringella [Scaringella 2007]: peak increase with bias in p-on-n
- Scaringella [Scaringella 2006]: peak constant with bias in p-on-n



- Scaringella [Scaringella Ph.D. Thesis, Chapt. 5, 2005]: peak decrease with bias in STFZ, DOFZ, p-on-n, irradiated by gamma-rays:



- Eremin [Eremin 2002], from Verbitskaya: increase filling current, more homogeneous electric field distribution and thus filling



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Experimental results: TSC and ZBTSC



Same for the CiOi and V2 peaks at higher temperature:

- negative SCR (p-type) \Rightarrow field reduction \Rightarrow peak delayed
- Positive SCR (n-type) \Rightarrow field increase \Rightarrow peak anticipated
- Hole emission:
 - negative SCR (p-type) \Rightarrow field reduction \Rightarrow peak delayed
 - Positive SCR (n-type) \Rightarrow field increase \Rightarrow peak anticipated

