

Proton decay in the minimal realistic $SO(10)$ GUT

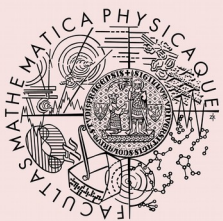
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SO(10)

(non-SUSY minimal renormalizable)

- Embedding of the Standard model

$$SU(3)_c \times SU(2)_L \times U(1)_Y \in SO(10)$$

- Charge quantization
 - Anomaly cancellation
 - Existence of the monopole
- All fermion fields from one generation + right-handed neutrino N_R in 16_F .

$$16_F = L_L \oplus \bar{d}_L \oplus Q_L \oplus \bar{u}_L \oplus \bar{e}_L \oplus N_L^c$$

SO(10)

(non-SUSY minimal renormalizable)

- Gauge fields in 45_G , it accommodates SM gauge fields + leptoquarks, diquarks (carry colour and flavour) with masses \sim symmetry breaking scale

$$45_G = G_\mu^b \oplus A_\mu^a \oplus B_\mu, Y_\mu \oplus (3, 1, \frac{2}{3}) \oplus (3, 2, -\frac{5}{6}) \\ \oplus (3, 2, \frac{1}{6}) \oplus (1, 1, 1) + h.c.$$

Mediate proton decay

SO(10)

(non-SUSY minimal renormalizable)

- Scalar sector is richer in comparison with the SM. In the minimal realistic setting it contains:
 - 45_S (SO(10) breaking, preserves rank)
 - 126_S (intermediate symmetry breaking, breaks rank, renormalizable Yukawa interaction)
 - 10_S (SM breaking, realistic fermionic spectrum)

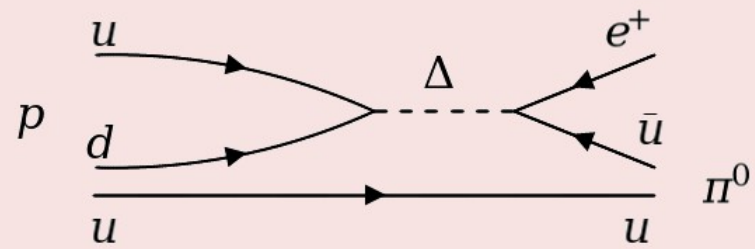
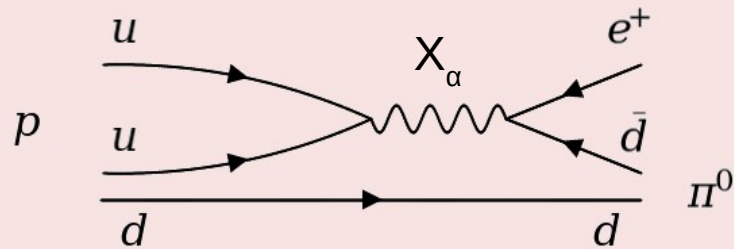
Neutrino mass generation – type I (N_R) and type II (triplet in 126_S)



Baryon number violation

Gauge & scalar leptoquarks lead to baryon number violation.

The most prominent BNV process is **proton decay**.

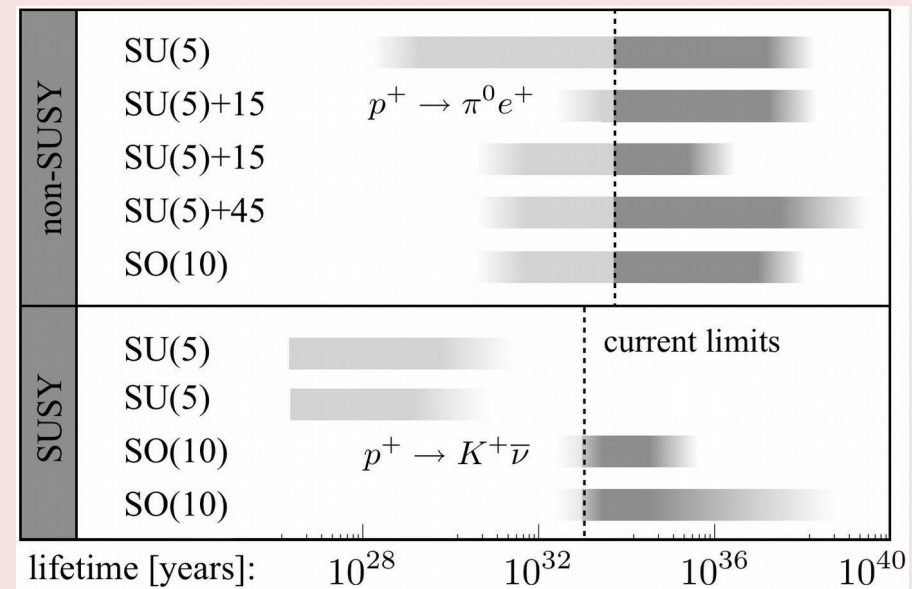


Under control only recently

$$\Gamma(p \rightarrow \pi^0 e^+) \propto |\langle \text{had.} \rangle|^2 \times \left(\frac{\alpha_G^2}{M_X^4} \right) \times$$

flavour \times eff. operator running

Proton lifetime prediction in the non-SUSY minimal renormalizable SO(10) is **robust with respect to the theoretical uncertainties**.

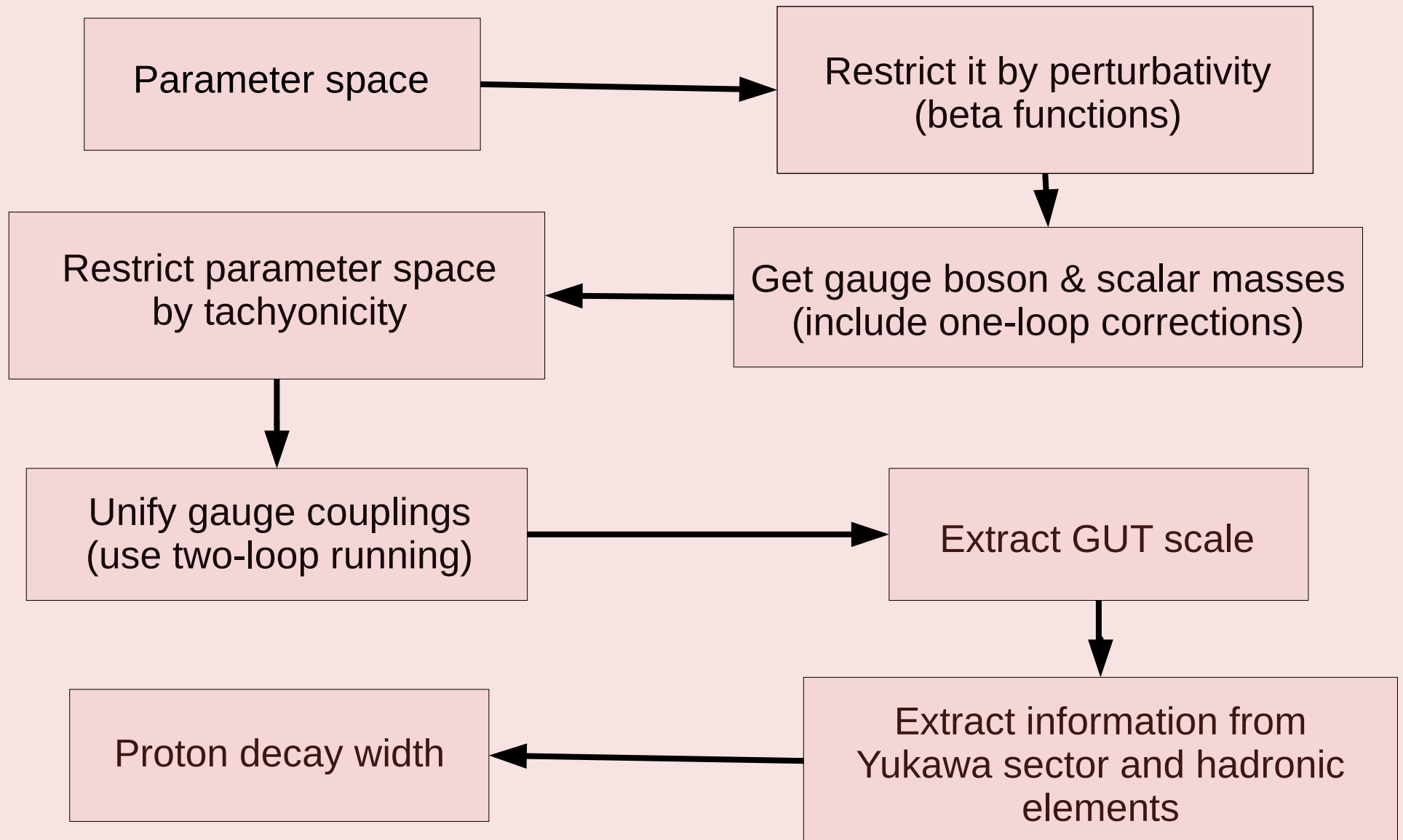


[Kolesova, Malinsky: Proceedings ICHEP 2014]

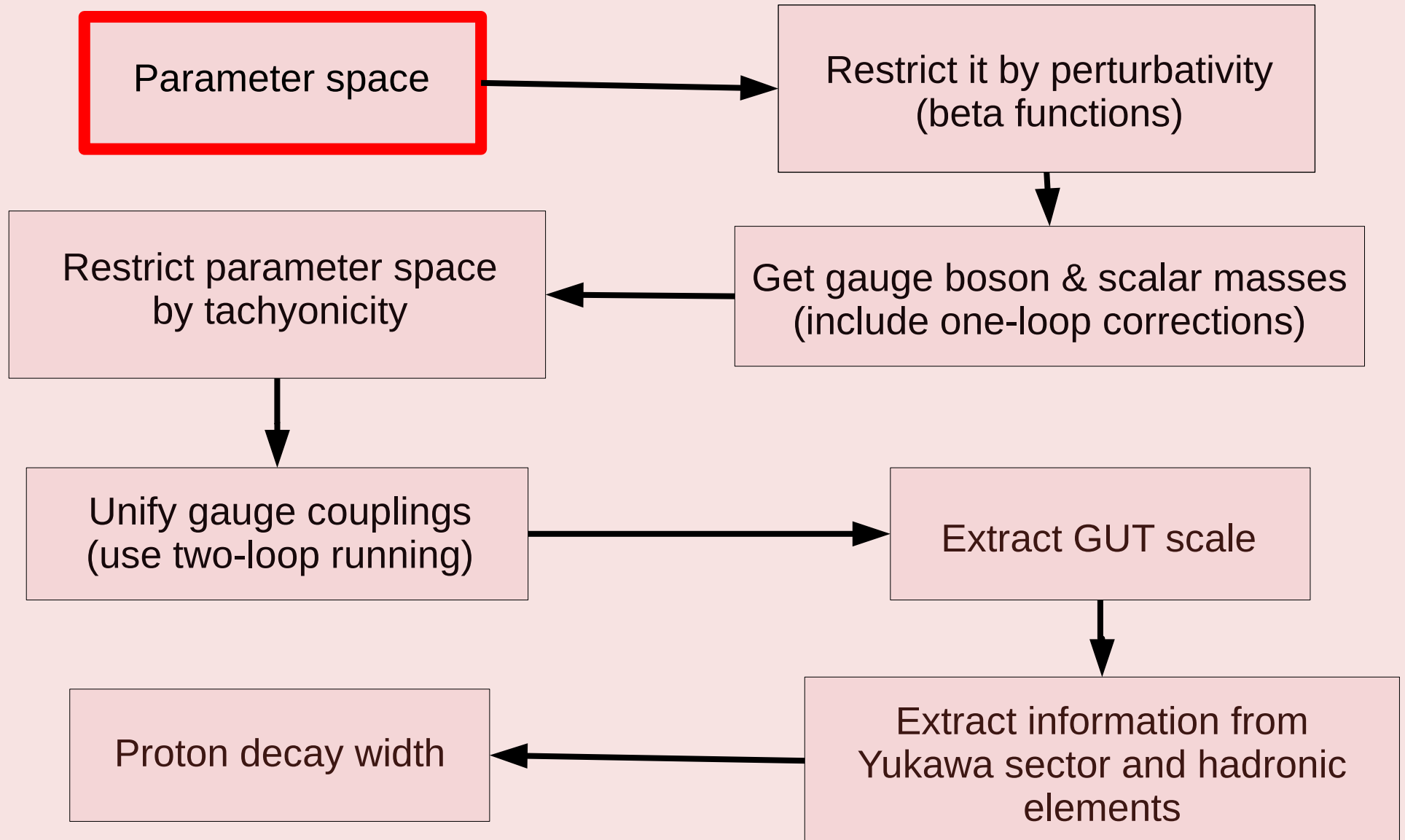
[Kolesova, Malinsky, Mede: AIP Conf. Proc. 2016]

Proton lifetime calculation

Proton decay cookbook



Proton decay cookbook



Parameter space

Scalar potential

$$\begin{aligned}
 V_0(\phi_{45}, \Sigma_{126}, \Sigma_{126}^*) = & -\frac{\mu^2}{4}(\phi\phi)_0 + \frac{a_0}{4}(\phi\phi)_0(\phi\phi)_0 + \frac{a_2}{4}(\phi\phi)_2(\phi\phi)_2 - \\
 & -\frac{\nu^2}{5!}(\Sigma\Sigma^*)_0 + \frac{\lambda_0}{(5!)^2}(\Sigma\Sigma^*)_0(\Sigma\Sigma^*)_0 + \frac{\lambda_2}{(4!)^2}(\Sigma\Sigma^*)_2(\Sigma\Sigma^*)_2 + \\
 & + \frac{\lambda_4}{(3!)^2(2!)^2}(\Sigma\Sigma^*)_4(\Sigma\Sigma^*)_4 + \frac{\lambda'_4}{(3!)^2}(\Sigma\Sigma^*)_{4'}(\Sigma\Sigma^*)_{4'} + \\
 & + \frac{\eta_2}{(4!)^2}(\Sigma\Sigma)_2(\Sigma\Sigma)_2 + \frac{\eta_2^*}{(4!)^2}(\Sigma^*\Sigma^*)_2(\Sigma^*\Sigma^*)_2 + \\
 & + \frac{i\tau}{4!}(\phi)_2(\Sigma\Sigma^*)_2 + \frac{\alpha}{2 \cdot 5!}(\phi\phi)_0(\Sigma\Sigma^*)_0 + \frac{\beta_4}{4 \cdot 3!}(\phi\phi)_4(\Sigma\Sigma^*)_4 + \\
 & + \frac{\beta'_4}{3!}(\phi\phi)'_4(\Sigma\Sigma^*)'_4 + \frac{\gamma_2}{4!}(\phi\phi)_2(\Sigma\Sigma)_2 + \frac{\gamma_2^*}{4!}(\phi\phi)_2(\Sigma^*\Sigma^*)_2
 \end{aligned}$$

where $\{\mu, \nu, \tau\} \leftrightarrow \text{VEVs } \{\sigma, \omega_{BL}, \omega_R\}$ and gauge coupling g

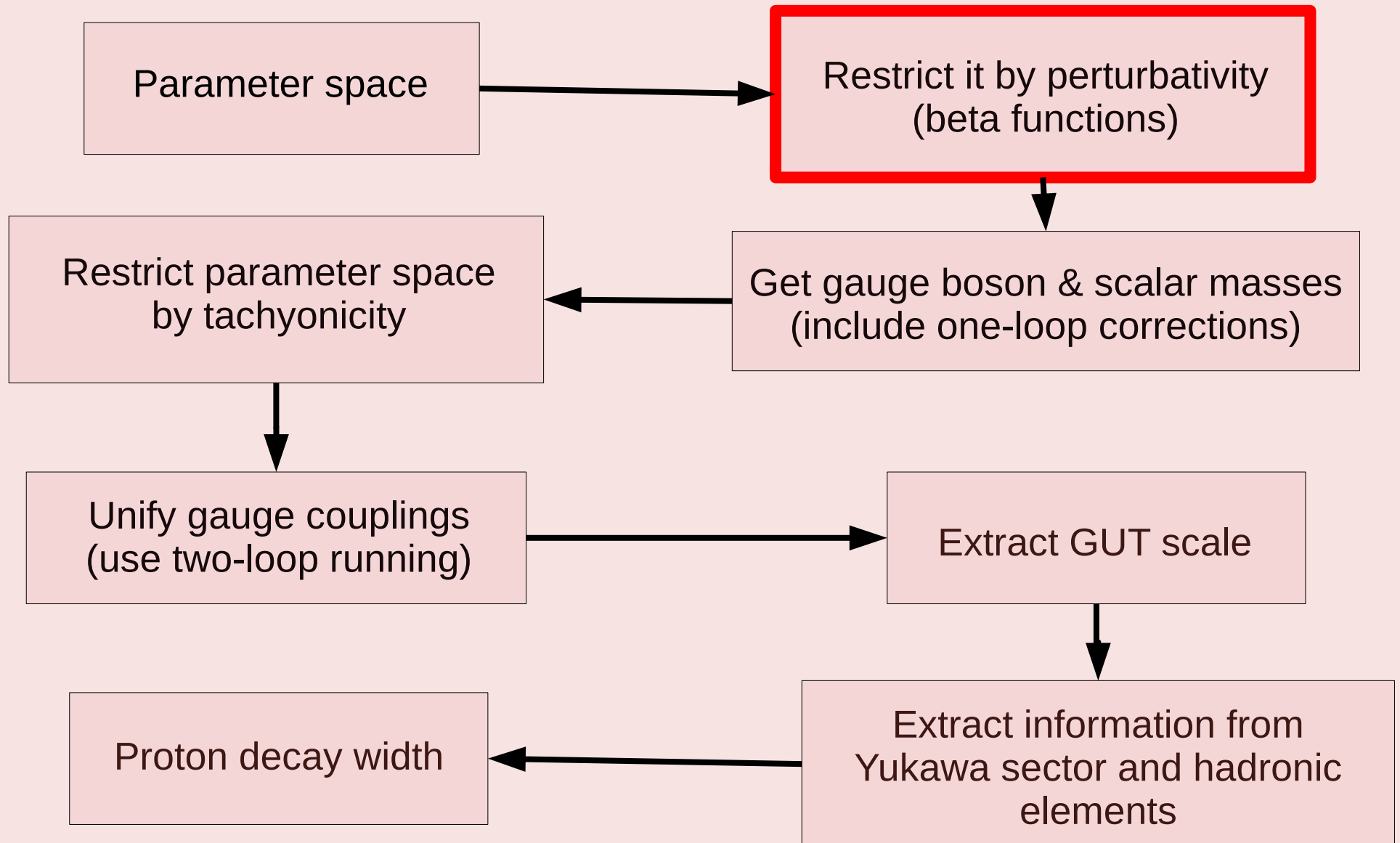
126_s , seesaw scale

45_s , GUT scale

17 degrees of freedom

excluding Yukawa sector

Proton decay cookbook



Perturbativity constraints

Vacuum stability:

- One-loop vacuum is not far from the tree-level vacuum (checked for different renorm. scales ← one-loop beta functions)

- Relation $\tau = P(\sigma, \omega_{BL}, \omega_R, \lambda) + Q(\lambda) \frac{\omega_{BL}\omega_R(\omega_{BL} + \omega_R)}{|\sigma|^2}$ $\leq \omega_{\max}$
P,Q ... polynomials

distinct cases: "almost one step breaking"

"(flipped) SU(5) interm. stage"

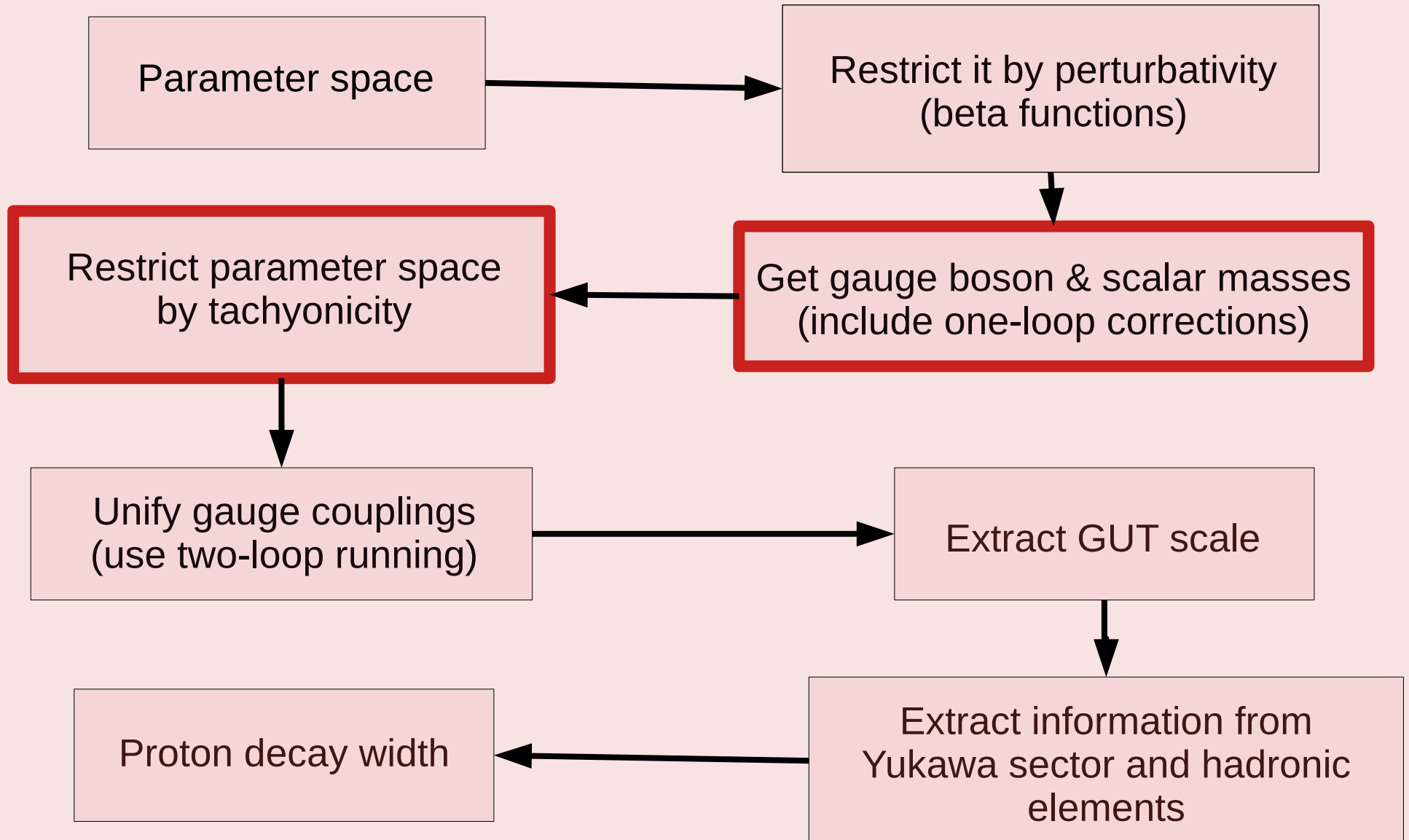
"SU(4)_c x SU(2)_L x U(1)_R interm. stage"

"SU(3)_c x SU(2)_L x SU(2)_R x U(1)_{B-L} interm. stage"

Mass stability:

- We require the maximal one-loop scalar mass correction to be smaller than the averaged tree-level mass

Proton decay cookbook



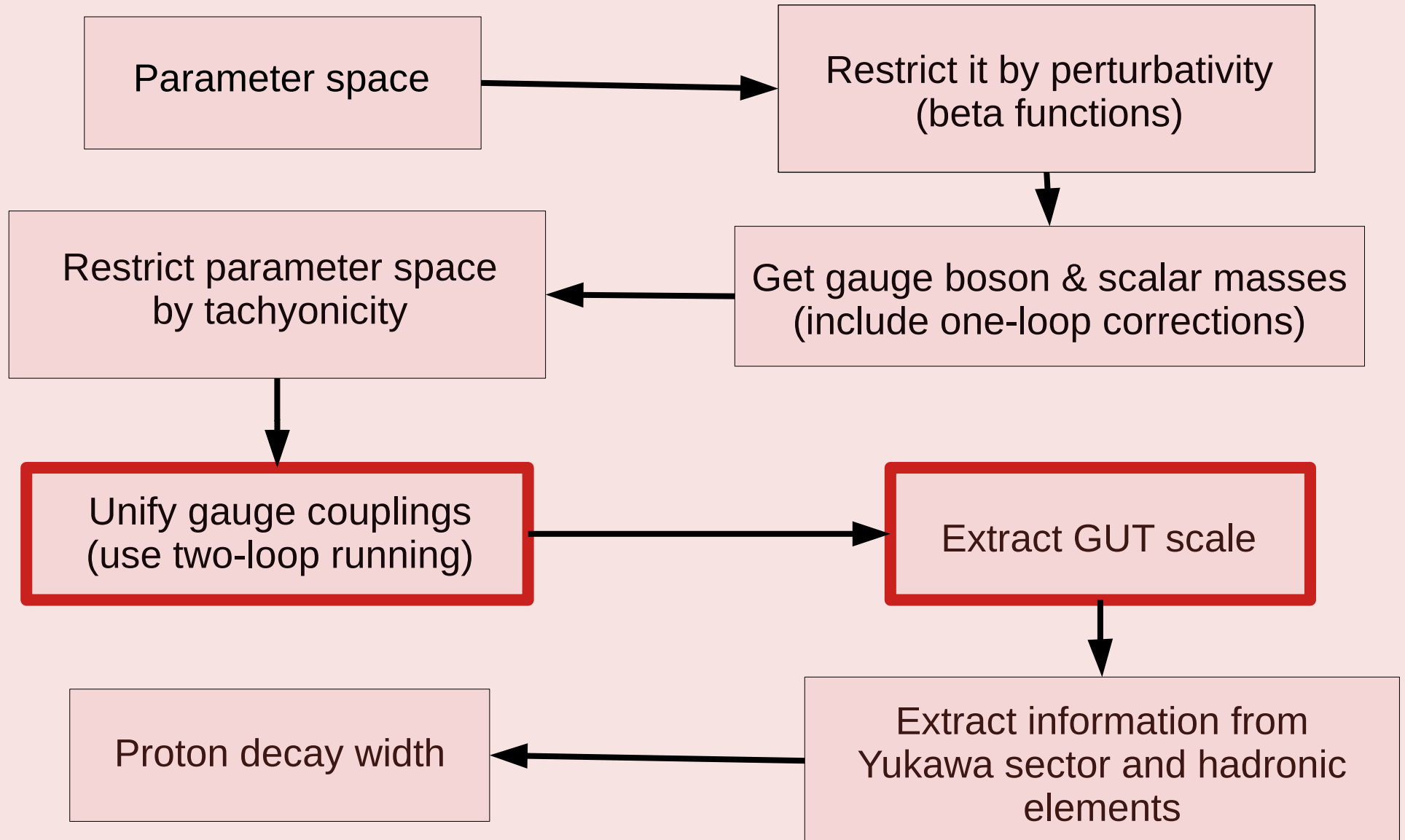
Tachyonicity constraints

The existence of non-tachyonic spectrum is the **crucial requirement** restricting viable parameter space. On top of that the realistic scenarios involve accidentally light pseudo-Goldstone scalars => one-loop calculations are necessary => inherently quantum model.



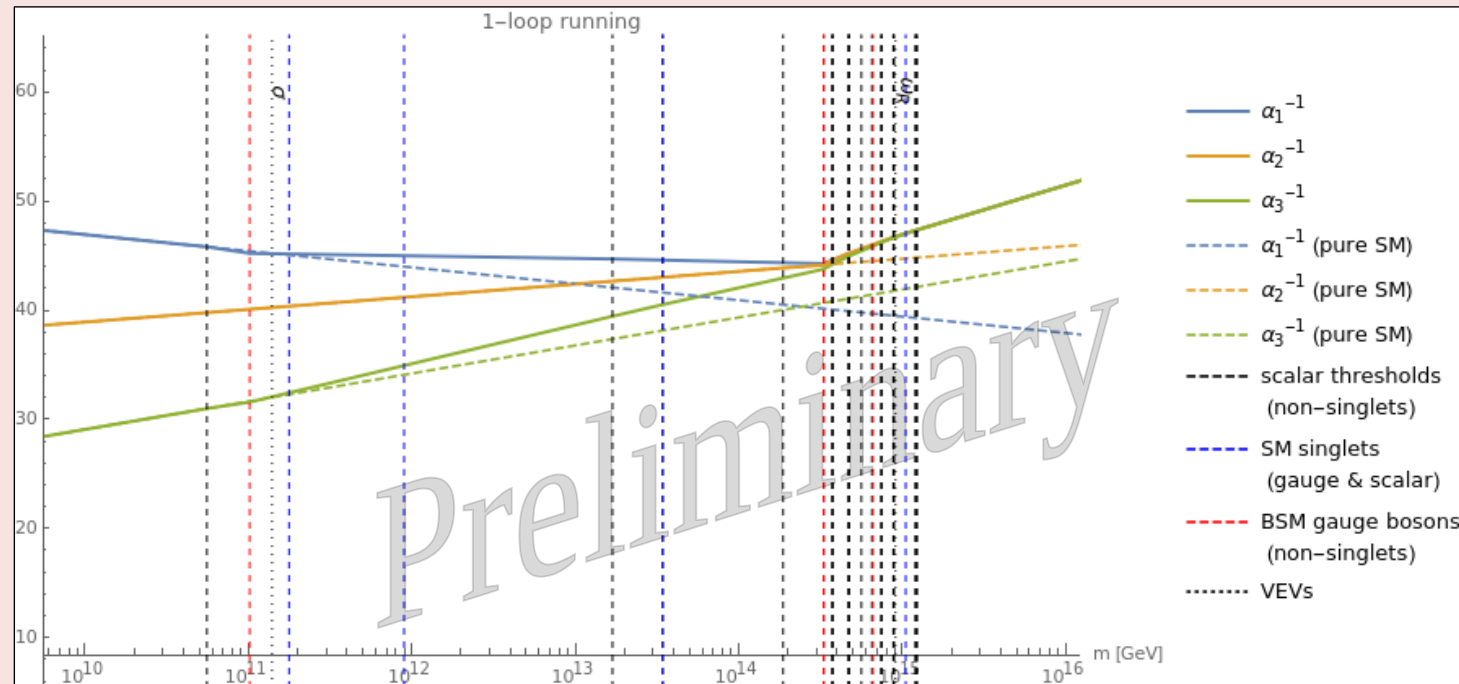
Complete calculation of full one-loop effective scalar mass corrections is implemented.

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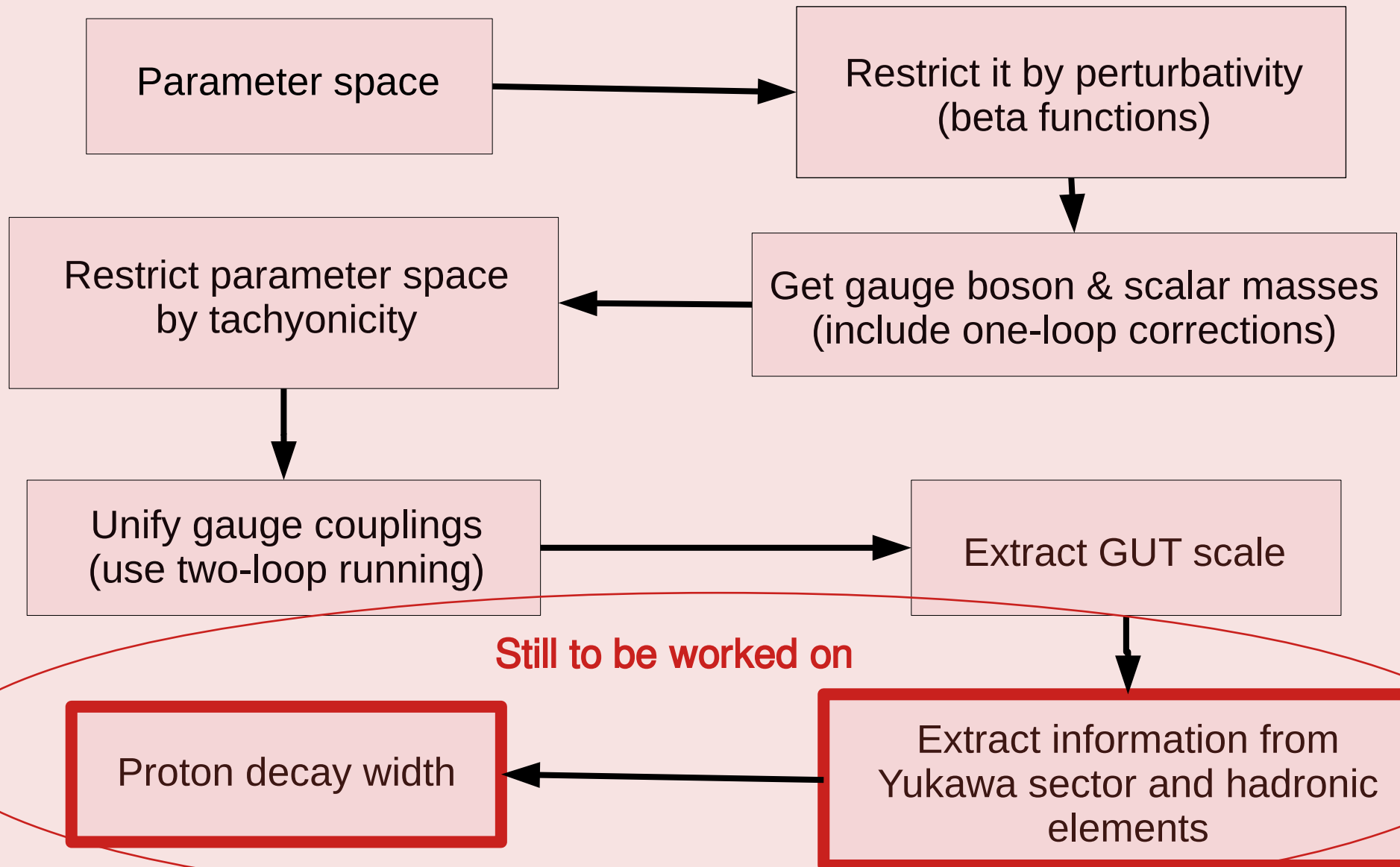


Gauge coupling unification

Robust proton lifetime estimate has to involve **two-loop gauge coupling running calculation** to extract the scale of unification. Successful unification also serves as a criterion to narrow down viable parameter space.



Proton decay cookbook



Summary

Minimal renormalizable non-SUSY **SO(10)**:

- Gauge bosons in 45_S
- All fermions from one generation (+ N_R) in 16_F
- Scalars in 45_S , 126_S (+ 10_S)

Proton lifetime estimate is robust with respect to theoretical uncertainties.

The **model analysis** involved:

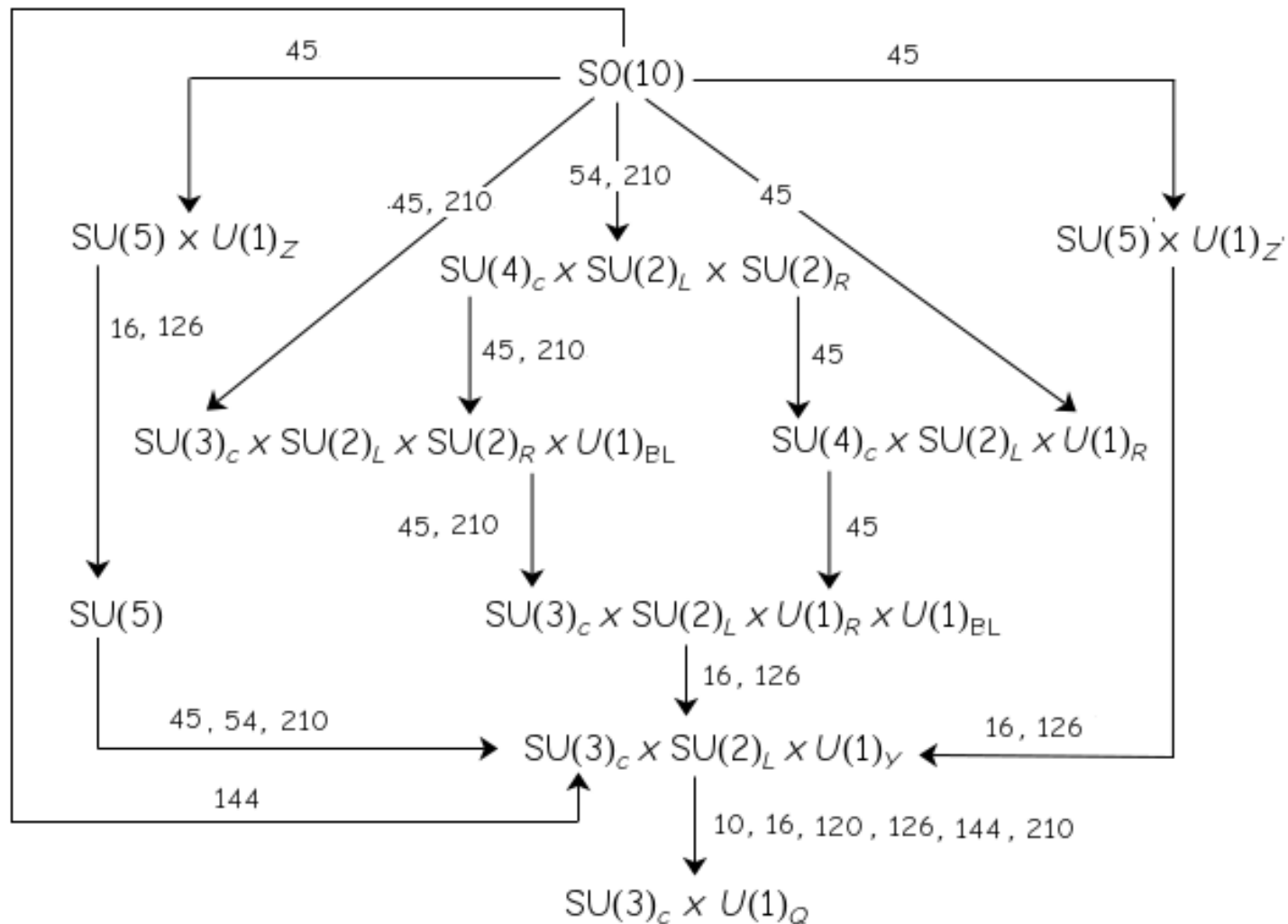
- Improved perturbativity considerations
- One-loop scalar masses calculation
- Two-loop gauge unification

Paper in preparation

Additional slides

SO(10) breaking chains

(non-SUSY minimal renormalizable)



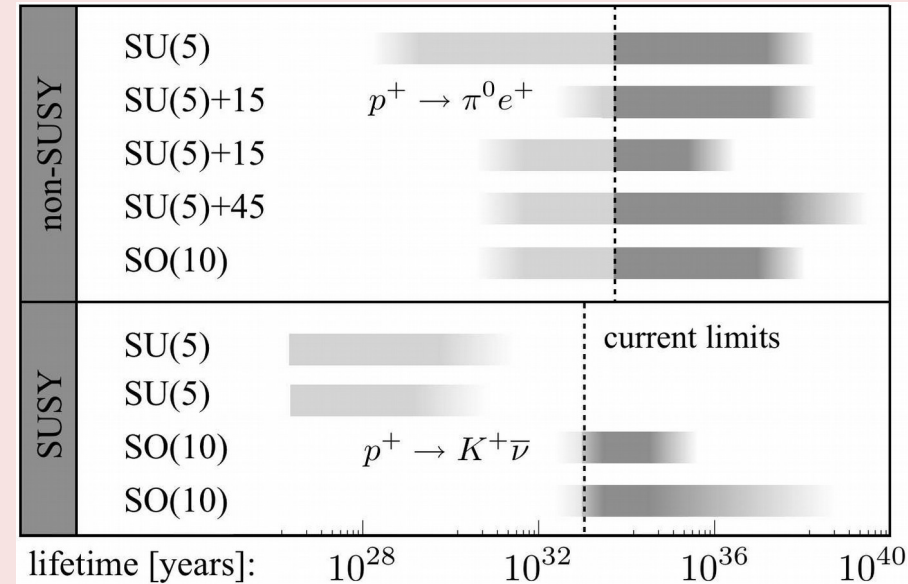
Theoretical uncertainties

Reducible:

- Finite order in perturbation theory

Irreducible:

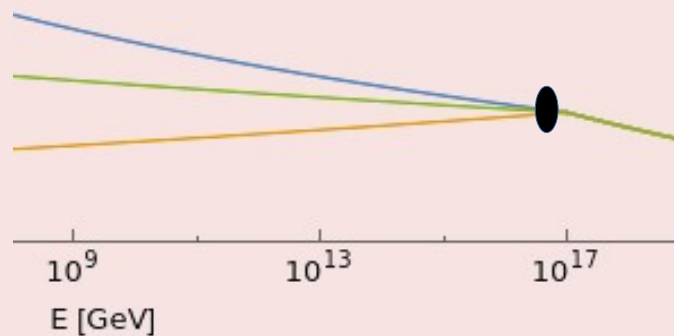
- Hadronic matrix elements
- Planck scale physics (dim 5 operators)



[Kolesova, Malinsky: Proceedings ICHEP 2014]

$$\frac{c}{M_{Pl}} F^{\mu\nu} \Phi F_{\mu\nu} \rightarrow c \frac{M_{GUT}}{M_{Pl}} F^{\mu\nu} F_{\mu\nu}$$

SO(10)



$$\frac{c}{M_{Pl}} f_i f_j H \Phi \rightarrow c \frac{M_{GUT}}{M_{Pl}} f_i f_j H$$

Suppressed in SO(10)

[Kolesova, Malinsky: Phys.Rev.D 99 (2019)]

[Kolesova, Malinsky, Mede: AIP Conf.Proc. 1743 (2016)]