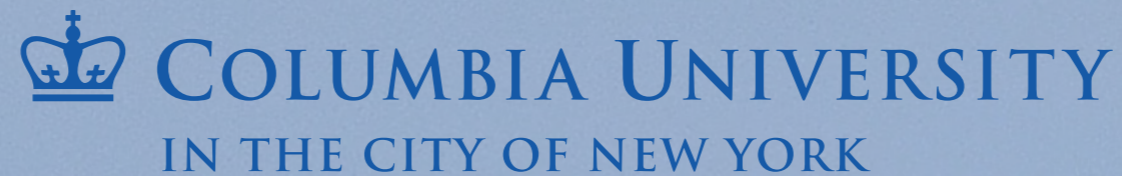


Chiral models of composite axions and accidental Peccei-Quinn symmetry

Alessandro Podo



Based on

R. Contino, AP, F. Revello - arXiv: 2108.xxxxx

SUSY 21 - 26 August 2021

The Standard Model paradigm

- Gauge dynamics
- Global symmetries are accidental
 - renormalizable lagrangian
 - higher-dimensional operators become irrelevant in the IR
 - global symmetries emerge as accidental symmetries
 - flavour symmetries and custodial $SO(3)$
 - Baryon number $U(1)_B$ and lepton number $U(1)_L$



+ [...]

- Fermions in chiral representations
 - bare fermion masses forbidden; generated dynamically

Bonus

- Unification of gauge couplings ?
 - fermions in complete GUT multiplets

Chiral gauge theories



the good

- Fermions in complex representations of the gauge group

No bare mass term $\delta\mathcal{L} = Q^\dagger \sigma^\mu i D_\mu Q - M \cancel{Q} Q$

- Dynamical generation of all energy scales
- Cancellation of gauge anomalies gives non-trivial constraints

Chiral gauge theories



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-
- Non-abelian simple gauge group: what happens at confinement?
 - No lattice simulations for chiral gauge theories

IR behaviour of theories with simple gauge group not understood



the bad

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the bad



the ugly

- Consider chiral product gauge group with confining vector-like factor

A calculable model with a chiral sector and SM interactions

Consider a massless vector-like theory

	SU(N _{DC})
ψ_1	\square
ψ_2	\square
χ_1	$\bar{\square}$
χ_2	$\bar{\square}$

Symmetry Breaking Pattern:

$$SU(2) \times SU(2) \times U(1)_V \longrightarrow SU(2)_V \times U(1)_V$$

spontaneous

A calculable model with a chiral sector and SM interactions

and consider the weak gauging of a subgroup $U(1)_D$

	SU(N _{DC})	U(1) _D
ψ_1	\square	+1
ψ_2	\square	-1
χ_1	$\bar{\square}$	- a
χ_2	$\bar{\square}$	+ a

chiral theory for $0 \leq a < 1$

Harigaya, Nomura - PRD 94 (2016) 035013

Co, Harigaya, Nomura - PRL 118 (2017) 101801

Symmetry Breaking Pattern:

$$SU(2) \times SU(2) \times U(1)_V \longrightarrow SU(2)_V \times U(1)_V$$

explicit

$$\longrightarrow U(1)_{3V} \times U(1)_V$$

$U(1)_D$ acquires a mass from the bilinear condensate $m_{\gamma_D} \sim f e_D$

A calculable model with a chiral sector and SM interactions

- chiral fermions should come in complete GUT multiplets

	SU(N _{DC})	U(1) _D	SU(5) _{GUT}
ψ_1	\square	+1	\square
ψ_2	\square	-1	\square
χ_1	$\bar{\square}$	-a	$\bar{\square}$
χ_2	$\bar{\square}$	+a	$\bar{\square}$

Contino, AP, Revello - JHEP 02 (2021) 091

99 NGBs :

pseudo eaten by γ_D

$$24^\pm \oplus 24^0 \oplus 24^{0'} \oplus 1^\pm \oplus 1^0$$

GUT breaking

$$\supset 1^0 \oplus 1^{0'} \quad \text{SM singlets}$$

one (combination) of the two singlets has anomalous couplings to SM

→ axion-like particle

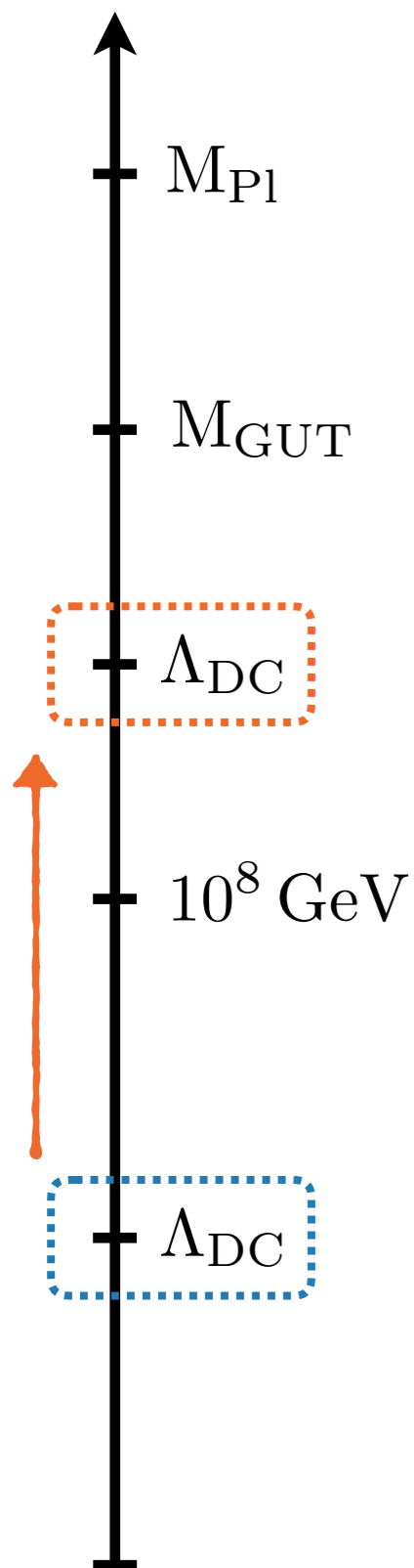
$$\delta m_{\tilde{a}}^2 \sim \frac{g_{\text{GUT}}^2}{16\pi^2} \frac{\Lambda_{\text{DC}}^4}{M_{\text{GUT}}^2}$$

GUT contribution

✗ not QCD axion

Accidental Peccei-Quinn symmetry

- raise the dynamical scale to PQ scale
- add SM singlets to obtain an exact QCD axion



	$SU(N_{DC})$	$U(1)_D$	$SU(5)_{GUT}$	$U(1)_{PQ}$
ψ_1	\square	$+1$	\square	$+1$
ψ_2	\square	-1	\square	$+1$
ψ_3	\square	$+1$	$\mathbf{1}$	-3
ψ_4	\square	-1	$\mathbf{1}$	-3
χ_1	$\overline{\square}$	$-a$	$\overline{\square}$	$+1$
χ_2	$\overline{\square}$	$+a$	$\overline{\square}$	$+1$
χ_3	$\overline{\square}$	$-a$	$\mathbf{1}$	-3
χ_4	$\overline{\square}$	$+a$	$\mathbf{1}$	-3

- accidental
- anomalous
- spont. broken

composite QCD axion

cfr.
 Kim - PRD 31 (1985) 1733
 Randall - PLB 284 (1992) 77
 Dobrescu - PRD 55 (1997) 5826
 + [...]

The axion quality problem

UV effects : PQ breaking operator with dimension Δ_{PQ}



IR effect

$$\Delta\theta \approx |c| \phi_{CP} \left(\frac{M_{\text{Pl}}}{\Lambda_{\text{QCD}}} \right)^4 \left(\frac{f_a}{M_{\text{Pl}}} \right)^{\Delta_{PQ}}$$

- neutron EDM experiment: $\Delta\theta \lesssim 10^{-10}$ nEDM - PRL 124 (2020) 8, 081803

for $f_a \lesssim 10^{12}$ GeV $\rightarrow \Delta_{PQ} \gtrsim 12$

- in the model we just defined: $\mathcal{O} = \psi_1 \psi_2 \chi_1 \chi_2$

Selection rules on PQ violating operators

- Not every PQ violating operator is dangerous!
- A generic PQ violating operator generates a potential only if it has non vanishing matrix element with a state containing axions:

$$\langle \psi_a | \mathcal{O}_{PQ} | 0 \rangle \neq 0$$

- The operator must be an interpolating operator for the axion with vanishing vectorial charges.

→ \mathcal{O}_{PQ} polynomial in $(\psi_r \chi_{\bar{r}})$, $(\psi_r \chi_{\bar{r}})^*$, $(\psi_r^\dagger \psi_r)$ and $(\chi_{\bar{r}}^\dagger \chi_{\bar{r}})$

caveat!

- It can be a composite operator built from the insertion of N local operators

$$d_{\text{eff}} = \sum_{i=1}^N d_i - 4(N - 1)$$

High quality accidental PQ

Can we modify our construction to have an high quality composite axion? **Yes!**

	$SU(N_{\text{DC}})$	$U(1)_D$	G_{SM}	$U(1)_{\text{PQ}}$
ψ_1	\square	p_1	r_1	α
ψ_2	\square	p_2	r_2	β
ψ_3	\square	p_3	r_3	γ
χ_1	$\bar{\square}$	q_1	\bar{r}_1	α
χ_2	$\bar{\square}$	q_2	\bar{r}_2	β
χ_3	$\bar{\square}$	q_3	\bar{r}_3	γ

G_{SM}	r_1	r_2	r_3
$SU(5)_{\text{GUT}}$	1	5	10
	1	$\bar{5}$	10
$SU(3)_c$	1	3	6
	1	$\bar{3}$	6

- confining $SU(N)$
- perturbativity of SM up to MPI
- accidental PQ

with appropriate choices of $U(1)$ charges PQ is protected up to dim 12

dim 12 operators always exist in these models as a consequence of anomaly cancellation

High quality axion - GUT model

Robust model: high quality irrespectively of GUT scalar sector

	SU(N_{DC})	U(1) _D	SU(5) _{GUT}	U(1) _{PQ}	
ψ_1	\square	+2	1	+5	
ψ_2	\square	+3	$\bar{5}$	+1	
ψ_3	\square	-5	10	-1	

χ_1	$\bar{\square}$	+3	1	+5	+ others!
χ_2	$\bar{\square}$	-6	5	+1	
χ_3	$\bar{\square}$	+6	$\bar{10}$	-1	

$$N_{\text{DC}} = 5$$

dim 12 PQ violating operator: $\mathcal{O}_1 = \psi_2 \chi_2 (\psi_3 \chi_3)^3$

QCD axion low energy couplings

- Axion low energy couplings
 - well-predicted but not distinctive low energy couplings
 - common to “hadronic axion” models

$$m_a = 5.70(7) \left(\frac{10^{12} \text{ GeV}}{f_a} \right) \mu\text{eV},$$

$$g_{a\gamma\gamma} = \frac{\alpha_{em}}{2\pi f_a} \left(\frac{E}{N} - 1.92(4) \right),$$

$$c_p = -0.47(3), \quad c_n = -0.02(3),$$

GUT models have fixed $E/N = 8/3$

Grilli di Cortona, Hardy, Vega, Villadoro - JHEP 05 (2016) 104

Cosmological evolution

- Cosmological evolution:

- PQ breaking after the end of inflation is disfavoured for composite models

- accidentally stable heavy resonances + domain walls



- PQ broken during inflation

- axion populated with misalignment mechanism (possibly DM)



- in the GUT scenario, parametrically lighter metastable NGBs

$$\delta m_{\tilde{a}}^2 \sim \frac{g_{\text{GUT}}^2}{16\pi^2} \frac{\Lambda_{\text{DC}}^4}{M_{\text{GUT}}^2}$$

- can leave imprints in cosmological observables

High-quality axion summary

Summary:

- Chiral model with dynamical generation of scales
- Strong CP solved by the QCD axion
- PQ is an high quality accidental symmetry through gauge protection
- Possibly DM
- Compatibility with SU(5) unified dynamics

Outlook:

- Observational distinctive signatures from pseudo-NGB ?

Backup Slides

High quality axion models classification

GUT n_{max}	$N_{DC} = 5$			$N_{DC} = 4$		
	10	15	20	10	15	20
AC solutions	77	189	341	77	189	341
HQ axions, $(\mathbf{1}, \bar{\mathbf{5}}, \mathbf{10})$	22	68	150	9	31	82
Robust HQ axions, $(\mathbf{1}, \bar{\mathbf{5}}, \mathbf{10})$	4	16	47	0	5	22
HQ axions, $(\mathbf{1}, \mathbf{5}, \mathbf{10})$	14	44	99	2	12	33
Robust HQ axions, $(\mathbf{1}, \mathbf{5}, \mathbf{10})$	3	16	36	0	7	21
No $d \leq 8$ operators	0	1	10	0	0	0

QCD n_{max}	$N_{DC} = 3$		
	10	15	20
AC solutions	16	40	96
HQ axions, $(\mathbf{1}, \bar{\mathbf{3}}, \mathbf{6})$	3	14	50
HQ axions, $(\mathbf{1}, \mathbf{3}, \mathbf{6})$	4	17	50
No $d \leq 8$ operators	0	1	8