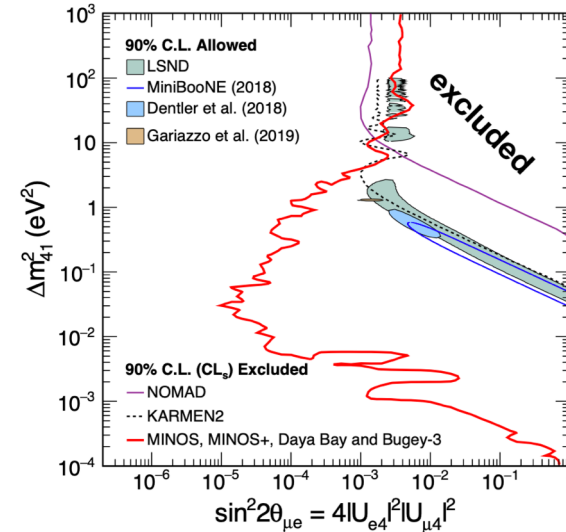
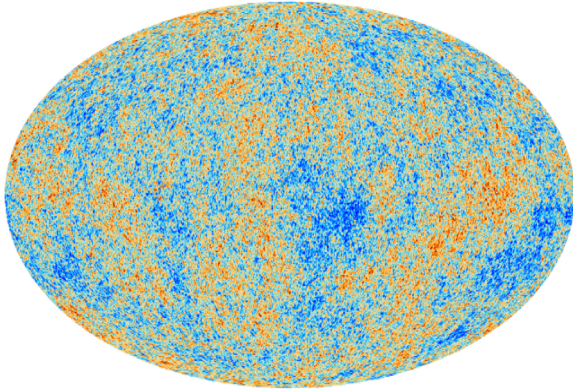


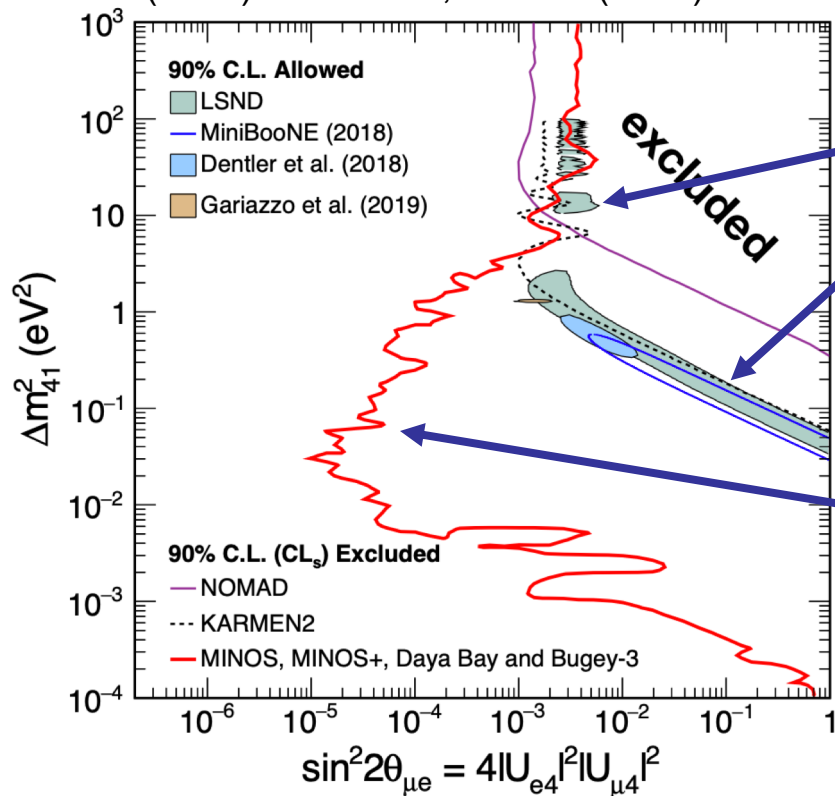
# Direct comparison of sterile neutrino constraints from cosmological data and oscillation data in a 3+1 model



Matthew Adams, Fedor Bezrukov, Jack Elvin-Poole, **Justin Evans**,  
Pawel Guzowski, Brían Ó Fearraigh, Stefan Söldner-Rembold

# Sterile neutrino tensions

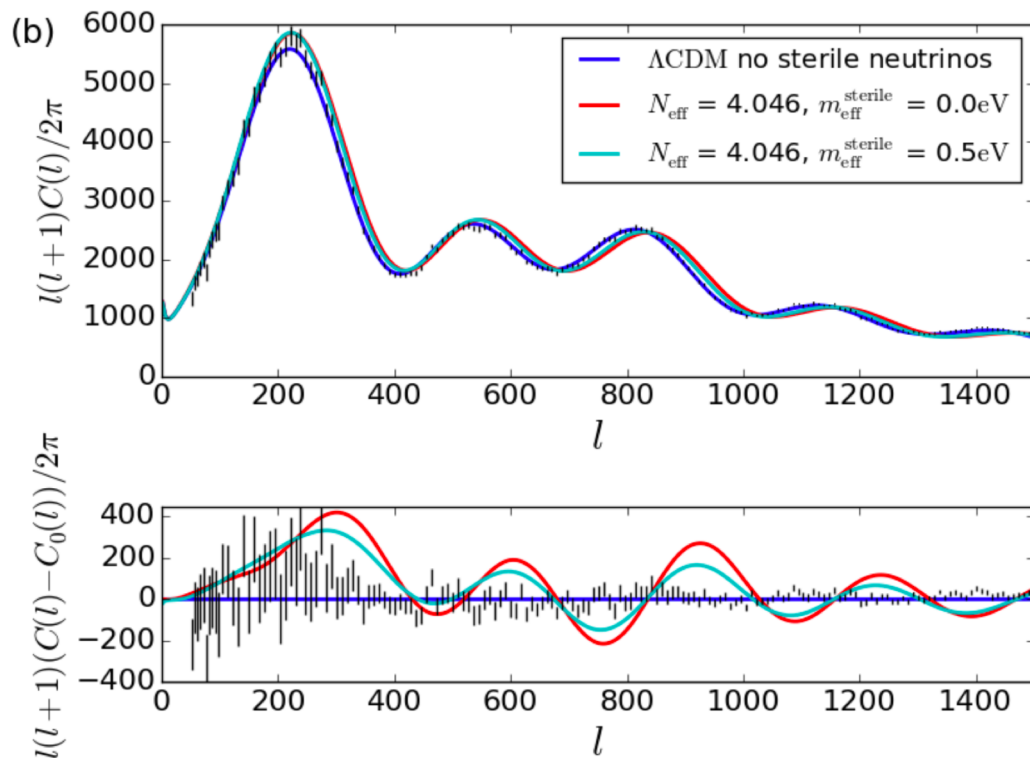
arXiv:2002.00301, PRD **95**, 072006  
(2017) & PRL **122**, 091803 (2019)



$\nu_\mu \rightarrow \nu_e$  appearance seen by LSND and MiniBooNE consistent with a sterile neutrino

MINOS, Daya Bay and Bugey exclude most of the allowed parameter space with  $\nu_\mu$  and  $\nu_e$  disappearance

# Cosmology and sterile neutrinos



Thermalised sterile neutrinos add an extra relativistic degree of freedom observable in the CMB power spectrum

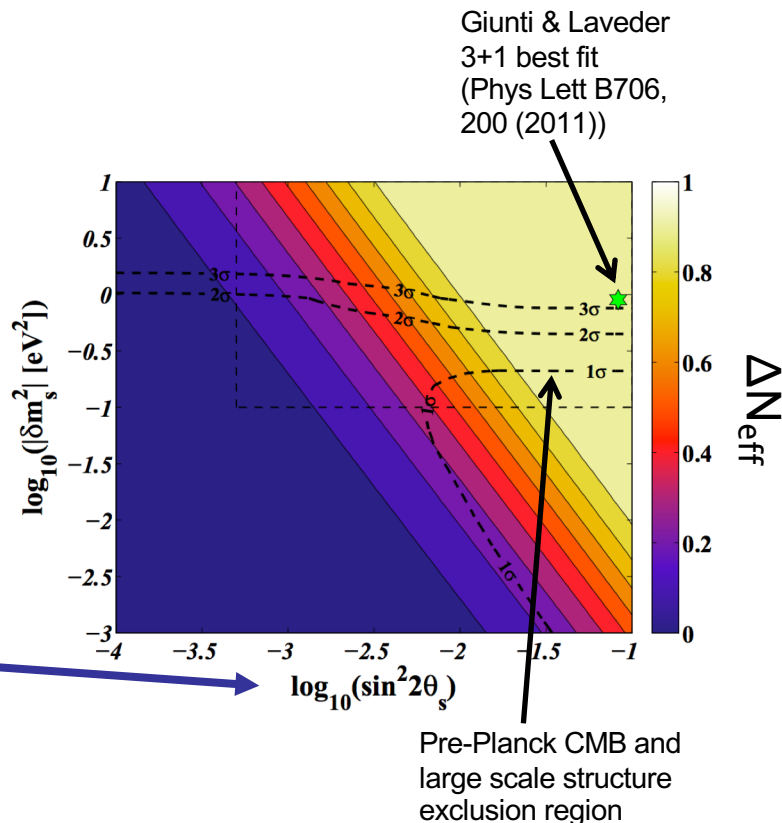
$$m_{\text{eff}}^{\text{sterile}} = \left(\frac{T_s}{T_\nu}\right)^3 m_4^{\text{thermal}} = (\Delta N_{\text{eff}})^{3/4} m_4^{\text{thermal}}$$

# Relating particle physics to cosmology

Graph from S. Hannestad *et al.*, Cosmol Astropart. Phys. **2012**, 025 (2012), arXiv:1204.5861  
Code available as LASAGNE

Larger sterile mass splitting  
increases temperature of  
thermalization  $\Rightarrow$  greater  $\Delta N_{\text{eff}}$

Larger mixing angle allows  
a higher thermalisation rate  
 $\Rightarrow$  greater  $\Delta N_{\text{eff}}$



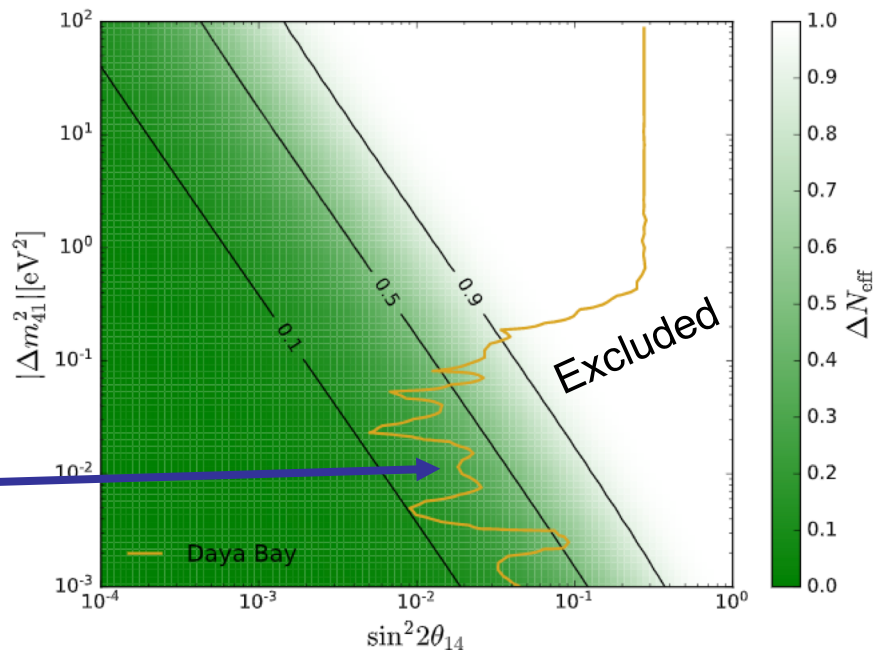
# $\bar{\nu}_e$ disappearance

1+1 model.

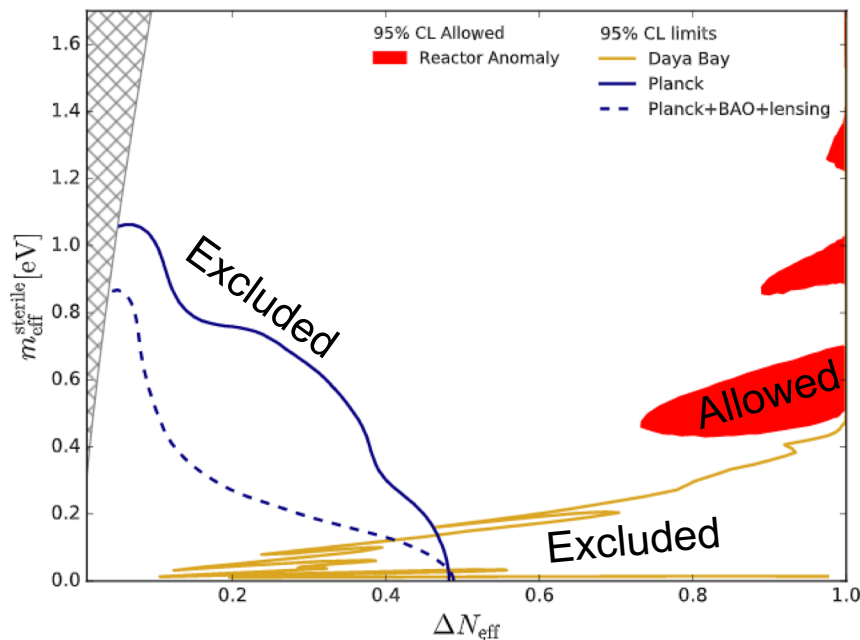
Only  $\nu_e$  flavour mixes into  $m_4$ .

Only non-zero mixing angle is  $\theta_{14}$ .

Daya Bay reactor  $\bar{\nu}_e$   
disappearance search



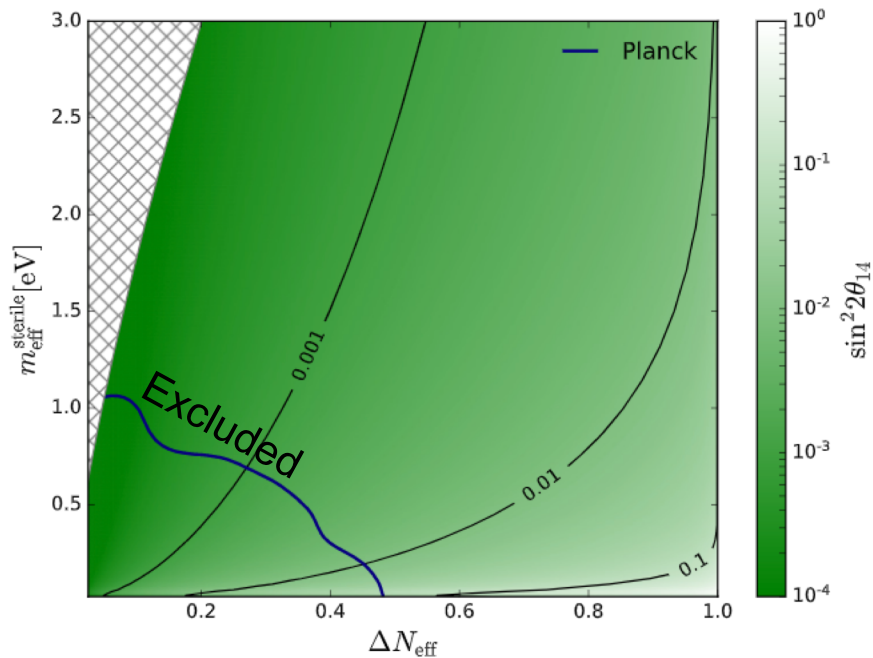
# $\bar{\nu}_e$ disappearance



Gallium anomaly allowed region  
and Daya Bay exclusion region  
expressed in cosmological  
parameter space

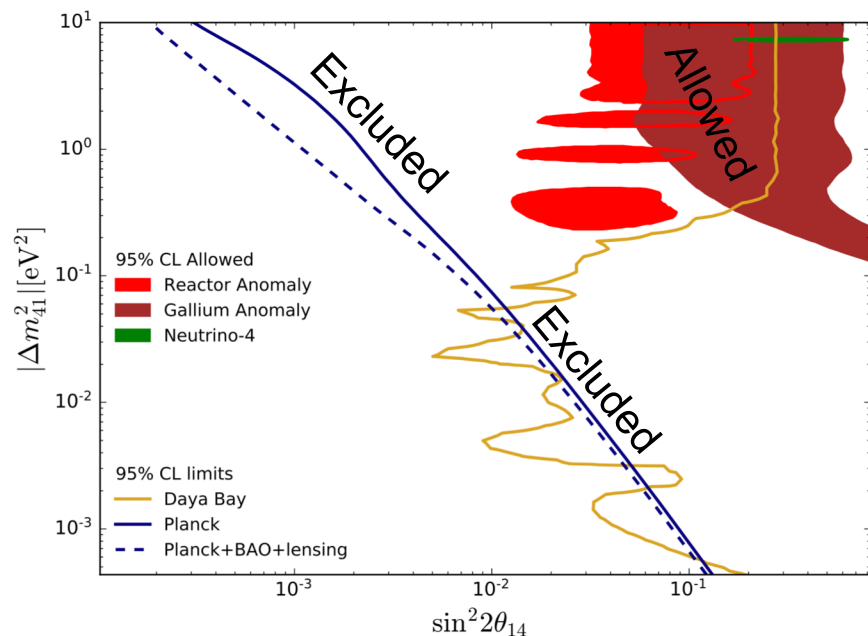
Compared to Planck exclusion

# $\bar{\nu}_e$ disappearance



Can also convert the Planck exclusion into neutrino-oscillation parameter space

# $\bar{\nu}_e$ disappearance

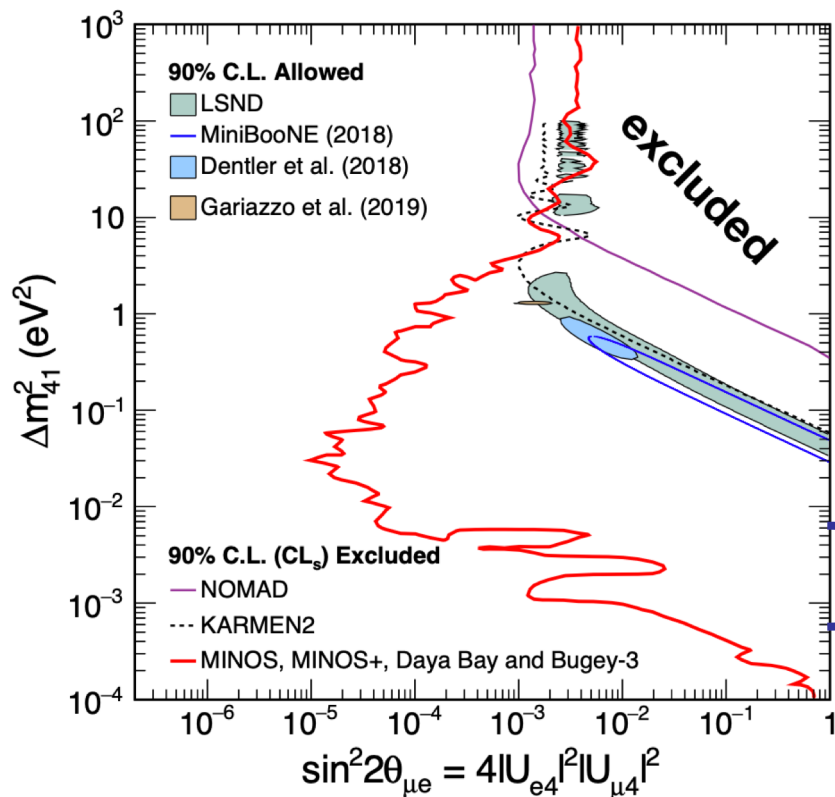


Can also convert the Planck exclusion into neutrino-oscillation parameter space

Including here the Neutrino-4 and gallium anomaly allowed regions



# A 3+1 model for $\nu_e \rightarrow \nu_\mu$ appearance

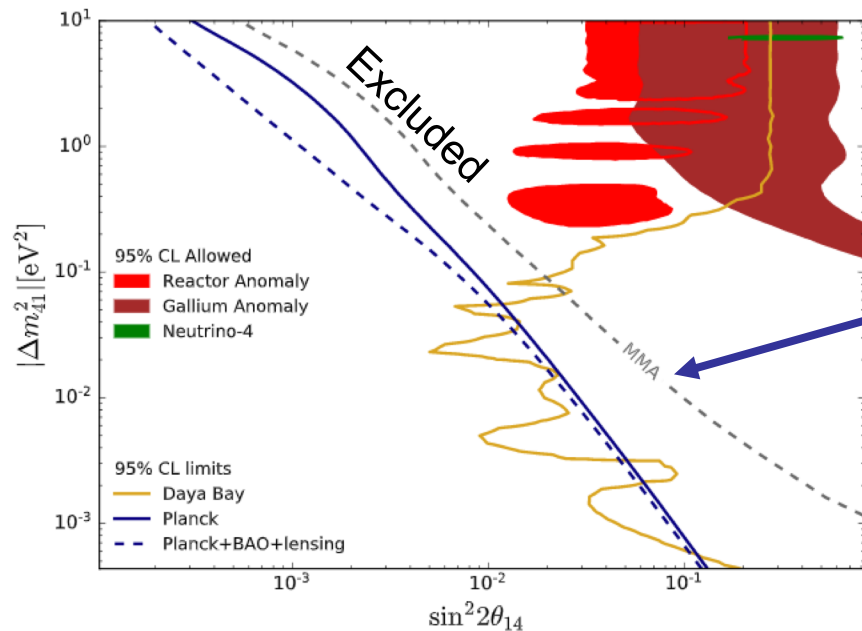


$$\sin^2(2\theta_{\mu e}) \equiv \sin^2(2\theta_{14})\sin^2(\theta_{24})$$

Requirement to solve quantum kinetic equations in four flavours requires simplification: mean momentum approximation

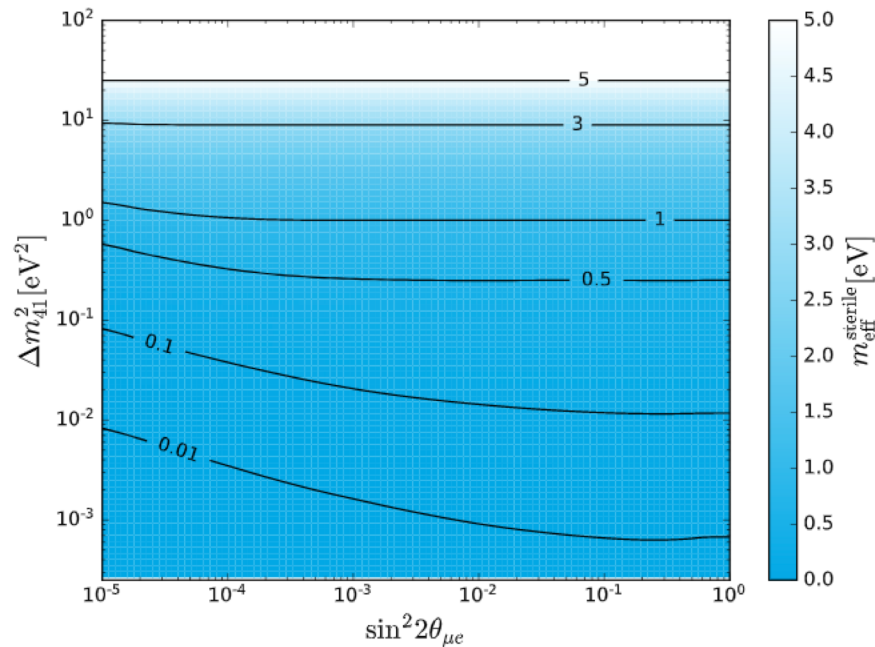
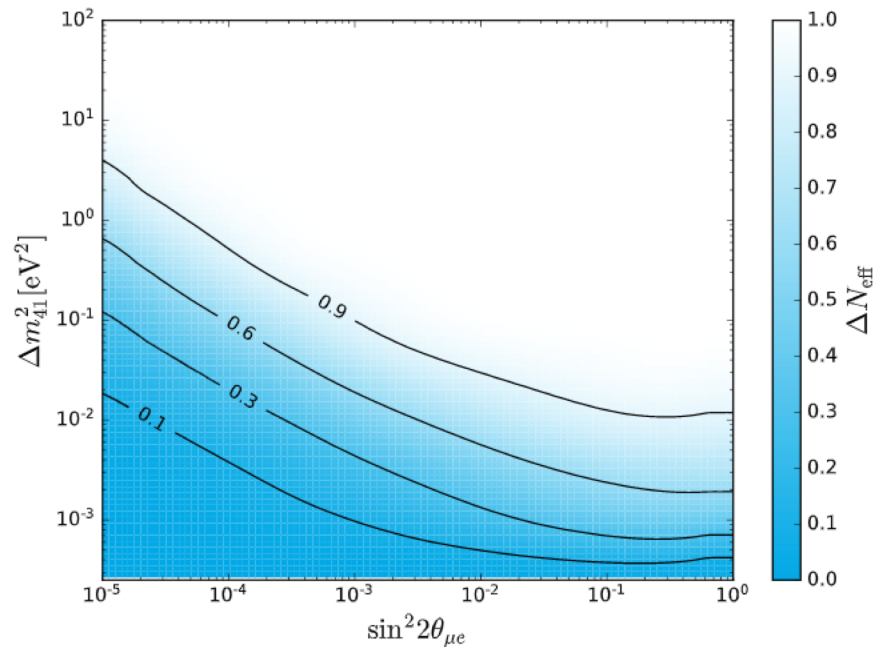
Degenerate region where  $\Delta m^2_{41} \sim \Delta m^2_{31}$

# Validating the mean momentum approximation



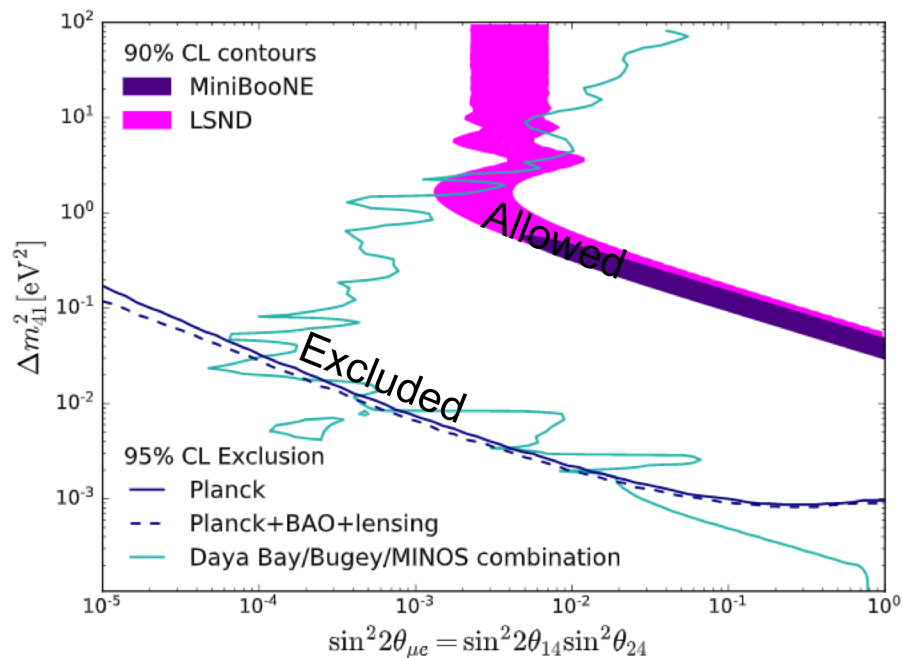
Planck+BAO+lensing with mean momentum approximation

# 3+1 model for $\nu_\mu \rightarrow \nu_e$ analysis



Cosmological parameters in the oscillation space of  $\nu_\mu \rightarrow \nu_e$

# 3+1 model for $\nu_\mu \rightarrow \nu_e$ analysis



Daya Bay/Bugey/MINOS combination  
from PRL **117** 151801 (2016)

# Conclusion

Cosmological and particle physics searches for sterile neutrinos can be compared in the same parameter space

## $\nu_e$ disappearance

- Cosmological limits are strongest above  $\Delta m_{41}^2 \sim 0.1 \text{ eV}^2$  and  $m_{\text{eff}}^{\text{sterile}} \sim 0.2 \text{ eV}$  but are model-dependent
- Daya Bay is more constraining at lower masses

## $\nu_\mu \rightarrow \nu_e$ appearance

- Cosmological limits are strongest at mass splittings above  $\Delta m_{41}^2 \sim 5 \times 10^{-2} \text{ eV}^2$  and  $m_{\text{eff}}^{\text{sterile}} \sim 0.2 \text{ eV}$  but are model-dependent
- Daya Bay / Bugey / MINOS combination is more constraining at lower masses

<https://inspirehep.net/literature/1781295>