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# Gravitational Waves from String Cosmology

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Physics Department, McGill University, Canada

SUSY 2021

# Outline

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# Map of the Cosmic Microwave Background (CMB)

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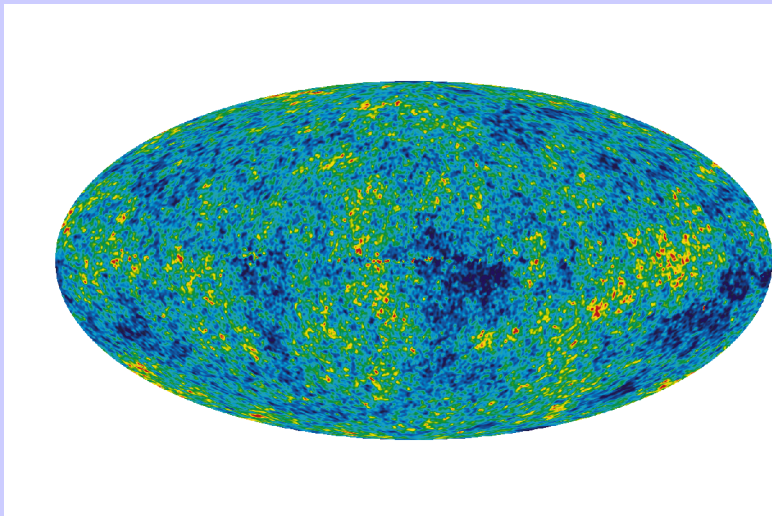
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Credit: NASA/WMAP Science Team

# Angular Power Spectrum of CMB Anisotropies

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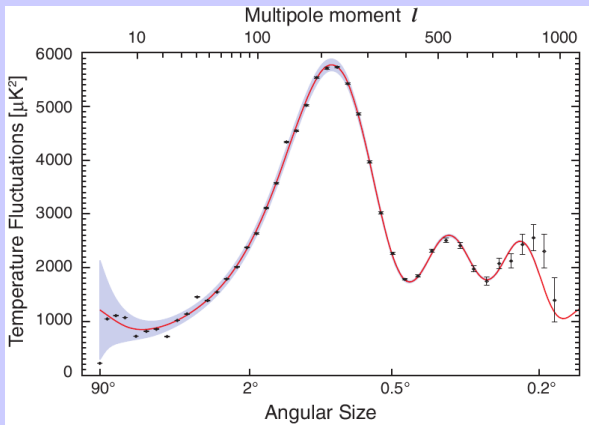
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1970Ap&SS...7....3S

SMALL-SCALE FLUCTUATIONS OF RELIC RADIATION

9

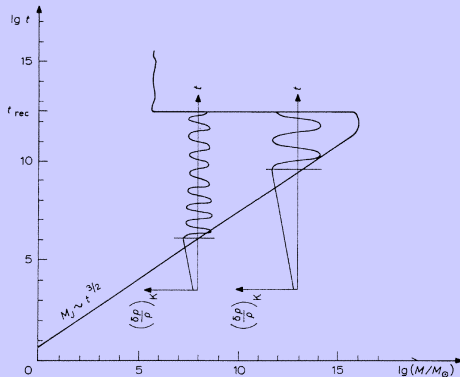


Fig. 1a. Diagram of gravitational instability in the 'big-bang' model. The region of instability is located to the right of the line  $M_J(t)$ ; the region of stability to the left. The two additional lines of the graph demonstrate the temporal evolution of density perturbations of matter: growth until the moment when the considered mass is smaller than the Jeans mass and oscillations thereafter. It is apparent that at the moment of recombination perturbations corresponding to different masses correspond to different phases.

# Key Realization

R. Sunyaev and Y. Zel'dovich, *Astrophys. and Space Science* **7**, 3 (1970); P. Peebles and J. Yu, *Ap. J.* **162**, 815 (1970).

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- Given a **scale-invariant power spectrum of adiabatic fluctuations** on "super-horizon" scales before  $t_{rec}$ , i.e. standing waves.
- → "correct" power spectrum of galaxies.
- → **acoustic oscillations in CMB angular power spectrum.**

# Early Work

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1970arXiv:1701.02736

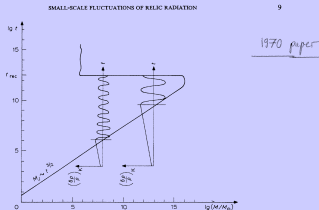


Fig. 1a. Diagram of gravitational instability in the "big-bang" model. The region of instability is located to the right of the line  $M_J(t)$ ; the region of stability to the left. The two additional lines of the graph demonstrate the temporal evolution of density perturbations of matter: growth until the moment when the considered mass is smaller than the Jeans mass and oscillations thereafter. It is apparent that at the moment of recombination perturbations corresponding to different masses correspond to different phases.

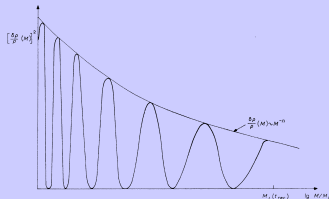


Fig. 1b. The dependence of the square of the amplitude of density perturbations of matter on scale. The fine line designates the usually assumed dependence  $(\delta\rho/\rho)_M \sim M^{-3}$ . It is apparent that fluctuations of relic radiation should depend on scale in a similar manner.

R. Sunyaev & Ya Zeldovich, Astrophysic and Space Science 7

3-19 (1970)



# Predictions from 1970

R. Sunyaev and Y. Zel'dovich, *Astrophys. and Space Science* **7**, 3 (1970); P. Peebles and J. Yu, *Ap. J.* **162**, 815 (1970).

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- → "correct" power spectrum of galaxies.
- → **acoustic oscillations in CMB angular power spectrum**.
- → **baryon acoustic oscillations in matter power spectrum**.

# Key Challenge

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## How does one obtain such a spectrum?

- Inflationary Cosmology is the first scenario based on causal physics which yields such a spectrum.
- But it is not the only one.

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# Hubble radius vs. Horizon

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- **Horizon**: forward light cone of a point on the initial Cauchy surface.
- Horizon carries information about causality.
- **Hubble radius**:  $l_H(t) \equiv H^{-1}(t)$ .
- Hubble radius crucial for the evolution of fluctuations.

# Criteria for a Successful Early Universe Scenario

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- **Horizon**  $\gg$  **Hubble radius** in order for the scenario to solve the “horizon problem” of Standard Big Bang Cosmology.
- Scales of cosmological interest today **originate inside the Hubble radius at early times** in order for a causal generation mechanism of fluctuations to be possible.
- Mechanism for producing a **scale-invariant spectrum of curvature fluctuations** on super-Hubble scales.

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# Inflation as a Solution

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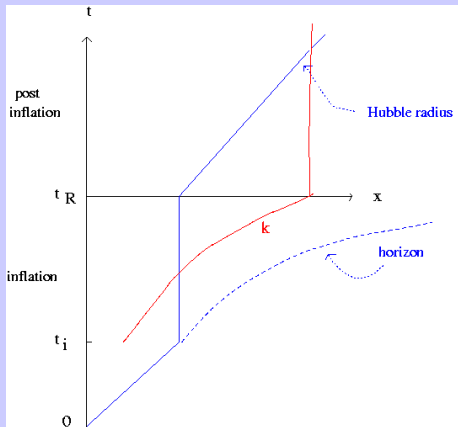
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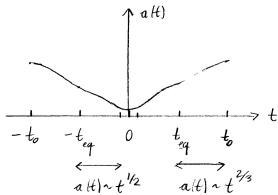
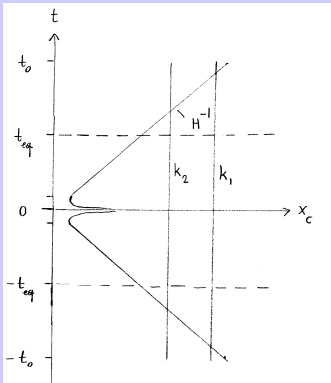
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# Emergent Universe

R.B. and C. Vafa, *Nucl. Phys. B*316:391 (1989)

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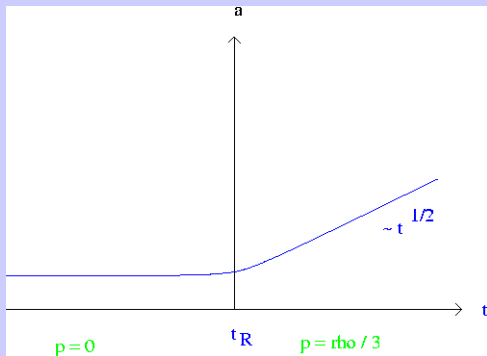
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# Emergent Universe as a Solution

A. Nayeri, R.B. and C. Vafa, *Phys. Rev. Lett.* 97:021302 (2006)

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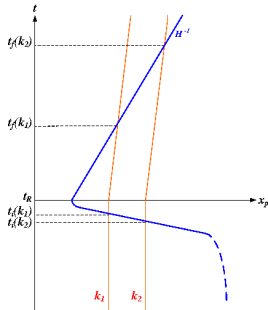
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# Ekpyrotic Bounce

J. Khoury, B. Ovrut, P. Steinhardt and N. Turok, Phys. Rev. D, 2001

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- Space-time described by Einstein-Hilbert action.
- Idea: **Slow contraction** given by matter with equation of state  $w \gg 1$ .
- Obtained by assuming that matter is dominated by a **scalar field  $\varphi$  with a negative exponential potential**.
- Anisotropies diluted, creates spatial flatness
- **Global attractor** in initial condition space (A. Ijjas et al, arXiv:2103.00584)
- Note: **Negative exponential potentials are ubiquitous in string theory.**

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J. Khoury, B. Ovrut, P. Steinhardt and N. Turok, Phys. Rev. D, 2001

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# S-Brane and Ekpyrosis

RB and Z. Wang, arXiv:2001.00638, arXiv:2004.06437

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## Challenges for Ekpyrotic Cosmology:

- How do we get the bounce?
- How do we obtain a scale-invariant spectrum of curvature fluctuations?
- Can we obtain a spectrum of gravitational waves relevant to current observations?

Adding an S-Brane to the EFT action can solve all three problems, and leads to two consistency relations for cosmological observables.

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RB and Z. Wang, arXiv:2001.00638, arXiv:2004.06437

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# Action

**Idea:** At the string scale a new tower of string states becomes comparable in mass to the usual low energy degrees of freedom;  
→ they must be included in the low energy effective action.

Included as an **S-Brane**.

$$S = \int d^4x \sqrt{-g} [R + \frac{1}{2} \partial_\mu \varphi \partial^\mu \varphi - V(\varphi)] \\ - \int d^4x \kappa \delta(\tau - \tau_B) \sqrt{\gamma},$$

$$\kappa \equiv N \eta_S^3,$$

**Note:** The S-brane has  $\rho = 0$  and  $p < 0$  → can mediate the transition between contraction and expansion.

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# Gravitational Waves from S-Brane Ekpyrosis

RB and Z. Wang, arXiv:2004.06437

## Computation of curvature fluctuations and **gravitational waves**:

- Begin with quantum vacuum fluctuations early in the contracting phase.
- Propagate using the usual theory of cosmological fluctuations until  $t = 0$ .
- Matching perturbations across the S-brane using the Israel matching conditions.
- → **scale-invariant spectrum** of curvature fluctuations and **gravitational waves**.

## Consistency relations:

$$\begin{aligned}n_t &= 1 - n_s \\ r &\sim (1 - n_s)^2\end{aligned}$$

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# Reheating from S-Brane Ekpyrosis

RB, K. Dasgupta and ZW, arXiv:2007.01203

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## Reheating after Ekpyrosis

Coupling of the S-brane to Standard Model gauge fields  $\rightarrow$   
gauge field production during S-brane decay.

# Superstrings in the Early Universe

R.B. and C. Vafa, *Nucl. Phys. B316:391 (1989)*

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Based on **new degrees of freedom** and **new symmetries** of string theory.

## New Degrees of Freedom

Assumption: All spatial dimensions toroidal, radius  $R$ .

String states:

- momentum modes:  $E_n = n/R$
- winding modes:  $E_m = mR$
- oscillatory modes:  $E$  independent of  $R$



# Superstrings in the Early Universe

R.B. and C. Vafa, *Nucl. Phys. B316:391 (1989)*

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# New Symmetries: T-Duality

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## T-Duality

- Momentum modes:  $E_n = n/R$
- Winding modes:  $E_m = mR$
- **Duality:**  $R \rightarrow 1/R$   $(n, m) \rightarrow (m, n)$
- Mass spectrum of string states unchanged
- Symmetry of vertex operators
- Symmetry at non-perturbative level  $\rightarrow$  existence of D-branes

# Absence of a Temperature Singularity in String Cosmology

R.B. and C. Vafa, *Nucl. Phys. B*316:391 (1989)

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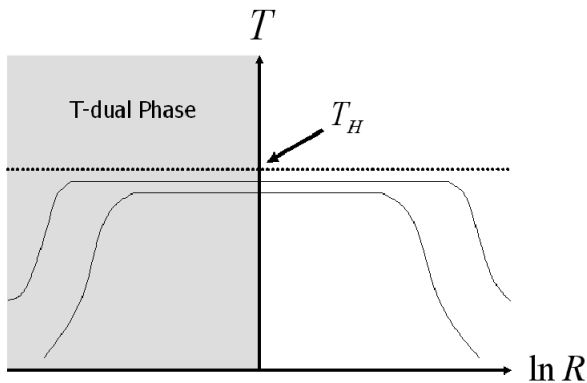
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## Temperature-size relation in string gas cosmology



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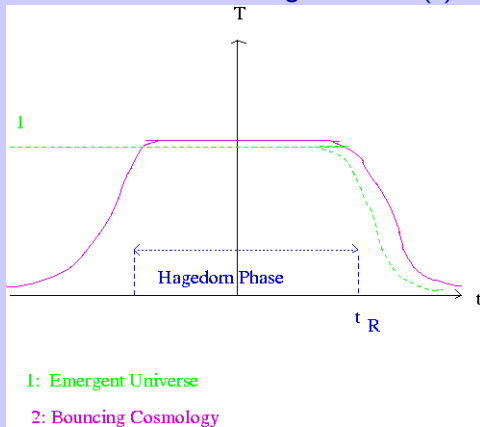
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Assume some action gives us  $R(t)$



# Position Operators

R.B. and C. Vafa, *Nucl. Phys. B*316:391 (1989)

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Position operators (dual to momenta)

$$|x\rangle = \sum_p \exp(ix \cdot p) |p\rangle$$

Dual position operators (dual to windings)

$$|\tilde{x}\rangle = \sum_w \exp(i\tilde{x} \cdot w) |w\rangle$$

# Position Operators

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# Heavy vs. Light Modes

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- $R \gg 1$ : momentum modes light.
- $R \ll 1$ : winding modes light.
- $R \gg 1$ : length measured in terms of  $|x \rangle$ .
- $R \ll 1$ : length measured in terms of  $|\tilde{x} \rangle$
- $R \sim 1$ : both  $|x \rangle$  and  $|\tilde{x} \rangle$  important.

**Conclusion:** At string scale densities usual effective field theory (EFT) based on supergravity will break down.

**Conclusion:** If an effective field theory description is valid, it must be an EFT in 18 spatial dimensions.

# Heavy vs. Light Modes

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- $R \gg 1$ : momentum modes light.
- $R \ll 1$ : winding modes light.
- $R \gg 1$ : length measured in terms of  $|x \rangle$ .
- $R \ll 1$ : length measured in terms of  $|\tilde{x} \rangle$
- $R \sim 1$ : both  $|x \rangle$  and  $|\tilde{x} \rangle$  important.

**Conclusion:** At string scale densities usual effective field theory (EFT) based on supergravity will break down.

**Conclusion:** If an effective field theory description is valid, it must be an EFT in 18 spatial dimensions.



# Physical length operator

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$$l_p(R) = R \quad R \gg 1$$
$$l_p(R) = \frac{1}{R} \quad R \ll 1$$

**Conclusion:** Resolution of the Cosmological Singularity.

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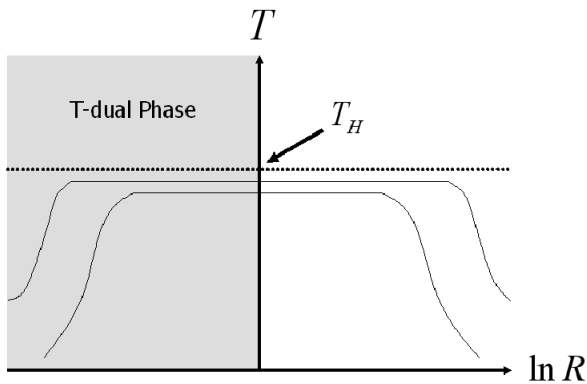
$$l_p(R) = R \quad R \gg 1$$
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**Conclusion:** Resolution of the Cosmological Singularity.

# Absence of a Temperature Singularity in String Cosmology

R.B. and C. Vafa, *Nucl. Phys. B*316:391 (1989)

## Temperature-size relation in string gas cosmology



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# Emergent Universe from String Cosmology

A. Nayeri, R.B. and C. Vafa, *Phys. Rev. Lett.* 97:021302 (2006)

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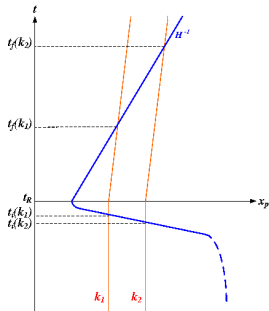
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# Fluctuations from String Cosmology

A. Nayeri, R.B. and C. Vafa, *Phys. Rev. Lett.* 97:021302 (2006), R.B., A. Nayeri, S. Patil and C. Vafa, *Phys. Rev. Lett.* 98:231302 (2007)

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Conclusions

- Perturbations arise as **thermal fluctuations** of the string gas.
- Thermodynamic quantities have **holographic scaling**.
- → scale invariant spectrum of **curvature fluctuations**, red tilt.
- → scale invariant spectrum of **gravitational waves**, blue tilt.
- **Consistency relation:  $n_t \simeq 1 - n_s$**
- **Consistency relation:  $r \sim (1 - n_s)^2$**

# Some Key Steps in the Calculation

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$$\langle |\Phi(k)|^2 \rangle = 16\pi^2 G^2 k^{-4} \langle \delta T_0^0(k) \delta T_0^0(k) \rangle,$$

$$\langle \delta T_0^0(k) \delta T_0^0(k) \rangle = \frac{T^2}{R^6} C_V.$$

$$C_V \simeq \frac{R^2 / l_s^3}{T(1 - T/T_H)}.$$

$$\langle |h(k)|^2 \rangle = 16\pi^2 G^2 k^{-4} \langle \delta T_j^i(k) \delta T_j^i(k) \rangle, \quad i \neq j,$$

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- **Swampland conjectures:** Inflation hard to realize in superstring theory.
- **Trans-Planckian Censorship Conjecture:** causality and unitarity problem for inflation.

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- There is a vast **landscape** of **effective field theories**.
- Any space-time dimension, and number of fields, any shape of the potential, any field range.
- **Superstring theory** is very **constraining**.
- Only a **small subset** of all EFTs is consistent with string theory.
- The rest lie in the **swampland**.

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# String Cosmology

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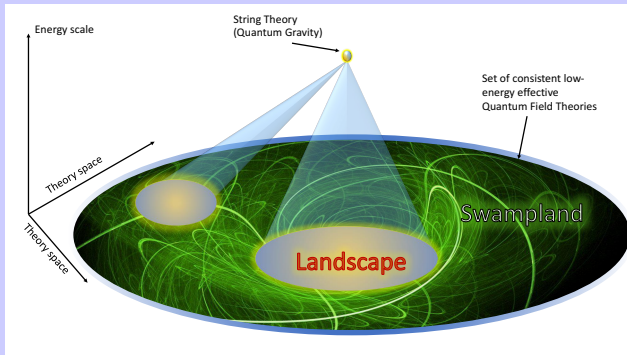
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# Swampland Conjectures

H. Ooguri and C. Vafa, hep-th/0605264; G. Obied, H. Ooguri, L. Spodyneiko and C. Vafa, arXiv:1806.08362; S. Garg and C. Krishnan, arXiv:1807.05193; H. Ooguri, E. Palti, G. Shiu and C. Vafa, arXiv:1810.05506.

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## What are conditions for habitable islands sticking out from the swamp?

- The effective field theory is only valid for  $\Delta\varphi < dm_{pl}$  (field range condition).
- The potential of  $\varphi$  obeys (de Sitter conjecture)

$$\left| \frac{V'}{V} \right| m_{pl} \geq c_1 \text{ or}$$
$$\frac{V''}{V} m_{pl}^2 \leq -c_2$$

Note:  $d, c_1, c_2$  constants of order 1.

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H. Ooguri and C. Vafa, hep-th/0605264; G. Obied, H. Ooguri, L. Spodyneiko and C. Vafa, arXiv:1806.08362; S. Garg and C. Krishnan, arXiv:1807.05193; H. Ooguri, E. Palti, G. Shiu and C. Vafa, arXiv:1810.05506.

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- No canonical single field inflation.
- No eternal inflation.
- No bare positive  $\Lambda$ .
- Dark Energy is not a bare cosmological constant.
- Quintessence dark energy is constrained (L. Heisenberg et al, arXiv:1808.02877).

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# Trans-Planckian Problem

J. Martin and R.B., *Phys. Rev. D63*, 123501 (2002)

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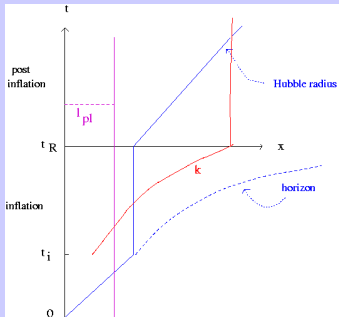
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- **Success of inflation:** At early times scales are inside the Hubble radius  $\rightarrow$  causal generation mechanism is possible.
- **Problem:** If time period of inflation is more than  $70H^{-1}$ , then  $\lambda_p(t) < l_{pl}$  at the beginning of inflation.
- $\rightarrow$  breakdown of effective field theory; new physics **MUST** enter into the calculation of the fluctuations.

# Trans-Planckian Censorship Conjecture (TCC)

A. Bedroya and C. Vafa., arXiv:1909.11063

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No trans-Planckian modes exit the Hubble horizon.

$$ds^2 = dt^2 - a(t)^2 dx^2$$

$$H(t) \equiv \frac{\dot{a}}{a}(t)$$

$$\frac{a(t_R)}{a(t_i)} \Big|_{pl} < H(t_R)^{-1}$$

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A. Bedroya and C. Vafa., arXiv:1909.11063

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A. Bedroya and C. Vafa., arXiv:1909.11063

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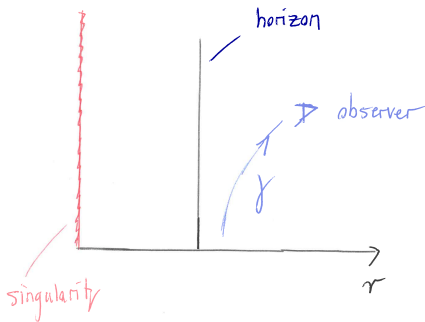
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# Justification

R.B. arXiv:1911.06056

## Analogy with Penrose's Cosmic Censorship Hypothesis:



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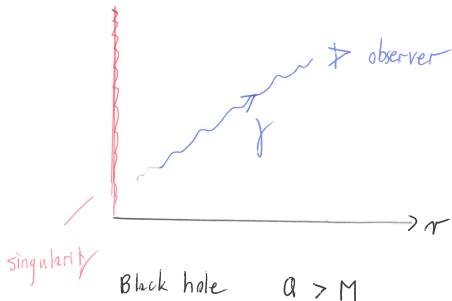
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R.B. arXiv:1911.06056

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- General Relativity allows for solutions with **timelike singularities**: super-extremal black holes.
- → Cauchy problem not well defined for observer external to black holes.
- Evolution **non-unitary** for external observer.
- Conjecture: ultraviolet physics → **external observer** shielded from the **singularity** and **non-unitarity** by **horizon**.

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R.B. arXiv:1911.06056

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## Translation

- Position space  $\rightarrow$  momentum space.
- Singularity  $\rightarrow$  trans-Planckian modes.
- Black Hole horizon  $\rightarrow$  Hubble horizon.

Observer outside of Hubble horizon must be shielded from the trans-Planckian modes.

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# Why Hubble Horizon?

R.B. arXiv:1911.06056; A. Bedroya and C. Vafa., arXiv:1909.11063

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- Recall: Fluctuations only oscillate on sub-Hubble scales.
- Recall: Fluctuations freeze out, become **squeezed states** and **classicalize** on super-Hubble scales.
- **Demand:** classical region be insensitive to trans-Planckian region.
- → no trans-Planckian modes ever exit Hubble horizon.



# Why Hubble Horizon?

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- Recall: **non-unitarity** of quantum field theory in an expanding universe (N. Weiss, Phys. Rev. D32, 3228 (1985)).
- $\mathcal{H}$  is the product Hilbert space of a harmonic oscillator Hilbert space for all comoving wave numbers  $k$
- Fixed  $k_{min}$ , time dependent  $k_{max}$  :  $k_{max}(t)a(t)^{-1} = m_{pl}$
- **Demand:** classical region be insensitive to non-unitarity.
- $\rightarrow$  no trans-Planckian modes ever exit Hubble horizon.

# Justification

R.B. arXiv:1911.06056; A. Bedroya and C. Vafa., arXiv:1909.11063

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# Application to Inflation

A. Bedroya, R.B., M. Loverde and C. Vafa., arXiv:1909.11106

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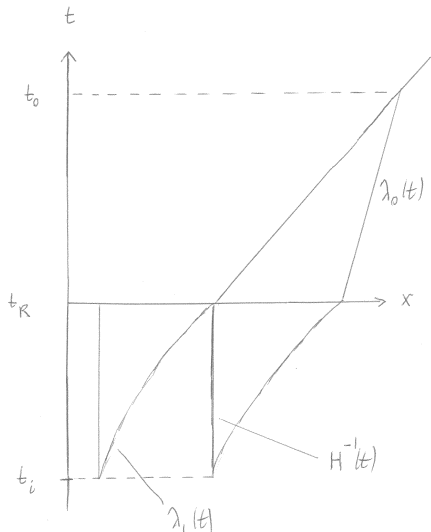
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A. Bedroya, R.B., M. Loverde and C. Vafa., arXiv:1909.11106

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TCC implies:

$$\frac{a(t_R)}{a(t_*)} |_{pl} < H(t_R)^{-1}$$

Demanding that inflation yields a causal mechanism for generating CMB anisotropies implies:

$$H_0^{-1} \frac{a(t_0)}{a(t_R)} \frac{a(t_R)}{a(t_*)} < H^{-1}(t_*)$$

# Application to Inflation

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# Implications

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**Upper bound** on the **energy scale of inflation**:

$$V^{1/4} < 3 \times 10^9 \text{GeV}$$

→ **upper bound** on the **primordial tensor to scalar ratio**  $r$ :

$$r < 10^{-30}$$

Note: Secondary tensors will be larger than the primary ones.

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**Upper bound** on the **energy scale of inflation**:

$$V^{1/4} < 3 \times 10^9 \text{GeV}$$

→ **upper bound** on the **primordial tensor to scalar ratio**  $r$ :

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Dark Energy cannot be a bare cosmological constant.

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- Data do not require an initial period of cosmological inflation.
- String theory → breakdown of effective field theory at string scale.
- String theory → constraints on low energy effective field theory (“Swampland” program).
- Inflation does **not naturally** emerge from superstring theory.
- Inflation is **highly fine tuned** assuming validity of the TCC.
- String Cosmology models which do not involve inflation and which are consistent with the Swampland Criteria and the TCC exist (e.g. String Gas Cosmology, S-Brane Ekpyrosis).
- **Superstring Cosmology** models yield a scale-invariant spectrum of gravitational waves with a **slight blue tilt**.

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