

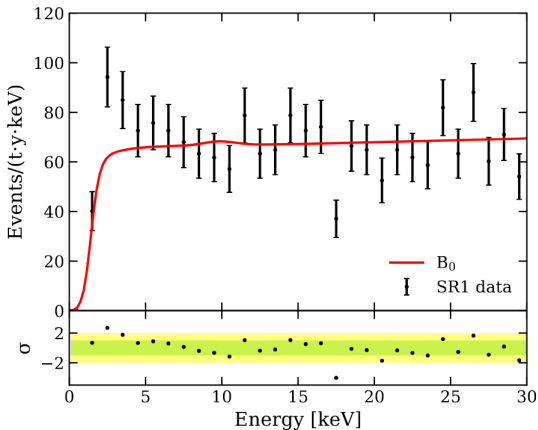
Anomaly-free leptophilic ALP and its LFV tests

Oscar Vives

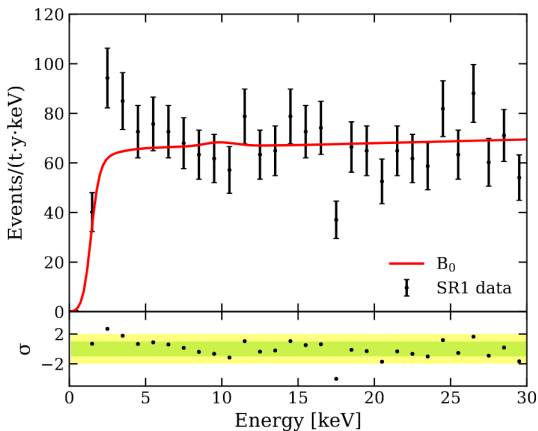
SUSY 2021, Beijing, 23 - 28 August



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Excess in e^- recoil energy between 2 and 3 keV.

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Best fit to data with large coupling to e^- , 3.4σ

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ALPs

Pseudo-Goldstone bosons from broken global symmetries with $m_a \simeq \mathcal{O}(\text{keV})$ and a weak coupling to e^- , could explain this excess

However, X-ray observations forbid anomalous coupling to photons for $m_a \gtrsim 0.1 \text{ keV}$



$U(1)_{\text{em}}$ anomaly-free ALP

SM particle content, only $B - L$ and L are $U(1)_{\text{em}}$ anomaly-free with family universal charges.

But ... breaking associated to N_R masses, too high in Seesaw models.



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- Global flavor-dependent $U(1)_\phi$ symmetry spontaneously broken
 \Rightarrow Axion Like Particle

Model I: $U(1)$ flavor symmetry

Scalar flavon field, ϕ , generates Yukawa couplings as function of

small vevs, $Y_{ij} = \left(\left(\frac{\langle \phi \rangle}{M} \right) \ll 1 \right)^n$. with n function of charges.

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- Angular component of flavon becomes ALP, if $U(1)_\phi$ symmetry global.
- Anomaly-free $U(1)_\phi$: $\sum_i Q_{L_i} = 0$ and $\sum_i Q_{e_i} = 0$.

Field	τ_L	μ_L	e_L	τ_R	μ_R	e_R	$N_{R,i}$	ϕ	H_1	H_2
$U(1)_\phi$	-1	0	1	1	0	-1	0	0	2	1
Z_2	+	+	+	-	-	-	-	+	-	+

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Soft-breaking of $U(1)_\phi \times Z_2$: $m^2 H_1 H_2$, with $m \sim \mathcal{O}(\text{EW})$.

$$\longrightarrow m_a \simeq m^2 / v_\phi$$

Yukawa couplings:

$$\mathcal{L}_Y \supset c_{ij}^e \epsilon^{n_{ij}^e} \bar{L}_i \tilde{H}_2 e_j + c_{ij}^\nu \epsilon^{n_{ij}^\nu} \bar{L}_i H_2 N_j + (M_R)_{ij} N_{R_i} N_{R_j}^c,$$

with $\epsilon \simeq 0.1$ and $n_{ij}^e = q_{L_i} - q_{e_j} + q_{H_2}$, $n_{ij}^\nu = q_{L_i} - q_{N_{R_j}} - q_{H_2}$,

$$n_{ij}^e = \begin{pmatrix} 4 & 3 & 2 \\ 3 & 2 & 1 \\ 2 & 1 & 0 \end{pmatrix}, \quad n_{ij}^\nu = \begin{pmatrix} 1 & 1 & 1 \\ 2 & 2 & 2 \\ 3 & 3 & 3 \end{pmatrix}.$$

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- After symmetry breaking: $\phi(x) = \frac{1}{\sqrt{2}} (v_\phi + s(x)) e^{i a(x)/v_\phi}$
- With couplings to charged leptons,

$$-\mathcal{L}_{ae} = i \frac{\partial_\mu a}{2f_a} \bar{e}_i \gamma^\mu (V_{ij}^e + \gamma^5 A_{ij}^e) e_j.$$

$$V_{ij}^e = \frac{1}{2} \left(U_R^{e\dagger} Q_e U_R^e + U_L^{e\dagger} Q_L U_L^e \right), \quad A_{ij}^e = \frac{1}{2} \left(U_R^{e\dagger} Q_e U_R^e - U_L^{e\dagger} Q_L U_L^e \right)$$

Model II: general mixing

Generalize the previous structure to allow for larger mixings determined by a larger symmetry \mathcal{F} , with global $U(1)_\phi$ part of \mathcal{F}

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- PMNS-like mixing in charged-leptons. Then:

$$V_{ij}^e, A_{ij}^e = \frac{1}{2} U_{\text{PMNS}}^{e\dagger} (Q_e \pm Q_L) U_{\text{PMNS}}^e$$

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Too large LFV ???

LFV and Astrophysics Constraints

$$\text{BR}(e_i \rightarrow e_j a) = \frac{m_{e_i}^3}{16\pi\Gamma(e_j)} \frac{|C_{ij}^e|^2}{4f_a^2} \left(1 - \frac{m_a^2}{m_{e_i}^2}\right)^2.$$

with $|C_{ij}^e|^2 = |V_{ij}^e|^2 + |A_{ij}^e|^2$

Lepton decay	BR limit	Projection
$\text{BR}(\mu \rightarrow e a)$	$< 2.6 \cdot 10^{-6}$ Jodidio et al.	$< 1.3 \cdot 10^{-7}$ MEGII-fwd
$\text{BR}(\mu \rightarrow e a)$	$< 2.1 \cdot 10^{-5}$ TWIST	$< 7.3 \cdot 10^{-8}$ Mu3e
$\text{BR}(\mu \rightarrow e a \gamma)$	$< 1.1 \cdot 10^{-9}$ Crystal Box	
$\text{BR}(\tau \rightarrow e a)$	$< 2.7 \cdot 10^{-3}$ ARGUS	$< 8.4 \cdot 10^{-6}$ Belle-II
$\text{BR}(\tau \rightarrow \mu a)$	$< 4.5 \cdot 10^{-3}$ ARGUS	$< 1.6 \cdot 10^{-5}$ Belle-II

Coling of white dwarfs and red giants:

$$f_a \gtrsim 2.3 \times 10^9 |C_{11}^e| \text{ GeV}, \quad f_a \gtrsim 1.2 \times 10^9 |C_{11}^e| \text{ GeV}.$$

Results

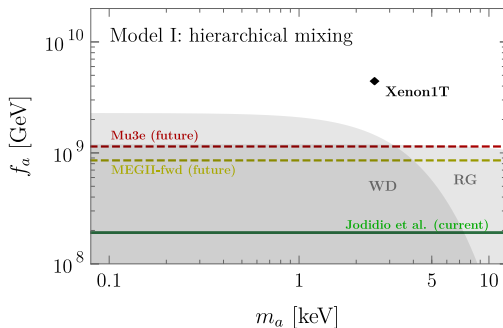
Analysis of *Takahashi-Yamada-Yin*, *Xenon1T* explained by ALP
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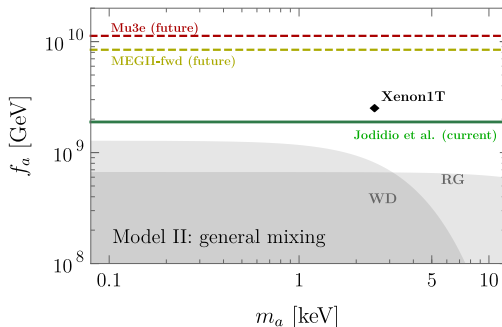
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Conclusions

- **Xenon1T** excess can be explained by flavored **ALP** with minimal particle content.
- Simultaneously, LFV **ALP** interactions could be detected at low-energy experiments.
- **ALP** flavor couplings depend on charged-lepton **mixings** and **charges**.
- Flavor models of small **mixings** (CKM-like) produce too small effects for proposed experiments.
- Models of PMNS-like charged-lepton **mixings** explaining **Xenon1T** could be detectable in LFV experiments.