

New measurement of the neutron electric dipole moment

Benoit Clément (on behalf of the nEDM collaboration)

ICHEP 2020

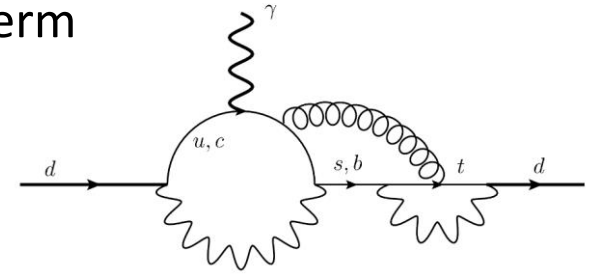


Electric dipole moments

For the theoritician : **CP**-violating term

...in the EM current : $\mathcal{L}_{int} = i \frac{d}{2} \bar{\psi} F_{\mu\nu} \sigma^{\mu\nu} \gamma^5 \psi$

Important for Baryon asymmetry !



In standard model, only through 3 (quarks) or 4 (leptons) loops EW diagrams

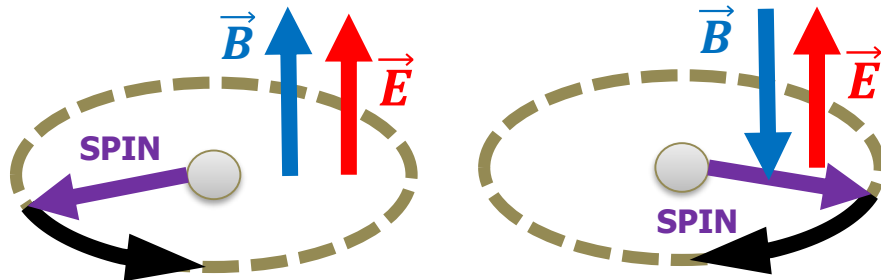
$$\Rightarrow d_n \approx 10^{-31} e.cm$$

BSM could contribute through lower order loops : $d_n \sim 10^{-29}$ to $10^{-26} e.cm$

For the experimentalist : **T**-violating effect

At low energy (non relativistic limit): $\mathcal{L}_{int} = -\vec{d} \cdot \vec{E} = d \hat{J} \cdot \vec{E}$

Coupling between the **Spin** and an external **Electric Field**



>> play >>

<< rewind <<

$$\mu|B| \gg d|E|$$

Change in Larmor frequency if B and E fields are (anti)parallel

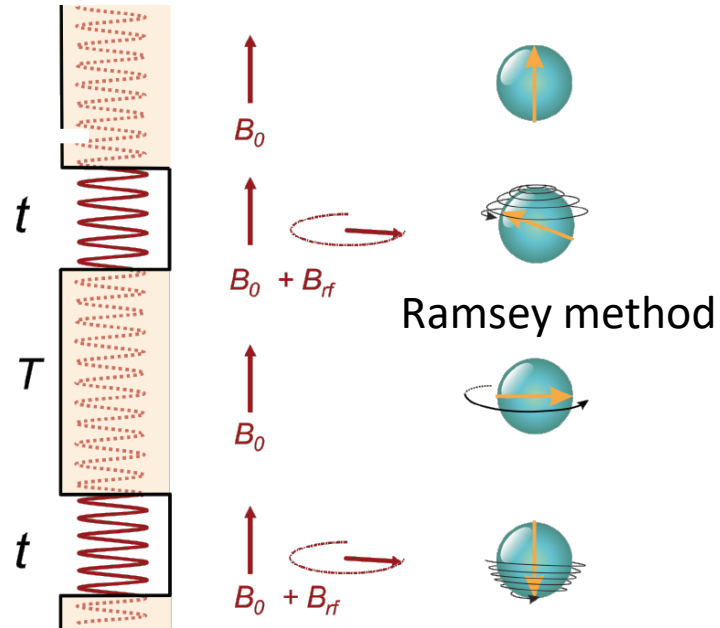
$$2\pi f = \frac{2\mu}{\hbar} |B| \pm \frac{2d}{\hbar} |E|$$

$$f(\uparrow\uparrow) - f(\uparrow\downarrow) = -\frac{2}{\pi\hbar} d|E|$$

$$d_n < 3 \times 10^{-26} e.cm @ 90\% CL$$

J.M. Pendlebury *et al*, Phys. Rev. D **92** 092003 (2015)

How to measure nEDM



Count neutrons in \uparrow and \downarrow spin states
Relate **asymmetry** to **precession frequency**

$$\text{Statistical sensitivity: } \sigma_{d_n} = \frac{\hbar}{2 \alpha E T \sqrt{N}}$$

1 cycle : $T \approx 180s, N = 5000$ to 15000

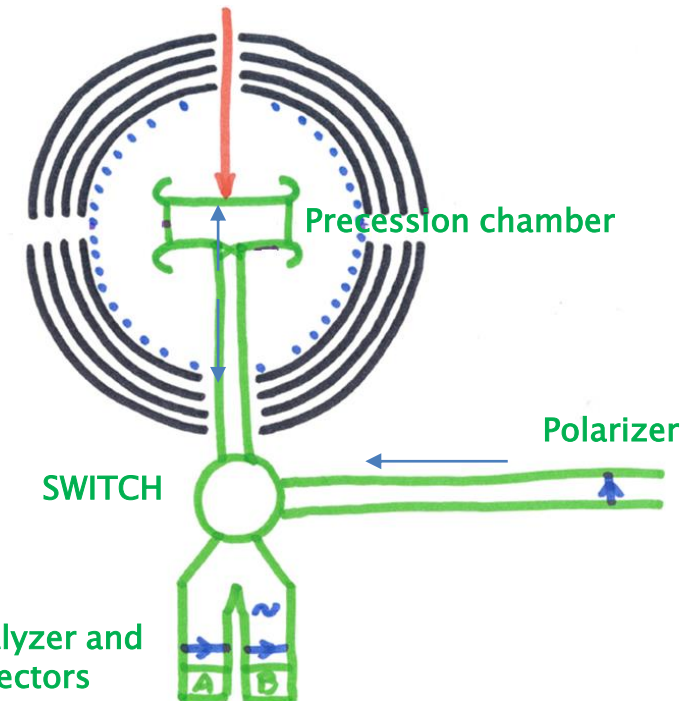
Ultracold neutrons :

- $E_{kin} \sim 100\text{neV}, v \sim 4\text{m. s}^{-1}$
- Reflected at all incidences
- Sensitive to gravity

Can be :

- guided through simple tubes
- stored in material « bottle »

Allow large storage times



The nEDM experiment

nEDM collaboration : ~ 50 physicists from 8 countries and 16 labs



Data taken over two years (2015 and 2016) at PSI

Using the apparatus used at ILL by the

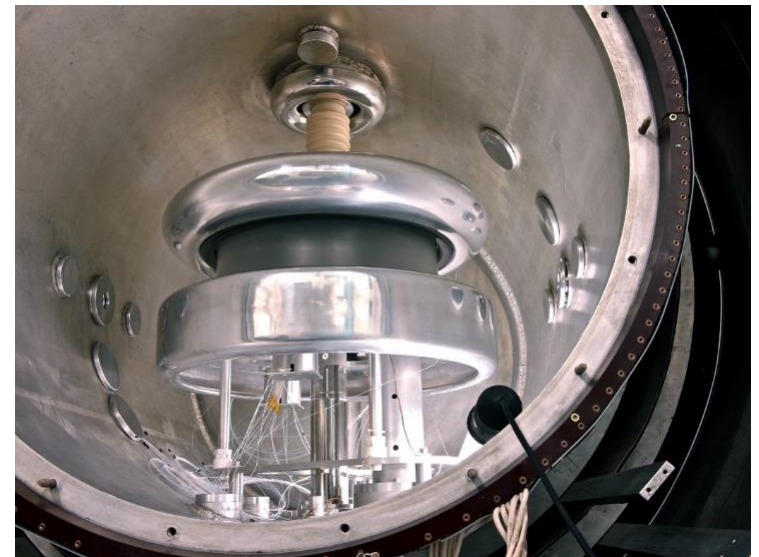
RAL/sussex collaboration

54 068 cycles of 180 s in 99 magnetic configurations (sequence)

~ 11 400 UCN per cycle

Electric field: $E \sim 11 \text{ kV/cm}$

Magnetic field: $B_0 \sim 1 \mu\text{T}$



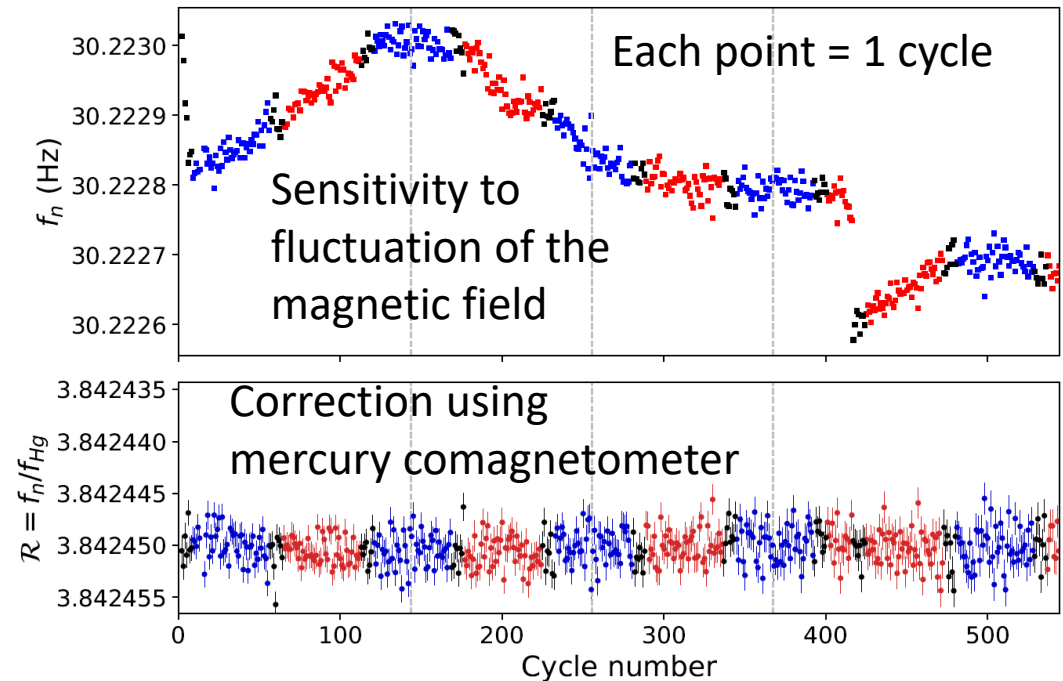
$$\sigma_{d_n} = \frac{\hbar}{2ET\alpha\sqrt{N_{\text{tot}}}}$$

Statistical sensitivity $\rightarrow 1, 1 \times 10^{-26} e \text{ cm}$

Magnetic management

Controlling the magnetic field and gradients:

- By reducing the ambient magnetic field :
 - Passive mu-metal shielding,
 - Active compensation coils,
 - Magnetic free environment.
- By measuring the field :
 - Cs magnetometers outside the chamber : field gradients,
 - Hg comagnetometer : measure (**almost**) the same field as seen by neutrons.



But any effect on Hg measurement linearly dependent on \vec{E} will induce a false Edm

Using magnetic gradients in the analysis

Fake Hg edm combines :

- Relativistic vxE
- Magnetic inhomogeneities (« gradients »)

$$d_{n\leftarrow\text{Hg}}^{\text{false}} = \frac{\hbar \gamma_n \gamma_{\text{Hg}} R^2}{8 c^2} (\mathbf{G}_{\text{grav}} + \hat{\mathbf{G}})$$

\mathbf{G}_{Grav} : Due to gravity, there is an average height difference of gaseous Hg and ultracold neutrons $\langle z \rangle = -0,39\text{cm}$ and feel slightly different fields. Both frequency ratio and false edm depends on this gradient

$\hat{\mathbf{G}}$: higher order gradients

Analysis strategy

- Create magnetic gradients building various magnetic configurations

- For each configuration: fit $\mathcal{R}(E) = \frac{f_n}{f_{\text{Hg}}} = \mathcal{R}_0 - \frac{2d_n}{\hbar|\gamma_n|B_0} E$

$$\mathcal{R}_0 = \left| \frac{\gamma_n}{\gamma_{\text{Hg}}} \right| \left(1 + \frac{\langle z \rangle}{B_0} \mathbf{G}_{\text{grav}} + \frac{\langle B_T^2 \rangle}{2B_0^2} \right) \quad d_n = d_n^{\text{true}} + \frac{\hbar \gamma_n \gamma_{\text{Hg}} R^2}{8 c^2} (\mathbf{G}_{\text{grav}} + \hat{\mathbf{G}})$$

- Use maps of the magnetic field to measure the gradients $\hat{\mathbf{G}}$ and the transverse field inhomogeneities $\langle B_T^2 \rangle$

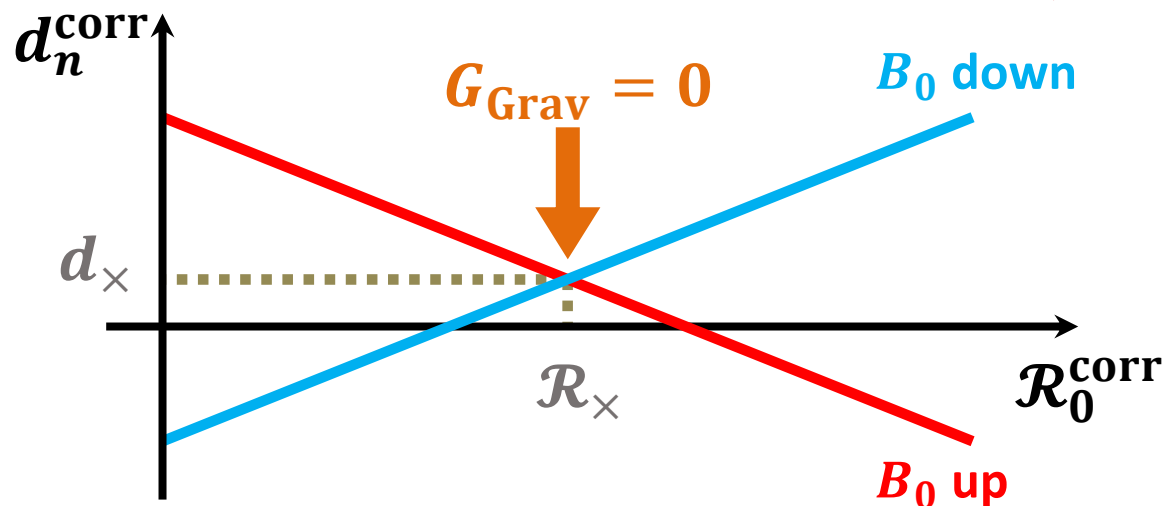
The crossing point

- After correction, we got 2 quantities that depend linearly on G_{grav} with a known slope, $\pm\eta$ (sign depends on magnetic field sign)

$$d_n^{\text{corr}} = d_n - \frac{\hbar \gamma_n \gamma_{\text{Hg}} R^2}{8 c^2} \hat{G} \quad \mathcal{R}_0^{\text{corr}} = \mathcal{R}_0 - \left| \frac{\gamma_n}{\gamma_{\text{Hg}}} \right| \frac{\langle B_T^2 \rangle}{2B_0^2}$$

- The crossing of the two lines correspond to $G_{\text{grav}} = 0$ cancelling the false Edm term

$$d_x = d_n - \frac{\hbar \gamma_n \gamma_{\text{Hg}} R^2}{8 c^2} \hat{G} = \boxed{d_n^{\text{true}}} + \frac{\hbar \gamma_n \gamma_{\text{Hg}} R^2}{8 c^2} \hat{G}_{\text{grav}}$$



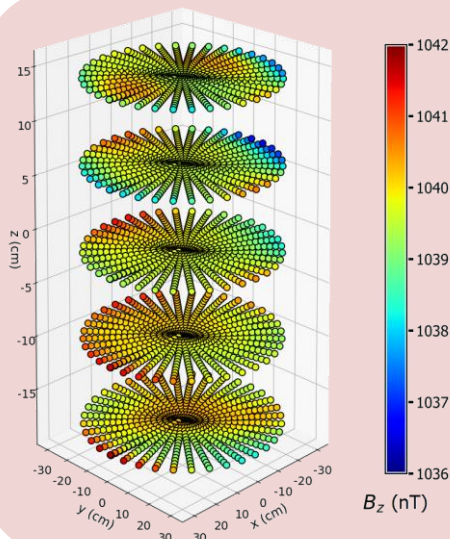
Systematics overview

		Effect	shift error $\times 10^{-28}$ e.cm	
False Hg EDM	}	Error on $\langle z \rangle$	-	7
		Higher order gradients \hat{G}	69	10
		Transverse field correction $\langle B_T^2 \rangle$	0	5
Other effects	}	Hg EDM[8]	-0.1	0.1
		Local dipole fields	-	4
		$v \times E$ UCN net motion	-	2
		Quadratic $v \times E$	-	0.1
		Uncompensated G drift	-	7.5
		Mercury light shift	-	0.4
		Inc. scattering ^{199}Hg	-	7
		TOTAL	69	18

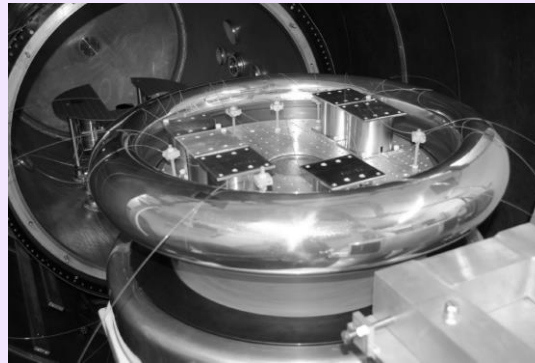
Final systematic :

$$\pm 0.2 \times 10^{-26} \text{ e.cm}$$

Field mapping



Cesium magnetometers

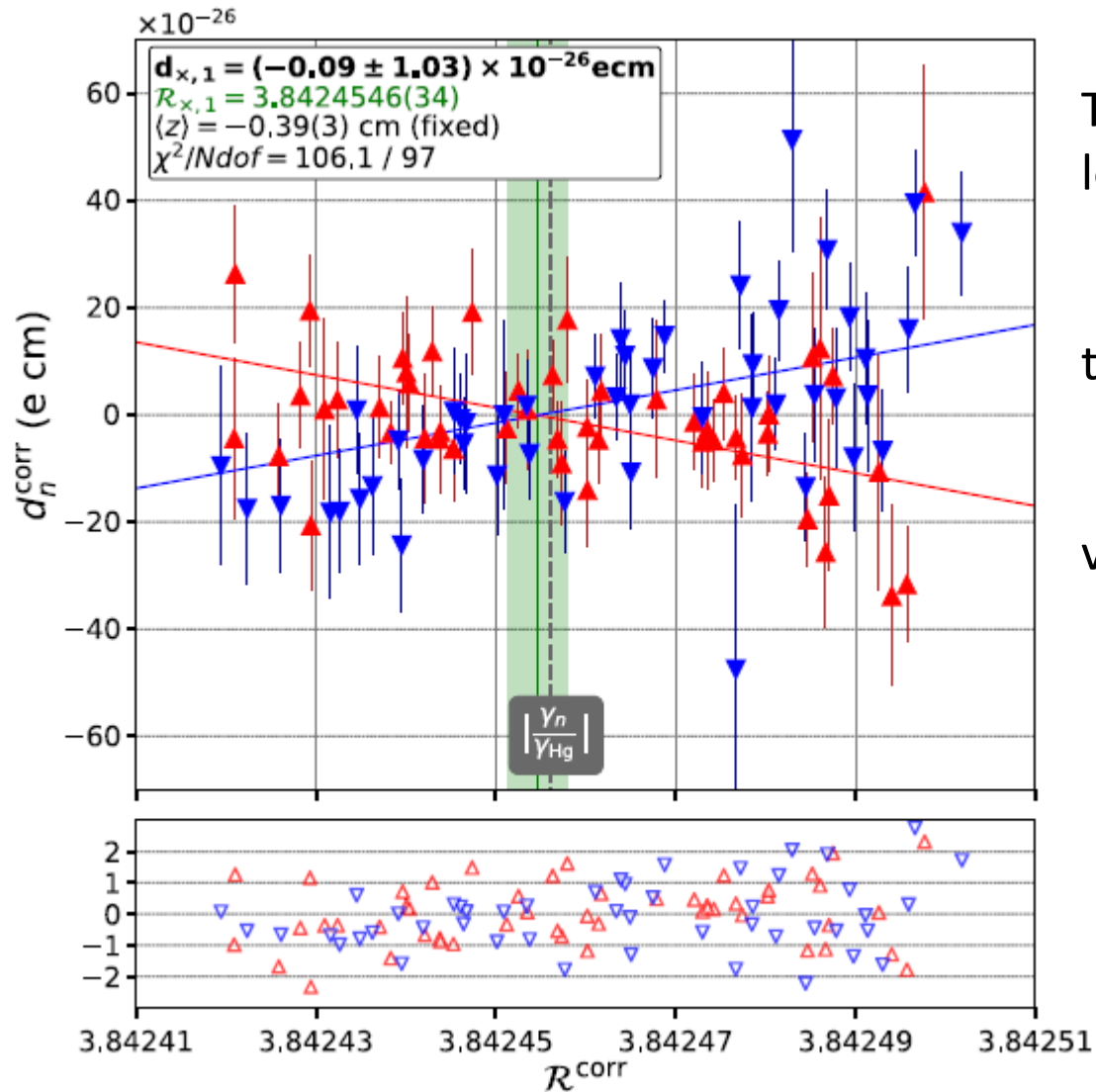


*Optically Pumped Cs Magnetometers
Enabling a High-Sensitivity Search for
the Neutron EDM,
submitted to PRA*

Scans for magnetic contaminations at PTB Berlin



Final result



Two analysis teams working on two level blinded data :

- optimize the selection and debug the analysis code independently

- **First unblinding** : check if central value are consistent between teams

- **Second unblinding** : Final results !

$$\text{I} \rightarrow (-0,09 \pm 1,03) \times 10^{-26} \text{ e cm}$$

$$\text{II} \rightarrow (+0,15 \pm 1,07) \times 10^{-26} \text{ e cm}$$

C. Abel *et al.* Phys. Rev. Lett. **124**, 081803 –
Published 28 February 2020

$$d_n = (0, 0 \pm 1, 1_{\text{stat}} \pm 0, 2_{\text{syst}}) \times 10^{-26} \text{ e cm}$$

Previous : $d_n = (-0.2 \pm 1.5_{\text{stat}} \pm 1.0_{\text{syst}}) \times 10^{-26} \text{ e cm}$

Coming next : n2EDM project

The nEDM collaboration has demonstrated its capacity to control systematic effects at a few 10^{-27} ecm

New spectrometer : 2022-...

Control of the magnetic field homogeneity and measurement : improved systematics

- Dual chamber apparatus : simultaneous measurement of both field configuration
- Colossal magnetic shield (6 mu-metal layers, $5 \times 5 \times 5 \text{ m}^3$), shielding factor 100000
- Advanced magnetometry (external and comagnetometry)

Increase the electric field

- 180 kV

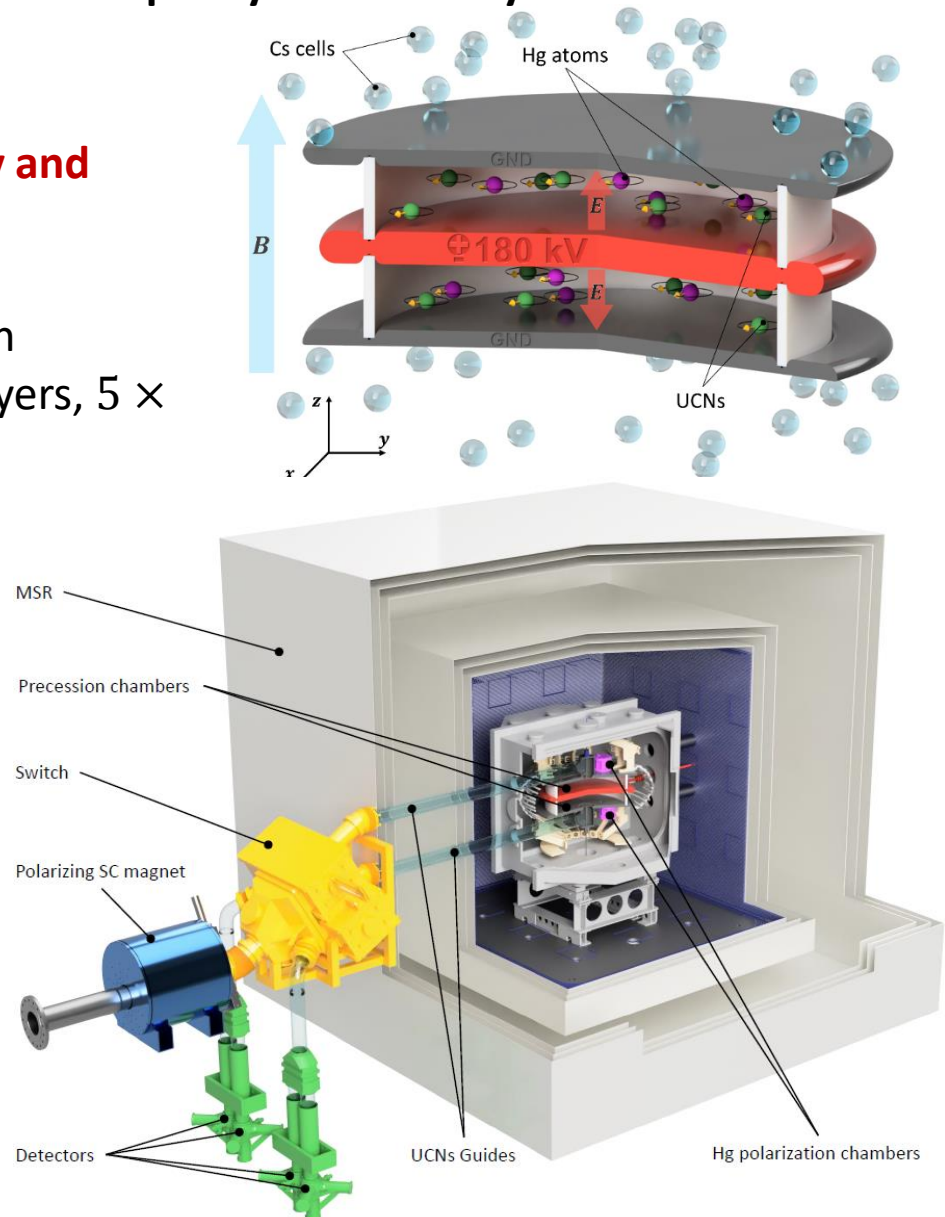
Increase the statistics

- 5 years data taking
- Larger chamber (Volume x 7)

Possible further improvement:

« magic field » to compensate Systematic $v \times E$ effect in comagnetometer

G.Pignol 1812.01420



Conclusion

nEDM experiment result

New result on the neutron EDM by the nEDM collaboration (PSI):

$$d_n = (0, 0 \pm 1, 1_{\text{stat}} \pm 0, 2_{\text{syst}}) \times 10^{-26} \text{ e cm}$$

Still (unfortunately) compatible with zero, translate into a limit:

$$d_n < 1.8 \times 10^{-26} \text{ e cm @ 90\% CL}$$

Improve the limit by a **factor 2**.

Systematics are reduced by a **factor 5**: Show a good understanding and control of the systematics effects.

n2EDM apparatus

- design includes our understanding of systematics to further improve the sensitivity.

- reduced statistical sensitivity (higher volume, higher E field)

GOAL : reach 10^{-27} e.cm in 2027

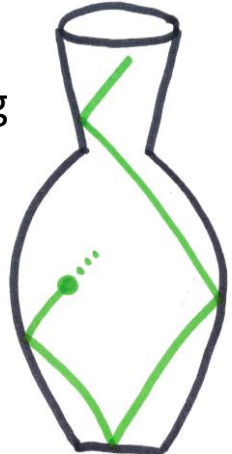
Extra material

Based on **ultra cold neutron storage** (all but one)

Neutron production : spallation or nuclear reactor + thermalisation + inelastic cooling (He-II or solid D₂)

Room temperature experiments : UCN need to be produced (source) then extracted/guided to the experimental chamber

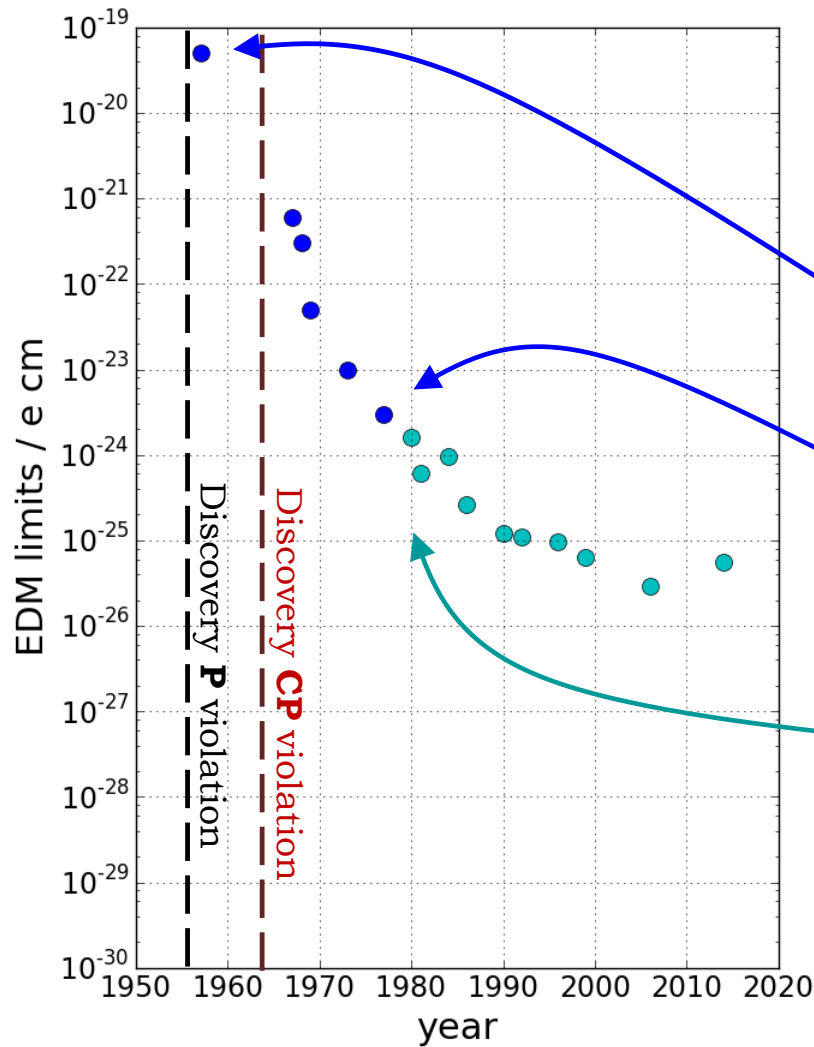
Cryogenic experiments : UCN production and EDM measurement takes place at the same place (in superfluid helium at 1K...)



Place	Neutron source	Concept	Stage/Readiness
SNS	Spallation + UCN production in situ	Cryogenic double chamber with helium comagnetometers	« large scale integration » phase
PSI	Spallation + sD2 UCN source	n2EDM: double Ramsey chamber with mercury comagnetometers	Source running, experiment under construction
LANL	Spallation + sD2 UCN source	double Ramsey chamber with mercury comagnetometers	Source running, experiment in design phase
TRIUMF	Spallation + superfluid He UCN source	double Ramsey chamber	Source under construction, experiment in conceptual phase
ILL	Reactor + superfluid He UCN source	panEDM: double Ramsey chamber, no comagnetometers	Source and experiment under construction
PNPI	WWR-M reactor + inpile superfluid He UCN source	Getting a really high density of UCNs	Source under construction
ESS	Spallation + Cold neutron beam	100m double beam + time of flight	Demonstration phase, small prototype operational @ ILL

Room temperature UCN experiments aiming at a precision of $\approx 1 \times 10^{-27} e \text{ cm}$

Neutron EDM history

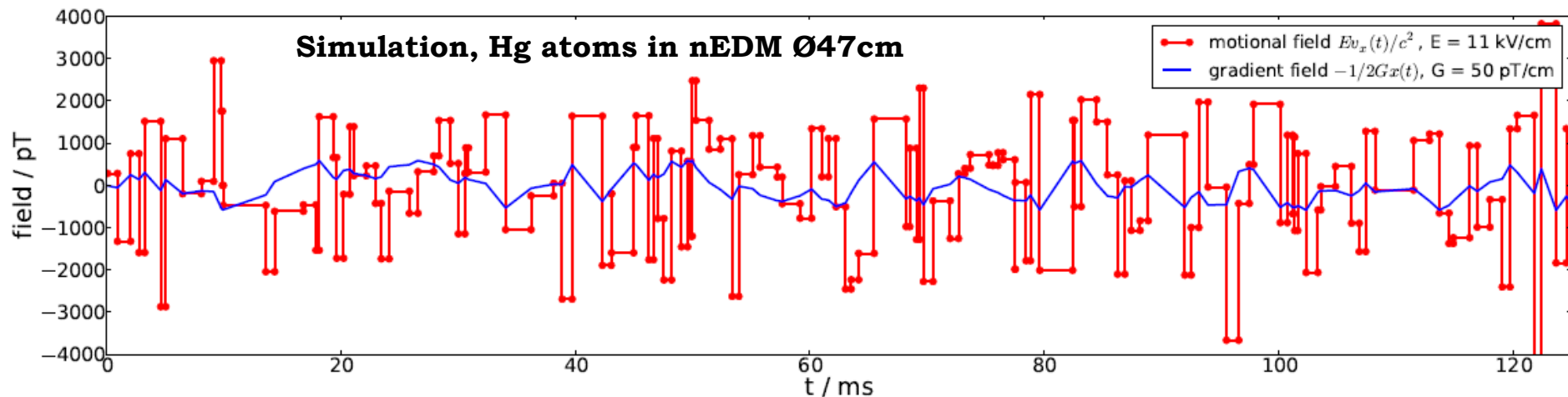


First experiment
by Smith, Purcell and Ramsey
with a **beam of thermal neutrons**,
interrogation time $T \approx 1$ ms

Best « beam » experiment in Grenoble
with **slower neutrons**,
interrogation time $T \approx 20$ ms

« Modern » experiments with
ultracold neutrons started.
interrogation time $T \approx 100$ s

The co-magnetometer problem: vxE/c^2



Frequency shift from a transverse magnetic noise \underline{B}

$$\delta f = \frac{\gamma^2}{4\pi} \int_0^\infty d\tau \operatorname{Im} e^{-i\omega\tau} \langle \underline{B}(0) \underline{B}^*(\tau) \rangle$$

Linear-in-E frequency shift (low B limit)

$$d_{n\leftarrow\text{Hg}}^{\text{false}} = -\frac{\hbar |\gamma_n \gamma_{\text{Hg}}|}{2c^2} \langle x B_x + y B_y \rangle$$

B-field uniformity is critical

$$G = 1 \text{ pT/cm} \leftrightarrow d_{n\leftarrow\text{Hg}}^{\text{false}} = 4.4 \times 10^{-27} \text{ ecm}$$

From nEDM to n2EDM @ PSI

$$d_n < 3 \times 10^{-26} \text{ e.cm}$$

$$d_n = (-0.2 \pm 1.5_{stat} \pm 1.0_{syst}) \times 10^{-26} \text{ ecm}$$

1998 1999 2000 2001 2002 2003 2004 2005 2006 2007 2008 2009 2010 2011 2012

ILL data production

RAL/Sussex collaboration



New apparatus : n2EDM

Transfer at PSI (new UCN source)
Various upgrade
New collaboration : nEDM



PSI data

n2EDM installation
and commissioning

n2EDM data...

2013 2014 2015 2016 2017 2018 2019 2020 2021 2022 2023 2024 2025 2026 2027

$$d_n = (0.0 \pm 1.1_{stat} \pm 0.2_{syst}) \times 10^{-26} \text{ ecm}$$

$$d_n < 1.8 \times 10^{-26} \text{ e.cm}$$

Abel et al., PRL 124 081803 (2020)

Beyond SM contribution

SM contribution through EW loops : $d_n \approx 10^{-33} e \text{ cm}$

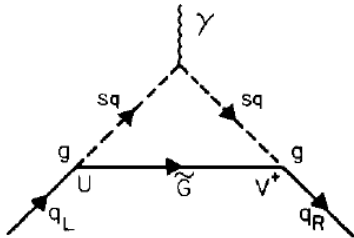
The QCD contribution : $\mathcal{L} = \frac{\alpha}{8\pi} \theta G^{\mu\nu} \widetilde{G}_{\mu\nu}$

$d_n \approx \theta \times 10^{-16} e \text{ cm}$
 $\rightarrow \theta < 10^{-10}$

No known mechanism to reduce θ to (nearly?) zero \rightarrow axions ?

« Strong CP problem »

One loop contribution : for example MSSM contains ~ 40 CP violating imaginary parameters...



$$d_n \approx e \frac{\alpha}{4\pi} \frac{m_q}{M_{CPV}^2} \approx \left(\frac{1 \text{ TeV}}{M_{CPV}} \right)^2 \times 10^{-25} e \text{ cm}$$

Fig. 2. One-loop diagram which may contribute to d_n in a softy broken susy model.

Ellis, Ferrara, Nanopoulos, *PLB* **114** (1982).

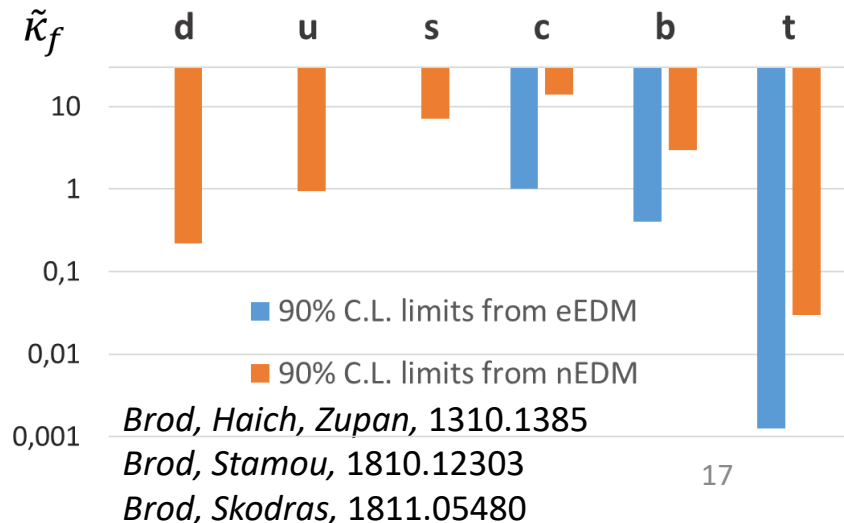
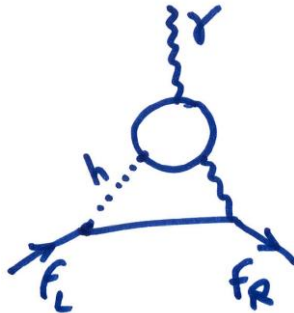
EDM induced by soft mass terms for squarks and gluinos

Two loops contribution :

Modified Higgs Yukawa coupling:

$$\mathcal{L} = -\frac{y_f}{\sqrt{2}} (\kappa_f \bar{f} f h + i \tilde{\kappa}_f \bar{f} \gamma_5 f h)$$

Barr, Zee, *PRL* **65** (1990)

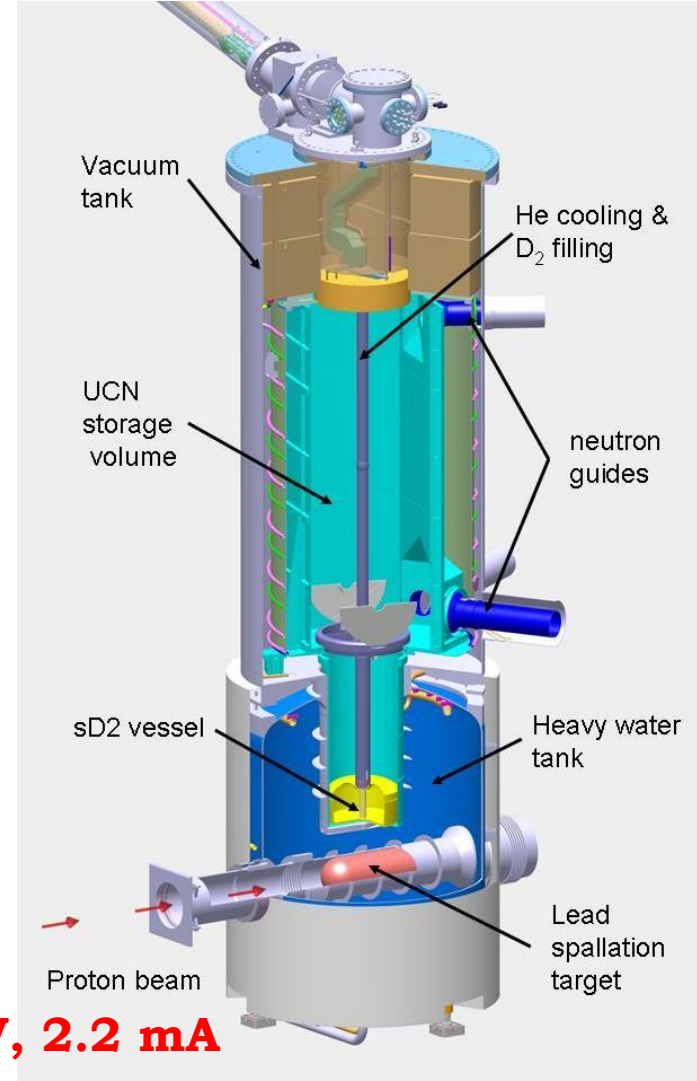


Ultracold neutrons @ PSI

UCN source at the Paul Scherrer Institute



pulsed UCN source
One kick per 5 min
online since 2011



600 MeV, 2.2 mA