



# Study of the normalized $p_T^W$ Distribution in $p\bar{p}$ Collisions at D0

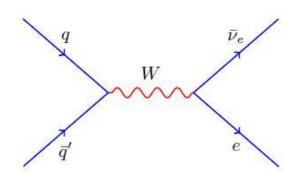
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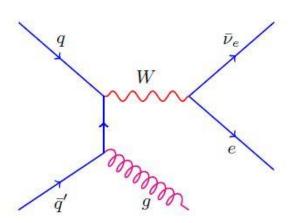
University of Science and Technology of China On Behalf of the D0 Collaboration

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#### Motivation

 $\triangleright p_T^V$  is described by QCD calculation



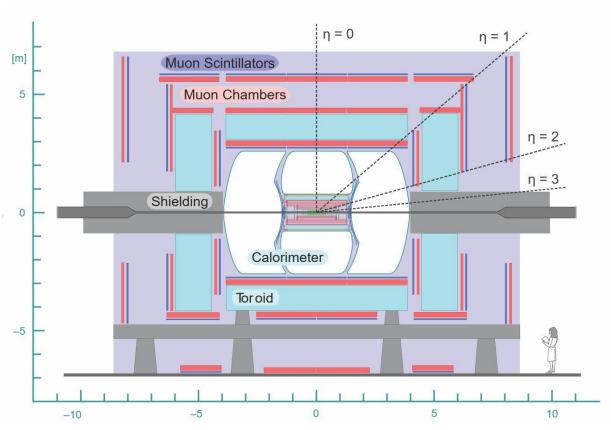


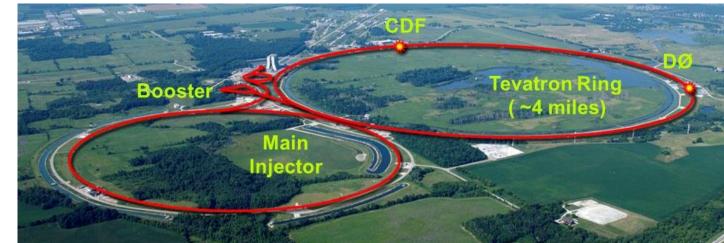
- $\triangleright$  Leading Order (LO):  $p_T^V = 0$
- $\triangleright$  Including higher order:  $p_T^V$  arise from the initial state parton emission
- > Test QCD predictions
- ightharpoonup In  $p\bar{p}$  collisions, the production dominated by valence quark
  - ➤ In the LHC experiments, it involves sea quarks
- $\triangleright$  Low  $p_T(V)$  region dominated by multiple soft gluon emissions
  - ➤ QCD predictions from a soft-gluon resummation formalism (CSS¹)
  - $\triangleright$  Using a form factor with 3 non-perturbative parameters,  $g_1$ ,  $g_2$  and  $g_3(BLNY^2)$
  - $\triangleright$   $g_1$ ,  $g_2$  and  $g_3$  fixed to previous measurement<sup>2</sup>
  - ➤ Constrain models of non-perturbative approaches
  - $\triangleright$  Benefit other related electroweak parameter measurements such as  $m_W$

- > Introduction
  - $\triangleright$  First Tevatron Run II  $p_T^W$  study
    - First  $p_T^W$  study at  $\sqrt{s} = 1.96 \text{ TeV}$
  - $\triangleright$  Based on the latest D0 published  $m_W$  measurement
    - ➤ Same data sample, 4.35 fb<sup>-1</sup> Run II Data
    - > Same background estimation strategy
    - > Same detector calibration methodologies
    - ➤ Same parametrized MC simulation (PMCS)
  - Focus on low  $p_T^W$  region (<15 GeV)
    - ➤ Sensitive to QCD non-perturbative parameters
  - > Provide reconstruction level results
    - ➤ A fast folding procedure for comparisons to other models

#### > D0 Detector

- Central tracking system
- Silicon Microstrip Tracker (SMT)
- Scintillating Central Fiber Tracker (CFT)
- ➤ 1.9 T Solenoid
- **Calorimeter**
- ► Liquid argon and uranium  $|\eta|$  < 4.2
- ➤ Electron energy measurement
- Hadronic recoil reconstruction





#### > Samples and selections

- ightharpoonup Data: Run II, 4.35 fb<sup>-1</sup>,  $\sqrt{s} = 1.96 \text{ TeV}$
- > Trigger requirement:
  - ➤ At least one electromagnetic (EM) cluster
  - > Transverse energy threshold: 25~27 GeV depending on instant luminosity
- > Offline selections:
  - > Electron candidate:

$$p_T^e > 25 \text{ GeV}, |\eta^e| < 1.05$$

Pass shower shape and isolation requirements

> W candidate:

At least one electron candidate

$$50 < m_T < 200 \text{ GeV}, p_T^{Missing} > 25 \text{ GeV}, u_T < 15 \text{ GeV}$$

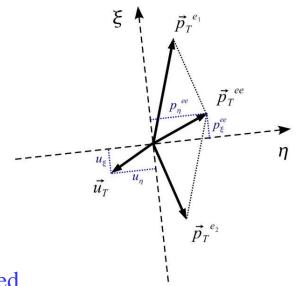
- ightharpoonup Hadronic Recoil  $\vec{u}_T = \sum \vec{E}_T^{calo}$ , represents  $p_T^W$ 
  - ➤ The vector sum of reconstructed energy clusters in the calorimeters excluding deposits from the lepton

$$ightharpoonup \vec{E}_T^{Missing} = -(\vec{u}_T + \vec{p}_T^e)$$
, represents  $p_T^v$ 

$$m_T = \sqrt{2p_T^e p_T^v (1 - \cos\Delta\phi)}$$

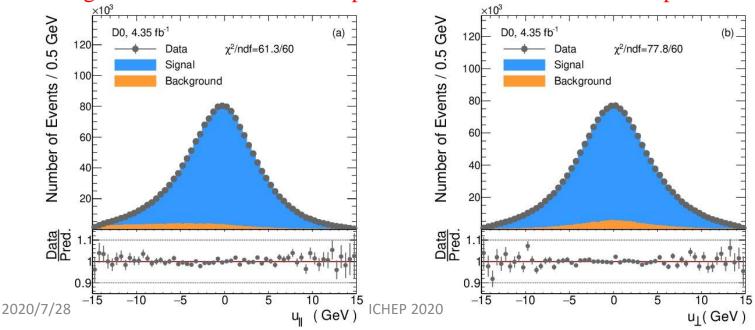
- Detector Calibration
  - Electron energy calibrated using Z mass
    - $\triangleright$  Two parameters:  $E_{corr} = \alpha E_{obs} + \beta$
  - ➤ Hadronic Recoil calibrated with Z candidates
    - $\triangleright \hat{\eta}$ : the direction bisecting the two electrons
    - $\succ$  Tuned by the imbalance in  $\hat{\eta}$  direction,  $\eta_{imb}$

$$\eta_{imb} = (\vec{u}_T + \vec{p}_T^{ee}) \cdot \hat{\eta}$$



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- ➤ In W candidates, only one charged lepton reconstructed
  - $\triangleright u_{\parallel}$  and  $u_{\perp}$ : the parallel and perpendicular components to the electron direction
  - > Tests the modeling of the hadronic recoil
- ➤ Good agreement between data and prediction on hadronic recoil response



#### ➤ Background Estimation

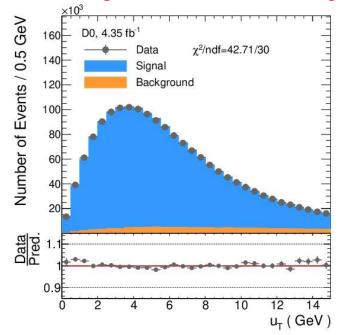
- $\blacktriangleright$  Three backgrounds: W  $\rightarrow \tau v \rightarrow e v v v$ , Z  $\rightarrow e e$ , Multi-Jet
  - $\triangleright$  W  $\rightarrow \tau v \rightarrow evvv$ : Estimated from MC simulation (PMCS)
  - $\gt Z \rightarrow ee$ : one electron escapes detection
  - ➤ Multi-Jet: one jet misidentified as one electron

Estimated from data

Background	W o  au v	Z  ightarrow ee	MJ
Fraction	$1.668\% \pm 0.004\%$	$1.08\% \pm 0.02\%$	1.018% + 0.065%

➤ Background less than 4%, uncertainty due to the background estimation is negligible

➤ Good agreement between data and prediction at the reconstruction level

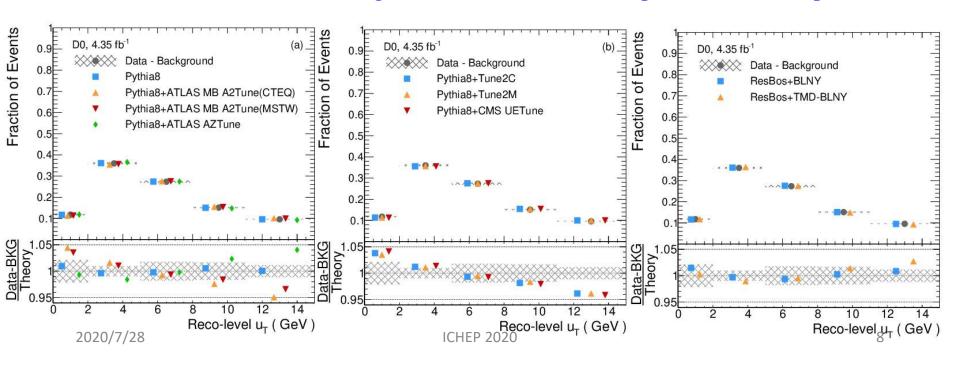


$u_T$ bin	0–2 GeV	2–5 GeV	5–8 GeV	8–11 GeV	11-15  GeV
Fraction of events in the $u_T$ bin	0.1181	0.3603	0.2738	0.1515	0.0963
Total uncertainty	0.0003	0.0005	0.0005	0.0004	0.0003

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#### > PMCS Reweighting

- $\triangleright p_T^W y^W$  2D distribution reweighted to other theory predictions
  - ➤ The default PMCS: ResBos+BLNY
  - Resummation: other non-perturbative functional form (TMD-BLNY)
  - ➤ Parton shower: different Pythia8 tunes from other collaborations
- > Systematic uncertainty estimated by changing parameters in PMCS
  - > Separately estimated with each model
  - > Dominated by the uncertainty due to the hadronic recoil calibration
  - ➤ Bin-by-bin correlation estimated
- > Fraction of events of the background-subtracted data compared to different predictions



- $\geq \chi^2$  calculation
  - $\triangleright$  The  $\chi^2$  value between the data and the reweighted PMCS are calculated
    - $\triangleright$  All 5  $u_T$  bins considered, n.d.f. equals to 4 due to the normalization
    - ➤ The bin-by-bin correlations are taken into account
    - > The uncertainties due to the resummation and the tune are ignored
    - ➤ The PDF uncertainty is negligible

#### **Conclusion**

- ➤ 2 models excluded: Pythia8+ATLAS MB A2Tune+CTEQ6L1
  Pythia8+CMS UE Tune CUETP8S1-CTEQ6L1+CTEQ6L1
- Other Pythia8 tunes except the default one are disfavored
- ➤ ResBos+BLNY and the default Pythia8 agree with the data very well

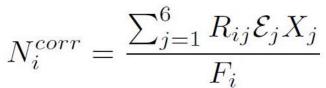
Generator/Model	$\chi^2/\mathrm{ndf}$	<i>p</i> -value	Signif.
RESBOS (Version CP 020811)+BLNY+CTEQ6.6	0.49	$7.41 \times 10^{-1}$	0.33
RESBOS (Version CP 112216)+TMD-BLNY+CT14HERA2NNLO	3.13	$1.39 \times 10^{-2}$	2.46
PYTHIA 8+CT14HERA2NNLO	0.32	$8.63 \times 10^{-1}$	0.17
PYTHIA 8+ATLAS MB A2Tune+CTEQ6L1	12.25	$5.84 \times 10^{-10}$	6.19
PYTHIA 8+ATLAS MB A2Tune+MSTW2008LO	6.17	$5.83 \times 10^{-5}$	4.02
PYTHIA 8+ATLAS AZTune+CT14HERA2NNLO	6.61	$2.60 \times 10^{-5}$	4.21
PYTHIA 8+Tune2C+CTEQ6L1	7.66	$3.61 \times 10^{-6}$	4.63
PYTHIA 8+Tune2M+MRSTLO	7.32	$6.89 \times 10^{-6}$	4.50
PYTHIA 8+CMS UE Tune CUETP8S1-CTEQ6L1+CTEQ6L1	8.80	$4.23 \times 10^{-7}$	5.06

- > Compare to other theory predictions
  - > Two approach to achieve this
    - > Provide an unfolded particle level result
    - > Provide a folding procedure
  - ➤ When the statistical uncertainty dominates
    - $\triangleright$  It can be proven that  $\chi^2_{reco}$  is equal to  $\chi^2_{unfold}$
    - > Same precision between the particle level and the reconstruction level comparisons
  - ➤ When the systematic uncertainty dominates
    - Linear extrapolation to the reconstruction level, non-linear to the particle level
    - ➤ Additional uncertainty on the particle level due to regularization procedure
    - ➤ The precision of the particle level comparison would be reduced
  - > The reduction of the precision would be
    - > Greater when the resolution of the distribution is worse
    - > Smaller when the bin width is enlarged
  - > This is why we chose to provide in our paper a folding procedure
    - ➤ Better precision than the unfolded results

# > Folding procedure

 $\triangleright$  The folded number of events in  $u_T$  bin i

# Response Matrix $R_{ii}$

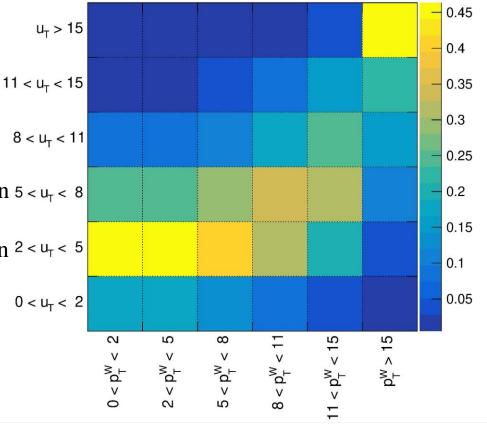


 $X_i$ : the number of events in the  $i^{th} p_T^W$  bin

 $\mathcal{E}_i$ : the efficiency correction in the  $i^{th} p_T^W$  bin 5 < u<sub>T</sub> < 8

 $F_i$ : the fiducial correction in the  $i^{th}$   $u_T$  bin

 $N_i^{corr}$ : the number of events in the  $i^{th}$   $u_T$  bin  $^{2 < u_T < 5}$ 



 $\triangleright$  Response Matrix  $R_{ij}$ :

The probability for the events in one  $p_T^W$  bin to be reconstructed into different  $u_T$  bins

The instruction to calculate the covariance matrix and details of the whole folding procedure are introduced in the appendix of the paper (arXiv:2007.13504)

https://www-d0.fnal.gov/Run2Physics/WWW/results/final/EW/E20A/

# > Summary

- First Tevatron Run II  $p_T(W)$  measurement at  $\sqrt{s} = 1.96$  TeV
  - $\triangleright$  Focus on low  $p_T(W)$  region
- ➤ The background subtracted data is compared to different predictions after PMCS simulation on the reconstruction level
  - > Uncertainty dominated by that due to the hadronic recoil calibration
  - ➤ Model uncertainty ignored
  - > PDF uncertainty negligible
- > Two models are excluded
  - ➤ Pythia8+ATLAS MB A2Tune+CTEQ6L1
  - ➤ Pythia8+CMS UE Tune CUETP8S1-CTEQ6L1+CTEQ6L1
- ➤ A folding procedure is provided for the comparison with other models
  - > Better precision than the unfolded results
  - ➤ Model dependence tested to be negligible
  - $\triangleright \chi^2$  difference smaller than the impact from the data fluctuation

#### Link to the paper:

https://www-d0.fnal.gov/Run2Physics/WWW/results/final/EW/E20A/ https://arxiv.org/abs/2007.13504

# **Backup**

➤ Collins-Soper-Sterman (CSS) resummation formalism

Production of a vector boson in the collision of two hadrons

$$\frac{d\sigma(h_1h_2 \to VX)}{dQ^2dQ_T^2dy} = \frac{1}{(2\pi)^2}\delta(Q^2 - M_V^2) \int d^2b \ e^{i\vec{Q}_T \cdot \vec{b}} \widetilde{W}_{j\bar{k}}(b, Q, x_1, x_2) + Y(Q_T, Q, x_1, x_2)$$

b: impact parameter

 $\triangleright$  the nonperturbative terms in the form of an additional factor  $\widetilde{W}_{j\bar{k}}^{NP}(b,Q,x_1,x_2)$ 

$$\widetilde{W}_{j\bar{k}} = \widetilde{W}_{j\bar{k}}^{pert} \widetilde{W}_{j\bar{k}}^{NP}$$

Brock-Landry-Nadolsky-Yuan form

$$\widetilde{W}_{j\bar{k}}^{NP}(b,Q,x_1,x_2) = \exp\left(-g_1 - g_2 \ln\left(\frac{Q}{2Q_0}\right) - g_1 g_3 \ln(100x_1x_2)\right) b^2$$

CSS: Nucl. Phys. B250, 199 (1985)

BLNY: Phys. Rev. D 67, 073016 (2003)

# ➤ Response Matrix *R*

 $\triangleright$  The probability for the events in one  $p_T^W$  bin to be reconstructed into different  $u_T$  bins

$$R_{ij} = P(\mathcal{N}_i | \mathcal{X}_j)$$

 $\mathcal{N}_i$ : the case that  $u_T$  is in the  $i^{th}$  bin  $\mathcal{X}_i$ : the case that  $p_T^W$  is in the  $i^{th}$  bin

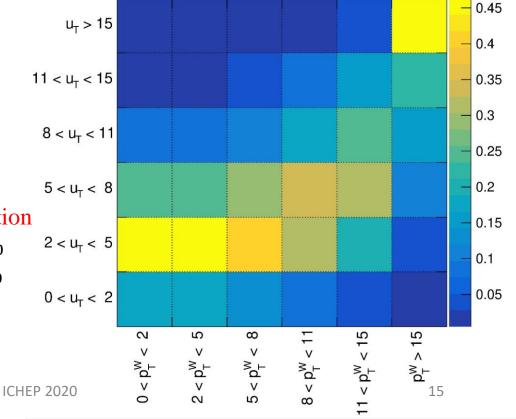
 $N_i$ : the number of events in the  $i^{th}$   $u_T$  bin  $X_i$ : the number of events in the  $i^{th}$   $p_T^W$  bin

$$N_i = \sum_j R_{ij} X_j$$

- $\triangleright$  Purity  $R_{ii}$ :
- The probability for the events in one  $p_T(W)$  bin to be reconstructed into the same  $u_T$  bin

Low purity caused by limited resolution Maximum Purity:  $max(R_{ii}) \sim 45\%$ 

Minimum Purity:  $min(R_{ii}) \sim 16\%$ 



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- Comparison to other theory models
  - ➤ Provide an unfolded particle-level result
    - > Directly compare to the theory prediction on the particle level
    - ➤ A common procedure widely used by other collaboration
  - > Provide a folding procedure
    - ➤ Account for the detector response and resolution effects
    - ➤ Compare to the background-subtracted data on the reconstruction level
  - > For these two approaches,
    - > Fiducial selection should be defined

$$p_T^e > 25 \text{ GeV}, |\eta^e| < 1.05$$
  
 $p_T^v > 25 \text{ GeV}, 50 < m_T < 200 \text{ GeV}$ 

- > Basic inputs are the same
- ➤ Basic inputs estimated from MC simulations
  - $\triangleright$  Fiducial Correction:  $u_T$  distribution within fiducial volume
  - Response Matrix: correct detector effects and migration
  - ➤ Efficiency Correction

- > Statistical uncertainty dominated situation
  - $\triangleright$  The  $\chi^2$  calculated on the reconstruction level:

$$\chi_{reco}^{2} = \left(N^{Data} - N^{Pred}\right)^{T} \Sigma^{-1} (N^{Data} - N^{Pred})$$

- $\triangleright$   $\Sigma$  is the bin-by-bin covariance matrix of the statistical uncertainty on the reconstruction level, which should be a diagonal matrix
- $\triangleright$  If the data and the prediction is rotated by a matrix M, which is the unfolding matrix

$$\chi_{unfold}^{2} = (MN^{Data} - MN^{Pred})^{T} \Sigma'^{-1} (MN^{Data} - MN^{Pred})$$

 $\triangleright \Sigma'$  is the bin-by-bin covariance matrix of the statistical uncertainty on the particle level

$$\Sigma' = M\Sigma M^T$$

$$\begin{split} \chi_{unfold}^2 &= \left(MN^{Data} - MN^{Pred}\right)^T \Sigma'^{-1} (MN^{Data} - MN^{Pred}) \\ &= \left(N^{Data} - N^{Pred}\right)^T M^T M^{T^{-1}} \Sigma^{-1} M^{-1} M \left(N^{Data} - N^{Pred}\right) \\ &= \left(N^{Data} - N^{Pred}\right)^T \Sigma^{-1} \left(N^{Data} - N^{Pred}\right) \\ &= \chi_{reco}^2 \end{split}$$

> Same precision between the particle level and the reconstruction level comparisons

- > Systematic uncertainty dominated situation
  - $\triangleright$  A simple unfolding procedure is to use  $R^{-1}$  as the unfolding matrix, M
    - $\triangleright$  The response matrix of the systematic variation,  $R' = R + \Delta R$
    - ➤ The uncertainty on the reconstruction level, a linear transformation

$$U_{reco} = R'X^{Pred} - RX^{Pred} = \Delta RX^{Pred}$$

➤ The uncertainty on the unfolding particle level, a non-linear transformation due to the inversion of the covariance matrix

$$U_{unfold} = R'^{-1}N^{Data} - R^{-1}N^{Data} = (R'^{-1} - R^{-1})N^{Data}$$

- ➤ The precision of the particle level comparison would be reduced
- > Unfolding method with a regularization scheme
  - ➤ No longer an unbiased estimation
  - > The reduction due to the non-linear transformation would be smaller
  - ➤ A model input or a regularization model required
  - ➤ Additional uncertainty due to the input model or the regularization
- ➤ The reduction of the precision:
  - > Greater when the purity is lower
  - > Smaller when the bin width is enlarged

# > Folding procedure

- > We choose to provide a folding procedure instead of an unfolded result, because
  - ➤ Better precision on the reconstruction level than the particle level
  - ➤ Not affected by the low purity problem
  - > The rise and hence the shape of the spectrum can be resolved
  - ➤ Avoid arbitrary definitions of the addition unfolding uncertainties and correlations
- $\triangleright$  The fraction of events in  $u_T$  bin i,  $\mathcal{N}_i$

$$\mathcal{N}_i = \frac{N_i^{corr}}{\sum_{j=1}^5 N_j^{corr}}, \quad N_i^{corr} = \frac{\sum_{j=1}^6 R_{ij} \mathcal{E}_j X_j}{F_i}$$

 $X_i$ : the number of events in the  $i^{th} p_T^W$  bin

 $\mathcal{E}_i$ : the efficiency correction factor in the  $i^{th}$   $p_T^W$  bin

 $F_i$ : the fiducial correction in the  $i^{th}$   $u_T$  bin

 $N_i^{corr}$ : the number of events after all the correction in the  $i^{th}$   $u_T$  bin

 $\triangleright \mathcal{N}_i$  is the folded result

- > The systematic uncertainty and its estimation
  - ➤ In total, 11 systematic variations provided
    - First 10 variations for the uncertainty due to the hadronic recoil calibration
    - > 5 positive change variations + 5 negative change variations
    - Last variation for the uncertainty due to the electron energy and efficiency
  - $\triangleright$  The covariance matrix of the systematic uncertainty,  $\Sigma^{(Syst.)}$ 
    - > Different variations uncorrelated from each other
    - Each variation, correlated bin-by-bin
    - > The positive change variations and negative change variations are averaged

$$\Sigma^{(Syst.)} = \frac{\sum_{k=1}^{10} \Sigma^{(k)}}{2} + \Sigma^{(11)}$$

 $\triangleright \Sigma^{(k)}$  is the covariance matrix of the  $k^{th}$  variation, its element  $\Sigma_{ij}^{(k)}$ 

$$\Sigma_{ij}^{(k)} = (\mathcal{N}_i^{(k)} - \mathcal{N}_i) \times (\mathcal{N}_j^{(k)} - \mathcal{N}_j)$$

 $\triangleright$  The covariance matrix used in the  $\chi^2$  calculation,  $\Sigma$ 

$$\Sigma = \Sigma^{(Data.\,Stat.)} + \Sigma^{(MC.\,Stat.)} + \Sigma^{(Syst.)}$$

- > The model dependence
  - > The folding inputs
    - $\triangleright$  The efficiency correction,  $\mathcal{E}_i$ , derived with  $p_T^W$  dependence, Model independent
    - $\triangleright$  The response matrix,  $R_{ij}$ , derived with  $u_T p_T^W$  dependence, Model independent
    - $\triangleright$  The fiducial correction,  $F_i$ , derived with  $u_T$  dependence, Model dependent
    - $\triangleright$  Check by changing the peak and the width of  $p_T^W$  distribution by 20%
    - > The impact is negligible compared to the total uncertainty of the folded result
  - ➤ The systematic uncertainty and its correlation
    - > Estimated by systematic variations

$$U_{reco} = R'X^{Pred} - RX^{Pred} = \Delta RX^{Pred}$$

- > The uncertainty and the covariance matrix should be model dependent
- ➤ The basic inputs from all the variations are provided

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