

# Comments on systematics of corrected MB distributions

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# For discussion (1)

- The event sample to normalize to
  - Sample defined by physics process
    - INEL (ALICE)
    - NSD (less well defined) (ALICE, CMS)
  - Sample defined by detector acceptance
    - INEL with at least one particle in some acceptance window
      - Published up to now:
      - ALICE in  $\eta$  window
      - ATLAS in both  $\eta$  and  $p_T$  window

# For discussion (2)

- The track sample to be normalized
  - Correct down to  $p_T = 0$  (used by ALICE and CMS)
  - Use a  $p_T$  cut-off defined by tracking efficiency (used by everybody)

# Sample defined by physics process

- Advantage

- Comparison with theoretical models, not necessarily a MC
- Comparison with other experiment

- Disadvantage

- Model dependent
  - One has to correct for the non-observed (triggered - selected) events. However, this extrapolation is not done just using the models: usually one uses constraint from ones own data and/or previous measurements.

# Sample defined by detector acceptance

- Advantage

- less model dependent, but not completely
- for finite tracking inefficiency  $\varepsilon$

$$N_{event} = \sum_M P_M \langle -\varepsilon^M \rangle$$

$$N_{track} = \sum_M P_M M \langle -\varepsilon \rangle$$

- $\langle N \rangle$  depends on true multiplicity distribution at low multiplicity ( $M$ )

# Sample defined by detector acceptance

- Disadvantage
  - This is suitable only for models formulated as Monte Carlo generators, but there are models which are not.  
Example: QGSM

# Correction down to $p_T=0$

- Advantage
  - Again comparison with non-MC models and with other experiments.
  - This correction can only be used when an experiment reaches low tracking efficiency for very low  $p_T$ 
    - The correction is small and very well-constrained (one measures the value on a already falling spectrum and it is constrained to zero at  $p_T=0$ )
    - in ALICE conservative error of  $\sim 0.5\%$  on mean multiplicity

# Correction down to $p_T=0$

- Disadvantage
  - One assumes that  $p_T$  dependence is falling down monotonically with  $p_T$  from 50MeV/c (ALICE) to 0
  - However, that seems reasonable and the model dependence is very small.



# Use a $p_T$ cut-off defined by tracking efficiency

- Advantage: well suited for MC model comparison,
  - However, there are still model dependent corrections (unless one goes to very large  $p_T \sim 1.5$  GeV).
  - At 500 MeV p and anti-p will have substantial energy loss (see Bethe-Bloch), one has to correct for this to get the momentum at primary vertex to do the  $p_T$  cut !
  - This correction is purely model dependent (relative amount of protons (even worse antiprotons)) their momentum spectra, and ultimately uncorrectable are  $dE/dx$  fluctuations).
  - Usually this uncertainty is much larger than the correction towards 0 (from a much lower cut-off).

# Use a $p_T$ cut-off defined by tracking efficiency

- Disadvantage
  - Comparison with non-MC models and other experiments practically impossible
  - Tuning of MC models in restricted range (1/3 of charged particles in  $p_T$  in the case of ATLAS)

# Conclusions

- Event selection
  - we have to do both, events samples defined by physics processes and detector acceptance
- Track selection
  - Correction or cut-off, depends on what is more suitable for a given detector and particular reconstruction method, i.e. which gives less systematic uncertainty and larger coverage.