

Dark Matter and Leptogenesis in the Type Ib Seesaw Model



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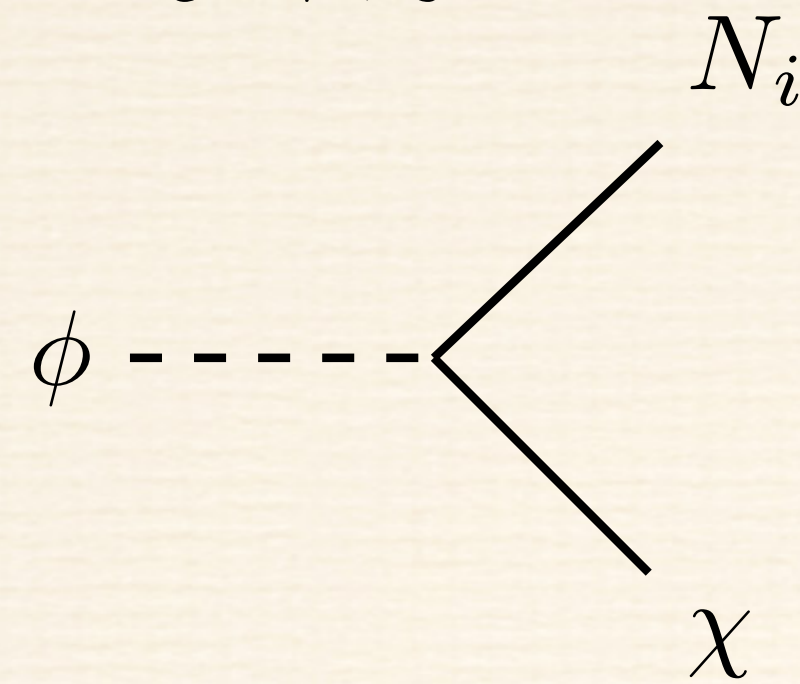
24 Aug 2021

Neutrino Portal

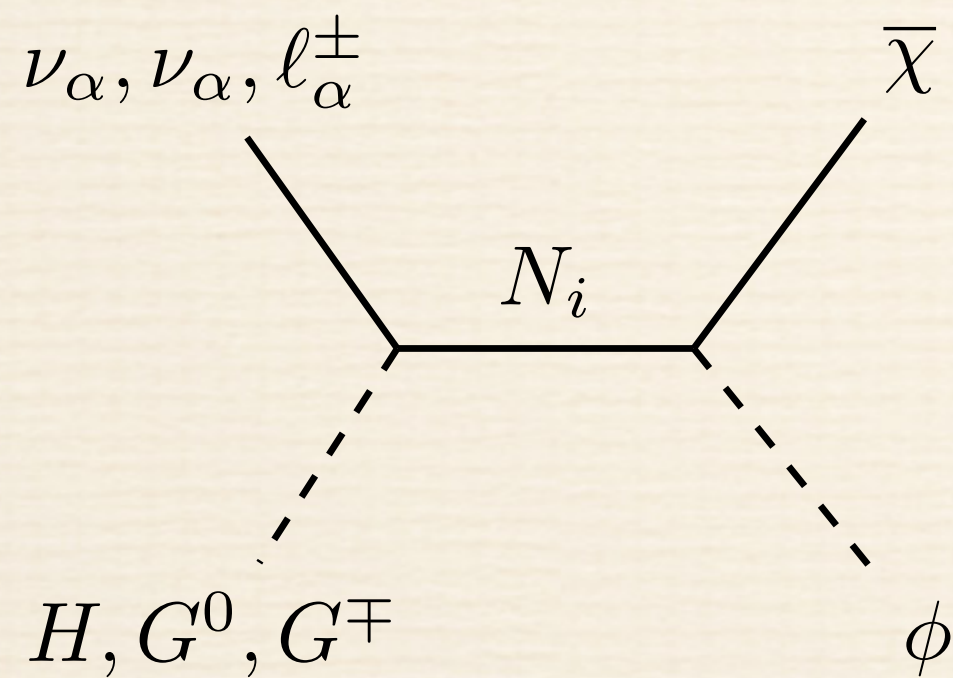
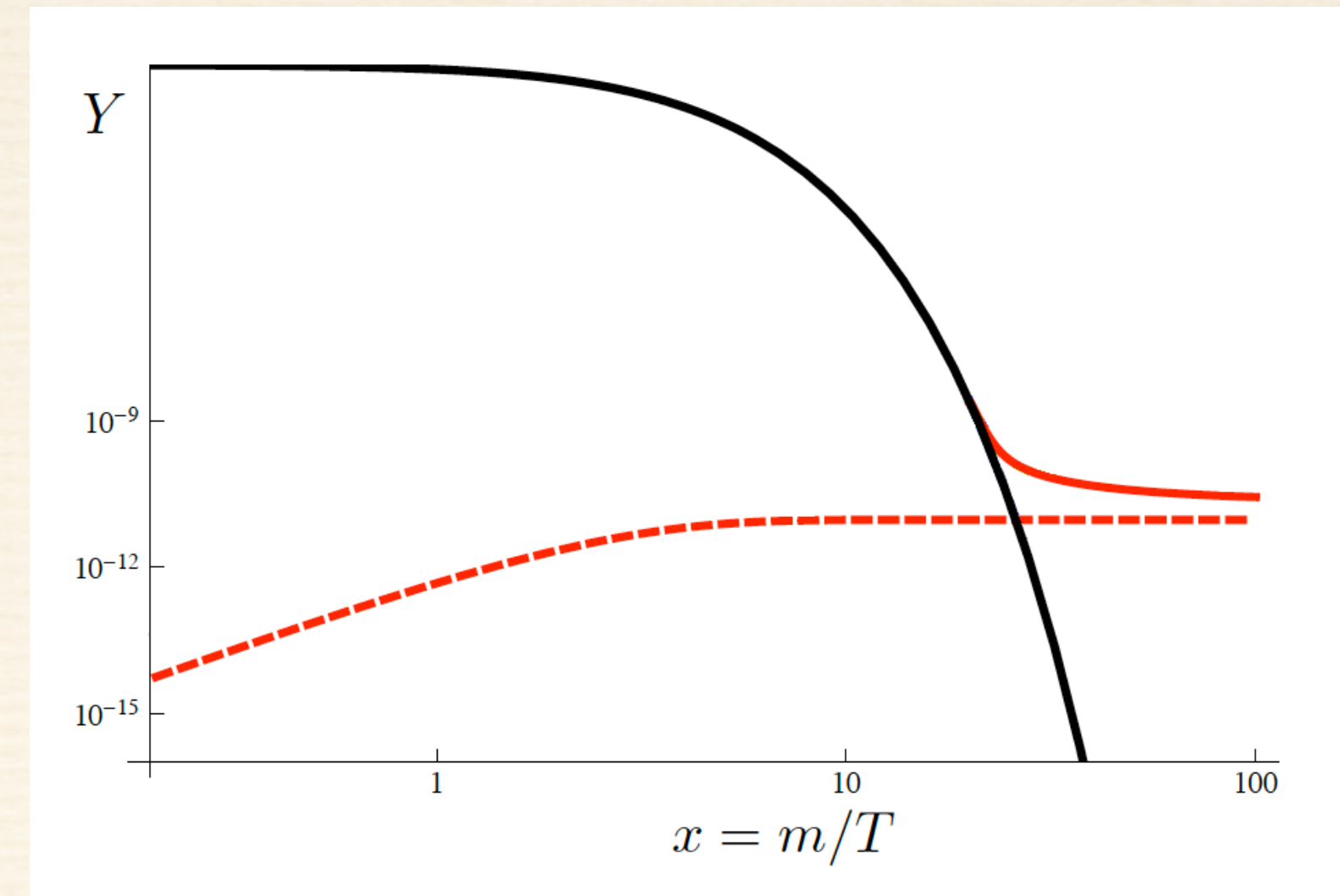
❖ General neutrino portal

$$y_i \phi \bar{\chi} N_i$$

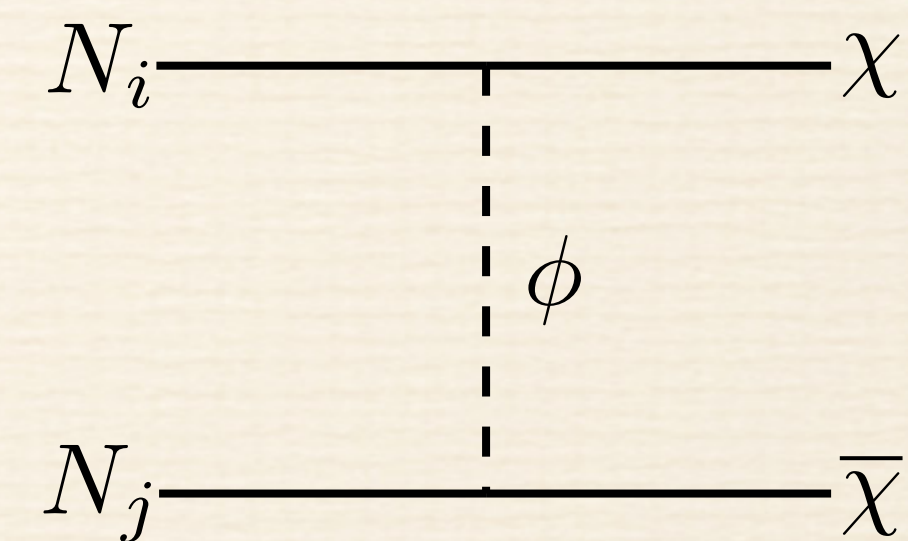
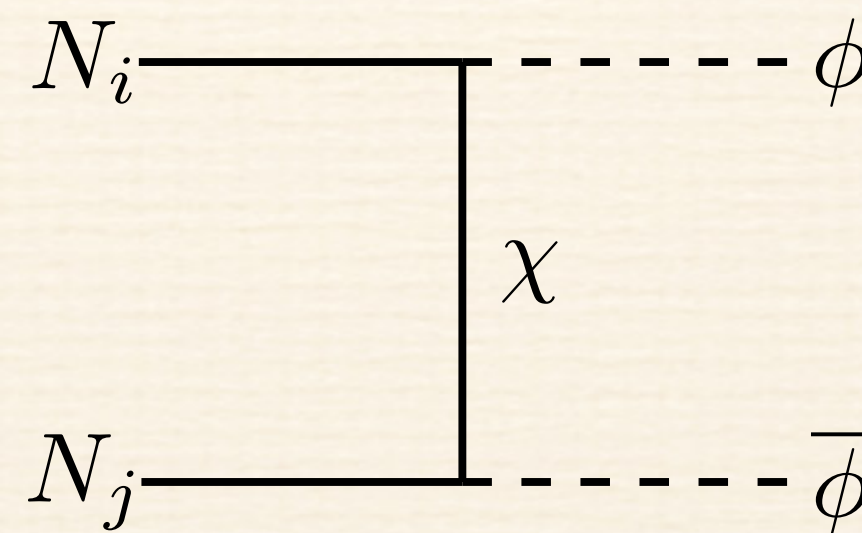
❖ heavy scalar scenario



❖ Dark matter producing process



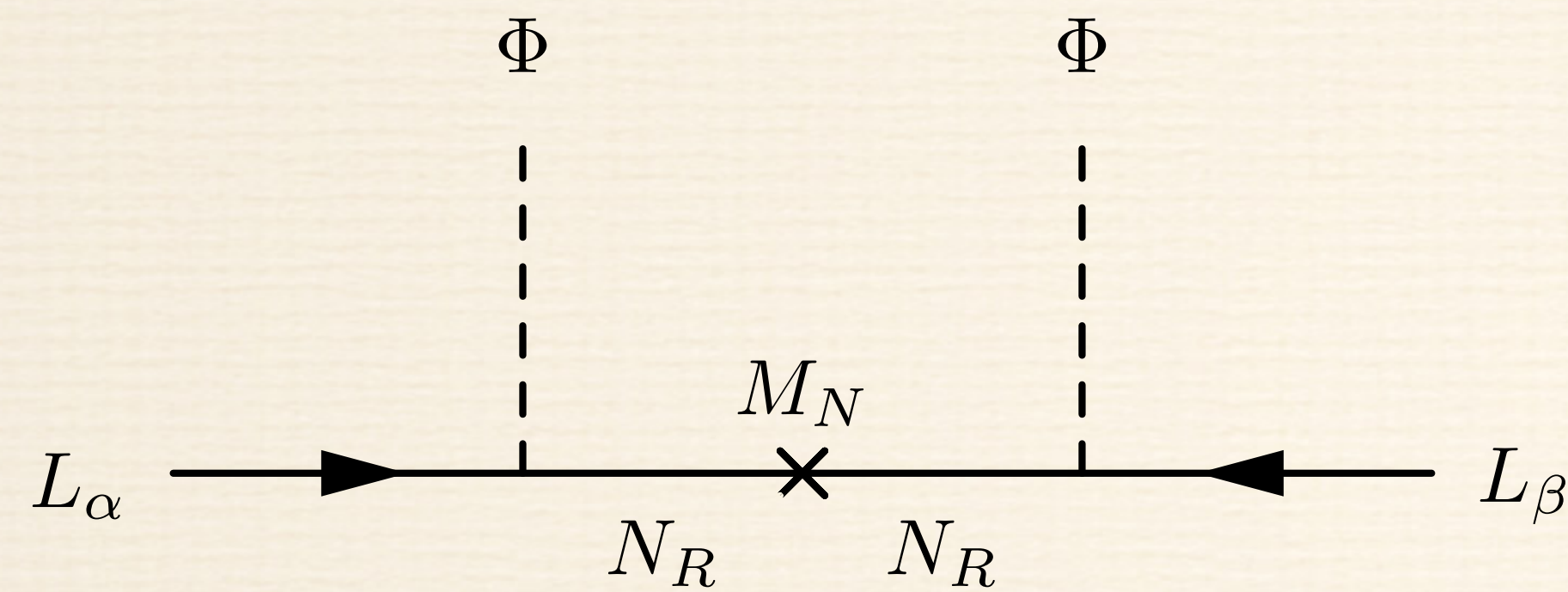
ν -Yukawa process $\propto Y^2 y^2$



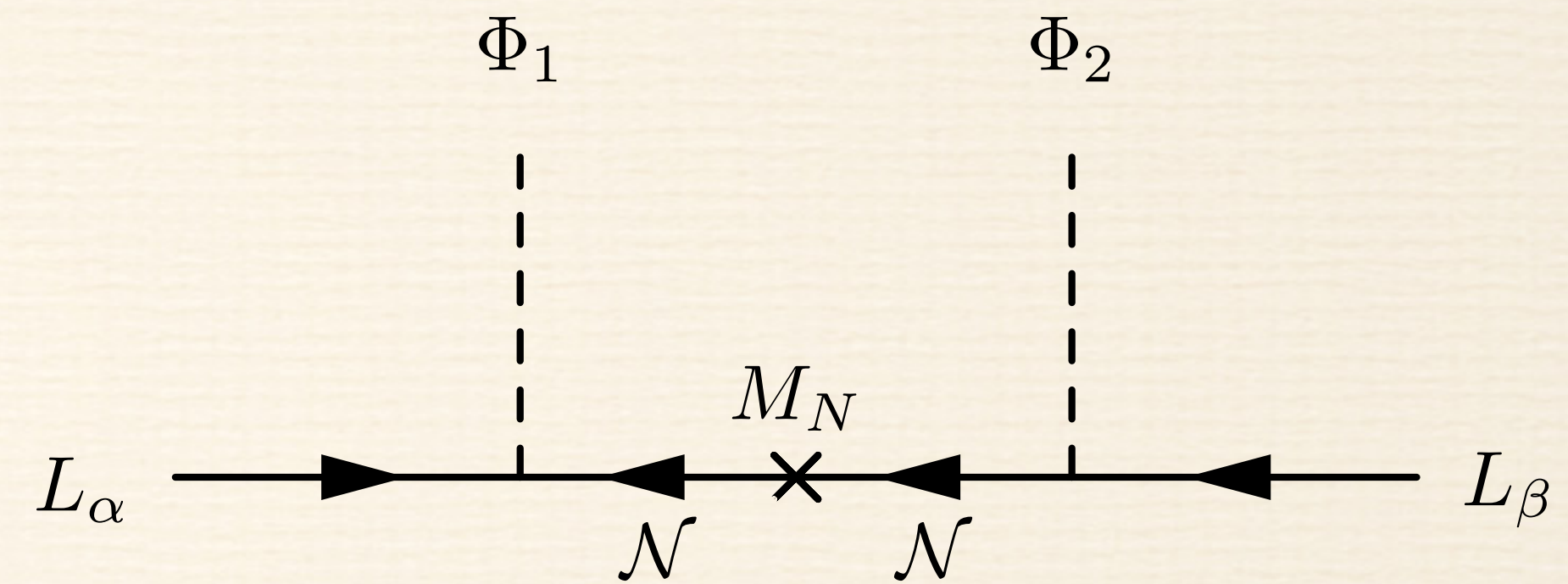
dark sector process $\propto y^4$

Why Type Ib Seesaw Model?

- ❖ Traditional type I seesaw mechanism (type Ia)
- ❖ Type Ib seesaw mechanism



$$m \propto \frac{Y^2 v^2}{M_N}$$



$$m \propto \frac{Y_1 Y_2 v_1 v_2}{M_N}$$

- ❖ $M_N \sim 1 \text{ GeV} \rightarrow Y \sim 10^{-7}$: too small to play roles in dark matter production
- ❖ One of Y_1, Y_2 can be small while the other one is large

Type Ib Seesaw Model with a Neutrino Portal

❖ Particles and symmetries

	Q_α	$u_{R\beta}$	$d_{R\beta}$	L_α	$e_{R\beta}$	Φ_1	Φ_2	N_{R1}	N_{R2}	ϕ	$\chi_{L,R}$
$SU(2)_L$	2	1	1	2	1	2	2	1	1	1	1
$U(1)_Y$	$\frac{1}{6}$	$\frac{2}{3}$	$-\frac{1}{3}$	$-\frac{1}{2}$	-1	$-\frac{1}{2}$	$-\frac{1}{2}$	0	0	0	0
Z_3	1	ω	ω	1	ω	ω	ω^2	ω^2	ω	ω	ω^2
Z_2	+	+	+	+	+	+	+	+	+	-	-

❖ Seesaw Lagrangian and neutrino portal

$$\mathcal{L}_{\text{seesawIb}} = -Y_{1\alpha} \bar{L}_\alpha \Phi_1 N_{R1} - Y_{2\alpha} \bar{L}_\alpha \Phi_2 N_{R2} - M_N \overline{N_{R1}^c} N_{R2} + \text{h.c.}$$

$$\mathcal{L}_{N_{R\text{portal}}} = y_1 \phi \overline{\chi_R} N_{R1}^c + y_2 \phi \overline{\chi_L} N_{R2} + \text{h.c.}$$

❖ The two RH neutrinos form a Dirac pair $\mathcal{N} = (N_{R1}^c, N_{R2})$

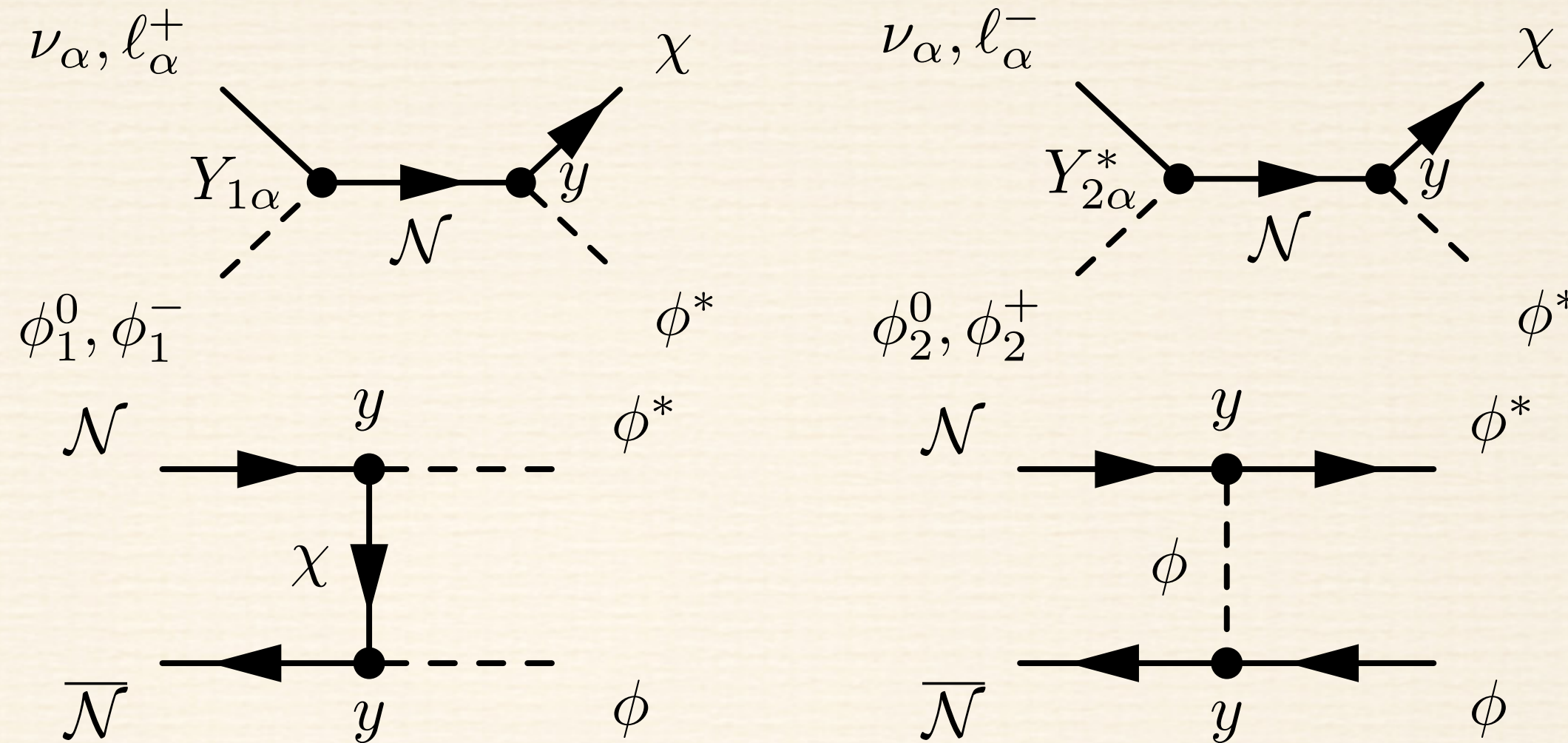
$$\mathcal{L}_{\text{seesawIb}} = -Y_{1\alpha}^* \bar{L}_\alpha^c \Phi_1^* \mathcal{N}_L - Y_{2\alpha} \bar{L}_\alpha \Phi_2 \mathcal{N}_R - M_N \overline{\mathcal{N}_L} \mathcal{N}_R + \text{h.c.}$$

$$\mathcal{L}_{N_{R\text{portal}}} = y \phi \overline{\chi} \mathcal{N} + \text{h.c.}$$

Dark Matter Production

❖ Dark matter producing processes

- ν -Yukawa processes



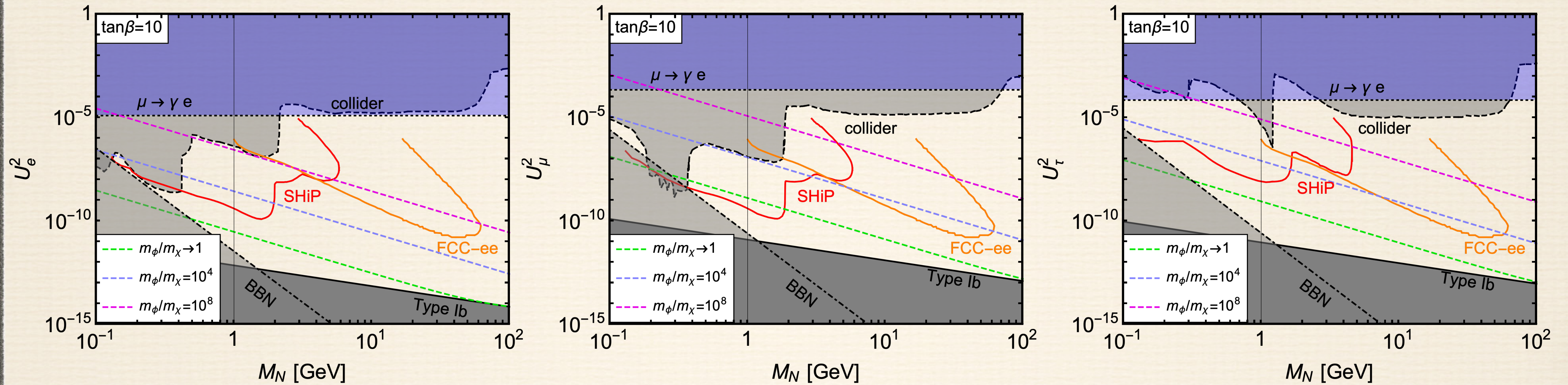
- dark sector processes

❖ If $m_\phi \gg m_\chi, M_N$, the dark matter production only depends on m_ϕ/m_χ

- ν -Yukawa dominance $y^2 (Y_1^2 + Y_2^2) \simeq 2.1 \times 10^{-22} \frac{m_\phi}{m_\chi}$ equal contribution
- dark sector dominance $y^4 \simeq 3.9 \times 10^{-23} \frac{m_\phi}{m_\chi}$ $(Y_1^2 + Y_2^2)^2 \simeq 5.5 \times 10^{-22} \frac{m_\phi}{m_\chi}$

❖ Condition for ν -Yukawa dominance $\frac{m_\phi}{m_\chi} \gtrsim 6.9 \times 10^{-9} \left(\frac{M_N}{1\text{GeV}} \right)^2 \frac{1}{\sin^2 2\beta}$ $\tan \beta = v_2/v_1$

Results



- ❖ ν -Yukawa dominance is allowed above the coloured dashed lines
- ❖ More constrained as m_ϕ/m_χ increases
- ❖ The strongest constraint is given by ν_μ mixing

Leptogenesis

Extended Type Ib Seesaw Mechanism

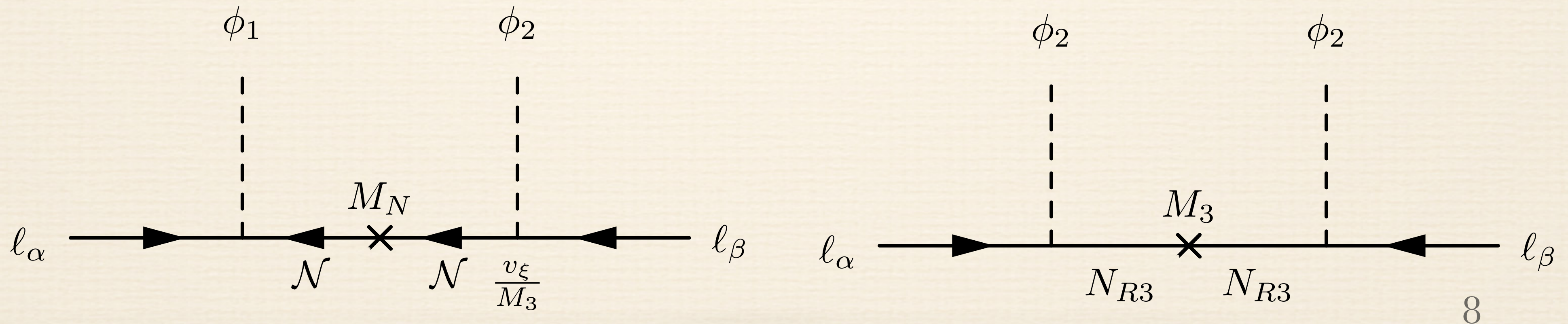
❖ Particles and symmetries

	q_α	$u_{R\beta}$	$d_{R\beta}$	l_α	$e_{R\beta}$	ϕ_1	ϕ_2	N_{R1}	N_{R2}	N_{R3}	ξ
$SU(2)_L$	2	1	1	2	1	2	2	1	1	1	1
$U(1)_Y$	$\frac{1}{6}$	$\frac{2}{3}$	$-\frac{1}{3}$	$-\frac{1}{2}$	-1	$-\frac{1}{2}$	$-\frac{1}{2}$	0	0	0	0
Z_4	1	ω^2	ω	1	ω	ω	ω^2	ω^3	ω	ω^2	ω

❖ Seesaw Lagrangian

$$\begin{aligned} \mathcal{L}_{\text{seesaw}} = & -Y_{1\alpha} \bar{l}_\alpha \phi_1 N_{R1} - Y_{3\alpha} \bar{l}_\alpha \phi_2 N_{R3} - 2Y_{13} \bar{\xi} \overline{N_{R3}^c} N_{R1} - 2Y_{23} \xi \overline{N_{R3}^c} N_{R2} \\ & - M \overline{N_{R1}^c} N_{R2} - \frac{1}{2} M_3 \overline{N_{R3}^c} N_{R3} + \text{h.c.} \end{aligned}$$

❖ Type Ib seesaw mechanism can be realised effectively



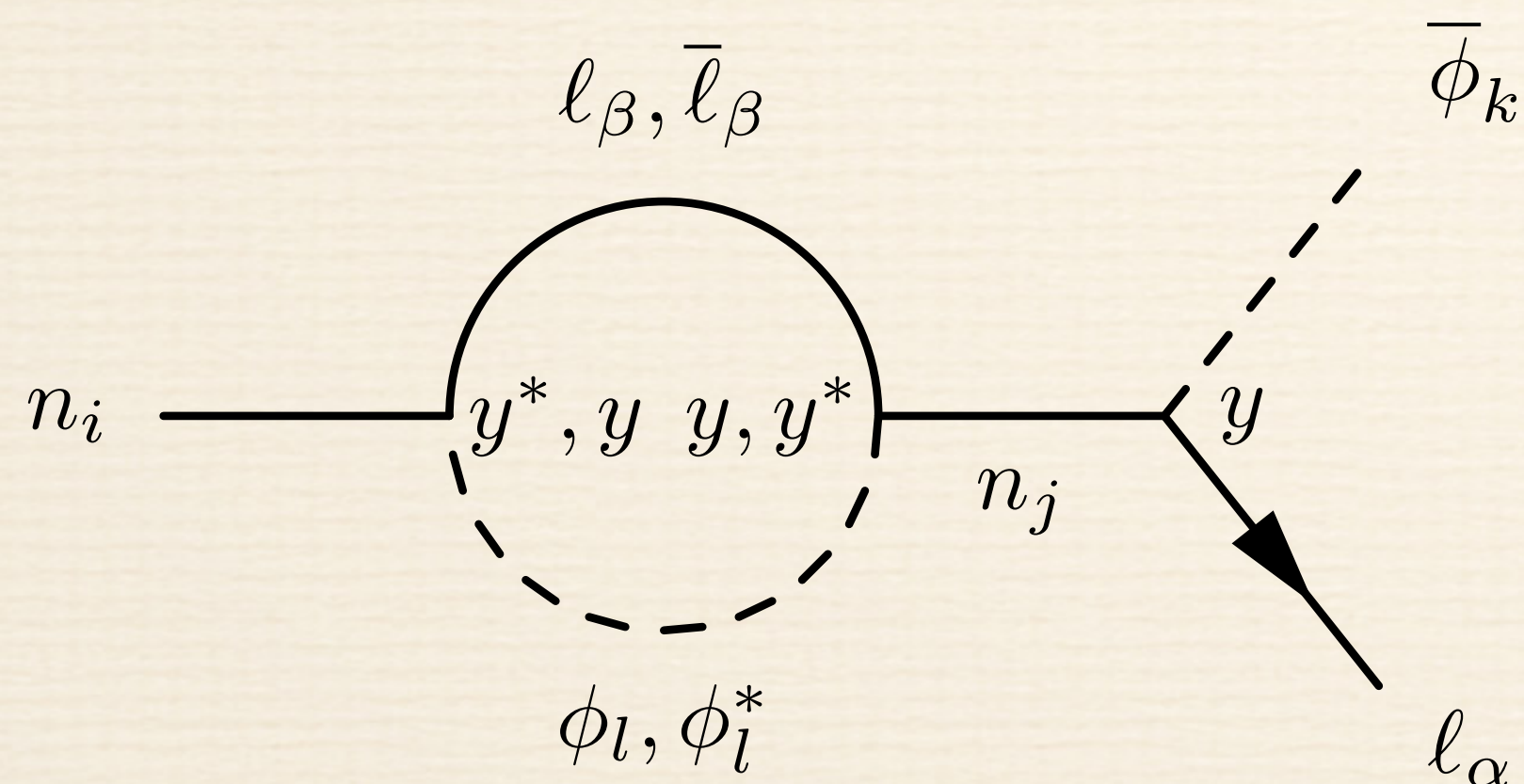
Resonant Leptogenesis

- ❖ N_{R1} and N_{R2} gain mass splitting through mixing with N_{R3}

$$\begin{pmatrix} N_{R1} \\ N_{R2} \\ N_{R3} \end{pmatrix} \rightarrow \begin{pmatrix} n_{R1} \\ n_{R2} \\ n_{R3} \end{pmatrix} = \begin{pmatrix} \frac{1}{\sqrt{2}} \left(1 - \frac{M_{13}^2 - M_{23}^2}{4MM_3}\right)^* & \frac{1}{\sqrt{2}} \left(1 + \frac{M_{13}^2 - M_{23}^2}{4MM_3}\right)^* & -\frac{M_{13}^* + M_{23}^*}{M_3} \\ \frac{i}{\sqrt{2}} \left(1 + \frac{M_{13}^2 - M_{23}^2}{4MM_3}\right)^* & -\frac{i}{\sqrt{2}} \left(1 - \frac{M_{13}^2 - M_{23}^2}{4MM_3}\right)^* & -i\frac{M_{13}^* - M_{23}^*}{M_3} \\ \frac{M_{13}}{\sqrt{2}M_3} & \frac{M_{23}}{\sqrt{2}M_3} & 1 \end{pmatrix} \begin{pmatrix} N_{R1} \\ N_{R2} \\ N_{R3} \end{pmatrix}$$

$$m_1 \simeq M - \frac{(M_{13} + M_{23})^2}{2M_3}, \quad m_2 \simeq M + \frac{(M_{13} - M_{23})^2}{2M_3}, \quad m_3 \simeq M_3$$

- ❖ Lepton number violation



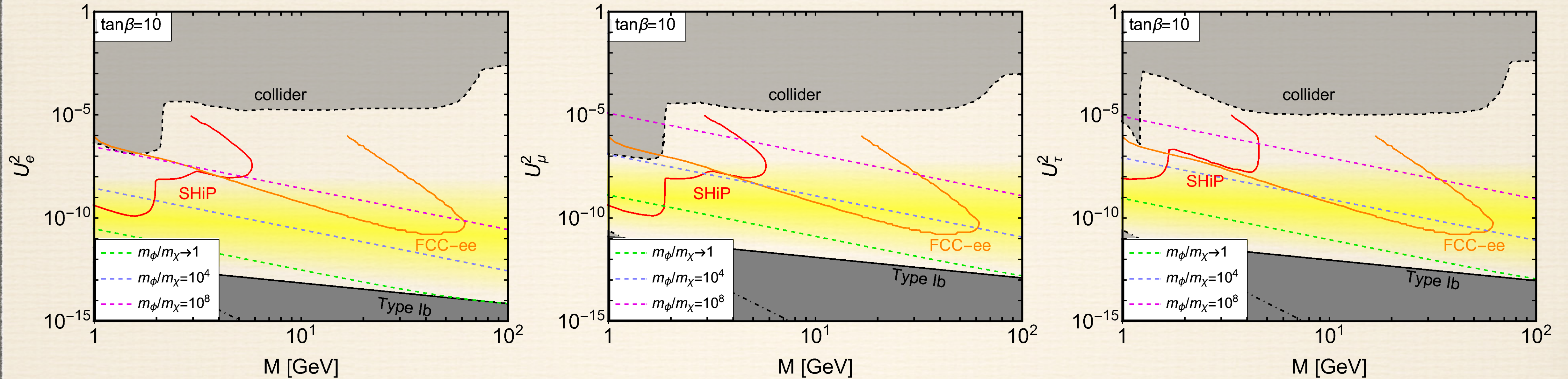
$$\left(\epsilon_{n_1}^{\text{wave}}\right)_{1\alpha} = \left(\epsilon_{n_2}^{\text{wave}}\right)_{1\alpha} \simeq \frac{Y_{1\alpha} Y_{1\alpha}^*}{2Y_1^2} \frac{\Im(M_{13}^2 + M_{23}^2)}{MM_3}$$

$$Y_{\Delta\alpha} \simeq \sum_i \epsilon_{i\alpha} Y_{n_i}^{\text{eq}} \frac{z_{\text{sph}}}{2C_2 \kappa_{i\alpha} |A_{\alpha\alpha}|}$$

$$C_2 \sim 0.1$$

$$z_{\text{sph}} \equiv M/T_{\text{sph}}$$

Results



❖ Theoretical and experimental constraints

- type Ib seesaw dominant
- pseudo-Dirac neutrino
- oscillation data
- observed baryogenesis

❖ Allowed region of mixing strength

$$1.0 \times 10^{-6} \cos^2 \beta \gg U_e^2 \gg 1.5 \times 10^{-14} \frac{v}{M} \cot \beta$$

$$7.3 \times 10^{-6} \cos^2 \beta \gg U_\mu^2 \gg 6.3 \times 10^{-13} \frac{v}{M} \cot \beta$$

$$5.7 \times 10^{-6} \cos^2 \beta \gg U_\tau^2 \gg 4.3 \times 10^{-13} \frac{v}{M} \cot \beta$$

Summary

- ❖ We discussed dark matter production in a minimal extension of type Ib seesaw model with a neutrino portal
- ❖ The DM- ν relation is constrained by experimental results when the heavy neutrino is 1-100 GeV
- ❖ The results can be further constrained by future experiments
- ❖ The type Ib seesaw mechanism can be realised effectively to produce BAU through resonant leptogenesis
- ❖ Leptogenesis and dark matter can be realised simultaneously in the type Ib seesaw model for GeV scale heavy neutrino

Thank You!