

# NLO VBF MONTE CARLO GENERATOR(S)

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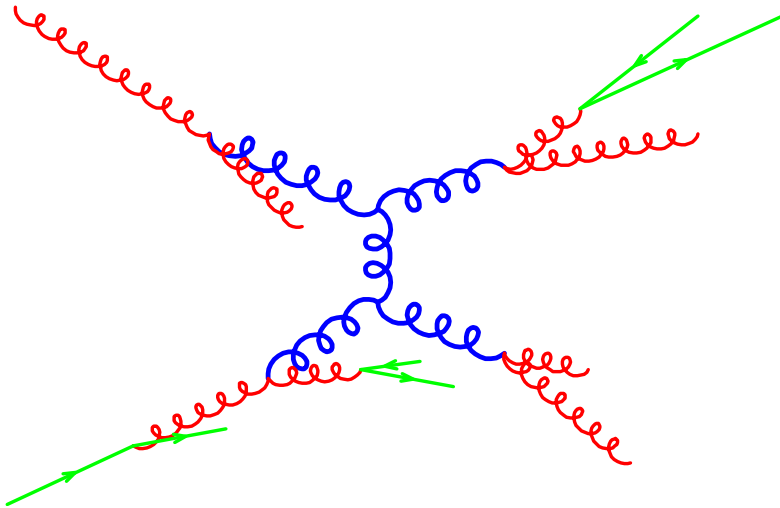
- Shower Monte Carlo + NLO
- POWHEG
- Higgs boson in VBF:  
POWHEG results
- Conclusions



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<sup>a</sup>On behalf of the VBF group: Ansgar Denner, Sinead Farrington, Christoph Hackstein and Daniela Rebutzi.

## Dominant corrections



Collinear-splitting processes in the initial and final state (always with **transverse momenta**  $> \Lambda_{\text{QCD}}$ ) are **strongly enhanced**. This is due to the fact that, in perturbation theory, the **denominators** in the propagators are **small**.

- The algorithms that evaluate (and **resum**) all these enhanced contributions are called **shower algorithms**.
- Shower algorithms give a description of a hard collision up to **distances of order**  $1/\Lambda_{\text{QCD}}$ .
- At larger distances, perturbation theory breaks down and we need to resort to **non-perturbative methods** (i.e. lattice calculations). However, these methods can be applied only to simple systems. The only viable alternative is to use **models of hadron formation**.

## SMC

- In high-energy collider physics not many questions can be answered without a Shower Monte Carlo (SMC).
- The name **shower** comes from the fact that we **dress** a **hard event** with **QCD radiation**.
- In general, the accuracy of SMC programs is **Leading Order + Leading Log**. They resum the largest logarithmic terms coming from the **collinear (and soft) regions**.
- Events are then characterized by a **small number** of **high- $p_T$** , **well-separated**, final-state partons + many collinear partons, whose collinear divergences have been correctly resummed.

# NLO

LO matrix elements are (often, but not always) good for shapes. Uncertain absolute normalization

$$\alpha_s(\mu) = \frac{\alpha_s(\mu_0)}{1 + b_0 \alpha_s(\mu_0) \log(\mu^2/\mu_0^2)}, \quad b_0 = \frac{11C_A - 4n_f T_F}{12\pi} \Big|_{n_f=5} \approx 0.6$$
$$\alpha_s^n(2\mu) \approx \alpha_s^n(\mu) [1 - b_0 \alpha_s(\mu) \log(4)]^n \approx \alpha_s^n(\mu) (1 - n\alpha_s(\mu))$$

For  $\mu = 100$  GeV,  $\alpha_s = 0.12$ , normalization uncertainty:

$W + 1J$	$W + 2J$	$W + 3J$
$\pm 12\%$	$\pm 24\%$	$\pm 36\%$

To improve on this, we need to go to NLO

- Positive experience with NLO calculations at LEP, HERA and Tevatron
- NLO results are cumbersome to compute: typically made up of an  $n$ -body (Born + virtual + soft and collinear remnants) and  $(n + 1)$ -body (real emission) terms, both divergent (finite only when summed up).
- Merging NLO with shower is a natural extension of present approaches
  - MC@NLO [Frixione, Webber]
  - POWHEG [Nason]

## POsitive-Weight Hardest Emission Generator

- ✓ it is a method for interfacing **NLO** calculations with **parton shower** programs [Nason, hep-ph/0409146]
- ✓ it generates the **hardest emission** first, with **NLO accuracy**. The produced events have **positive weights**. The acronym comes from this feature
- ✓ it is **independent** from **parton-shower** programs. It can be interfaced with any programs that comply with the LHE format: **PYTHIA**, **HERWIG** ... It is then possible to **compare** the **different outputs**
- ✓ **No need to implement new interfaces**  
Two possible ways to interface to shower Monte Carlo programs
  1. **Les Houches Event** format. The event is written on a file that is subsequently showered by HERWIG, PYTHIA...
  2. **on the fly**. We provide UPINIT and UPEVNT directly running in HERWIG and PYTHIA

## Higgs boson in VBF: analysis cuts

We have used the POWHEG BOX to generate events for Higgs boson in VBF

- $pp$  collider at 7 TeV
- jet clustered with  $k_T$  algo,  $R = 0.7$
- **jet** defined by

$$p_{Tj} > 20 \text{ GeV}, \quad |y_j| < 5$$

- **tagging jets**

$$p_T^{\text{tag}} > 30 \text{ GeV}, \quad |y_{j_1} - y_{j_2}| > 4.2, \quad y_{j_1} \cdot y_{j_2} < 0, \quad m_{jj} > 600 \text{ GeV}$$

$$\text{veto jet: } \min(y_{j_1}, y_{j_2}) < y_j < \max(y_{j_1}, y_{j_2})$$

- info on the **jet shape**. If  $\vec{p}^j$  is the momentum of the jet, then

$$p_T^{\text{rel},j} = \sum_{k^i \in p^j} \frac{|\vec{k}^i \times \vec{p}^j|}{|\vec{p}^j|}$$

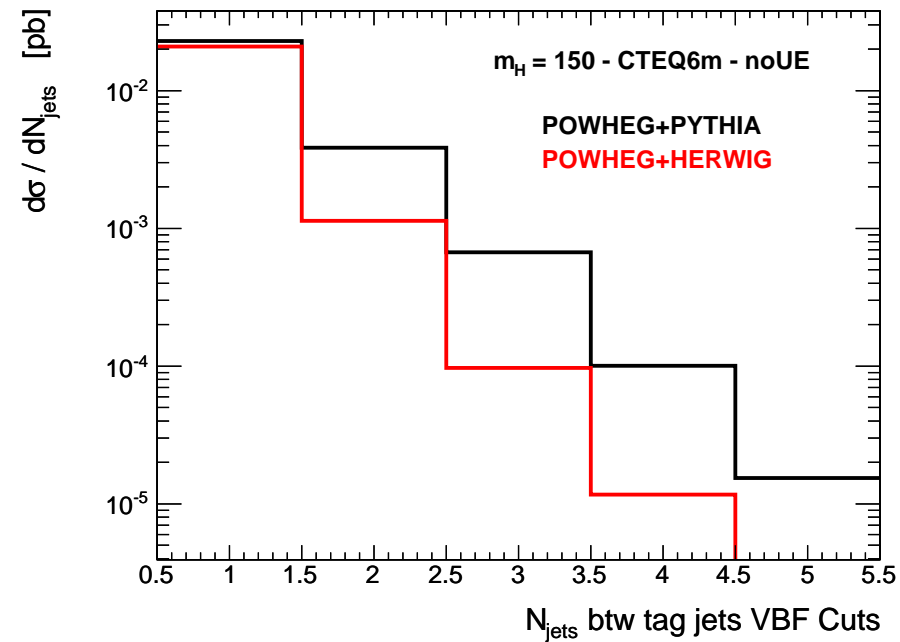
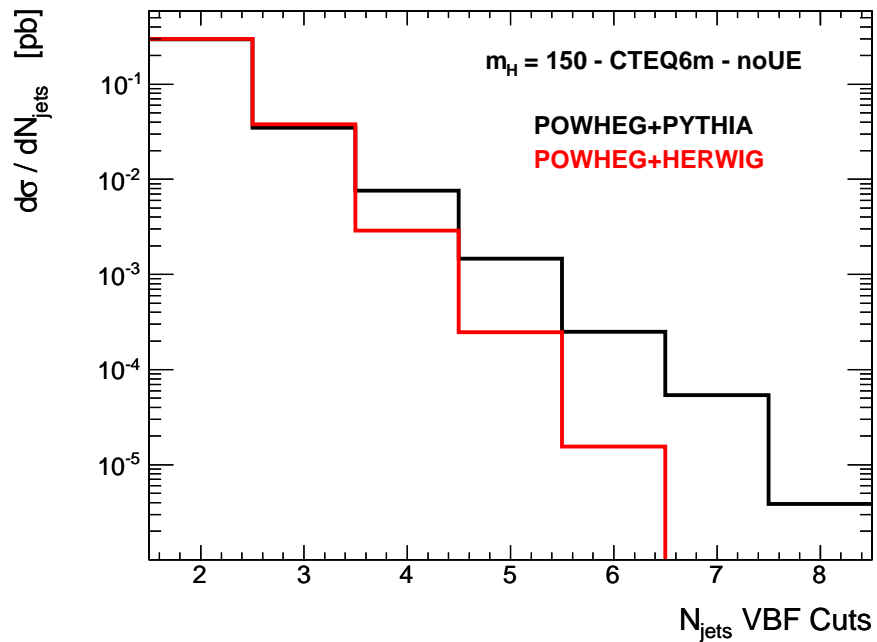
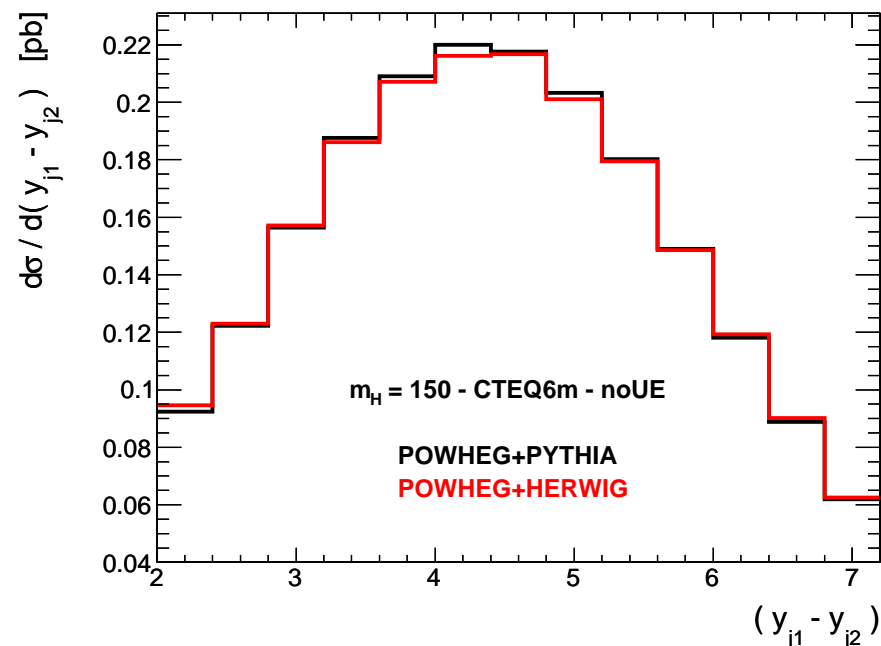
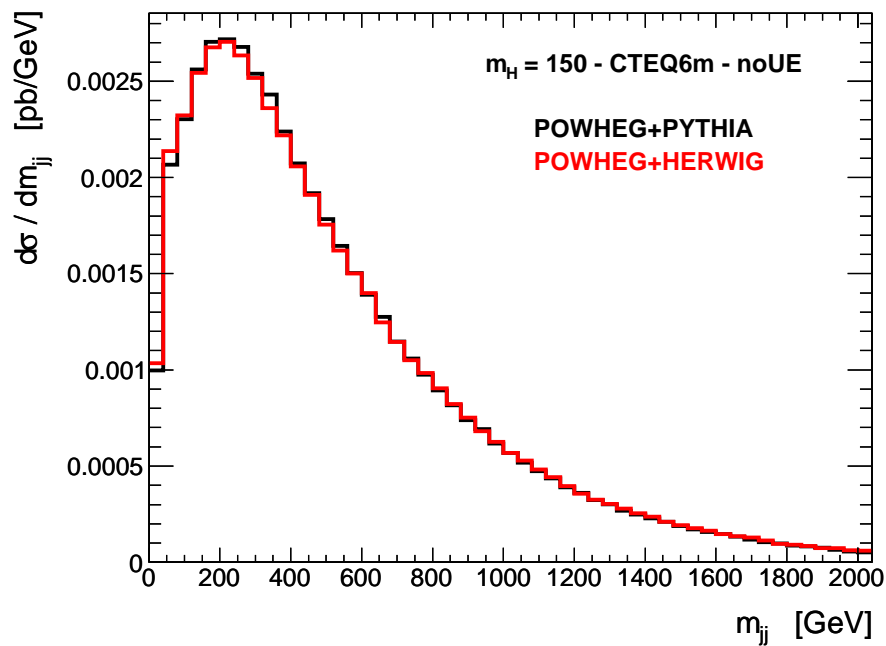
## Higgs boson in VBF: results

- comparisons among
  - POWHEG + HERWIG (ver 6.510)
  - POWHEG + PYTHIA (ver 6.4.21)
  - HERWIG alone
  - PYTHIA alone
- for  $m_H = 150, 200, 300, 400, 500$  GeV
- for two pdf sets: **CTEQ6m** and **MSTW2008NLO** (and the respective LO ones when used with HERWIG and PYTHIA alone)
- **without** and **with** underlying events
  - HERWIG  $\implies$  IPROC (no +10000)  
It contains only **soft** UE. **NO Multiple Interactions** (provided by the JIMMY package).
  - PYTHIA  $\implies$  mstp(81)=1 and **new  $p_T$ -ordered shower** (**MI included** in the default mode)

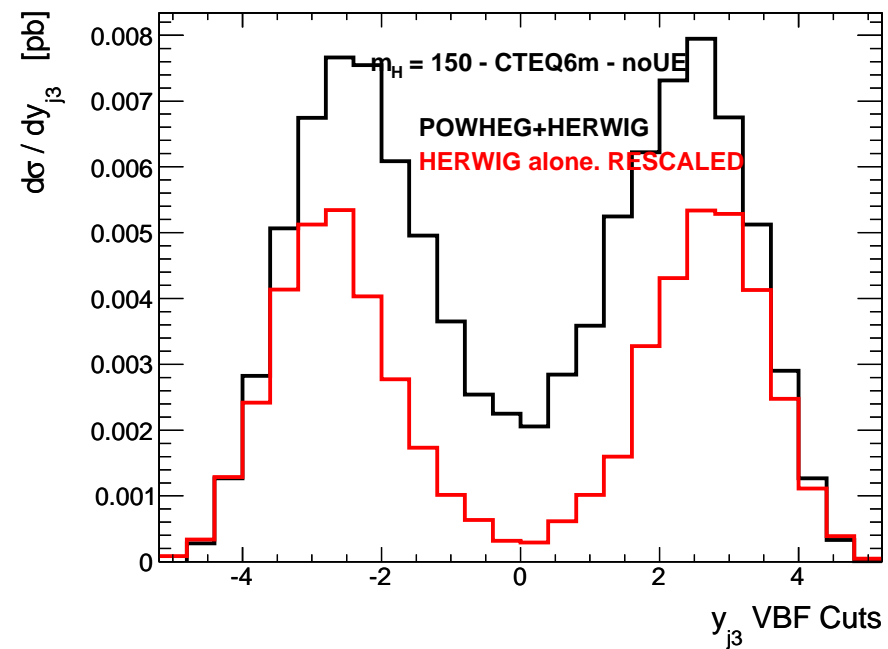
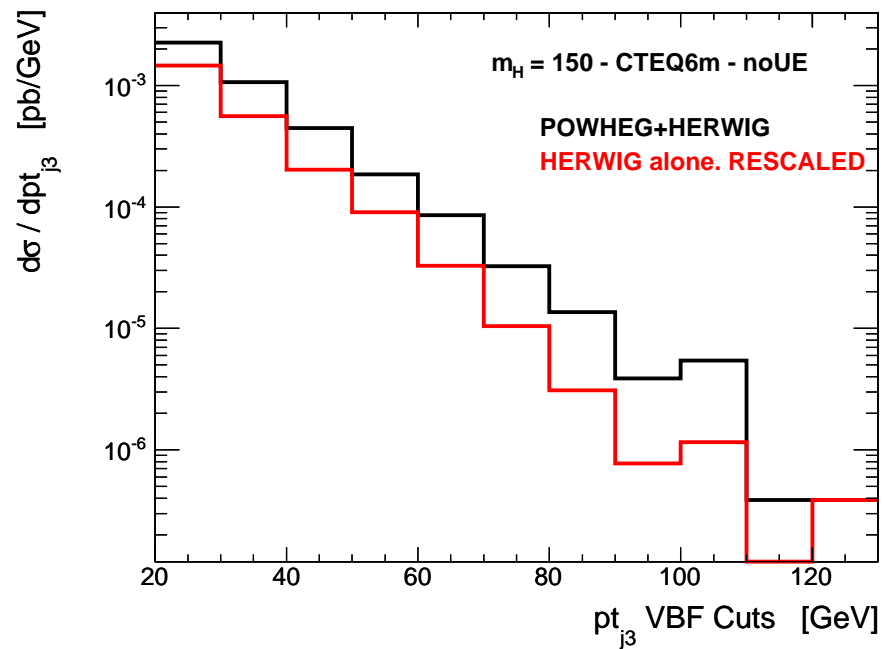
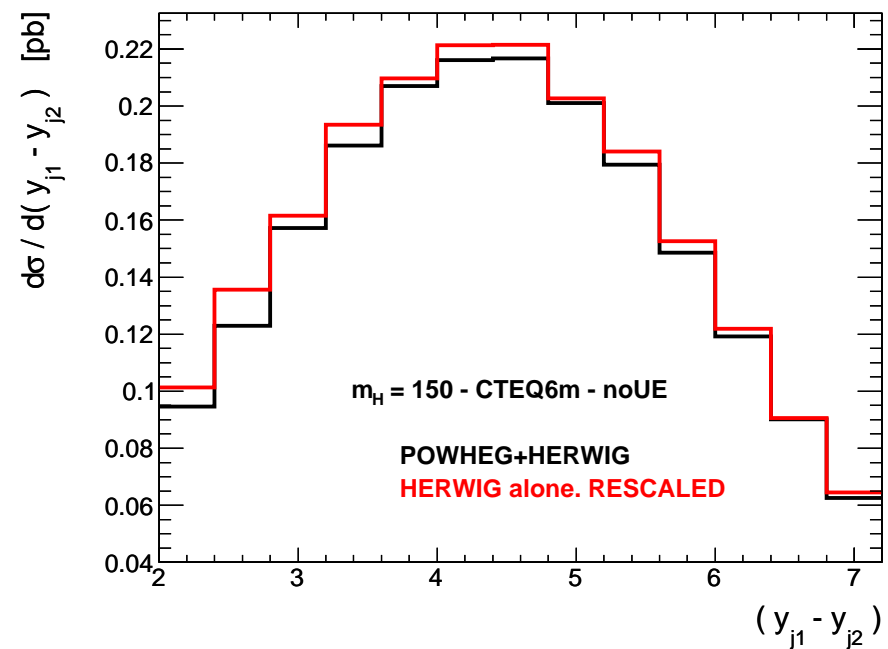
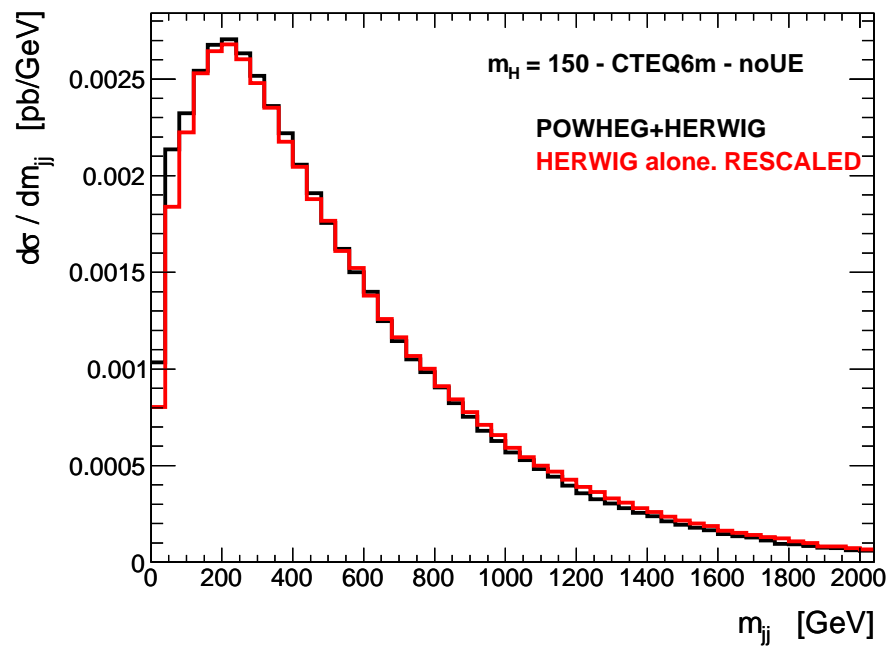
## Higgs boson in VBF: results

- We will present results mostly for  $m_H = 150$  GeV.
- **Same conclusions** for all the Higgs boson masses up to 500 GeV
- Results from the analysis of 500.000 events generated by the POWHEG BOX

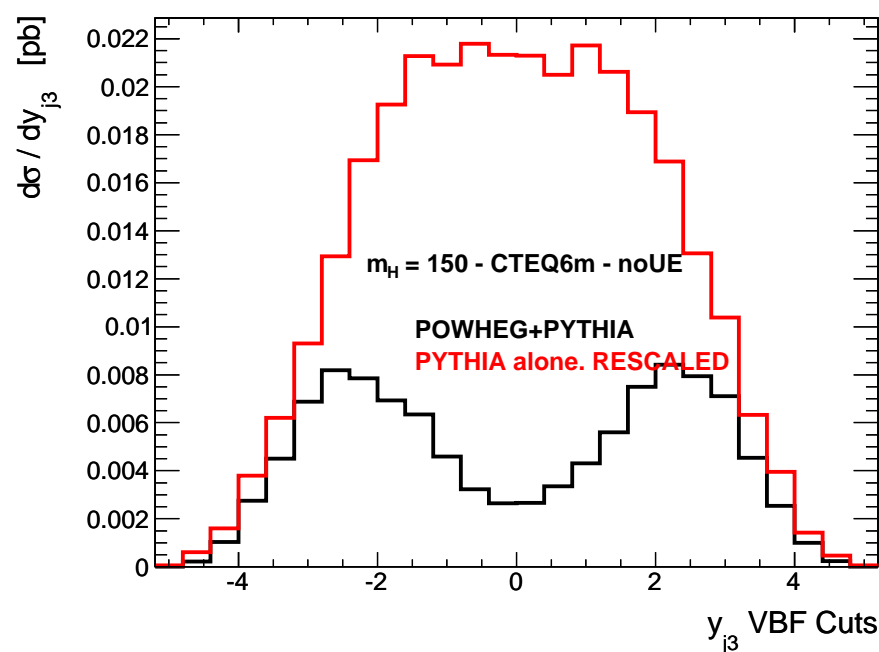
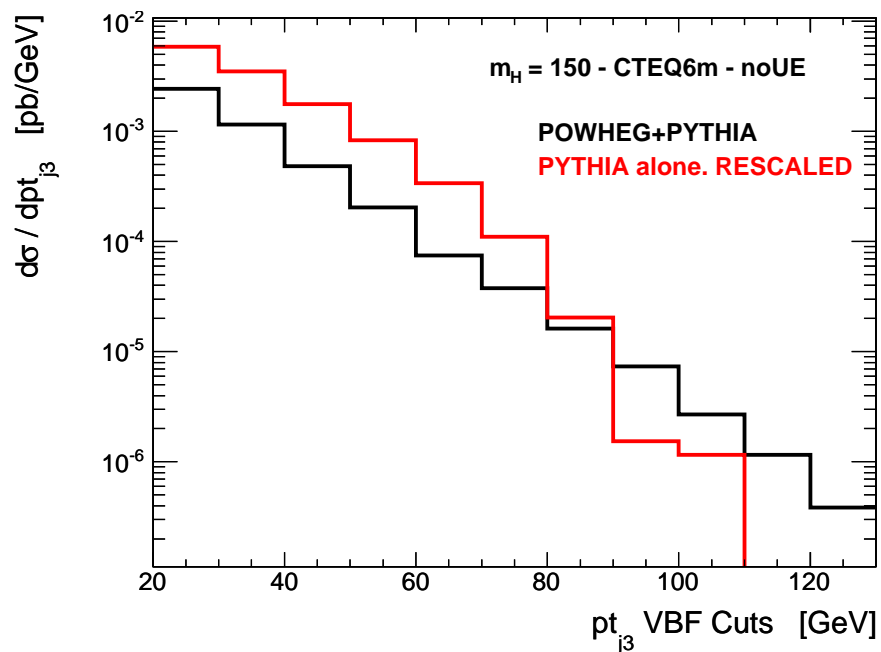
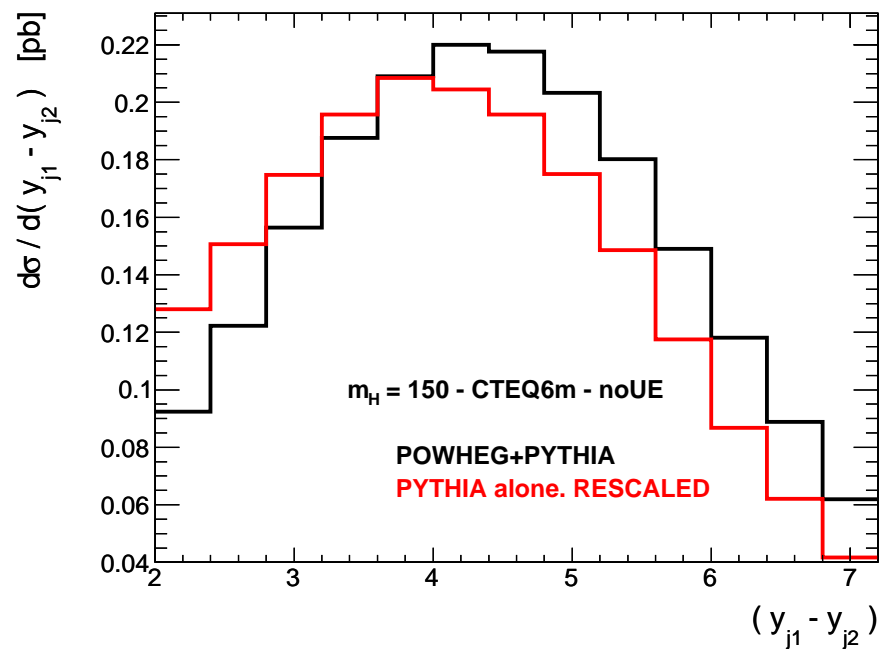
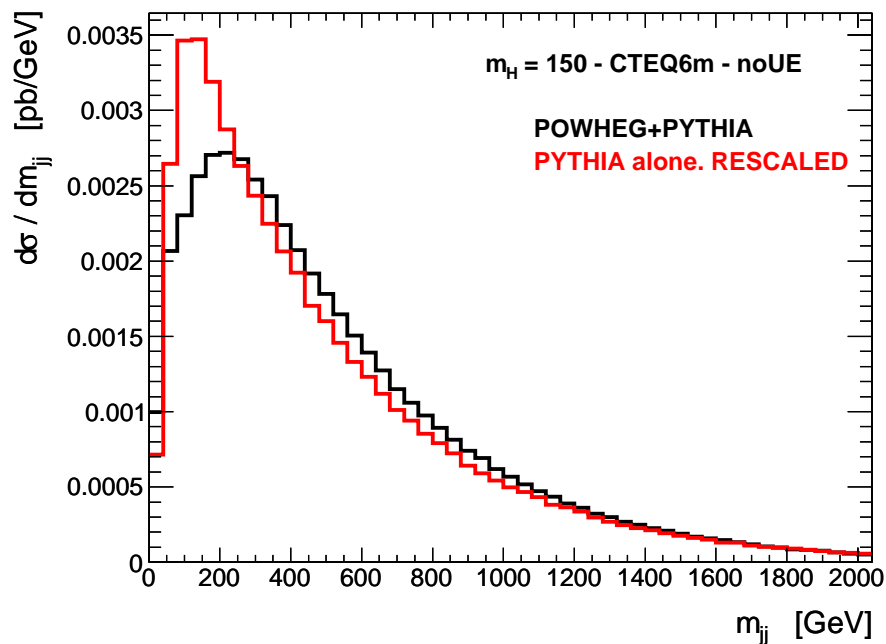




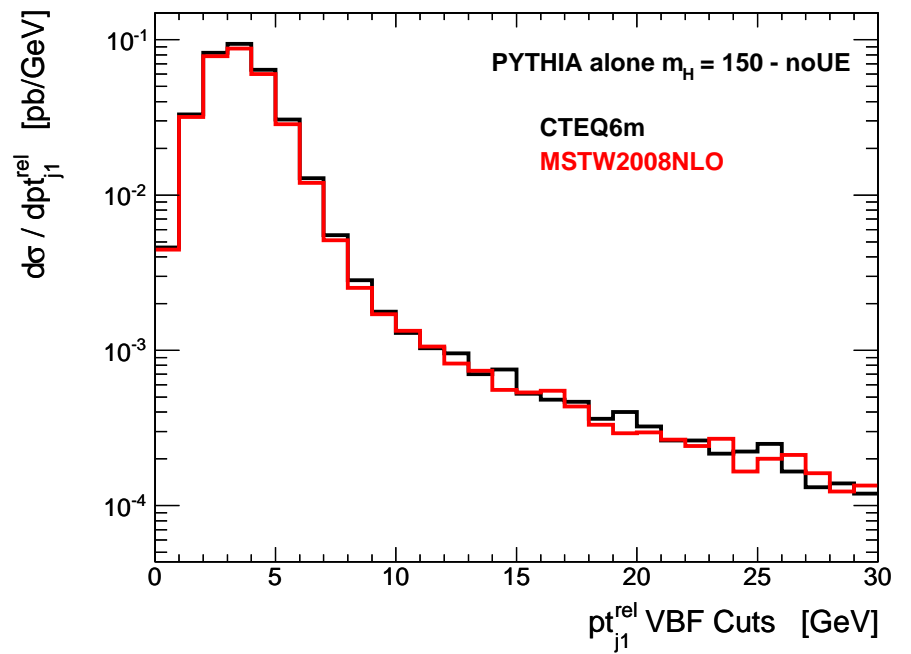
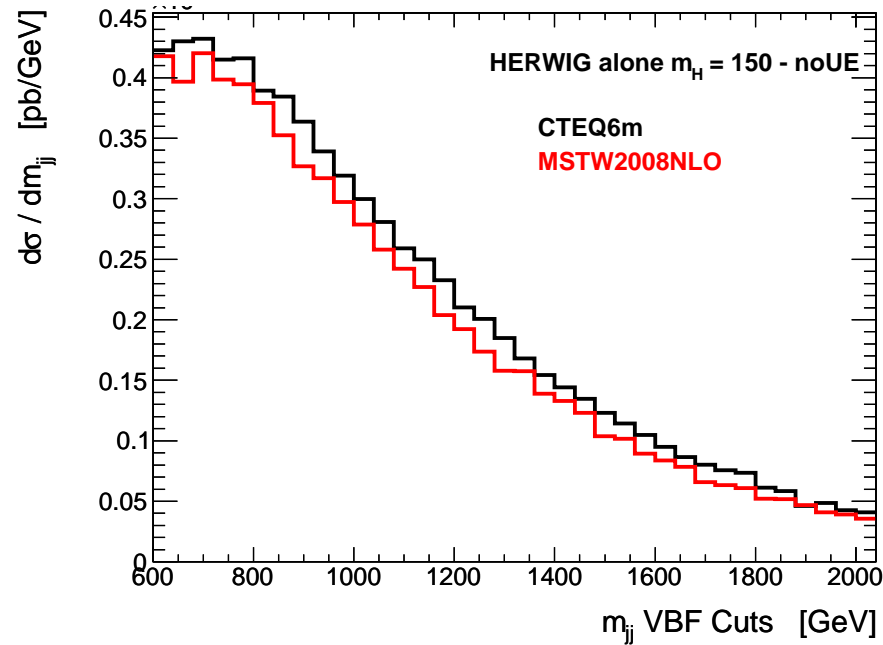
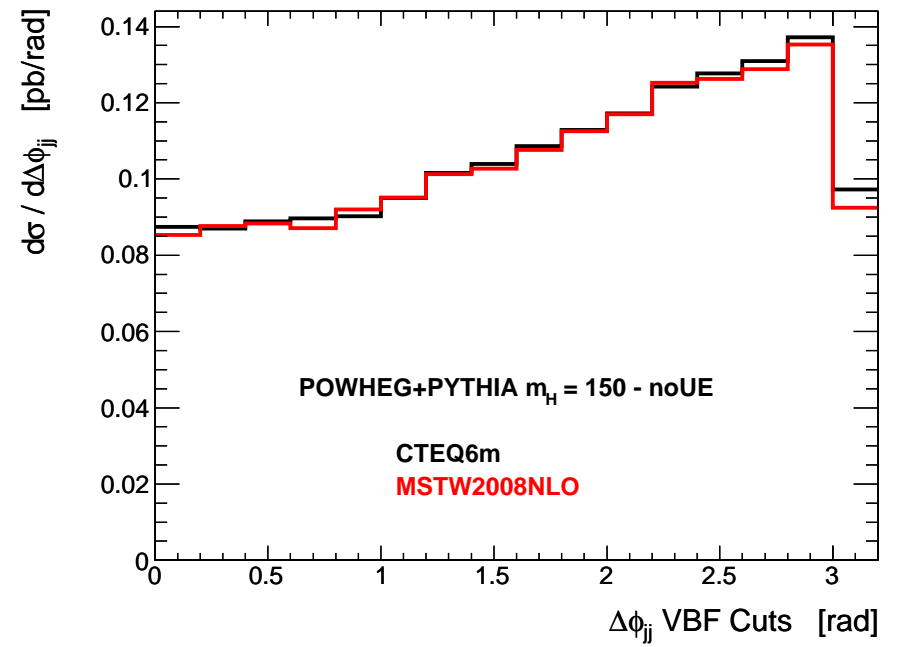
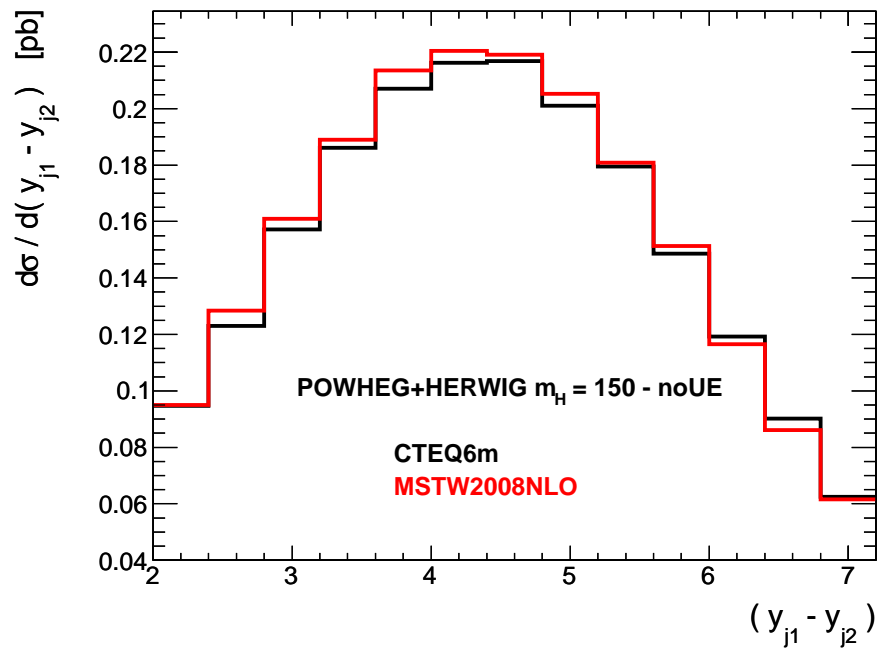
PYTHIA jets tend to be harder than HERWIG's ones.  $N_{jets}$  is correct up to 3 (NLO), since generated by POWHEG



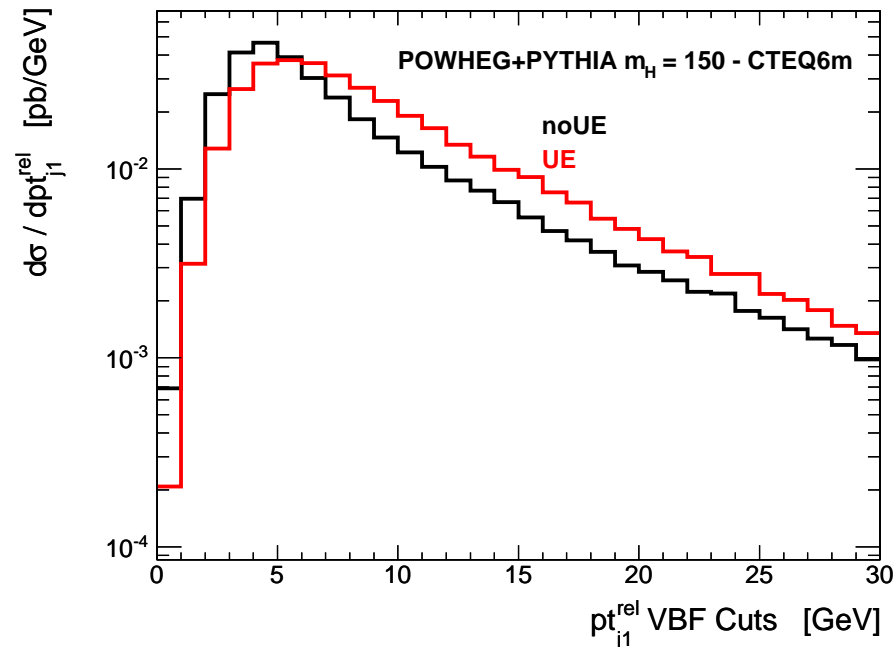
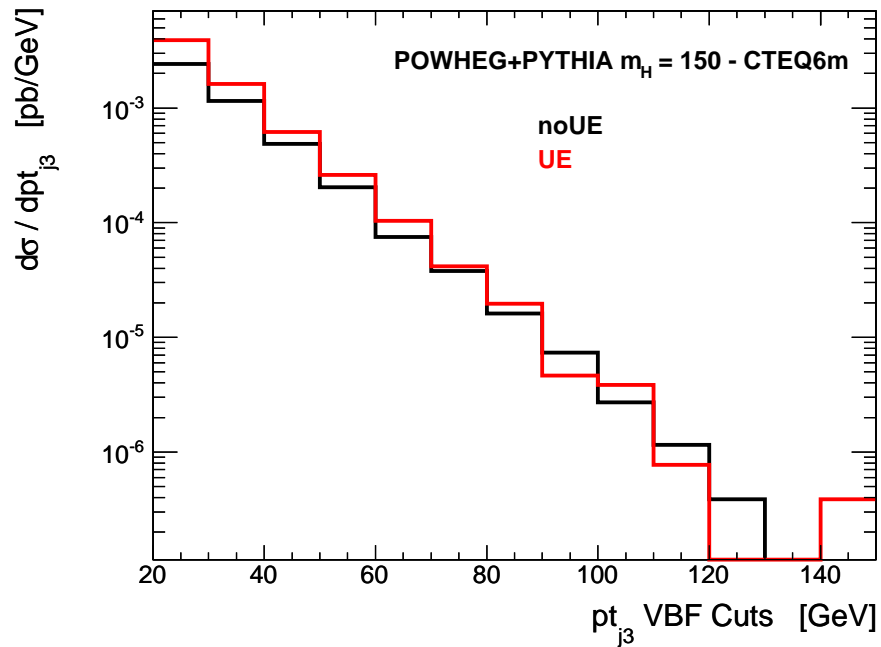
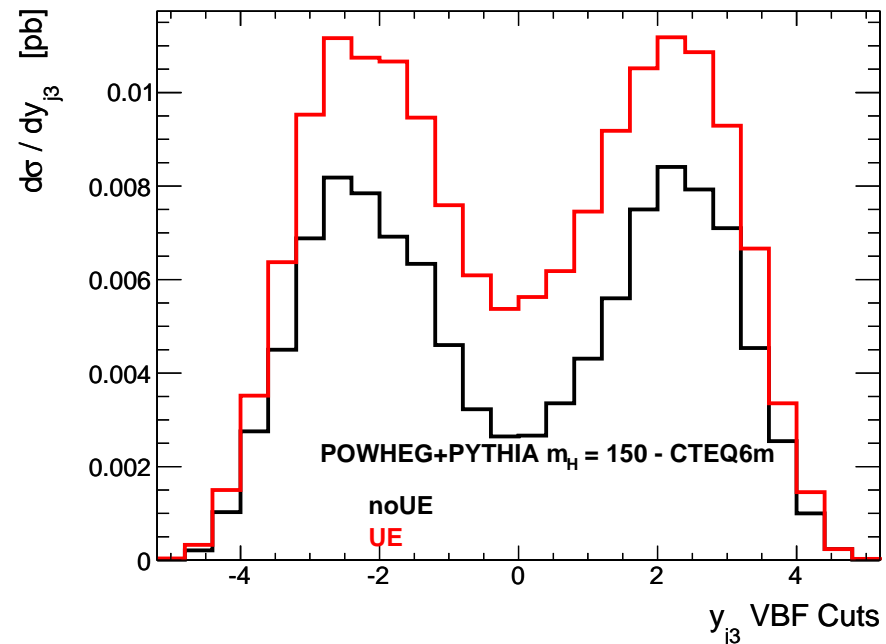
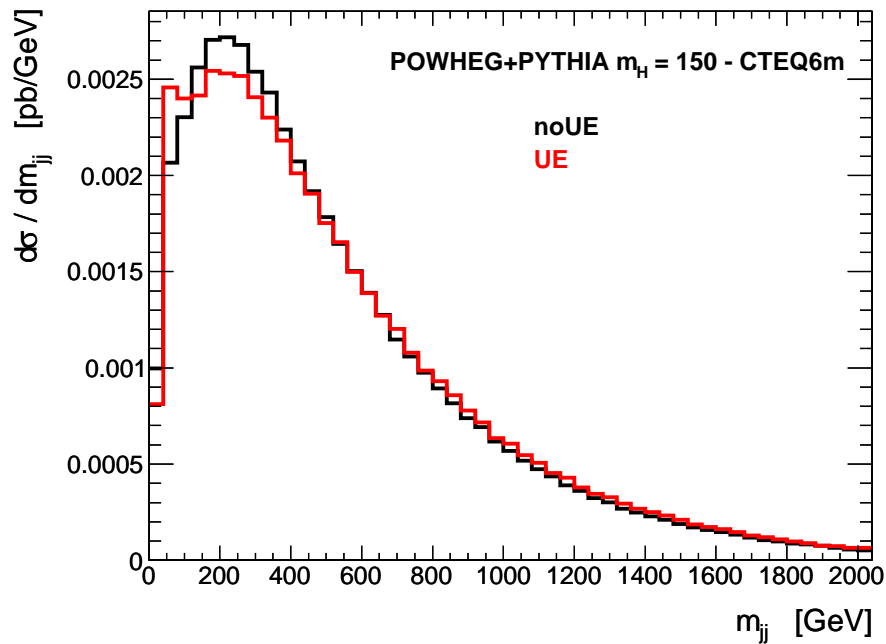
Neither HW alone nor PY alone are expected to describe correctly the 3rd jet. HW results **rescaled** to the NLO total Xsec



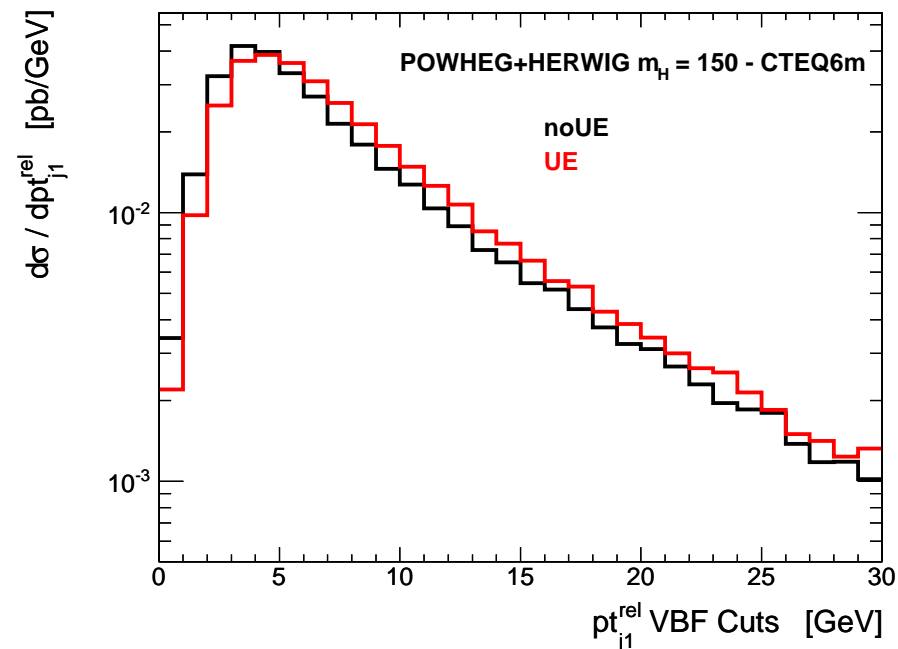
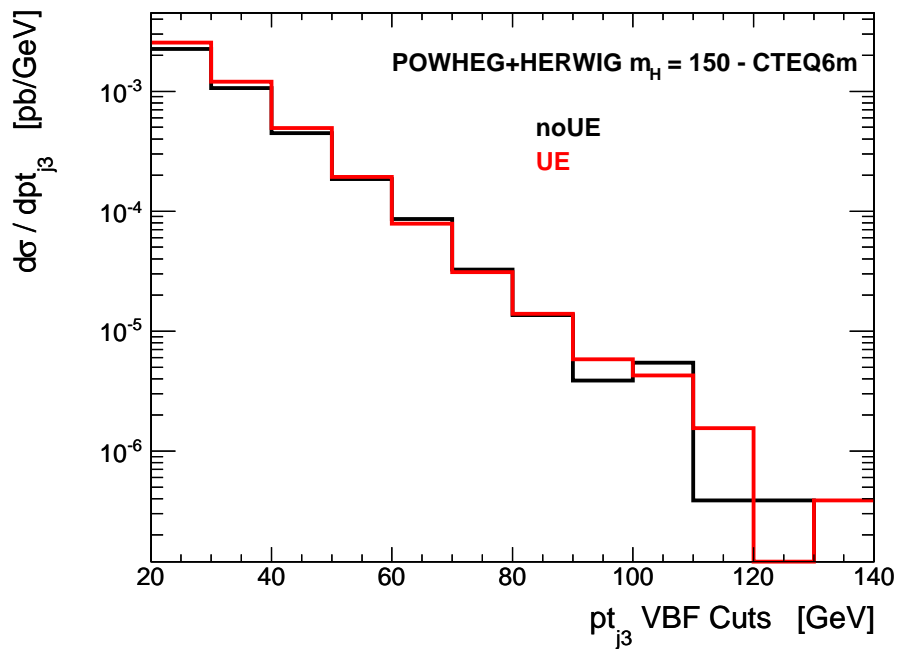
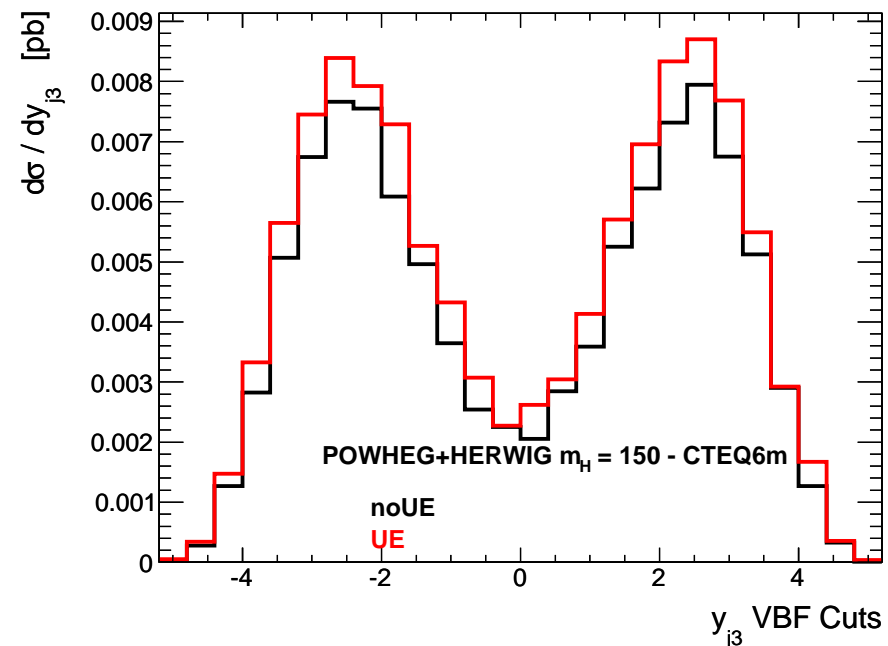
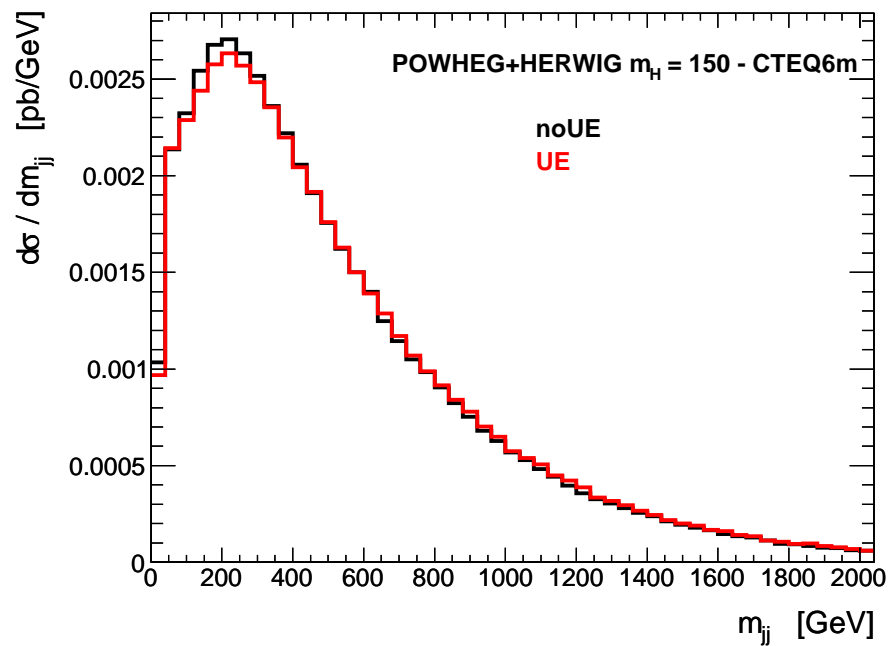
Harder PY jets “mimic” tagging jets. Wrong shapes for  $m_{jj}$  and  $(y_{j1} - y_{j2})$ . Same for  $y_{j3}$ . PY results **rescaled** to the NLO total Xsec



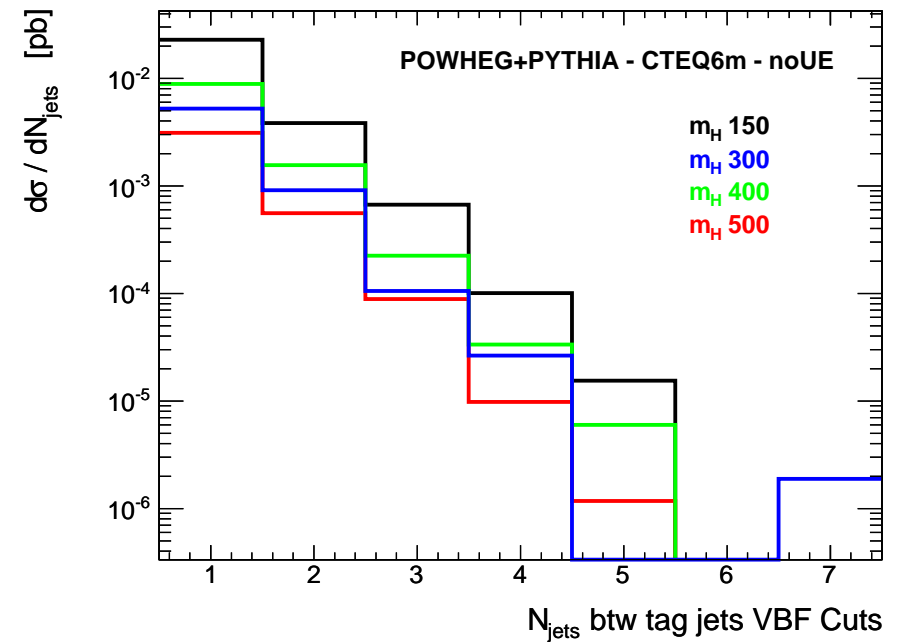
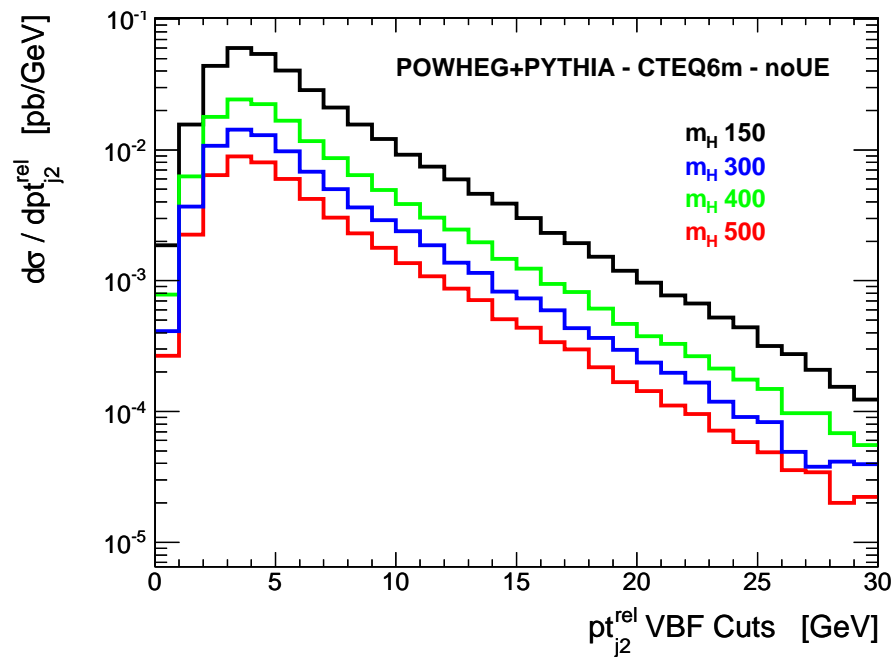
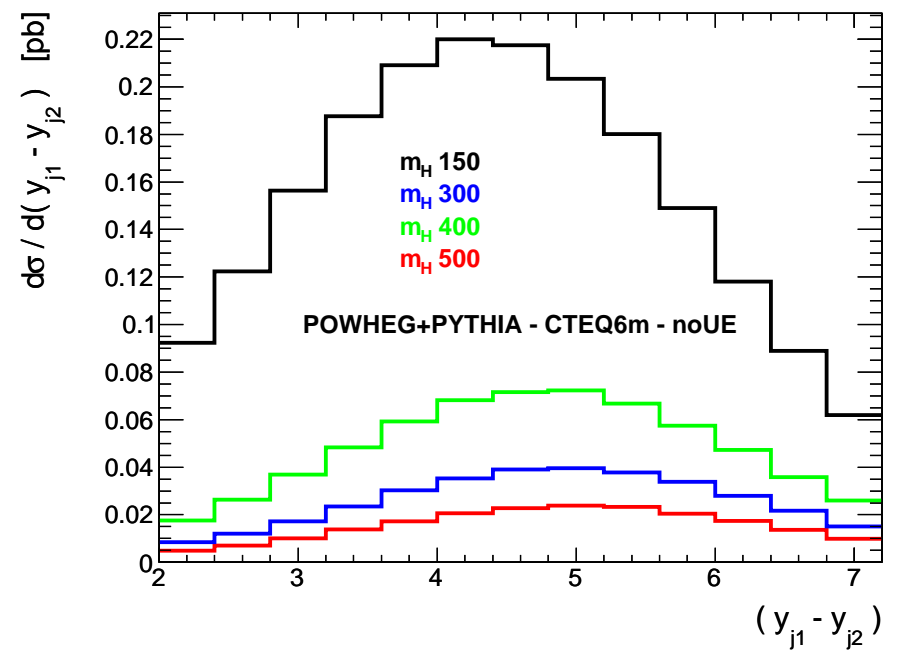
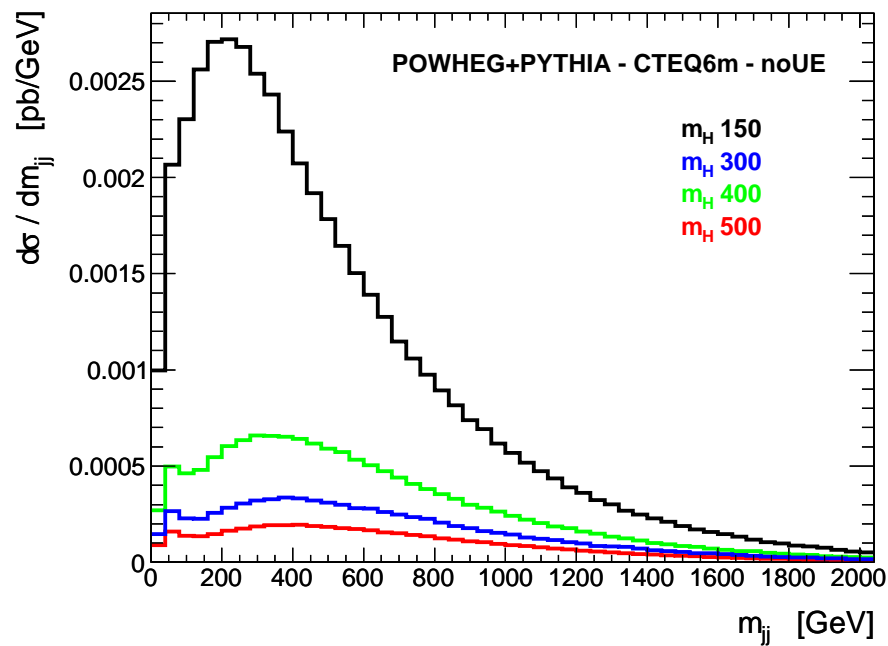
Mild dependence with POWHEG (only 1st emission). When run alone, larger differences with HW than with PY



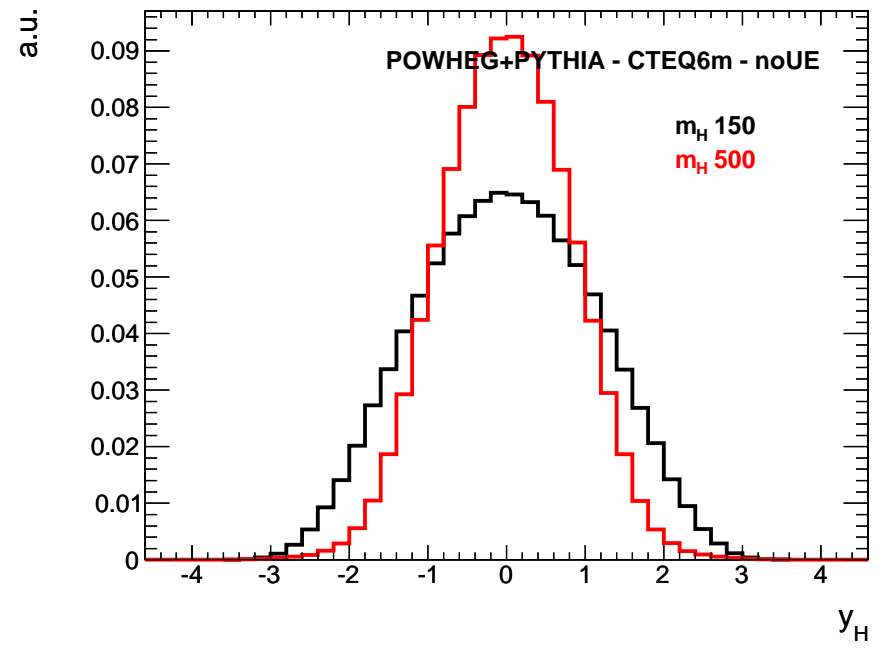
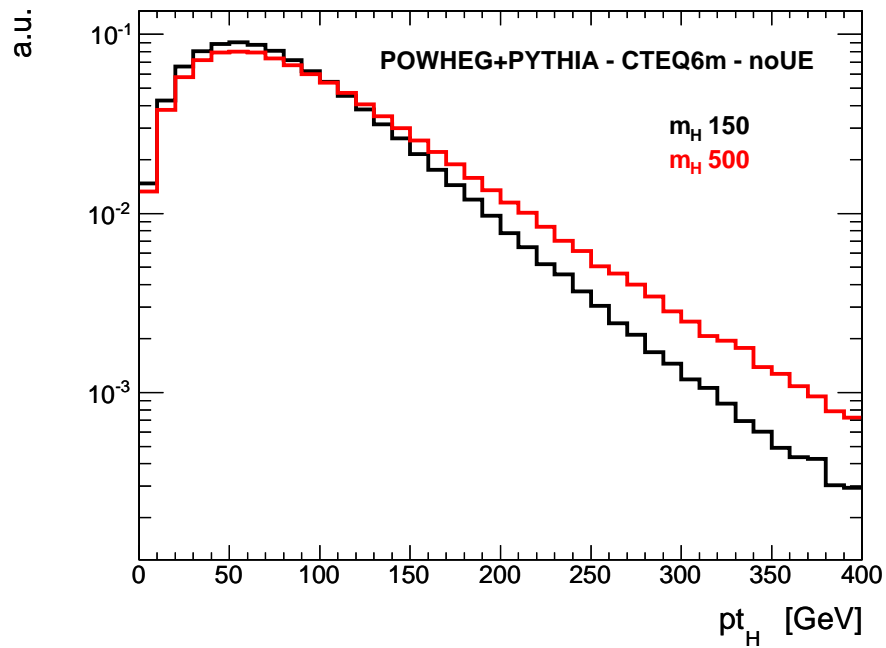
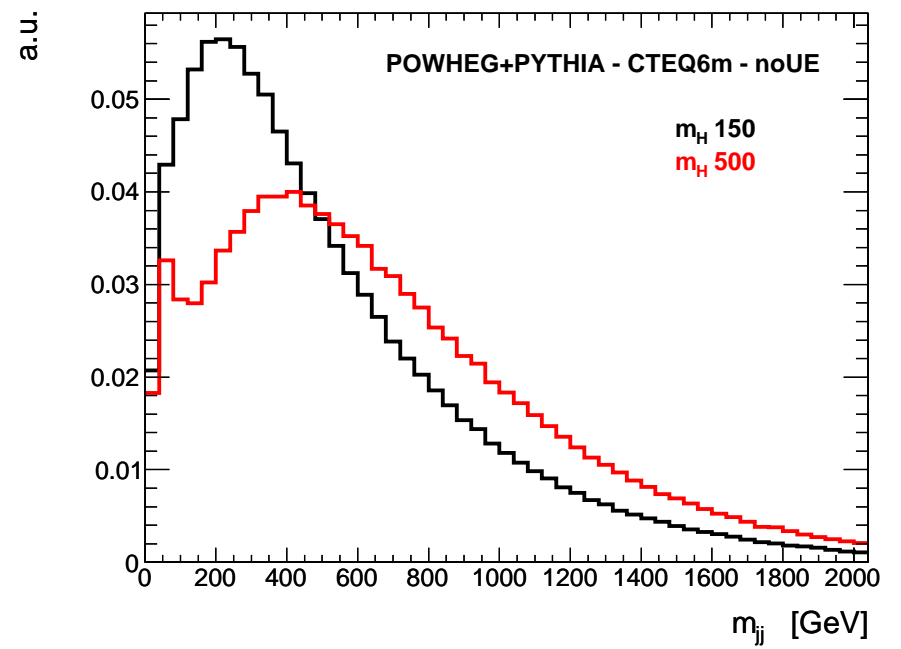
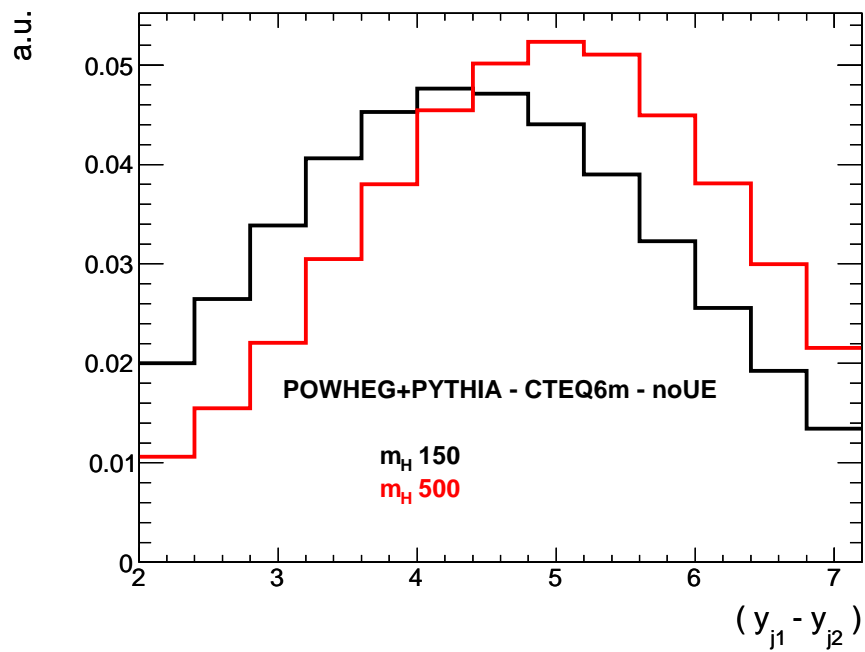
PYTHIA UE + MI events affect appreciably event shapes: low  $m_{jj}$  extra peak,  $y_{j_3}$ ,  $p_T^{rel}$



Do not forget that HERWIG only has soft UE. NO MI



Shift of  $m_{jj}$  and  $(y_{j1} - y_{j2})$  peak to higher values. Need to compare with  $ggH$  for new VBF cuts.  $H$  in VBF less promising



Plots normalized to 1 to show differences in shape



## Higgs boson in VBF

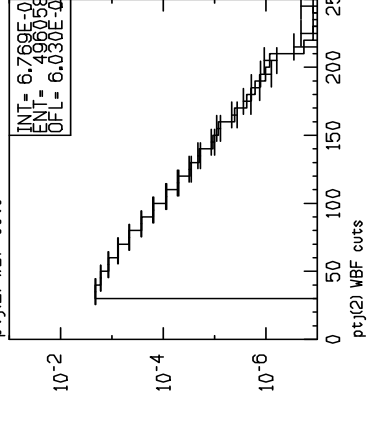
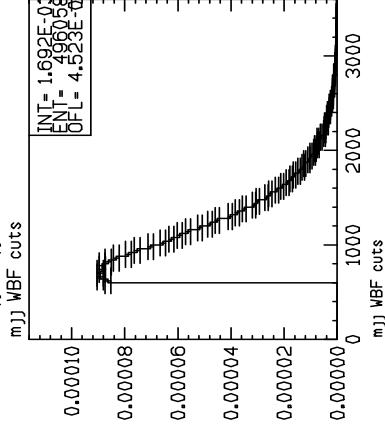
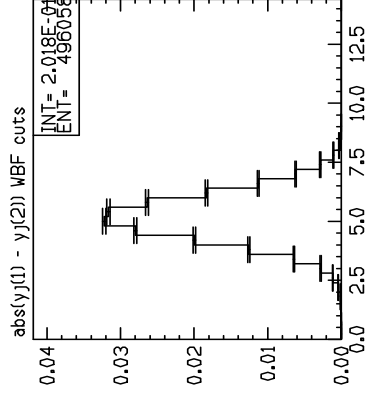
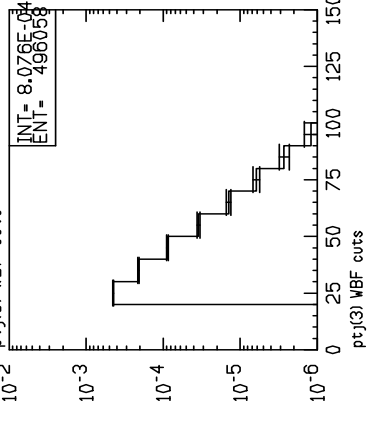
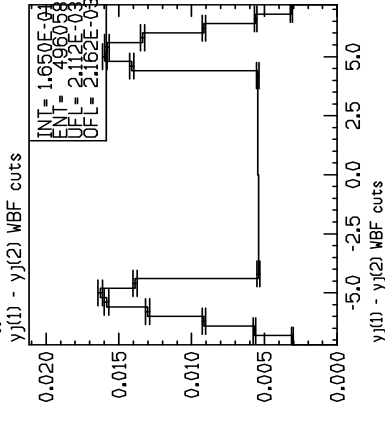
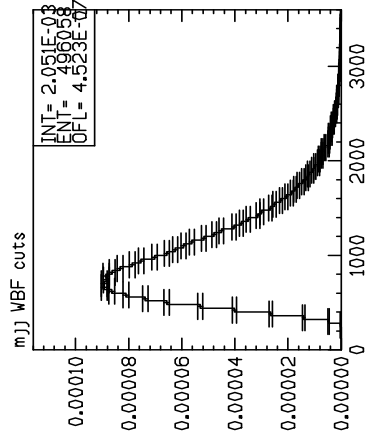
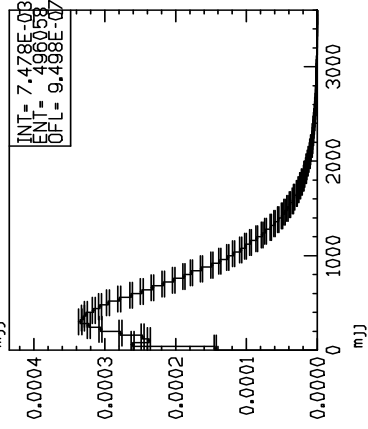
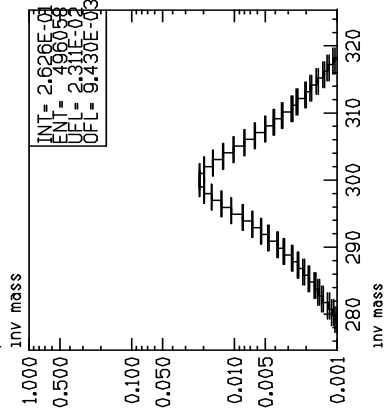
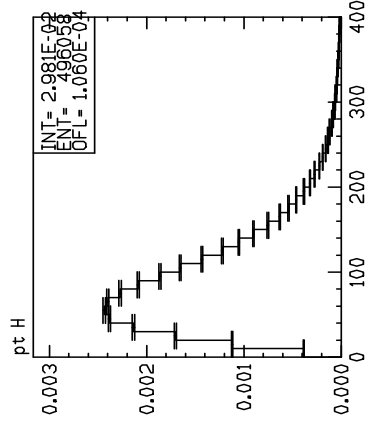
Plots for all the runs in

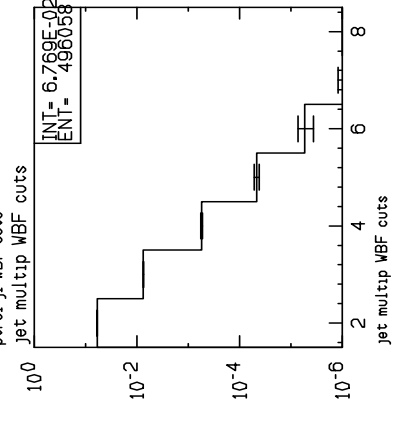
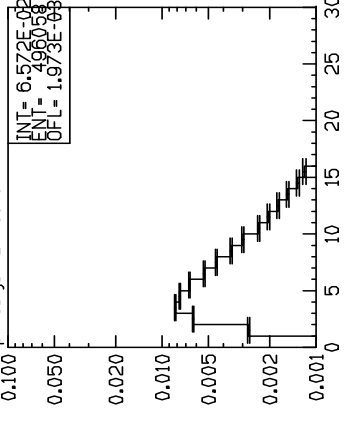
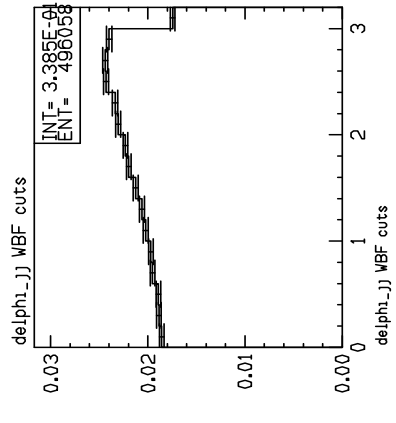
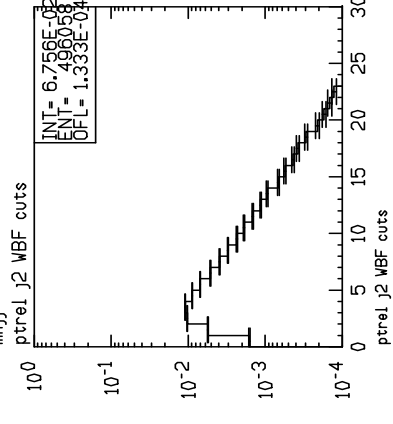
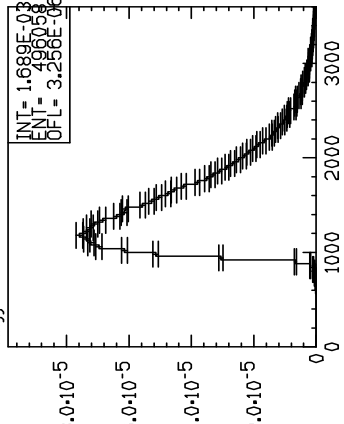
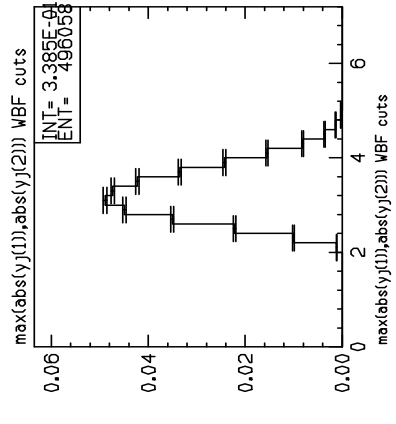
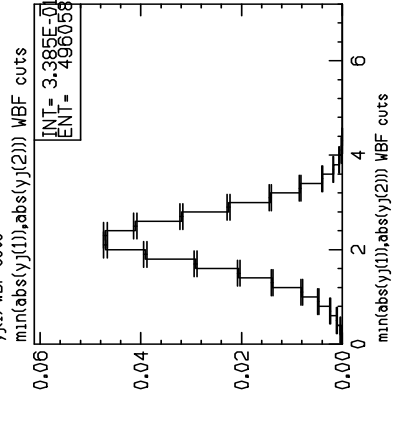
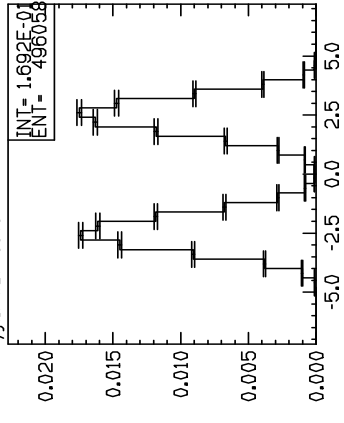
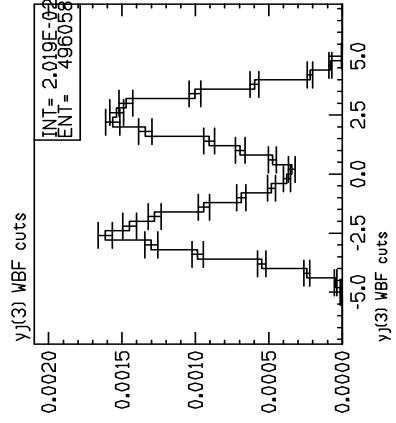
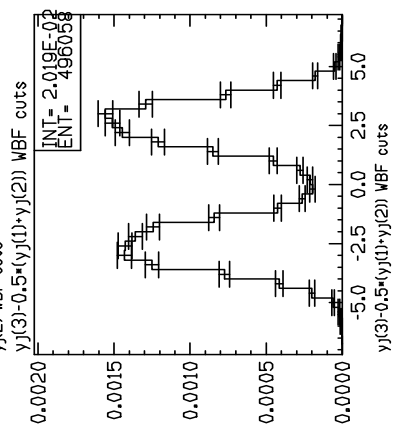
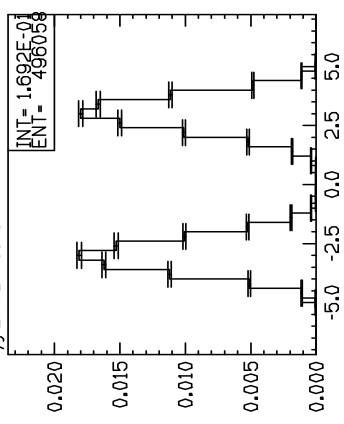
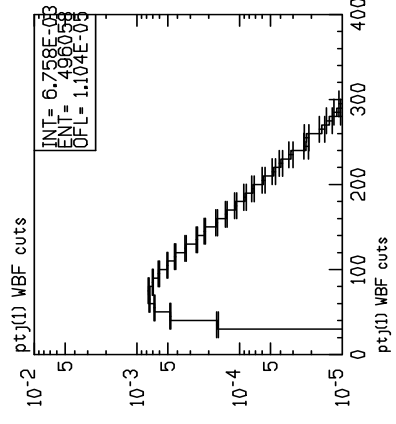
[http://virgilio.mib.infn.it/~oleari/public/VBF\\_H\\_Plots.tgz](http://virgilio.mib.infn.it/~oleari/public/VBF_H_Plots.tgz)

In the following, an example of the output result file (in topdrawer format) that can be found there for

POWHEG + HERWIG, no UE, MSTW2008NLO pdf,  $m_H = 300$  GeV

with error bars (HERWIG-lhef\_noUE\_MSTW2008NLO\_H300.top)





## Conclusions

- ✓ For the **experimentalists**: please use the most up-to-date tools. SMC programs, if not merged with hard matrix elements (ME), cannot describe hard jets beyond the LO ones.
- ✓ For the **theorists**: if you are interested in turning your “favorite” NLO calculation into POWHEG, an easy-to-use tool is available: the **POWHEG BOX**.

This is a **computer framework**, presented in [Alioli, Nason, Oleari and Re, arXiv: 1002.2581], that implements in practice the theoretical construction of the **POWHEG** formalism, for **generic NLO processes**, according to the general formulation of POWHEG. The only ingredients to be provided are: **lists of the Born and real processes**, the **Born phase space**, the **Born squared amplitudes**, the **color-correlated** and **spin-correlated** amplitudes, the **real squared amplitude** and the finite part of the **virtual** corrections.

- ✓ **Many thanks** to Andrea Benaglia, who has run the POWHEG BOX code for Higgs boson production in VBF and has produced all the plots.