

# New laser setup in Bari for MPGD gain measurements

RD51 mini-week

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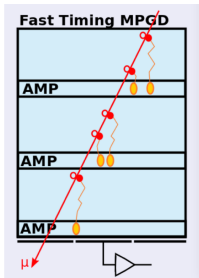


## Outline

- Laser setup requirements for the FTM
- The laser specifications and optical setup
- Characterization of a triple-GEM prototype in the laser box

# Requirements for the characterization of the FTM

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## Working principle and time resolution

$$\sigma_t = \lambda / N v_{\text{drift}}$$

$\lambda$  = ionization mean free path

$N$  = number of stages

1. Signal pickup by external readout: **only resistive electrodes**

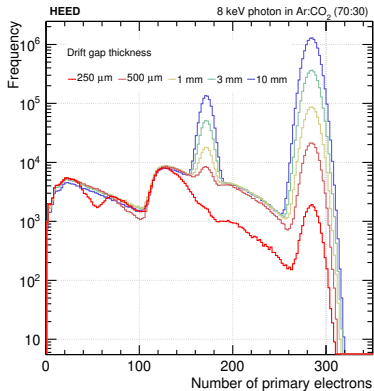
Gain calibration with non-monochromatic sources is made difficult: no copper electrode for fluorescence conversion

2. Tiny drift gaps (250  $\mu\text{m}$ )

Cannot perform gain calibration by conventional sources (X-rays)

New small-size ( $\sim 4 \times 4 \text{ cm}^2$ ) prototype currently under tests in Bari

# X-ray photon conversion in gas mixture



No distinct peak at gaps  $< 500 \mu\text{m}$   $\rightarrow$  **X-ray energy loss is subjected to large fluctuations**

# **Laser specifications and optical setup**

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# Principles of laser-gas interaction

- Photons in lasers have **too low energy** ( $\sim 4.7$  eV @ 266 nm) to ionize typical counting gas molecules (13-15 eV)
- Common mixtures contain some ppm impurity molecules with low ionization potential ( $\sim 9$  eV)  $\rightarrow$  laser ionization is possible by **multi-photon processes**:

$$\frac{R}{V} = N\sigma^{(n)}\phi^n$$

$R/V$  = ionization rate density

$N$  = molecule concentration

$\sigma^{(n)}$  = n-photon cross-section equivalent

$\phi$  = beam flux

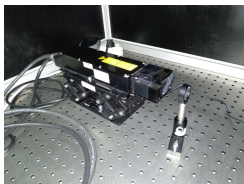
- At low intensities, two-photon ionization dominates:

$$\frac{R}{V} = \left(\frac{\lambda}{hc}\right)^2 N\sigma^{(2)}I^2$$

Primary current in detector is proportional to **square of laser pulse energy**

# Specifications of the laser setup

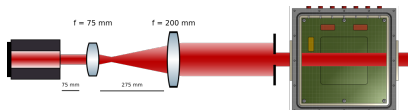
<b>Pulse energy</b>	51 $\mu$ J	can provide a MIP-like energy deposit
<b>Waist radius</b>	400 $\mu$ m	low angular divergence
<b>Wavelength</b>	266 nm/4.7 eV	two-photon ionization of hydrocarbons
<b>Pulse duration</b>	1 ns FWHM	lower than triple-GEMs time resolution
<b>Spatial mode</b>	TEM <sub>00</sub>	gaussian beam beam quality <1.5



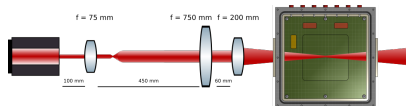


# Optical setup preparation

## Collimated setup



## Focused setup

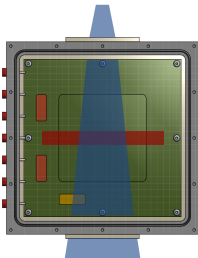
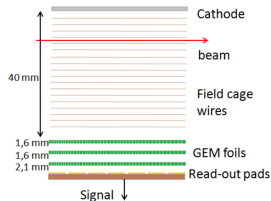
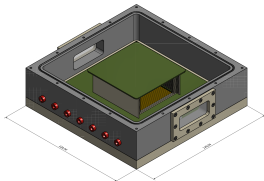


	Collimated	Focused
<b>Waist radius</b>	1500 $\mu\text{m}$	23.4 $\mu\text{m}$
<b>Angular divergence</b>	0.06 mrad	$\sim 5$ mrad
<b>Beam intensity</b>	34 $\mu\text{J}/\text{mm}^2$	$3 \times 10^4$ $\mu\text{J}/\text{mm}^2$
	Optical filter + Pinhole to reduce pulse energy	Point-like primary ionization

# **Characterization of a triple-GEM detector in the laser box**

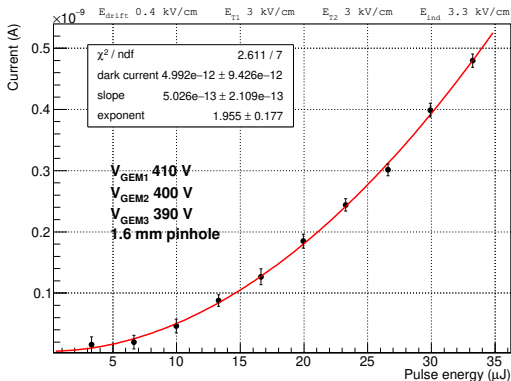
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# The time projection GEM prototype

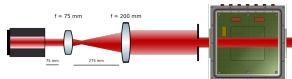


- 40 mm drift gap, suitable for benchmarking of the laser setup
- beam passes through **quartz windows** (transparent to UV)
- signal readout on 2 rows of 60 pads ( $6 \times 2 \text{ mm}^2$  each)
- the small instrumented area compared with the total gas volume complicates the gain calibration with X-rays

# Observation of multi-photon ionization



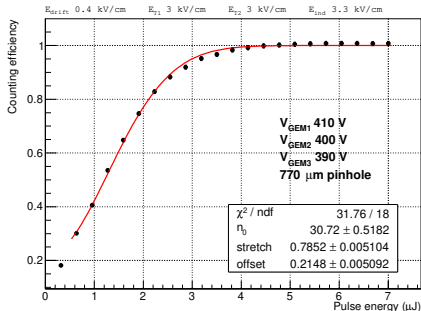
- collimated setup (low intensity)
- anode current  $\propto$  primary ionization rate
- ionization rate  $\propto$  (pulse energy) $^m$  for m-photon absorption



$m = 1.96 \pm 0.18$ , compatible with **two-photon absorption**

# Estimation of primary ionization rate

**Problem** How to determine the number of primaries created by a single laser pulse?



Counting efficiency scan vs laser pulse energy @ 100 Hz:

$$\epsilon = \frac{\text{anode signal rate}}{\text{laser pulse rate}}$$

Assuming Poisson fluctuations on primary electron number,

$$\epsilon = 1 - \sum_{n=0}^{n_{th}} \frac{\exp[-n_0(E/E_0)] n_0^n (E/E_0)^{2n}}{n!}$$

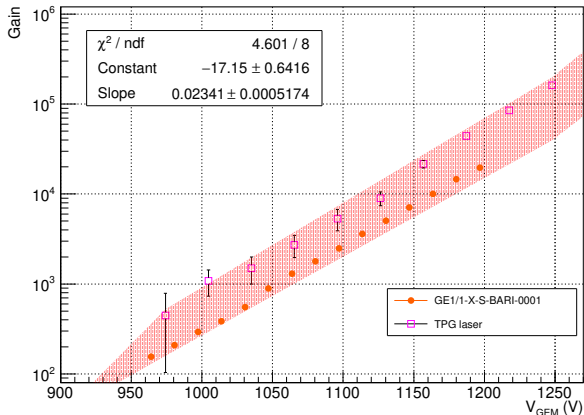
$E_0$  = reference pulse energy       $n_0$  = primary electrons per laser pulse at  $E_0$

$n_{th}$  = n. of primary electrons corresponding to the discriminator threshold

**$n_0 = 30.7 \pm 0.5$  electrons at  $10 \mu\text{J}$  in the active gas volume**

Waiting for confirmation from: primary ionization current, single-electron spectrum

# Gain curve measurement

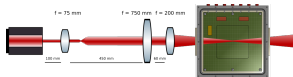
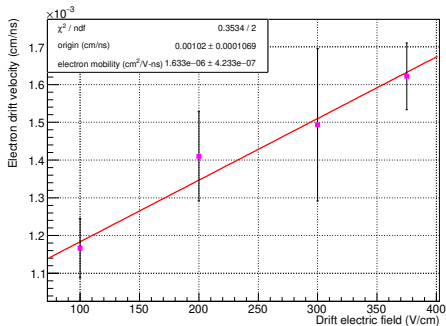


$$G_{\text{eff}} = \frac{\text{anode current}}{n_p \times q_e \times \text{laser rate}}$$

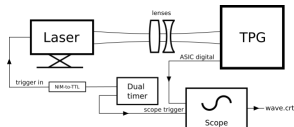
$$n_p = 30.7 @ 10 \mu\text{J pulse energy}$$

Compatible with the gain curves of CMS GE1/1 chambers tested at CERN and in Bari

# Timing measurements



$(1.63 \pm 0.42) \times 10^{-6} \text{ cm}^2/\text{V} \cdot \text{ns}$   
electron mobility in Ar:CO<sub>2</sub> (70:30)



- optical setup in focused configuration: **point-like ionization**
- laser pulse emission triggered by external clock
- detector digital signals acquired by a scope
- the average signal arrival time is plotted at different ionization positions in the drift gap
- **electron drift velocity** is given by a linear fit

$2.36 \times 10^{-6} \text{ cm}^2/\text{V} \cdot \text{ns}$   
from a Magboltz simulation

# Conclusions and summary

## Setup for the characterization of small-gap MPGDs

- Validation with triple-GEM chamber
- **Collimated** low intensity setup for gain calibration
- **Focused** high intensity setup for timing measurements

## Future development

- Comparison with other gain calibration techniques with lasers
  - direct **primary current measurement**
  - **single-electron** response
- The UV laser bench is **ready for the characterization of the FTM**
- Femtosecond laser for time resolution measurements on the FTM