Perspective of searching for axion-like particles in the mass range 10⁻⁷ - 10³ eV with stimulated photon-photon collider

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- 1. pseudo Nambu-Goldstone bosons as candidates for dark components in the Universe
- 2. Principle of stimulated resonant photon-photon scattering
- 3. Quick overview of experimental activities
- 4. Summary



pNGBs can be dark components of the Universe



Inflaton (ALP) mass and coupling to photons



Photon-photon collision systems



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S-matrix for two-body interactions in stimulated resonant scattering

K.Homma, Y. Kirita arXiv:1909.00983

 $S^{(2)} = \frac{i^2}{2} \int d^4x \int d^4y T[F_{\mu\nu}(x)F^{\mu\nu}\phi(x)F_{\rho\sigma}(y)F^{\rho\sigma}(y)\phi(y)]$ $N[F_{\mu\nu}(x)F^{\mu\nu}(x)F_{\sigma\rho}(y)F^{\sigma\rho}(y)\langle 0|T[\phi(x)\phi(y)]|0\rangle]$ $\propto a_{\vec{l}}^{\dagger}a_{\vec{l}}^{\dagger}a_{\vec{k}} a_{\vec{l}} \qquad \text{Scalar - propagator}$

Coherent state:

$$\begin{split} |N\rangle\rangle &\equiv \exp\left(-N/2\right)\sum_{n=0}^{\infty}\frac{N^{n/2}}{\sqrt{n!}}|n\rangle \quad |n\rangle = \frac{1}{\sqrt{n!}}\left(a^{\dagger}\right)^{n}|0\rangle \\ \langle\langle N|N\rangle\rangle &= 1 \quad \langle\langle N|\left(a^{\dagger}a\right)|N\rangle\rangle = N \quad a|N\rangle\rangle = \sqrt{N}|N\rangle\rangle, \quad \text{and} \quad \langle\langle N|a^{\dagger} = \sqrt{N}\langle\langle N|_{n}|_{n}\rangle\rangle \\ &= \sqrt{N}|N\rangle\rangle, \quad \text{and} \quad \langle\langle N|a^{\dagger} = \sqrt{N}\langle\langle N|_{n}|_{n}\rangle\rangle \\ &= \sqrt{N}|N\rangle\rangle, \quad \text{and} \quad \langle\langle N|a^{\dagger} = \sqrt{N}\langle\langle N|_{n}|_{n}\rangle\rangle \\ &= \sqrt{N}|N\rangle\rangle, \quad \text{and} \quad \langle\langle N|a^{\dagger} = \sqrt{N}\langle\langle N|_{n}|_{n}\rangle\rangle \\ &= \sqrt{N}|N\rangle\rangle, \quad \text{and} \quad \langle\langle N|a^{\dagger} = \sqrt{N}\langle\langle N|_{n}|_{n}\rangle\rangle \\ &= \sqrt{N}|N\rangle\rangle, \quad \text{and} \quad \langle\langle N|a^{\dagger} = \sqrt{N}\langle\langle N|_{n}|_{n}\rangle\rangle \\ &= \sqrt{N}|N\rangle\rangle \\ &= \sqrt{N}|N\rangle\rangle, \quad \text{and} \quad \langle\langle N|a^{\dagger} = \sqrt{N}\langle\langle N|_{n}\rangle\rangle \\ &= \sqrt{N}|N\rangle\rangle, \quad \text{and} \quad \langle\langle N|a^{\dagger} = \sqrt{N}\langle\langle N|_{n}\rangle\rangle \\ &= \sqrt{N}|N\rangle\rangle, \quad \text{and} \quad \langle\langle N|a^{\dagger} = \sqrt{N}\langle\langle N|_{n}\rangle\rangle \\ &= \sqrt{N}|N\rangle\rangle \\ &= \sqrt{N}|N\rangle\rangle \\ &= \sqrt{N}|N\rangle\rangle, \quad \text{and} \quad \langle\langle N|a^{\dagger} = \sqrt{N}\langle\langle N|_{n}\rangle\rangle \\ &= \sqrt{N}|N\rangle\rangle \\ &= \sqrt{N}|N\rangle\rangle$$

Transition amplitude: $1 + 1 \rightarrow 3 + 4$ $\ll N_1 | \ll N_4 | \ll 1_3 | S^{(2)} | N_1 \gg | N_4 \gg | 0 >$ $\propto \ll N_1 | \ll N_4 | \ll 1_3 | a_{\vec{l}}^{\dagger} a_{\vec{j}}^{\dagger} a_{\vec{k}} | N_1 \gg | N_4 \gg | 0 >$ $\propto \sqrt{N_1} \sqrt{N_1} \sqrt{N_4} \ll N_1 | N_1 \gg \ll N_4 | N_4 \gg < 0 | 0 >$

s-channel propagator cannot be implemented - creation and decay points are spatially apart -



CAST, Theopisti Dafni, 7th Patras Workshop, Mykonos 2011

Light Shining through a Wall (LSW)



Okun 1982, Skivie 1983, Ansel'm 1985, Van Bibber et al. 1987

s-channel scattering contains resonance



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The first search for scalar field with FWM



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PTEP

The first search for sub-eV scalar fields via four-wave mixing at a quasi-parallel laser collider

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A search for sub-eV scalar fields coupling to two photons has been performed via four-wave mixing at a quasi-parallel laser collider for the first time. The experiment demonstrates the novel approach of searching for resonantly produced sub-eV scalar fields by combining two-color laser fields in the vacuum. The aim of this paper is to provide the concrete experimental setup and the analysis method based on specific combinations of polarization states between incoming and outgoing photons, which is extendable to higher-intensity laser systems operated at high repetition rates. No significant signal of four-wave mixing was observed by combining a $0.2 \,\mu$ J/0.75 ns pulse laser and a 2 mW CW laser on the same optical axis. Based on the prescription developed for this particular experimental approach, we obtained the upper limit at a confidence level of 95% on the coupling–mass relation.

Run I at Kyoto-ICR



with atomic Four-Wave Mixing (FWM)

Time structures of the number of photons $POL{1}$ $POL{2}$ s С The number of photons/event The number of photons/event 0.015 0.015 0.01 0.01 Only C C+I 0.005 0.005 60 60 70 80 90 100 110 70 80 90 100 110 Time [ns] Time [ns] Ρ The number of photons/event The number of photons/event 0.015 0.015 Pedestal Only I 0.01 0.01 0.005 0.005 0 70 100 60 80 90 110 60 70 80 100 110 90 Time [ns] Time [ns]

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Pressure dependence of atomic FWM (corrected)



Quadratic pressure dependence is consistent with $\chi^{(3)}$ process

PTEP

Search for sub-eV scalar and pseudoscalar resonances via four-wave mixing with a laser collider

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Run II with 10 TW at Kyoto-ICR



Extreme-Light-Infrastructure (ELI)





ELI-NP facility (280M€) Comm. starts from 2019 10 PW x 2 1 PW x 2 0.1 PW x 2

0.2-19.5 MeV gamma beam produced by ~700 MeV e- + laset⁶

Sensitivity below sub-eV mass domain in Quasi-Parallel collision



Asymmetric head-on collision for high-mass pNGB search



Sensitivity in 0.1 eV–10 keV mass domain in Asymmetric Head-on collision



Entire scope by this method



Klystron as a 10⁻⁵ eV photon source

PEAK POWER CAPABILITY OF MICROWAVE AMPLIFIER TUBES







J.Plasma Fusion Res. Vol.86 (2010)

Can we reach dilaton ?

K.Homma, Y. Kirita arXiv:1909.00983



Summary

- Charged particle collisions discovered SM particles in extremely strong coupling and high-mass domains over 10⁵ 10¹¹ eV
- Stimulated photon-photon collider can provide windows in extremely weak coupling and lowmass domains over 10⁻⁷ – 10³ eV
- Current technologies can provide the ways to access gravitational coupling strength in elementary scattering processes !