

New results from the CUORE experiment

Andrea Giachero

University & INFN of Milano-Bicocca



Nuclear Double Beta Decay

Very rare radioactive decay where the charge of the nucleus changes by two units

- $(A, Z) \rightarrow (A, Z + 2) + 2e^- + 2\bar{\nu}_e \quad (2\nu\beta\beta)$

Allowed in the Standard Model for even-even nuclei and already observed with $T_{1/2} > 10^{18}$ y

- $(A, Z) \rightarrow (A, Z + 2) + 2e^- \quad (0\nu\beta\beta)$

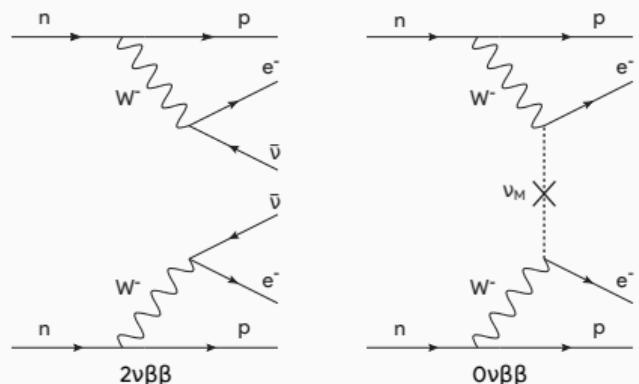
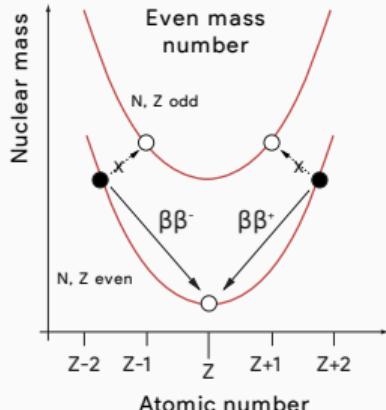
Forbidden by Standard Model ($\Delta L = 2$), expected with $T_{1/2} > 10^{25}$ y

- **L-violation**, i. e. creation of a pair of electrons

- the discovery of $0\nu\beta\beta$ decay would imply the violation of the SM predicted accidental symmetries

- Key tool to **study neutrinos**

- Majorana or Dirac nature ($\nu = \bar{\nu}$ or $\nu \neq \bar{\nu}$)
- ν mass scale and ordering (inverted or normal hierarchy)
- CP -violation in the lepton sector



- $0\nu\beta\beta$ is a nuclear process:

- 2nd order transition: $(A, Z) \rightarrow (A, Z + 2)$
- even-even nuclei: β -decay is forbidden

- half-life expression can be factorized as:

$$\left[t_{1/2}^{0\nu} \right]^{-1} = G_{0\nu} |\mathcal{M}|^2 F$$

- $G_{0\nu}$: Phase Space Factor (PSF, atomic physics)
 - \mathcal{M} : Nuclear Matrix Element (NME, nuclear physics)
 - F : New Physics Term
- ⇒ in case of 3 light neutrino exchange:

$$F = |m_{\beta\beta}|^2$$

where $m_{\beta\beta}$ is the Majorana mass

- $m_{\beta\beta}$ is the key quantity in the $0\nu\beta\beta$

- Absolute value of the ee-entry of the neutrino mass matrix

$$\bullet |m_{\beta\beta}| = \left| \sum_{i=1,2,3} e^{i\xi_i} |U_{ei}^2| m_i \right|$$

- U can be identified with the mixing matrix of the oscillation analysis

F. Capozzi, et al., arXiv:1804.09678 [hep-ph]

F. Capozzi, et al., Phys. Rev. D 95 (2017) 096014

$$\bullet U \equiv U|_{\text{osc.}} \cdot \text{diag} \left(e^{-i\xi_1/2}, e^{-i\xi_2/2}, e^{i\phi - i\xi_3/2} \right)$$

- 1 CP-violating + 3 Majorana phases

- Only two phases play a *physical* role

$$|m_{\beta\beta}| = \left| e^{i\alpha_1} \cos^2 \theta_{12} \cos^2 \theta_{13} m_1 + e^{i\alpha_2} \cos^2 \theta_{13} \sin^2 \theta_{12} m_2 + \sin^2 \theta_{13} m_3 \right|$$

Experimental search for $0\nu\beta\beta$

- **Detection:** energy (track) of the 2 emitted e^-
 - monochromatic peak at the Q -value of the decay ($Q_{\beta\beta}$)
 - smearing due to finite energy resolution
- **Observable:** decay half-life of the isotope, $t_{1/2}^{0\nu}$
 - in the case of a peak in the energy spectrum

$$t_{1/2}^{0\nu} = \ln 2 \cdot T \cdot \varepsilon \cdot \frac{N_{\beta\beta}}{N_{\text{peak}}} \quad \left(\frac{\delta t_{1/2}^{0\nu}}{t_{1/2}^{0\nu}} = \frac{\delta N_{\text{peak}}}{N_{\text{peak}}} \right)$$

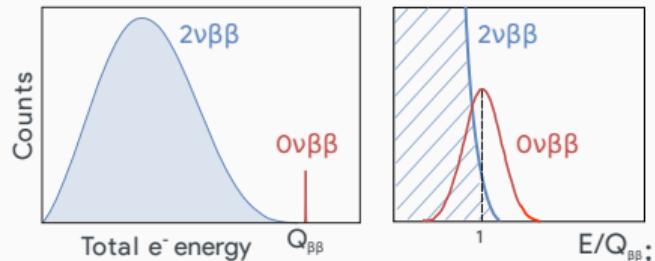
- if no peak is detected, the **sensitivity** corresponds to the maximum signal that can be hidden by the background fluctuations $n_B = \sqrt{MTB\Delta E}$

$$S_{1/2}^{0\nu} = \ln 2 \cdot T \cdot \varepsilon \cdot \frac{n}{n_\sigma \cdot \sigma_B} = \ln 2 \cdot \frac{1}{n_\sigma} \cdot \frac{xN_A}{M_A} \cdot \eta \cdot \varepsilon \sqrt{\frac{MT}{B\Delta E}}$$

($M T$ is the total exposure)

M = detector mass T = measuring time B = background level ΔE = energy resolution

ε = detection efficiency η = candidate isotopic abundance

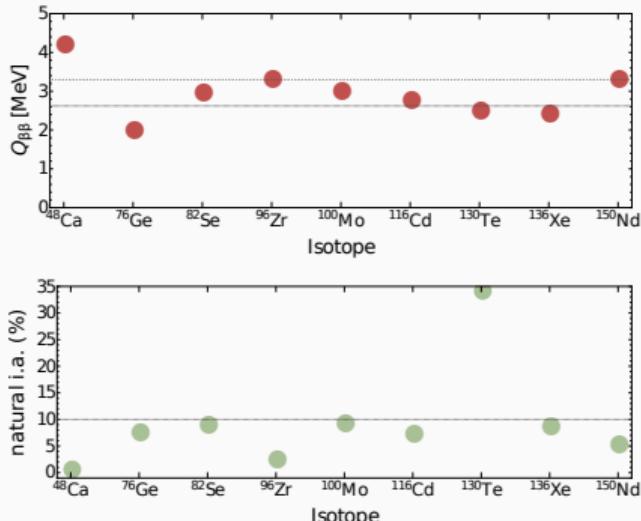


- **Good energy resolution ΔE**
 - only protection against $2\nu\beta\beta$ spectrum tail
 - $R_{0\nu/2\nu} \propto \left(\frac{Q_{\beta\beta}}{\Delta}\right)^6 \frac{t_{1/2}^{2\nu}}{t_{1/2}^{0\nu}}$
- **Very low background B**
 - underground location + shielding
 - radio-pure materials for detector and surrounding parts
 - analysis rejection techniques
- **Large isotope mass M**
 - present: some tens up to hundreds of kg
 - tonnes required to cover the IH region

Choice of the isotope

- Value of $Q_{\beta\beta}$ → influences the background
 - $> 2.6 (> 3.3)$ MeV end-point of main γ s (β s)
 - avoid radioactivity peak position
 - < 4.01 MeV starting point of main α s
- High isotopic abundance
 - ease of material enrichment
(technologically + economically)
- Availability of the isotope
 - tonnes required for future $0\nu\beta\beta$ experiments
→ high cost + large procurement time
- Compatibility with a detection technique

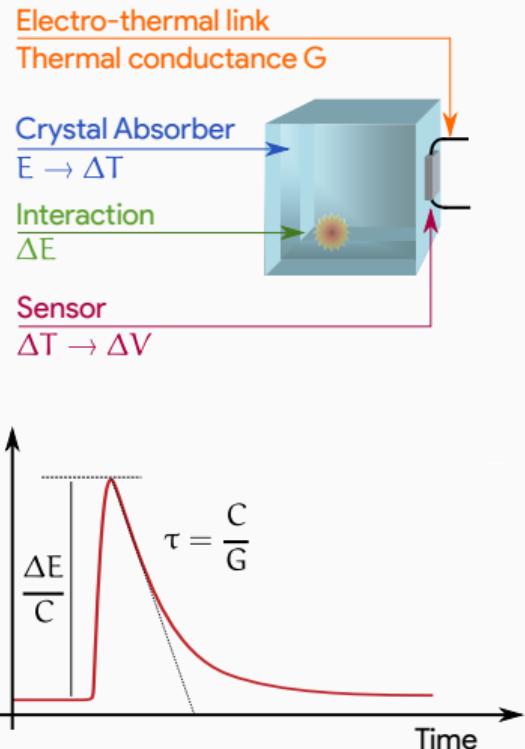
Most suitable **isotope + detector** combination



Promising isotopes

^{48}Ca	^{76}Ge	^{82}Se
^{96}Zr	^{100}Mo	^{116}Cd
^{130}Te	^{136}Xe	^{150}Nd

- Complete energy thermalization:
ionization, excitation \Rightarrow heat \Rightarrow calorimetry;
- $\Delta T = \Delta E/C$ where ΔE is the released energy and C the total thermal capacity;
 - Absorber with very low thermal capacity: $C \downarrow \Rightarrow \Delta T \uparrow$;
 - Debye law for superconductors below T_C and dielectric: $C \propto (T/\Theta_D)^3$;
 - A very low temperature is needed:
 $T \downarrow \Rightarrow C \downarrow \Rightarrow \Delta T \uparrow \Rightarrow (T = 10 \div 100 \text{ mK})$;
- Limit to energy resolution \Rightarrow statistical fluctuation of internal energy
 $\Delta E_{rms} = \sqrt{k_B T^2 C}$;
- $\Delta T(t) = \frac{\Delta E}{C} e^{-t/\tau}$ with $\tau = \frac{C}{G}$ and G thermal conductance.



Decay under study: $^{130}\text{Te} \rightarrow ^{130}\text{Xe} + \text{e}^-$

Absorber

Why ^{130}Te :

- High isotopic abundance: 34.167% \Rightarrow no need for enrichment;
- Q-value (2527.5 keV)
almost above the natural γ background (> 2615 eV)
below the natural α background (< 4.01 MeV);

Why TeO₂:

- Easy to grow big crystals with low radioactive contaminations;
- Good mechanical properties;
- Low heat capacity (dielectric and diamagnetic).

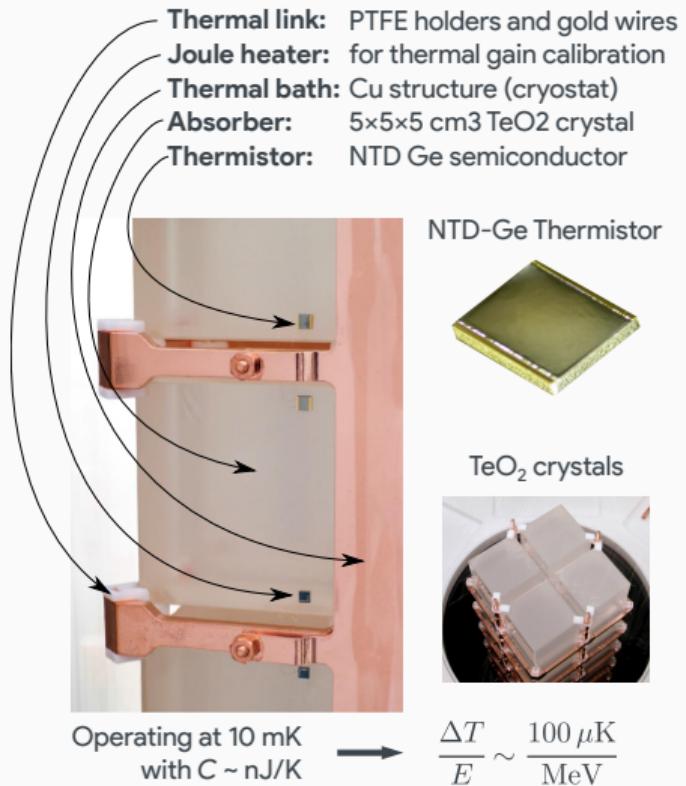
Sensor

NTD-Ge thermistor:

- Germanium crystal doped by thermal neutrons;
- Operate in VRH (Variable Range Hopping) regime;

$$R(T) = R_0 e^{\left(\frac{T_0}{T}\right)^{\frac{1}{2}}}$$

$$R_0 = 1\Omega \quad , \quad T_0 = 4\text{ K}$$



Cryogenic Underground Observatory for Rare Events

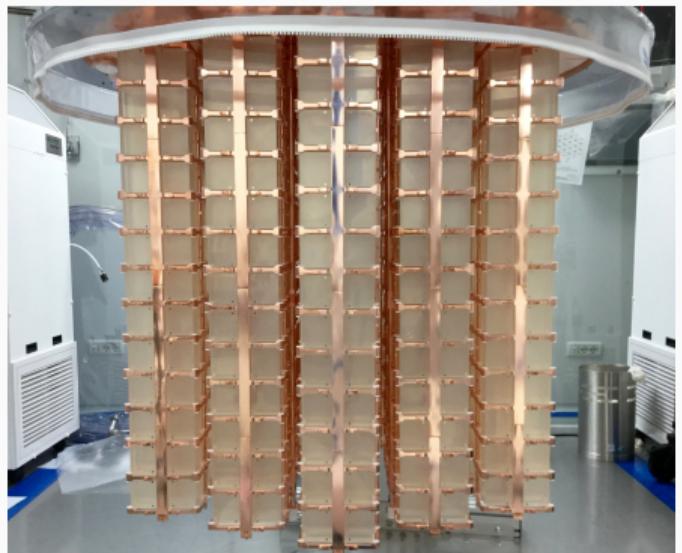
The larger thermal detector array ever built

CUORE Detector:

- Custom Cryostat to reach $T_0 \simeq 10$ mK;
- Array of 988 TeO₂ crystals;
- Total mass of 742 kg;
- 19 towers, 13 floors, 4 crystal for each floor;
- TeO₂ absorber $5 \times 5 \times 5$ cm³ (750 g each);
- Stringent radiopurity control on materials and assembly;
- Target energy resolution: 5 keV @ 2615 (FWHM);
- Background aim: 10^{-2} count/keV/kg/yr;
- $t_{1/2}^{0\nu}$ sensitivity in 5 years:

$$t_{1/2} \geq 9.0 \cdot 10^{25} \text{ yr} @ 90\% \text{ C.L}$$

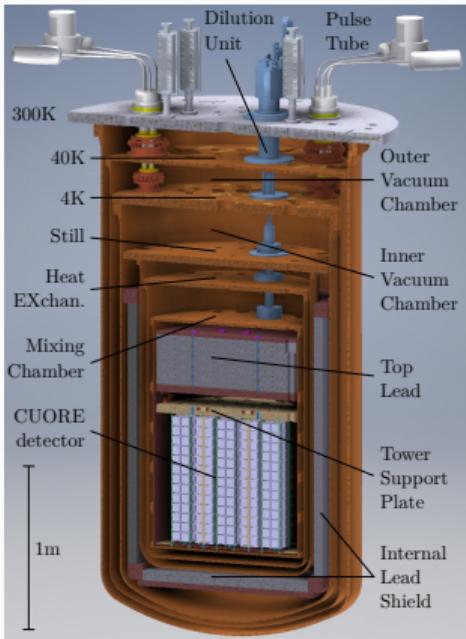
$$m_{\beta\beta} \leq (45 - 211) \text{ meV}$$



- All 19 towers installed between July-August 2016;
- First cool down in December 2016;
- Base temperature 7 mK in 26/01/2017;

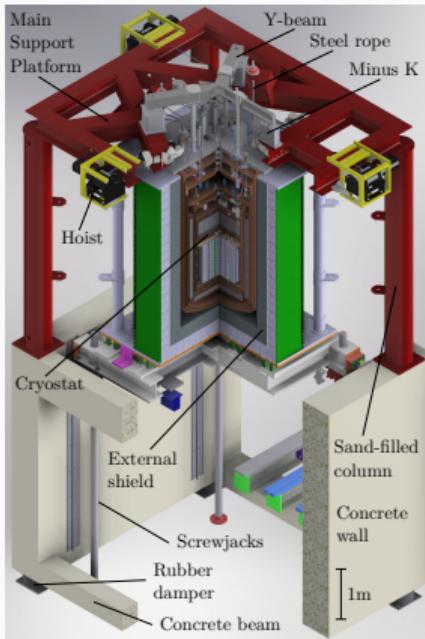
first event observed the same day;

Rendering of the CUORE cryostat



Cryogenics 102 (2019) 9-21
 Cryogenics 93 (2018) 56-65

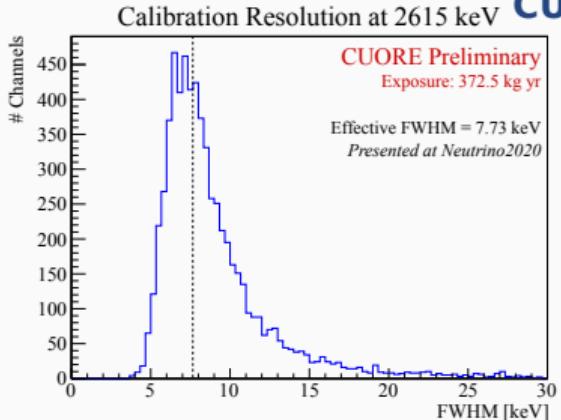
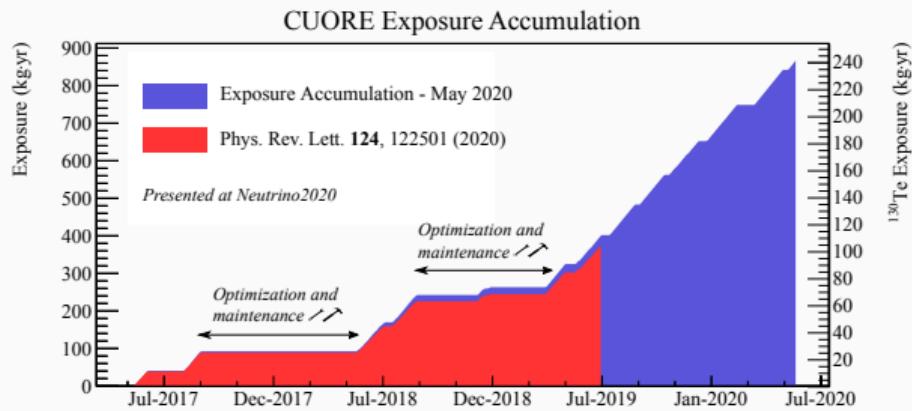
Rendering of the cryostat support structure



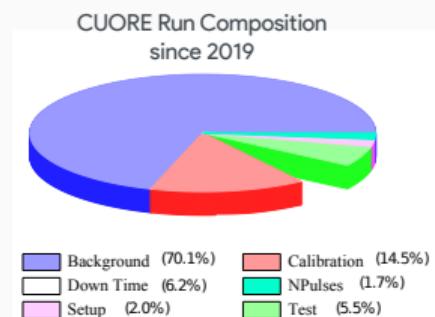
The coldest cubic meter in the known Universe

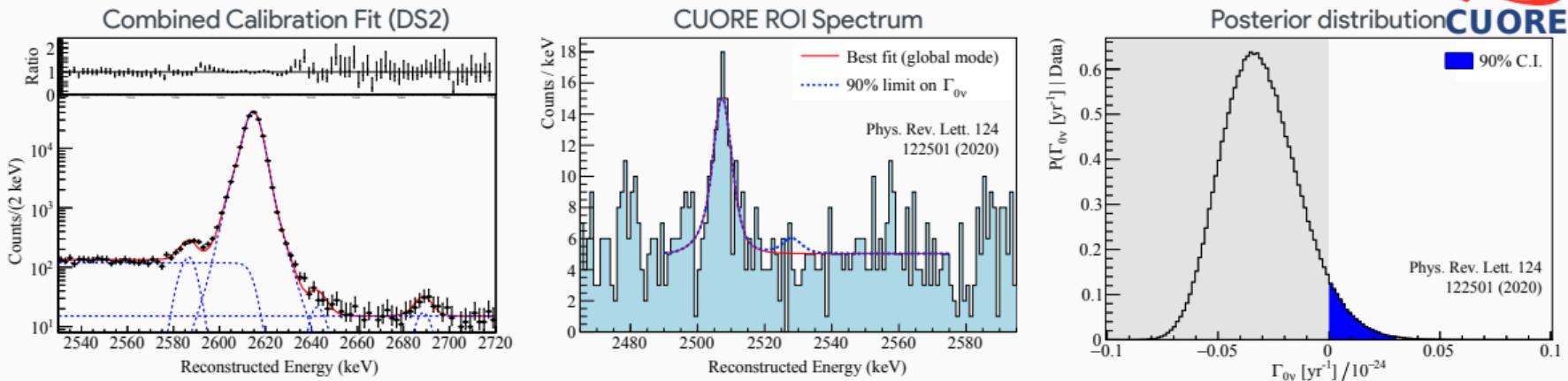
- Powerful multistage cryogen-free cryostat
 - Precooled by a fast cooling system;
 - Pulse Tubes (PTs),
 - $^3\text{He}/^4\text{He}$ dilution refrigerator
- cooling power: $5 \mu\text{W}$ at 10 mK
- Cooldown time around 1 month;
- Cryogenic vessels and shielding:
 - 13 tonnes $< 4 \text{ K}$
 - 5 tonnes $< 50 \text{ mK}$
 - 1500 kg @ 10 mK (detectors + materials)
- External Shielding:
 - 30 cm lead
 - 18 cm polyethylene + 2 cm borated material
- Mechanical vibration isolation;
- Active noise cancellation;

CUORE data taking



- 7 dataset from April 14, 2017;
- 900(min)-950(max) (over 988) valid calorimeters;
- 7.73 keV FWHM @ 2615 keV;
- 7.0 keV FWHM @ Q-value;
- 372.5 kg·yr of TeO₂ exposure accumulated;
- ~ 10 keV of trigger threshold;
- 40 keV of analysis threshold;





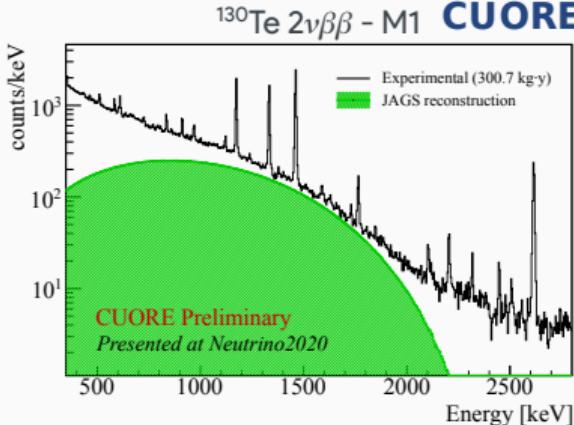
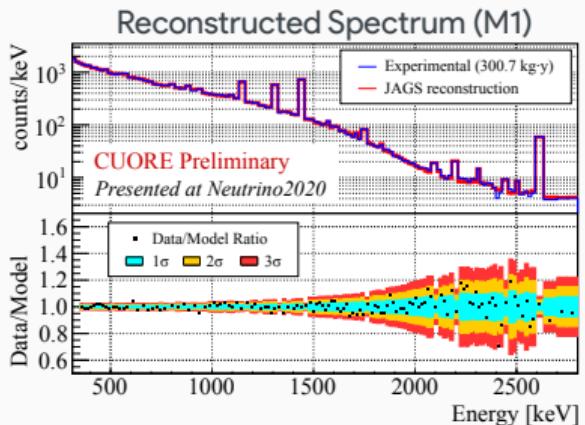
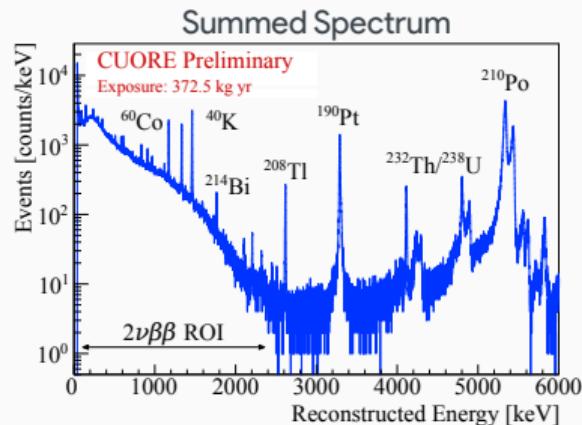
- Bayesian Analysis (with BAT)
- Likelihood model: flat continuum (BI), and posited peak for $0\nu\beta\beta$ (rate), peak for ^{60}Co (rate+position)
- Unbinned fit on physical range (rates non negative), uniform prior on $\Gamma_{0\nu}$;
- Systematics: repeat fits with nuisance parameters, allow negative rates (<0.4% impact on limit)

- No evidence for $0\nu\beta\beta$ decay of ^{130}Te ;
- Background index in the ROI:

$$B = (0.0138 \pm 0.0007) \text{ cnts/keV/kg/yr}$$
- $t_{1/2}^{0\nu}$ sensitivity (90% CL):

$$t_{1/2}^{0\nu} > 3.2 \cdot 10^{25} \text{ yr}$$
- Limit on the neutrino mass:

$$m_{\beta\beta} < (75 - 350) \text{ meV}$$



- Reconstruct CUORE continuum background;
- GEANT4 simulation + measured detector response function to produce expected spectra;
- 62 sources considered, Bayesian fit with flat priors (except for muons);
- Exploit coincidences and detector self-shielding to constrain location of sources

- CUORE-0 ([Eur. Phys. J. C77 \(2017\) 1, 13](#))

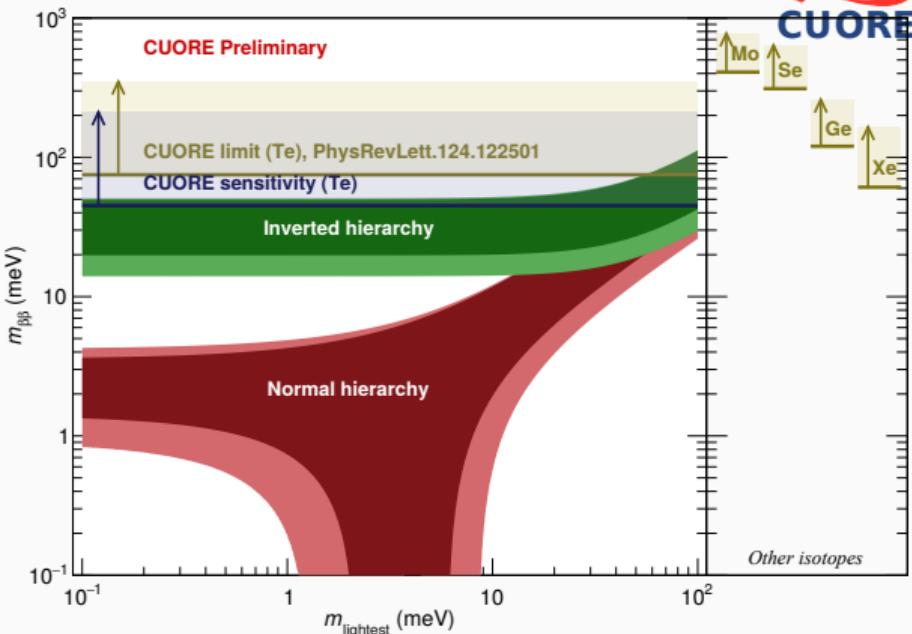
$$t_{1/2}^{2\nu} = [8.2 \pm 0.2_{(\text{stat})} \pm 0.6_{(\text{sys})}] \cdot 10^{20} \text{ yr}$$

- CUORE (preliminary)

$$t_{1/2}^{2\nu} = \left[7.71_{-0.06 \text{ (stat)}}^{-0.08} \quad -0.17_{-0.15 \text{ (stat)}} \right] \cdot 10^{20} \text{ yr}$$

- reduction of a factor 3 on both statistical and systematic uncertainties

CUORE Collaboration



- CUORE data taking continues smoothly
 - ⇒ stable conditions allowed continued data taking with minimal onsite activity during recent lockdowns;
- CUORE will reach the target sensitivity in a few years, approaching the Inverted Hierarchy Region
 - ⇒ next unblinding is foreseen at 1 tonne· year

- The study of frontier subjects in neutrino physics require the use of frontier detectors;
- Observation of $0\nu\beta\beta$ decay would prove that neutrinos are Majorana particle
- CUORE accumulated over 400 kg·yr exposure, new data is being acquired;
- ... and will approach the Inverted Hierarchy Region in a few years;
- Successfull performance of CUORE justifies future developments of cryogenics calorimeters;

⇒ CUPID program: Cuore Upgrade with Particle IDentinficaion

(see Davide Chiesa's talk just after this)