

Implications of rare kaon decays on lepton number violating interactions

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In collaboration with
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Based on
arXiv:2008.XXXXX

31.07.20 ICHEP 2020



Lepton number violation (LNV)

Neutrino masses are a clear sign of physics beyond the Standard Model (BSM)

Well motivated scenario: Majorana mass

$$\nu \equiv \nu^c$$

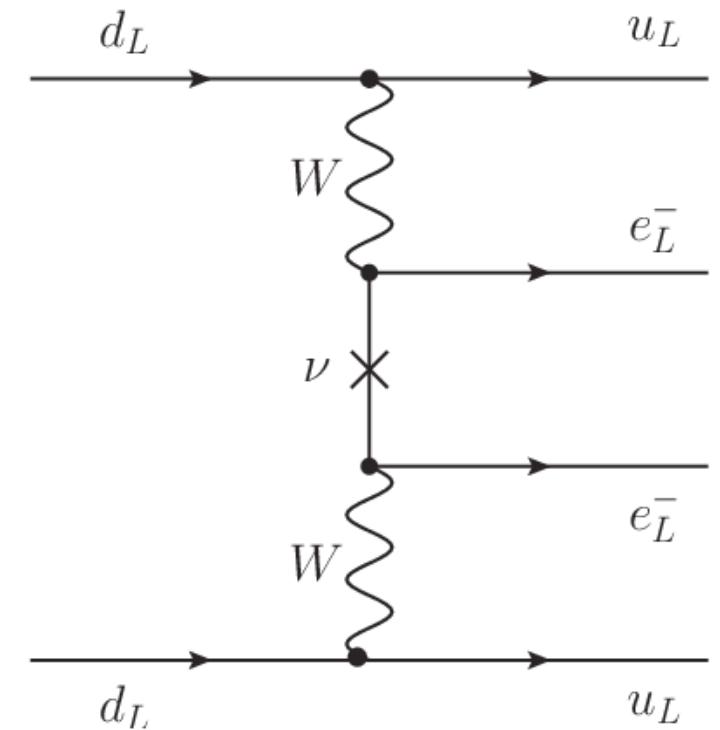
Neutrino Majorana mass implies
lepton number violation (LNV)

Interesting wrt Leptogenesis

Neutrinoless double beta decay ($0\nu\beta\beta$):

+Most sensitive probe of LNV

-Only sensitive to electron flavor



Any other place to look for LNV?

Rare kaon decays

$$\text{BR}(K^+ \rightarrow \pi^+ \nu \bar{\nu})_{\text{SM}} = (8.4 \pm 1.0) \times 10^{-11}$$

$$\text{BR}(K_L \rightarrow \pi^0 \nu \bar{\nu})_{\text{SM}} = (3.4 \pm 0.6) \times 10^{-11}$$

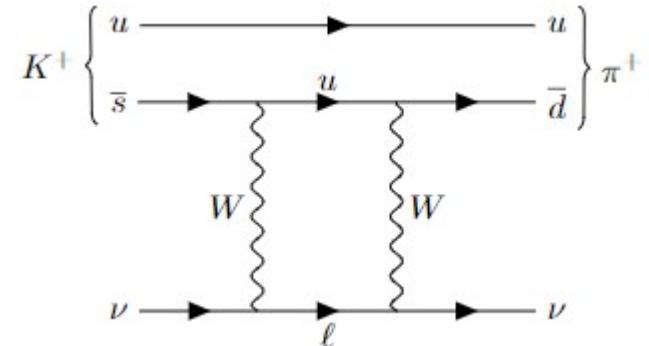
Buras et al, JHEP 2015

$$\text{BR}(K^+ \rightarrow \pi^+ \nu \bar{\nu})_{\text{NA62}} < 1.78 \times 10^{-10}, \quad \text{at 90\% CL}$$

NA62 collaboration, arXiv:2007.08218

$$\text{BR}(K_L \rightarrow \pi^0 \nu \bar{\nu})_{\text{KOTO}} < 3.0 \times 10^{-9}, \quad \text{at 90\% CL}$$

KOTO collaboration, Phys. Rev. Lett. (2019)



See also other talks
 @ ICHEP 2020
 NA62 (Marchevski)
 KOTO (Shimizu)

SM: Neutrino+anti-neutrino LNC

Consequence of LNV in K decays?

BSM: two neutrinos/anti-neutrinos LNV

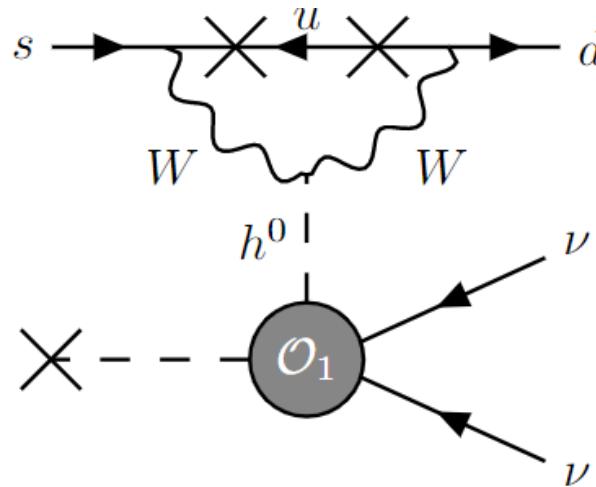
Observable?

Beyond the Standard Model

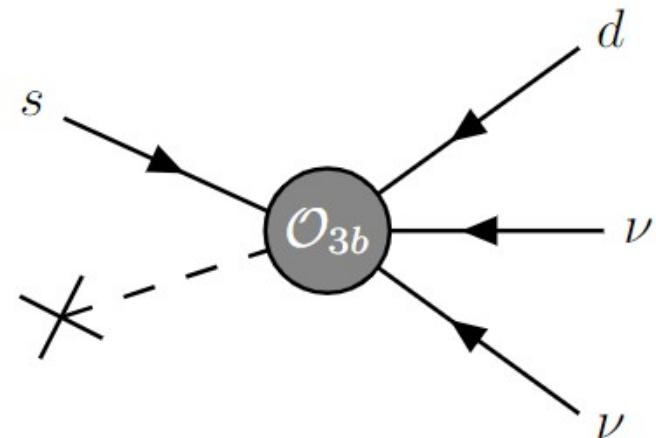
The SM decay is a lepton number conserving (LNC) dimension 6 process
 A BSM decay can be a LNV dimension 7 process

$$\mathcal{L}_{\text{EFT}} = \sum_i C_i \mathcal{O}_i + \text{h.c.} \quad \text{Wilson coefficient: } C_i \propto \frac{1}{\Lambda^{(D-4)}}, \Lambda = \text{New Physics (NP) scale}$$

$$\mathcal{O}_1 = L^\alpha L^\beta H^\rho H^\sigma \epsilon_{\alpha\rho} \epsilon_{\beta\sigma}$$



$$\mathcal{O}_{3b} = L^\alpha L^\beta Q^\rho d^c H^\sigma \epsilon_{\alpha\rho} \epsilon_{\beta\sigma}$$



-Weinberg operator is GIM suppressed

+One possible operator at mass dimension 7

Detection prospects

Search at E949:

SM: LNC, vector current

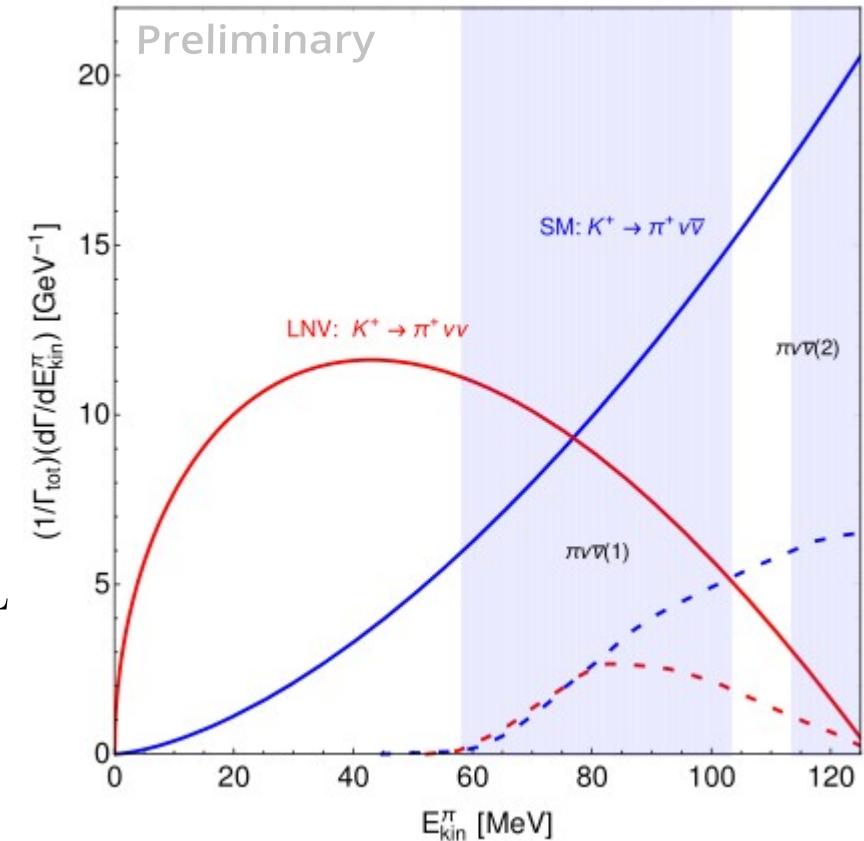
$$\mathcal{L}_{\text{SM}}^{K \rightarrow \pi \nu \bar{\nu}} = \frac{1}{\Lambda_{\text{SM}}^2} \sum_{i=1}^3 (\bar{\nu}_i \gamma^\mu \nu_i) (d \gamma_\mu s)$$

BSM: (can be) LNV, scalar current

$$\mathcal{L}_{\text{BSM}}^{K \rightarrow \pi \nu \nu} = \langle \pi \nu \nu | \mathcal{O}_{3b} | K \rangle = \frac{v}{\Lambda_{\text{BSM}}^3} \bar{s} d \nu_i \nu_j$$

$\text{BR}(K^+ \rightarrow \pi^+ \nu \bar{\nu})_{\text{E949}}^{\text{vector}} < 3.35 \times 10^{-10}$, at 90% CL

$\text{BR}(K^+ \rightarrow \pi^+ \nu \bar{\nu})_{\text{E949}}^{\text{scalar}} < 21 \times 10^{-10}$, at 90% CL

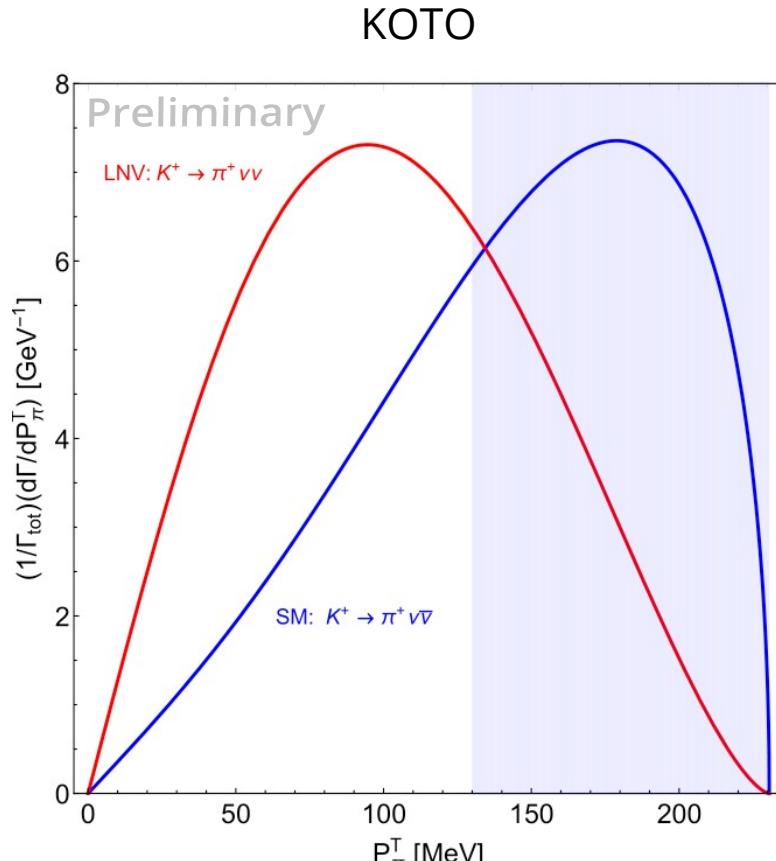


Different current → different distribution

Dashed lines: after experimental acceptance

Deppisch, KF, Harz arXiv:2008.XXXXX

Phase space: KOTO



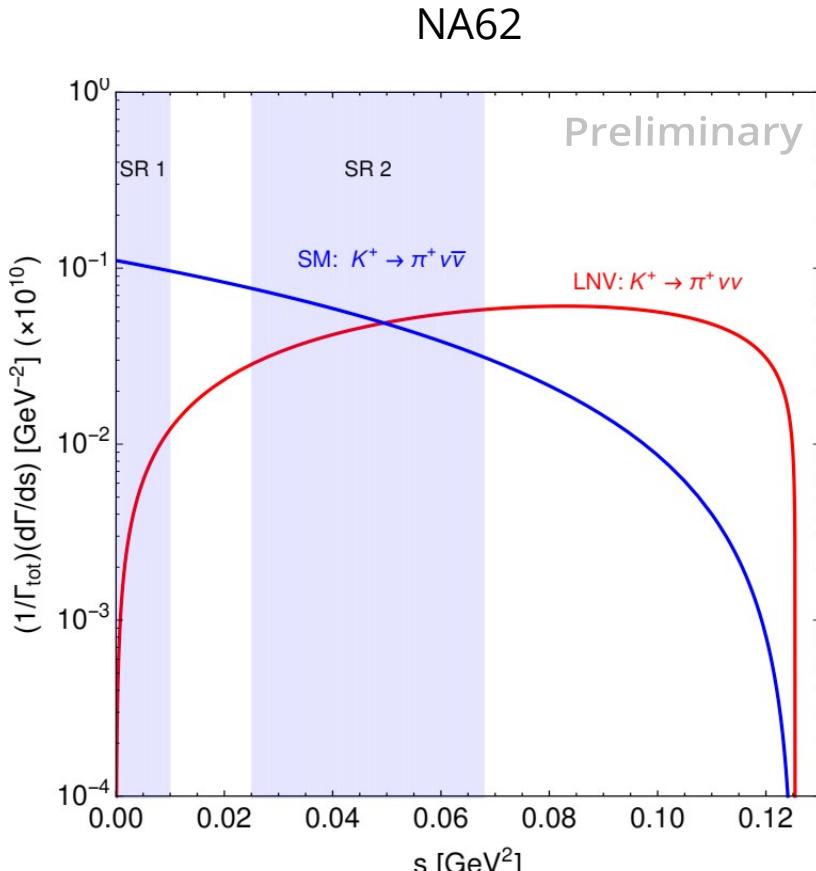
Deppisch, KF, Harz arXiv:2008.XXXXX

Experiment	SM (Vector)	LNV (Scalar)
NA62 SR 1	6%	0.3%
NA62 SR 2	17%	15%
E949 $\pi\nu\bar{\nu}(1)$	29%	2%
E949 $\pi\nu\bar{\nu}(2)$	45%	38%
KOTO	64%	30%

All three experiments are more sensitive to vector currents than to scalar currents

At E949 and NA62:
different # of events in each signal region

Phase space: NA62



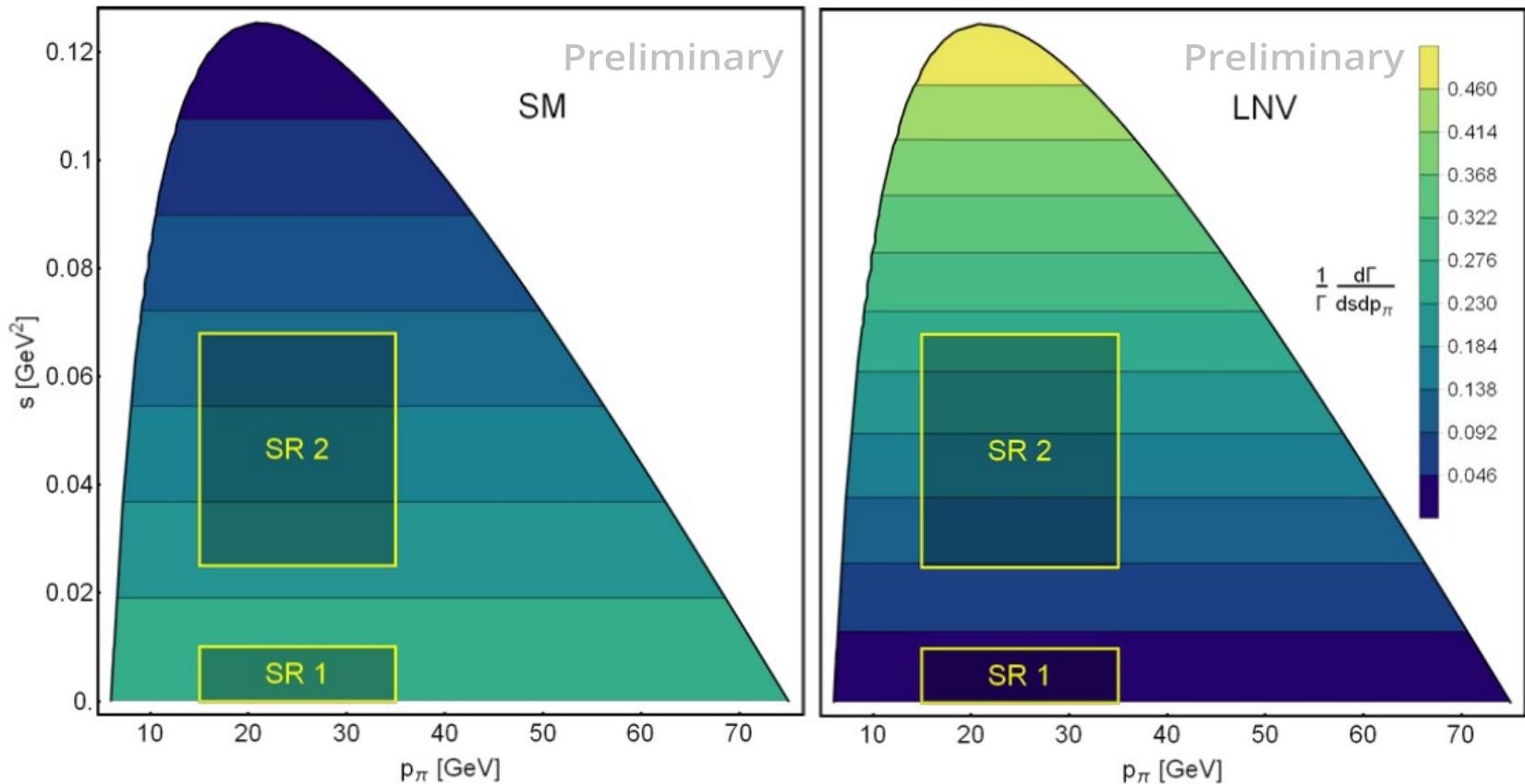
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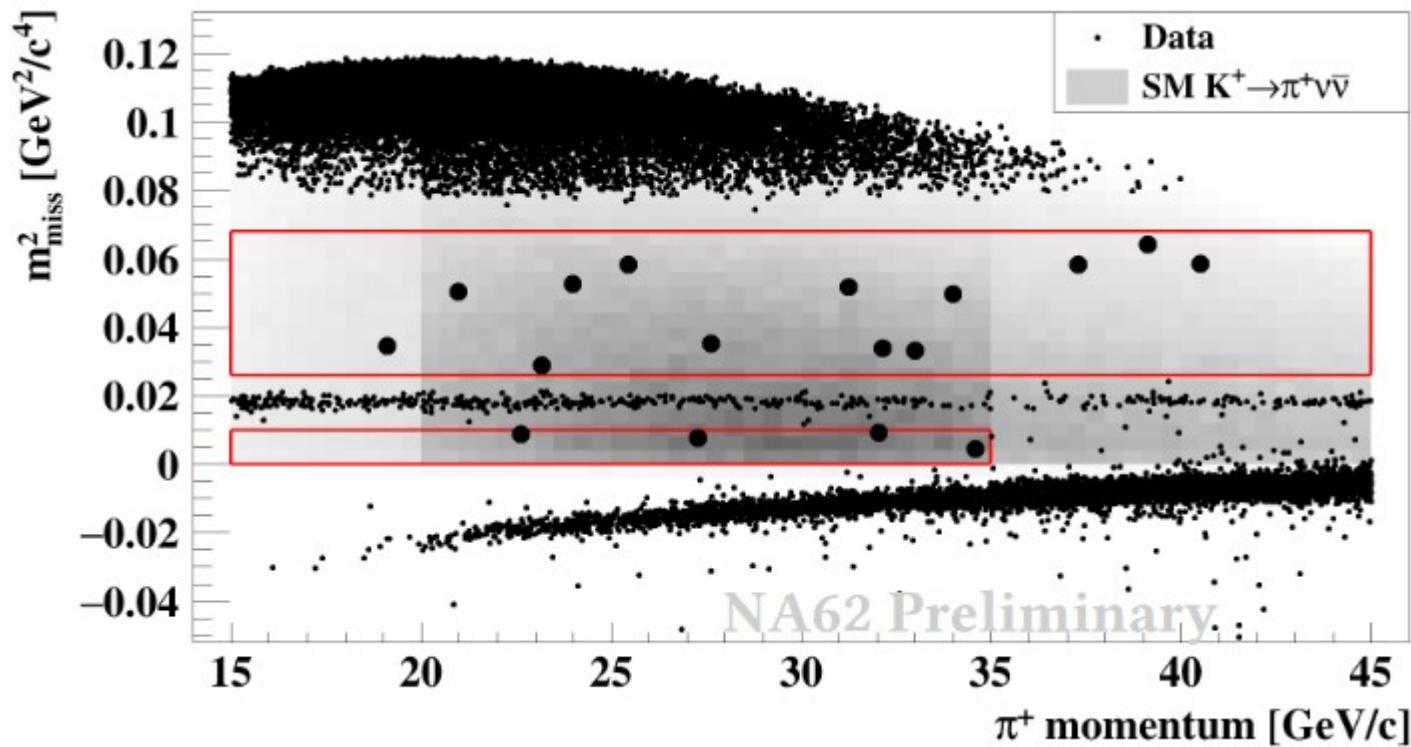
$$\text{Missing energy } s: \quad s = (E_K - E_\pi)^2$$

Deppisch, KF, Harz arXiv:2008.XXXXXX

For a given s, the distribution is independent of the pion momentum.

LNC: More events in SR 1, LNV: More events in SR 2

Reported @ ICHEP 2020

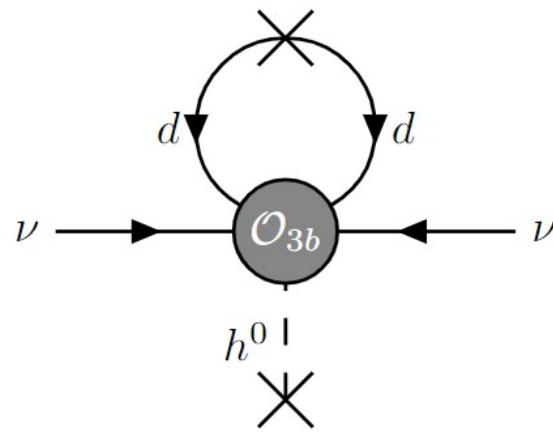


See talk by R. Marchevski on behalf of the NA62 collaboration

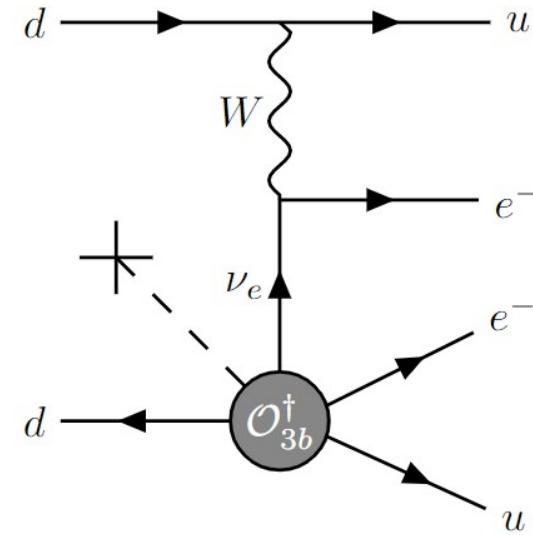
<https://youtu.be/kH9sg-9BXG4?t=10359>

Another side of dim-7 LNV

Radiative neutrino mass



Neutrinoless double beta decay ($0\nu\beta\beta$)



Radiative neutrino mass depends on the UV completion: non-trivial corrections

$0\nu\beta\beta$ is only sensitive to electron neutrinos and up/down quarks

Implications for leptogenesis

$$\frac{d\eta_{\Delta L}}{dz} \propto -\eta_{\Delta L} \frac{\Lambda_{\text{Pl}}}{\Lambda_{\text{BSM}}} \left(\frac{T}{\Lambda_{\text{BSM}}} \right)^{2D-9}$$

Lepton asymmetry can be translated into a baryon asymmetry via sphalerons

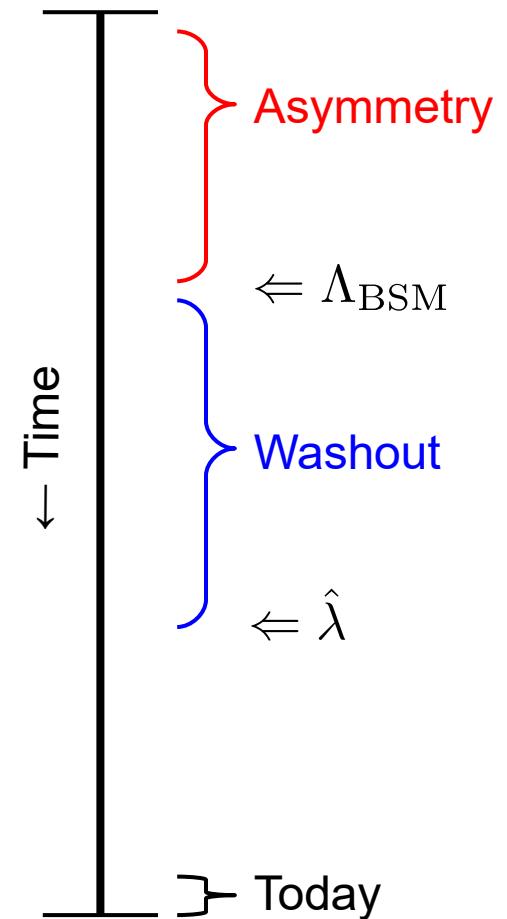
Washout: process that reduces asymmetry

Assuming preexisting asymmetry: washout acts to remove it

$$\hat{\lambda} = f(\eta_B^{\text{obs}}, \Lambda_{\text{BSM}}, \dots) \quad \eta_B^{\text{obs}} = (6.20 \pm 0.015) \times 10^{-10}$$

Planck collaboration (2018)

Strong washout: $\hat{\lambda} \lesssim T \lesssim \Lambda_{\text{BSM}}$



NP scales, washout scales

$$\text{BR}(K \rightarrow \pi \nu \bar{\nu})_{\text{tot}} = \text{BR}(K \rightarrow \pi \nu \bar{\nu})_{\text{SM}} + \left(\frac{A_{\text{LNV}}}{A_{\text{SM}}} \right) \times \sum_{i \leq j=1}^3 \text{BR}(K \rightarrow \pi \nu_i \nu_j)_{\text{LNV}}$$

Relative acceptance: $\left(\frac{A_{\text{LNV}}}{A_{\text{SM}}} \right)$

Strong washout: $\hat{\lambda} \lesssim T \lesssim \Lambda_{\text{BSM}}$

$$\begin{aligned} \mathcal{O}_{3b} : \quad K^+ \rightarrow \pi^+ \nu \nu: \quad & \Lambda_{\text{BSM}}^{\text{NA62}} > 17.2 \text{ TeV} & \Rightarrow \quad \hat{\lambda} > 0.196 \text{ TeV} \\ K_L \rightarrow \pi^0 \nu \nu: \quad & \Lambda_{\text{BSM}}^{\text{KOTO}} > 12.3 \text{ TeV} & \Rightarrow \quad \hat{\lambda} > 0.178 \text{ TeV} \end{aligned}$$

Rare kaon decays are probing NP scales of order 10 TeV

Conclusion

- Scalar/vector distribution difference is significant in kaon decays
- Distinguishing LNV from LNC in rare kaon decays may be plausible
- An observation of LNV in rare kaon decays would put high scale leptogenesis under tension

Thanks for listening

Backup: leptoquarks

$$\begin{aligned}\mathcal{L} = & \mathcal{L}_{\text{SM}} - \tilde{R}_2^{\dagger\alpha}(\square + m_{\tilde{R}_2}^2)\tilde{R}_{2\alpha} - S_1^*(\square + m_{S_1}^2)S_1 \\ & + \mu S_1 H^{\dagger\alpha}\tilde{R}_{2\alpha} - g_1^{ik}\bar{L}_{i\alpha}i\sigma_2^{\alpha\beta}\tilde{R}_{2\beta}^*d_k - g_2^{jn}\bar{Q}_n^{c\alpha}L_j^\beta\epsilon_{\alpha\beta}S_1 \\ & - g_3^{jn}\bar{u}_n^ce_jS_1 + \text{h.c.}\end{aligned}$$

$$\mathcal{L}_{7D} = \frac{\mu g_1^{ik}g_2^{jn}}{m_{\tilde{R}_2}^2 m_{S_1}^2} L_i^\alpha H^\beta d_k^c Q_n^\mu L_j^\nu \epsilon_{\alpha\beta} \epsilon_{\mu\nu} + \frac{\mu g_1^{ik}g_3^{jn}}{m_{\tilde{R}_2}^2 m_{S_1}^2} L_i^\alpha H^\beta d_k^c u_n^c e_j^c \epsilon_{\alpha\beta}$$

O. Cata and T. Mannel, arXiv:1903.01799

