

# PRODUCTION OF THERMAL AXIONS ACROSS THE EWPT

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Instituto de  
Física  
Teórica  
UAM-CSIC



Based on: 2007.06579 and more to appear, FAA, F. D'Eramo, L. Merlo, A. Notari, R. Zambujal Ferreira.

Fernando Arias Aragón  
fernando.arias@uam.es

# Outline

- Motivation
- Setup
- Results
- UV complete models
- Complementarity with XENON1T
- Summary

# Motivation

- Light or massless new DoF appear frequently in BSM

- Axions are particularly well motivated:

- Solution to the Strong CP Problem

Crewther, Di Vecchia, Veneziano & Witten, 1980

$$SM \rightarrow d_n \sim \bar{\theta} \cdot O(10^{-16}) e \cdot cm$$

$$EXP \rightarrow d_n < 0,30 \cdot 10^{-25} e \cdot cm \Rightarrow \bar{\theta} \leq 10^{-10}$$

Baker et al., 0602020 Afach et al., 1509.04411

- Possible DM candidate [D. J. E. Marsh, 1510.07633](#)
- Arise naturally from string theories [Aravnitaki et al., 0905.4720](#)
- **Possible explanation for XENON1T excess!** [XENON Collabotation, 2006.09721](#)

# Motivation

- Many strategies to search for the axion:

- Helioscopes
- Haloscopes
- Light shining through wall
- Dark matter recoil
- Star cooling

- In addition to this, there is a cosmological probe

$$\Delta N_{eff} \simeq 74.85 Y_a^{4/3}$$

- Probed in the future (CMB-S4, LiteBIRD, Simons Observatory)
- $\Delta N_{eff}$  depends on its production, thus on its couplings and its decay constant

# Setup

- General axion interactions

$$\mathcal{L}_{\text{axion-int}} \supset \frac{1}{f_a} \left[ a c_X \frac{\alpha_X}{8\pi} X^{a\mu\nu} \tilde{X}_{\mu\nu}^a + \partial_\mu a c_\psi \bar{\psi} \gamma^\mu \psi \right]$$

$$X = \{G, W, B\}$$
$$\psi = \{Q_L, u_R, d_R, E_L, e_R\}$$

- Three classes of processes

- $X + X \rightarrow X + a$

- $\psi + X \rightarrow \psi + a, \psi + \bar{\psi} \rightarrow X + a$

- $\psi + H \rightarrow \psi + a, \psi + \bar{\psi} \rightarrow H + a$

- Production is efficient when production rate  $\Gamma$  exceeds Hubble rate, with  $H \propto T^2$

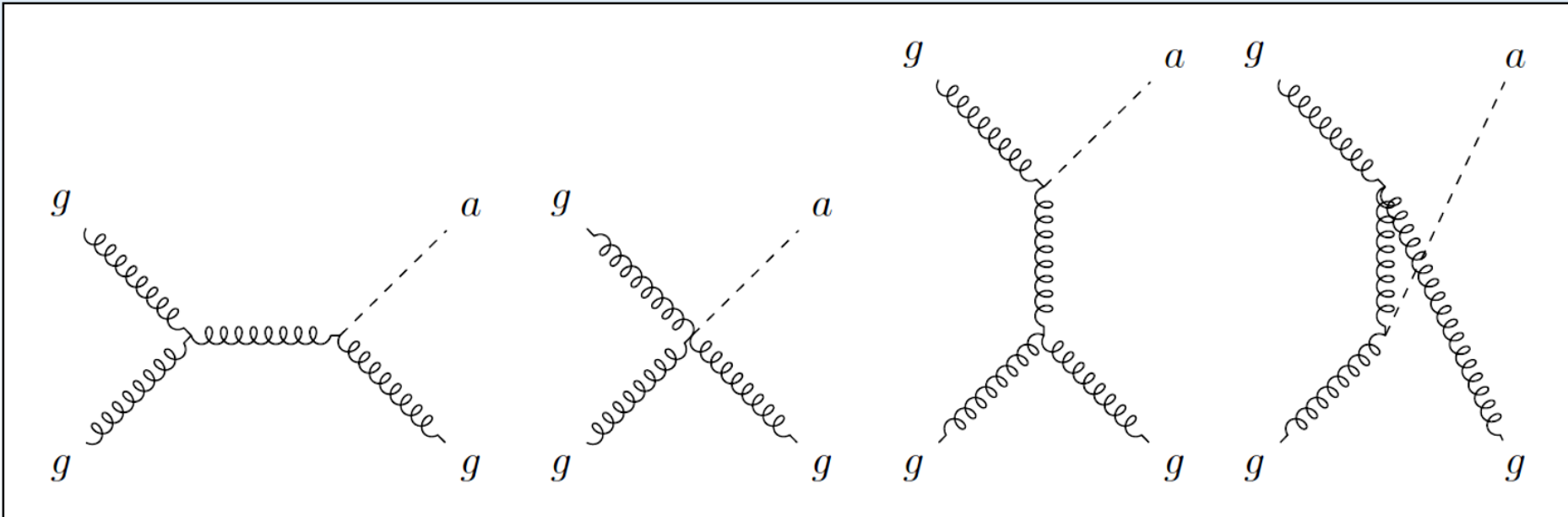
# Setup

- $X + X \rightarrow X + a$

- Existing both above and below the EWPT

- Production rate:  $\Gamma \simeq \alpha_X^3 c_X^2 \frac{T^3}{f_a^2}$

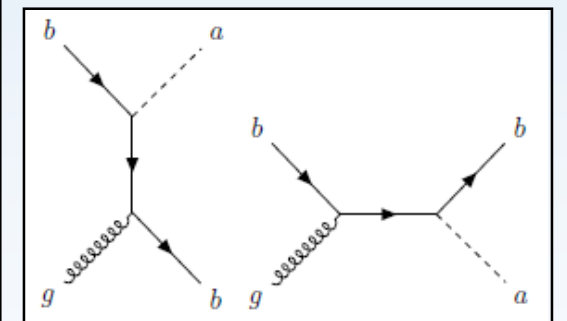
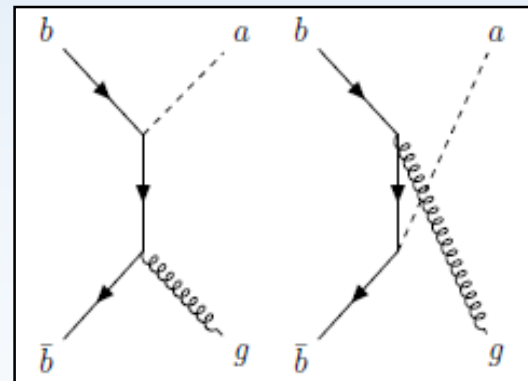
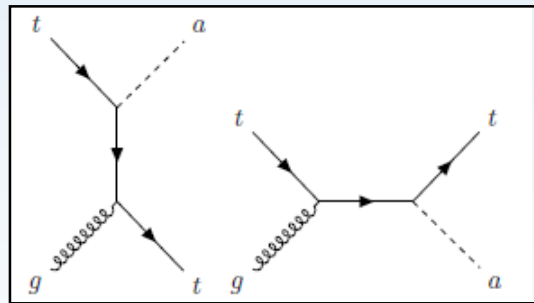
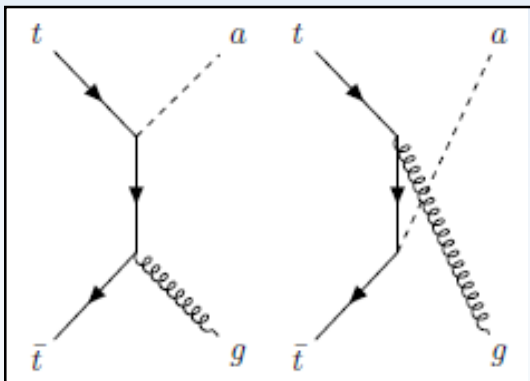
- Gluons dominate. Mostly relevant at high T, where  $\Delta N_{eff} \simeq 0,027$  if there is thermalization



# Setup

- $\psi + X \rightarrow \psi + a, \quad \psi + \bar{\psi} \rightarrow X + a$

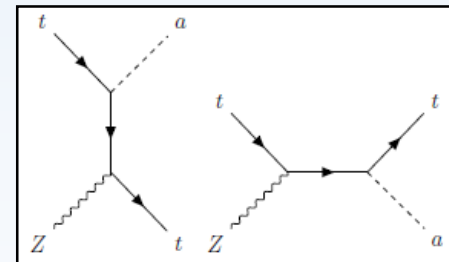
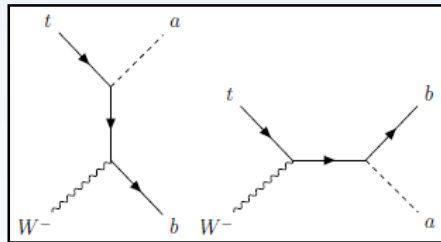
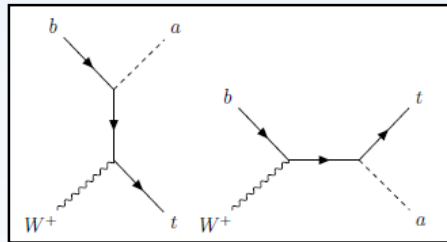
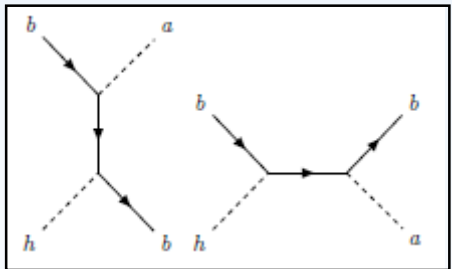
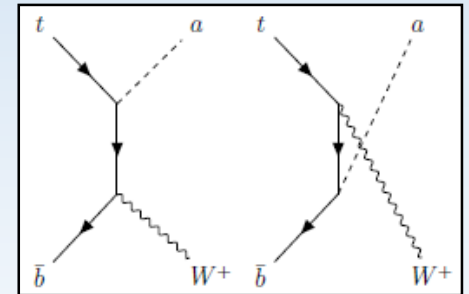
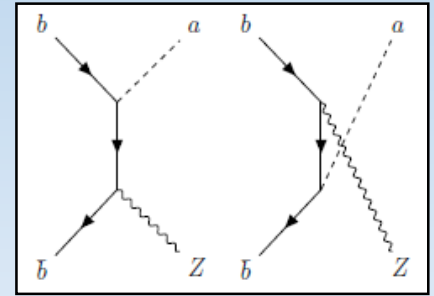
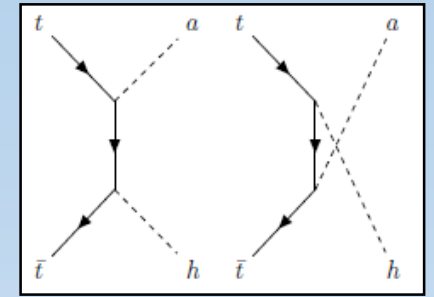
- Existing only below the EWPT
- Boltzmann suppressed for  $T < m_\psi$
- Production rate for  $T > m_\psi: \Gamma \simeq \alpha_X c_\psi^2 m_\psi^2 \frac{T}{f_a^2}$
- Gluons and heavy fermions dominate.  $\Gamma/H$  peaks at  $T \sim m_\psi$



# Setup

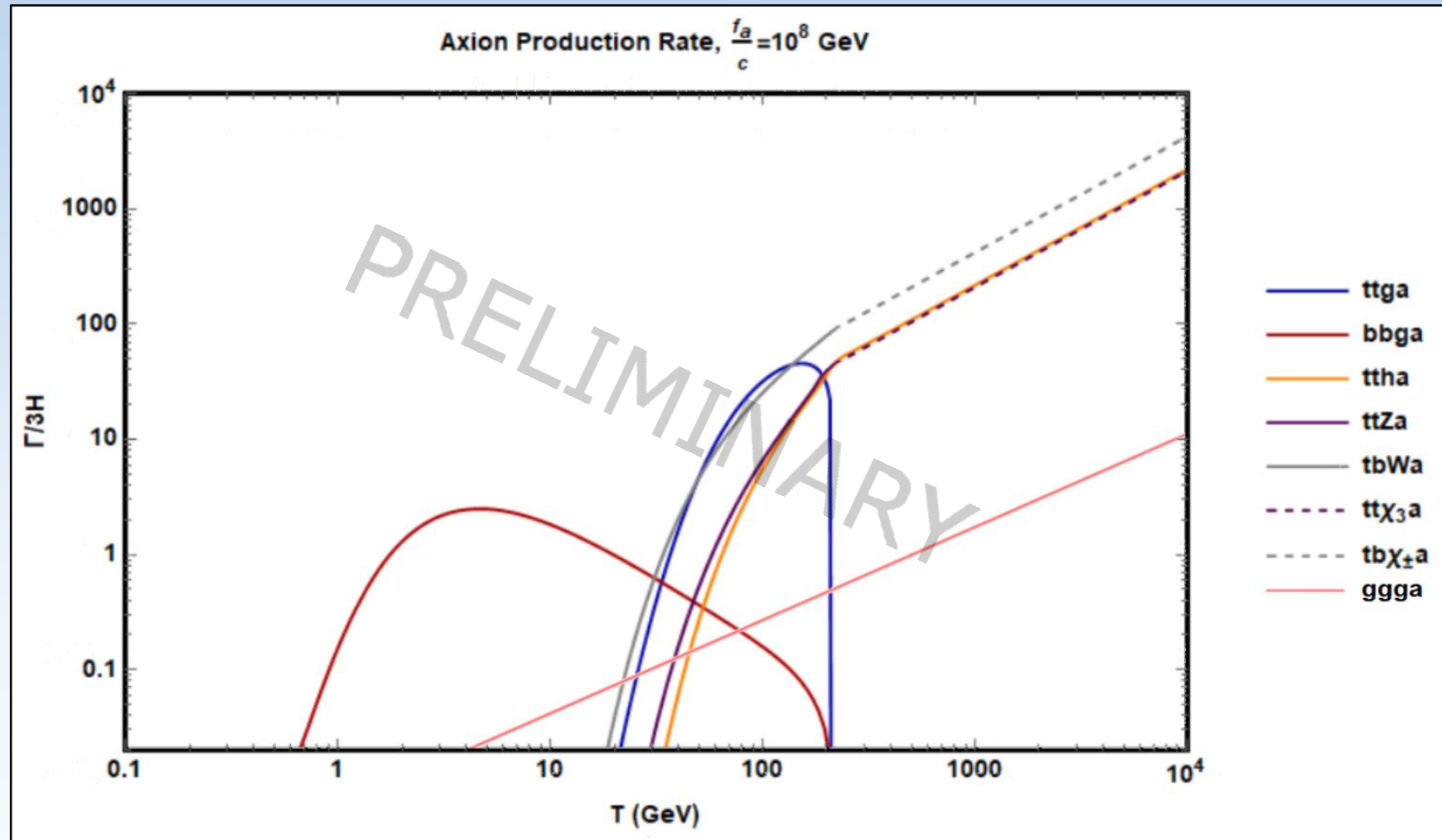
- $\psi + H \rightarrow \psi + a, \psi + \bar{\psi} \rightarrow H + a$

- Above EWSB: any of the four real components of the Higgs doublet
- Below EWSB: physical Higgs  $h$  and the longitudinal components of  $W$  and  $Z$  bosons
- Boltzmann suppressed for  $T < m_\psi$
- Production rate for  $T > m_\psi: \Gamma \simeq c_X^2 y_\psi^2 \frac{T^3}{f_a^2}$
- Heavy fermions dominate. Can thermalize at high  $T$ .





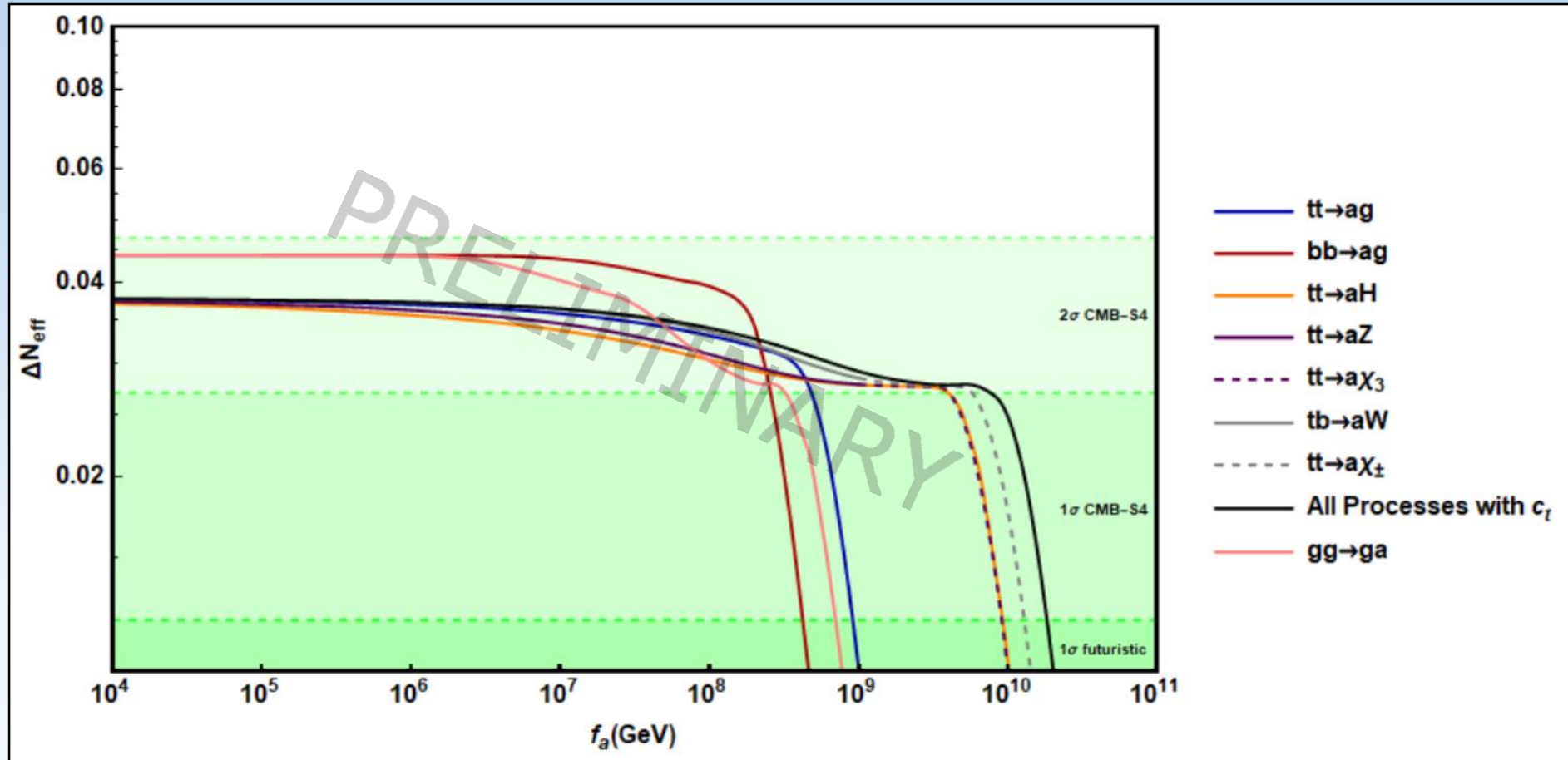
# Setup



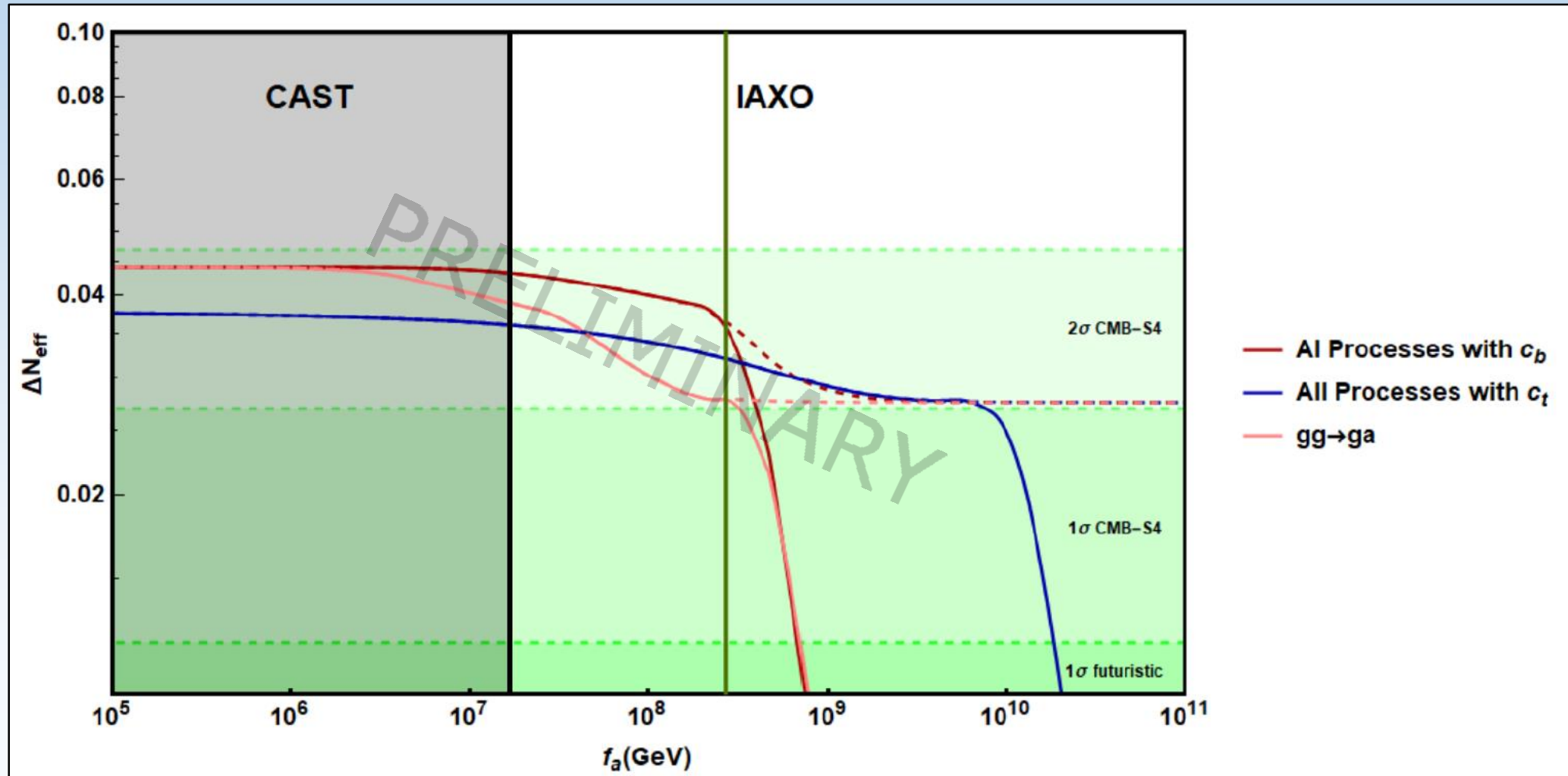
# Results

- Model independent approach
- Operator-by-operator analysis: assume only one axion coupling at a time
- Solve the Boltzmann equation and find  $\Delta N_{eff}$  for each set of processes
- Stop the equation at 1 GeV to keep strong interactions perturbative
- Result sensitive to Early Universe conditions:
  - Initial condition: zero or thermal abundance
  - Reheating temperature

# Results



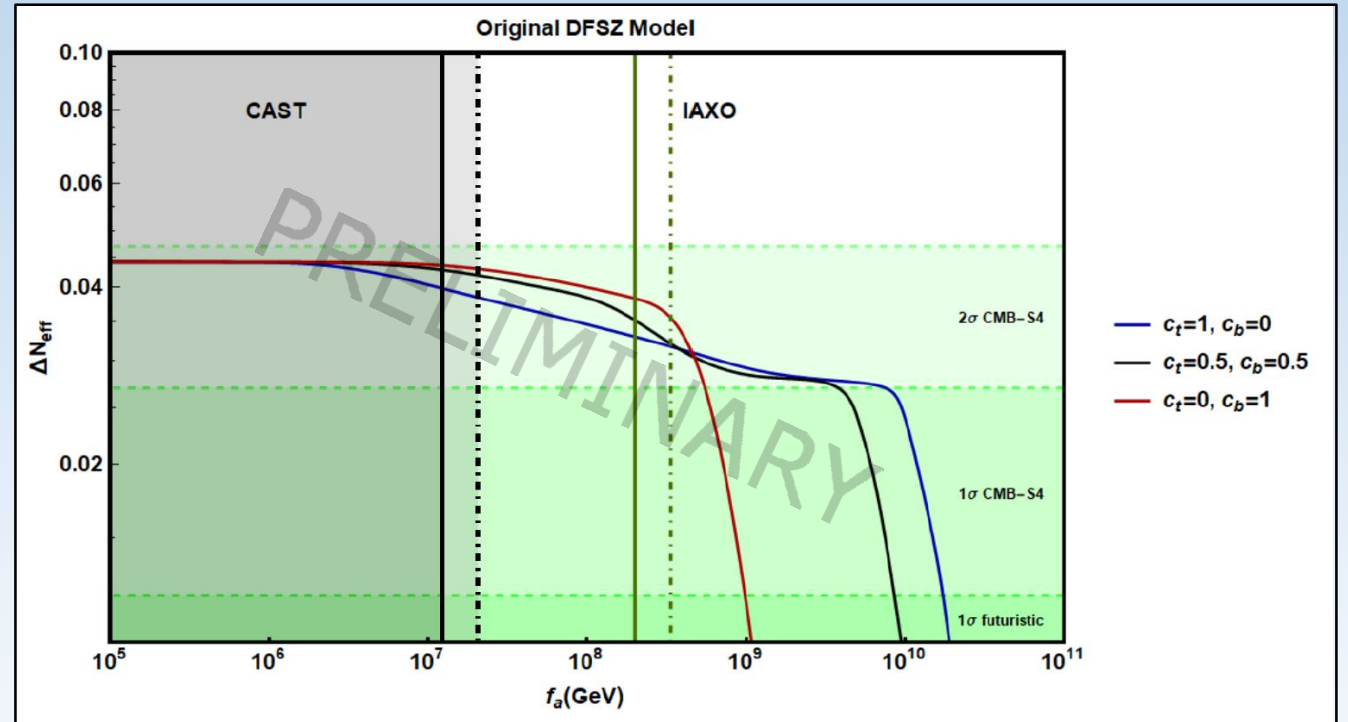
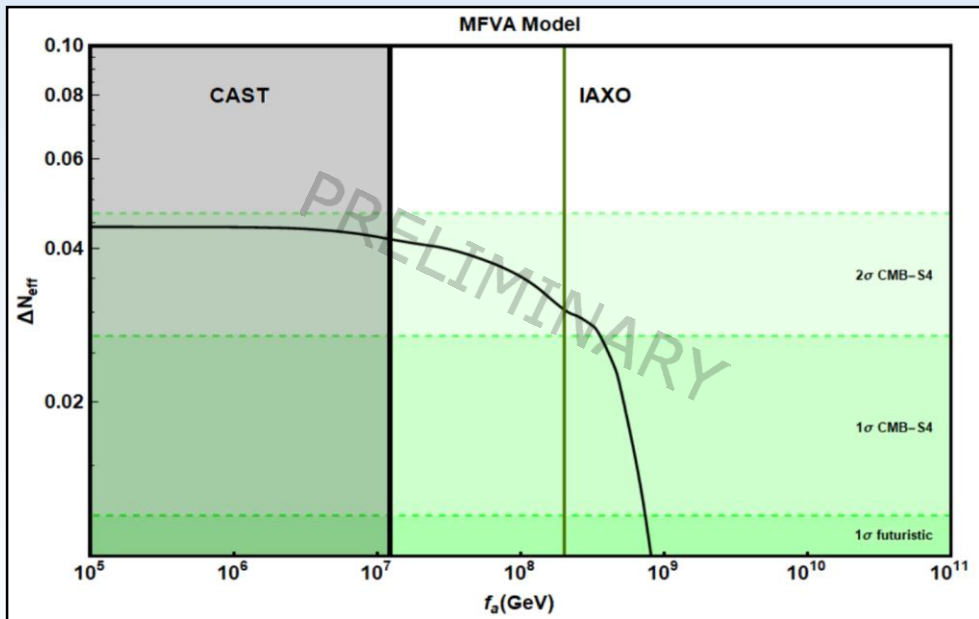
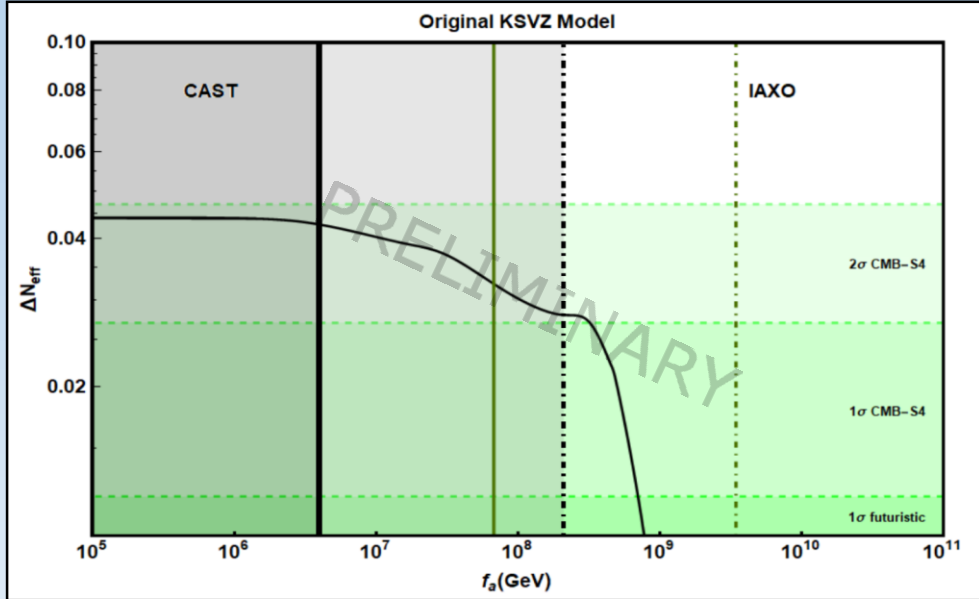
# Results



# UV complete models

- Specific models give a single prediction for  $\Delta N_{eff}(f_a)$
- Two classical invisible axion scenarios:
  - DFSZ.  $c_t + c_b = 1$ ,  $E/N$  has two possible values  
M. Dine, W. Fischler, and M. Srednicki, Phys. Lett.B104.  
A. R. Zhitnitsky, Sov. J. Nucl. Phys.31(1980) 260
  - KSVZ. Only gluon process, many values for  $E/N$   
J. E. Kim, Phys. Rev. Lett. 43 (1979) 103  
M. A. Shifman, A. I. Vainshtein, and V. I. Zakharov, Nucl. Phys.B166(1980) 493–506
- An example of a flavourful axion model:
  - The Minimal Flavour Violating Axion.  $c_t = 0$ ,  $E/N = 8/3$   
FAA & Merlo, 1709.07039

# UV complete models



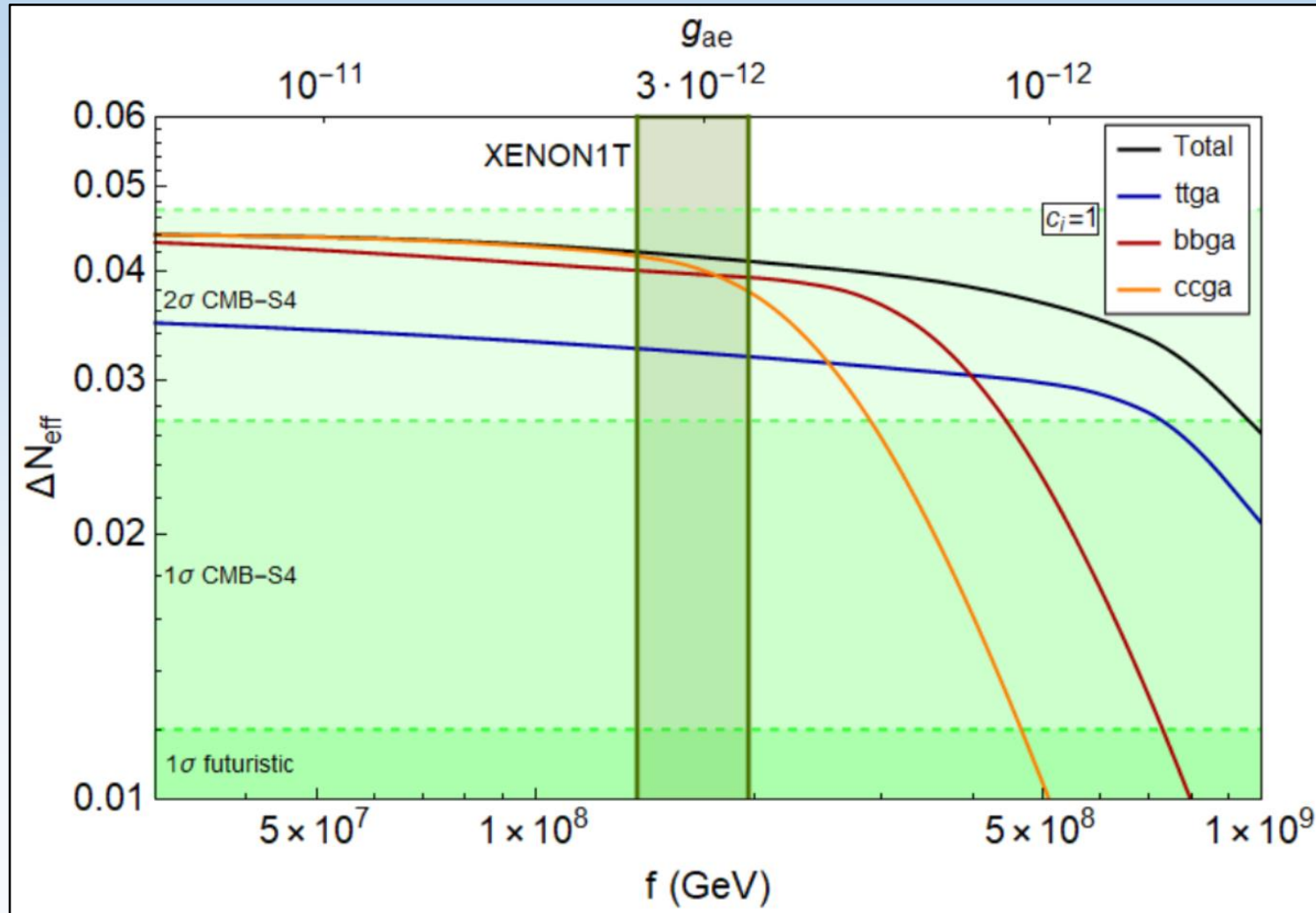
# Complementarity with XENON1T

FAA, F. D'Eramo, R. Z. Ferreira, L. Merlo, A. Notari, 2007.06579

- An excess has been measured by XENON1T, compatible with solar axions
- They could also leave an imprint on the Early Universe through  $N_{eff}$
- Using the expected axion-electron coupling, we look at the expected  $\Delta N_{eff}$  produced by these light relics in three scenarios:
  - Democratic non-anomalous ALP, with all couplings  $c_\psi = 1$
  - Non-anomalous ALP with only one coupling at tree level, inducing the rest at 1 loop
  - DFSZ axion

# XENON1T – Democratic

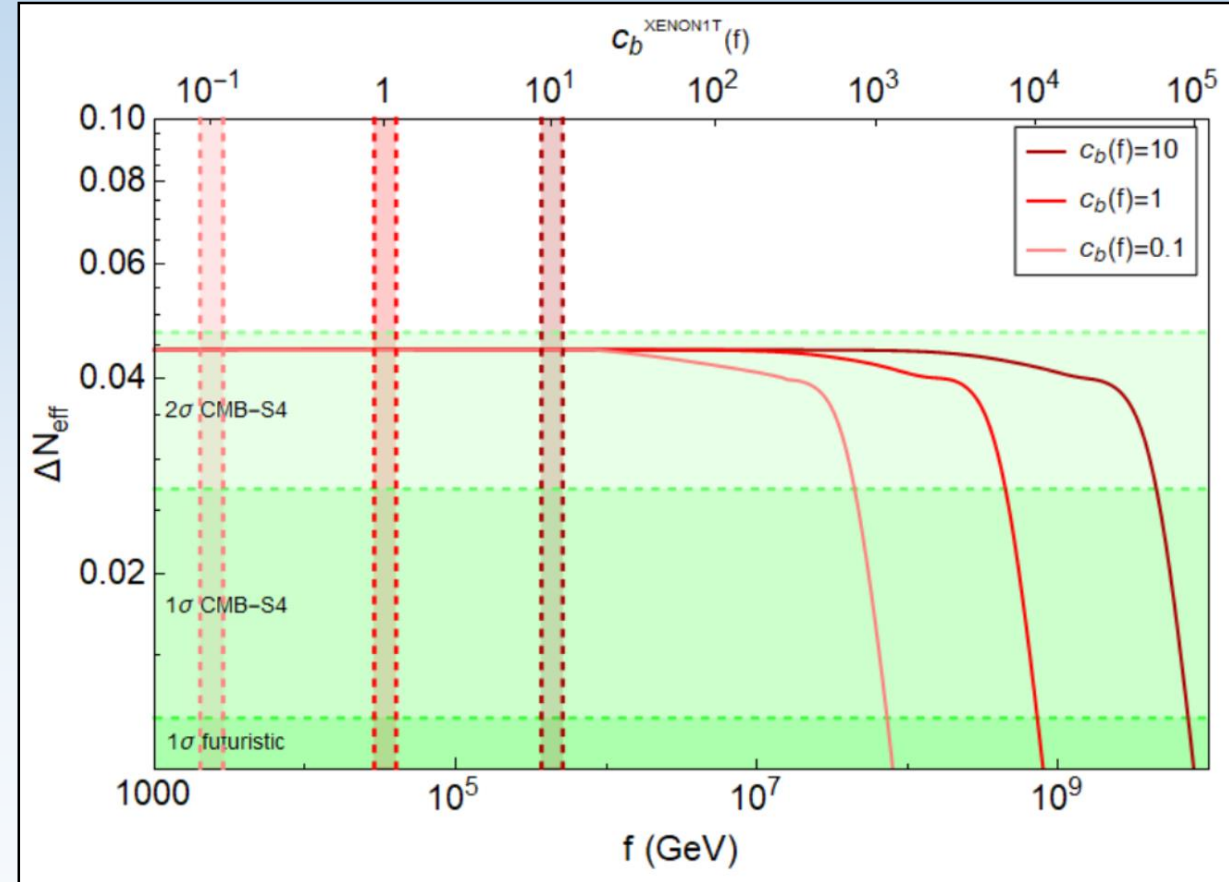
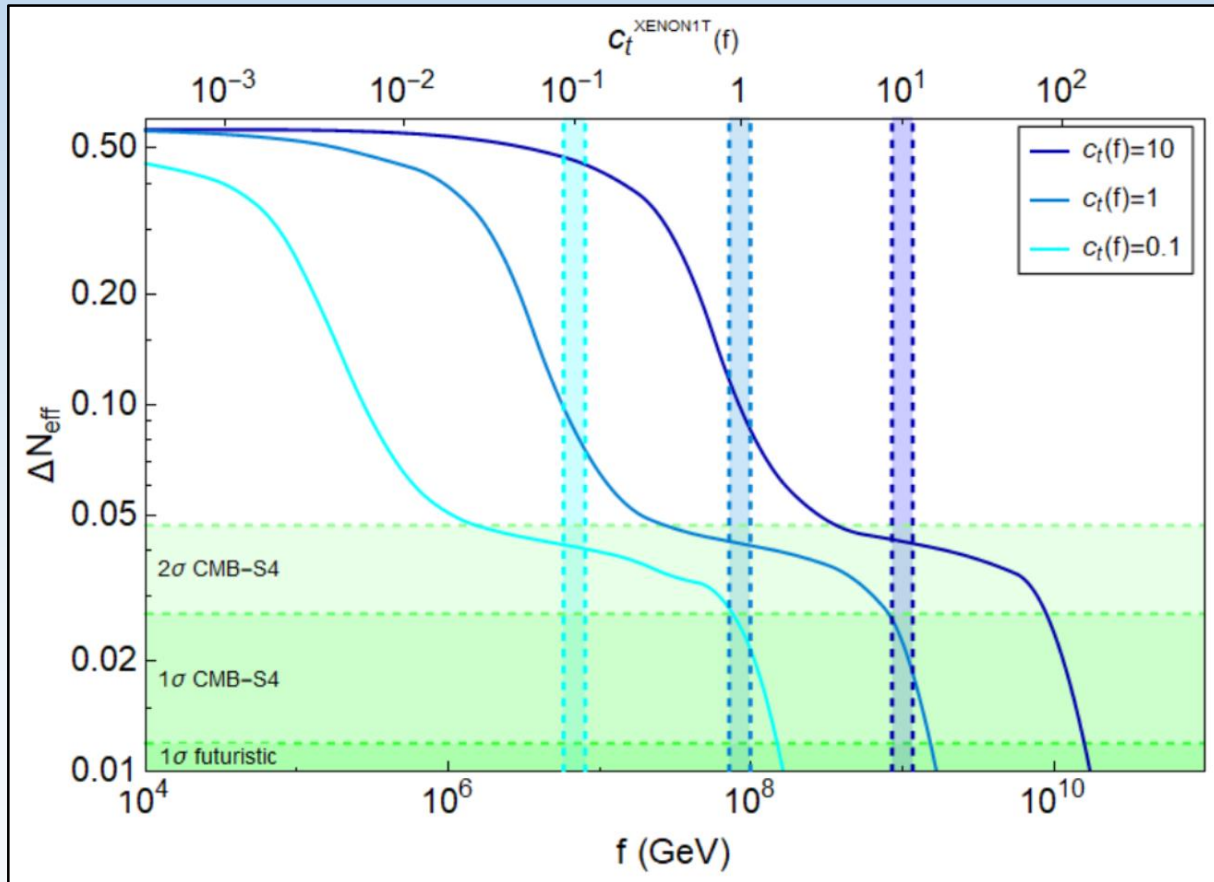
FAA, F. D'Eramo, R. Z. Ferreira, L. Merlo, A. Notari, 2007.06579





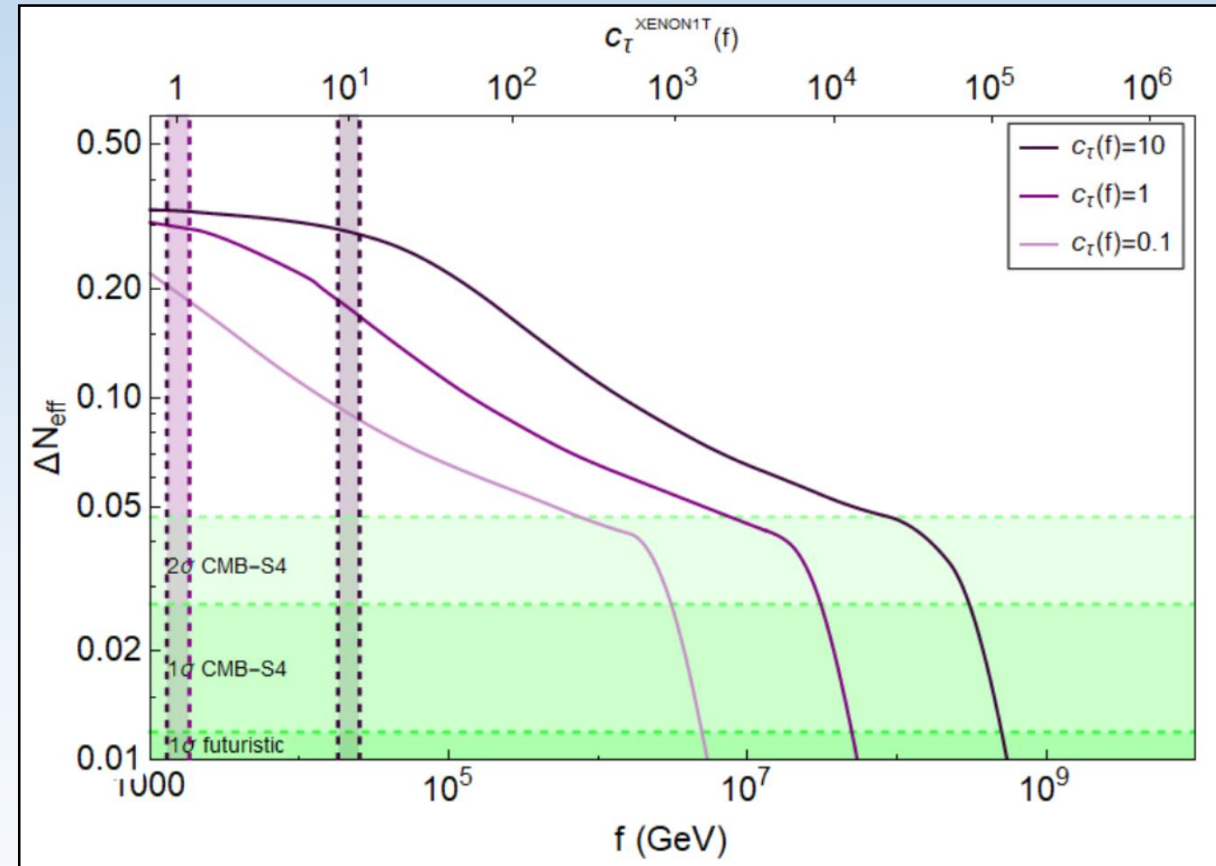
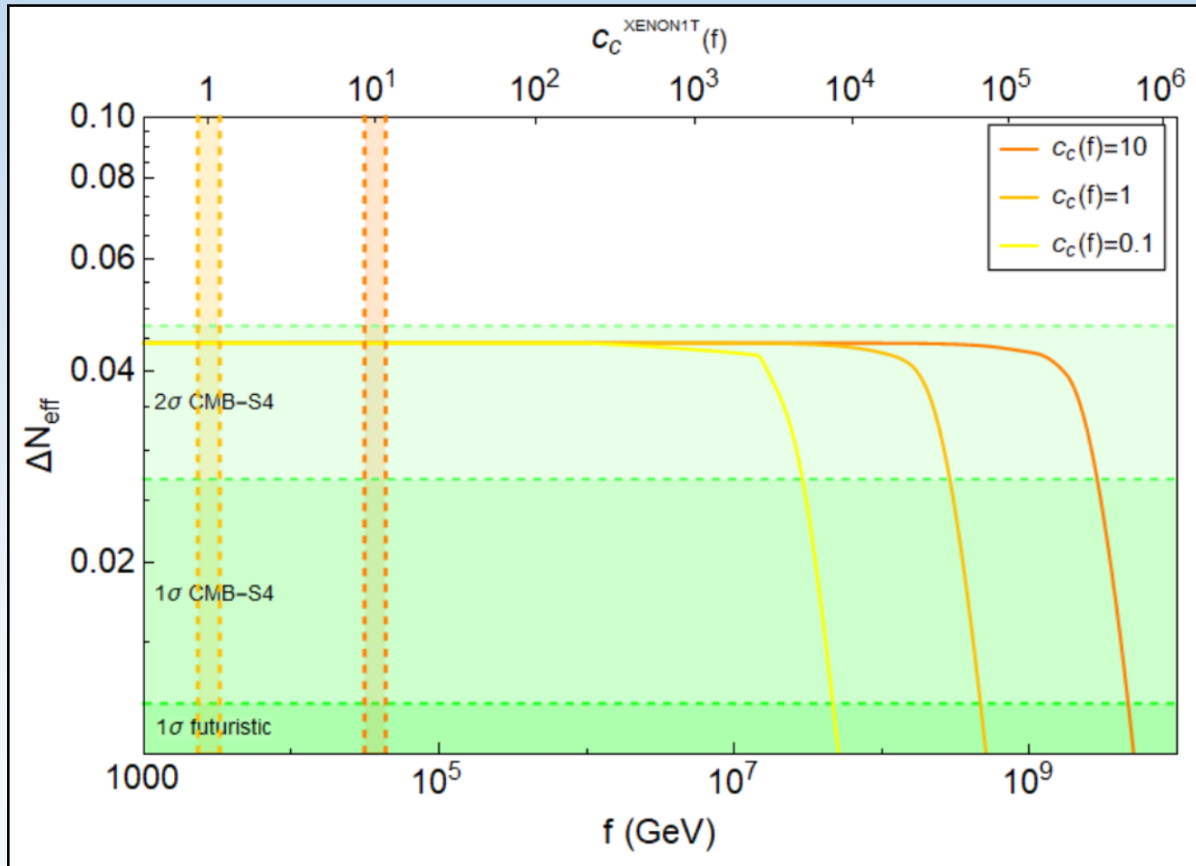
# XENON1T – Loop-induced $c_e$

FAA, F. D'Eramo, R. Z. Ferreira, L. Merlo, A. Notari, 2007.06579



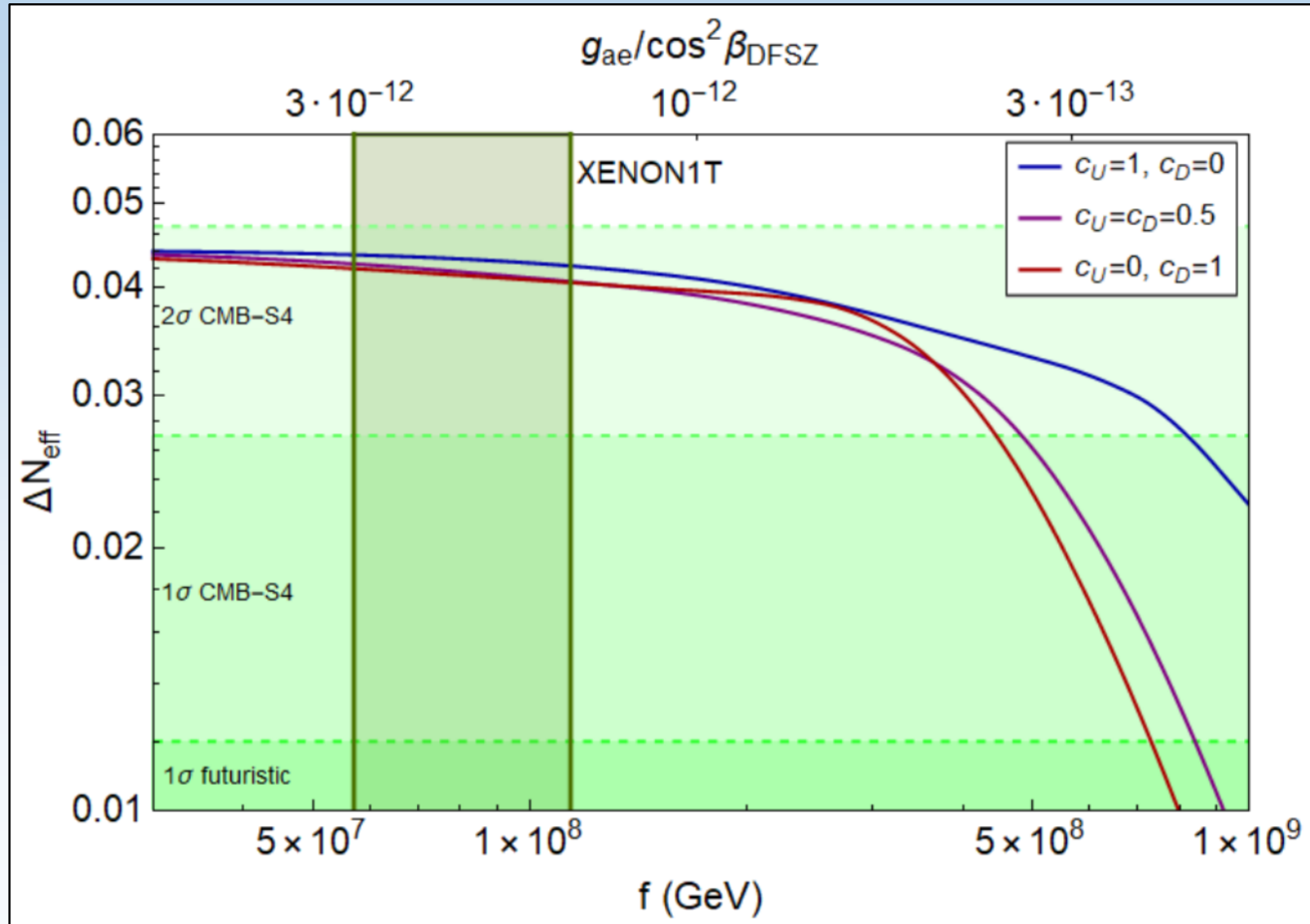
# XENON1T – Loop-induced $c_e$

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# XENON1T – DFSZ

FAA, F. D'Eramo, R. Z. Ferreira, L. Merlo, A. Notari, 2007.06579



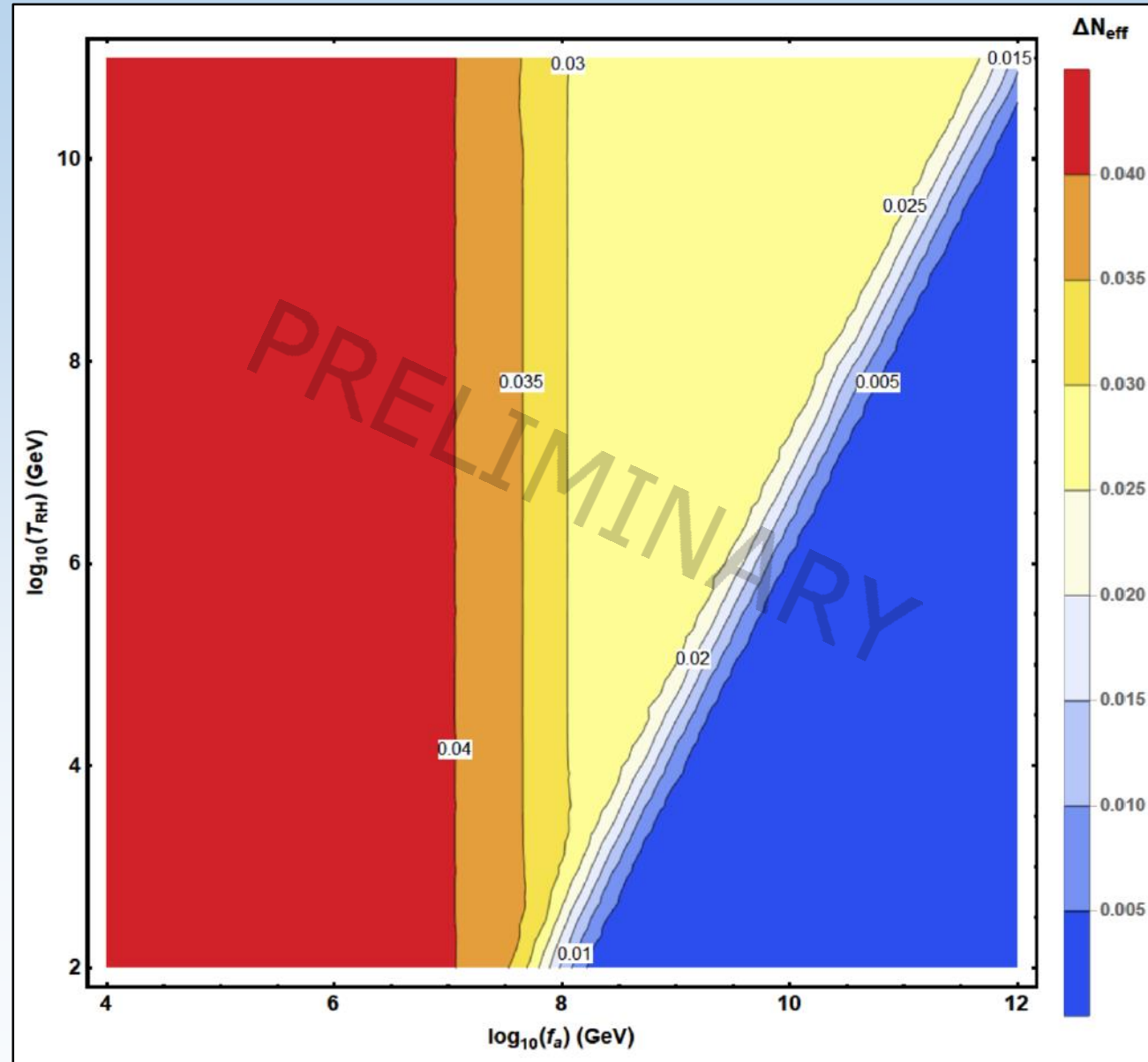
# Summary

- Axions are a very well motivated and useful extension to the SM
- $N_{eff}$  is an interesting observable, complementary to other searches
- It can reward us with knowledge about the axion couplings and/or Early Universe details
- If measured with enough precision, it can help distinguish different models
- Measurable at, at least,  $2\sigma$  with the CMB-S4 experiment if XENON1T is confirmed to be an axion!

THANK YOU FOR  
YOUR ATTENTION

# BACKUP SLIDES

# Results



# UV complete models – DFSZ

- Two Higgs doublets: one for up and another for down quarks
- Same axion couplings to all three generations
- Axion-quark couplings satisfy this relation:

$$c_t + c_b = 1$$

- Charged leptons can couple to either of the Higgs doublets
- Two possible axion-photon couplings
- Different  $\Delta N_{eff}(f)$  predictions depending on specific  $c_t$  and  $c_b$



# UV complete models – KSVZ

- Axion does not couple to SM particles at tree level
- Axion couples to heavy exotic quarks through a new EW singlet scalar
- Many axion-photon couplings depending on heavy quark representations
- $\Delta N_{eff}(f)$  driven solely by gluon process

# UV complete models – MFVA

R. S. Chivukula and H. Georgi, PLB **188**, 99-104 (1987)

- Minimal Flavour Violation: only the SM Yukawas break flavour
- PQ symmetry identified within the MFV symmetry group

(D'Ambrosio et al, 020736; Cirigliano et al, 0507001)

$$\mathcal{G}_F = SU(3)_{Q_L} \times SU(3)_{u_R} \times SU(3)_{d_R} \times SU(3)_{L_L} \times SU(3)_{e_R} \\ \times U(1)_B \times U(1)_L \times U(1)_Y \times U(1)_{PQ} \times U(1)_{e_R}$$

- Flavour conserving non-universal couplings, with  $c_t = 0$

$$\mathcal{L}_Y = \left(\frac{\Phi}{\Lambda_\Phi}\right)^{x_d - x_Q} \bar{Q}_L H d_R \mathcal{Y}_d + \left(\frac{\Phi}{\Lambda_\Phi}\right)^{x_u - x_Q} \bar{Q}_L \tilde{H} u_R \mathcal{Y}_u + \left(\frac{\Phi}{\Lambda_\Phi}\right)^{x_e - x_L} \bar{L}_L H e_R \mathcal{Y}_e + h.c.$$

$$\Phi = \frac{\rho + v_\Phi}{\sqrt{2}} e^{ia/v_\Phi}$$

$$Y_u = \left(\frac{v_\Phi}{\sqrt{2}\Lambda_\Phi}\right)^{x_u - x_Q} \mathcal{Y}_u, \quad Y_d = \left(\frac{v_\Phi}{\sqrt{2}\Lambda_\Phi}\right)^{x_d - x_Q} \mathcal{Y}_d, \quad Y_e = \left(\frac{v_\Phi}{\sqrt{2}\Lambda_\Phi}\right)^{x_e - x_L} \mathcal{Y}_e$$