Physics at Future Circular Colliders

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The open questions in HEP

• Data driven:

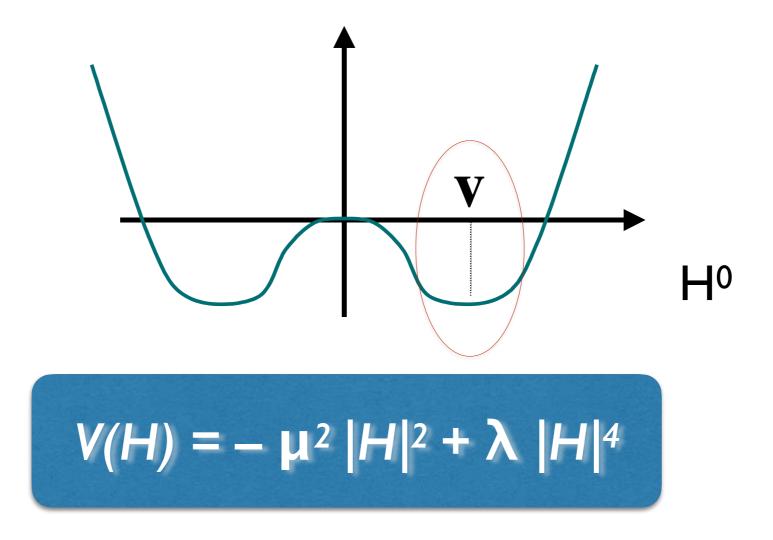
- DM
- Neutrino masses
- Matter vs antimatter asymmetry
- Dark energy
- ...

• Theory driven:

- The hierarchy problem and naturalness
- The flavour problem (origin of fermion families, mass/mixing pattern)
- Quantum gravity
- Origin of inflation
- ...

We have no guarantees as to where answers to these questions will come from, and what are the experiments that will eventually answer them.

But there is one question that can only be addressed by colliders, and future collider efforts must focus on its thorough exploration



Where does this come from?

- The search for the origin of the Higgs and EW symmetry breaking is justified independently of prejudice on the relevance of theoretical puzzles like the hierarchy problem
- It is reasonable to expect that the dynamics underlying the Higgs phenomenon sits nearby the EW scale, justifying the yet unfulfilled hope that new physics should be seen by LHC...
- .. thus many theoretical ideas are emerging, postponing to much higher energies or to alternative scenarios the framework to understand the origin of the weak scale
- The detailed experimental investigation of Higgs properties remains nevertheless a sine qua non condition to make progress no matter what is our bias

The importance of the in-depth exploration of the Higgs properties was acknowledged by the 2020 update of the European Strategy for Particle Physics:

> "An electron-positron Higgs factory is the highest-priority next collider"

- The precision measurement of Higgs properties must be a guaranteed deliverable of all future colliders
- Whether the measurements will challenge or confirm the SM properties, these measurements are a key ingredient in exploration of physics beyond the SM.
- Should they show deviations from the SM, the hint to BSM will be explicit, and the correlations among the various deviations will guide the interpretation of their origin
- Should they agree with the SM, the more accurate the measurements, the more constraining their power in identifying the microscopic origin of possible BSM effects observed in other parts of the programme
 - The LEP precision measurements are still today an essential constraint in evaluating BSM models proposed whenever some anomaly is detected in the data

Other important open issues on the Higgs sector

- Is the Higgs the only (fundamental?) scalar field, or are there other Higgslike states (e.g. H[±], A⁰, H^{±±}, ..., EW-singlets,) ?
 - Do all SM families get their mass from the **<u>same</u>** Higgs field?
 - Do I₃=1/2 fermions (up-type quarks) get their mass from the <u>same</u> Higgs field as I₃=-1/2 fermions (down-type quarks and charged leptons)?
 - Do Higgs couplings conserve flavour? $H \rightarrow \mu \tau$? $H \rightarrow e \tau$? $t \rightarrow Hc$?
- Is there a deep reason for the apparent metastability of the Higgs vacuum?
- Is there a relation among Higgs/EWSB, baryogenesis, Dark Matter, inflation?
- What happens at the EW phase transition (PT) during the Big Bang?
 - what's the order of the phase transition?
 - are the conditions realized to allow EW baryogenesis?
 - the Higgs discovery does not close the book, it opens a whole new chapter of exploration, based on precise measurements of its properties, which can only rely on a future generation of colliders

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Key question for the future developments of HEP: Why don't we see the new physics we expected to be present around the TeV scale ?

- Is the mass scale beyond the LHC reach ?
- Is the mass scale within LHC's reach, but final states are elusive to the direct search ?

These two scenarios are a priori equally likely, but they impact in different ways the future of HEP, and thus the assessment of the physics potential of possible future facilities

Readiness to address both scenarios is the best hedge for the field:

- precision \Rightarrow higher statistics, better detectors and experimental conditions
- sensitivity (to elusive signatures) \Rightarrow ditto
- •extended energy/mass reach ⇒ higher energy

From ESPP 2020:

"Europe, together with its international partners, should investigate the technical and financial feasibility of a future hadron collider at CERN with a centre-of-mass energy of at least 100 TeV and with an electronpositron Higgs and electroweak factory as a possible first stage. "

Answer to these challenges: Future Circular Collider

LHC

http://cern.ch/fcc

Switzerland

France

100km tunnel

FCC 100 km circumference

FCC-ee: e+e- @ 91, 160, 240, 365 GeV
FCC-hh: pp @ 100 TeV
FCC-eh: ебодек ръотек @ 3.5 TeV

What the future <u>circular</u> collider can offer

• <u>Guaranteed deliverables</u>:

 study of <u>Higgs</u> and <u>top</u> quark properties, and exploration of <u>EWSB</u> phenomena, with the best possible precision and sensitivity

• Exploration potential:

- exploit both direct (large Q²) and indirect (precision) probes
- enhanced mass reach for direct exploration at 100 TeV
 - E.g. match the mass scales for new physics that could be exposed via indirect precision measurements in the EW and Higgs sector
- <u>Provide firm Yes/No answers</u> to questions like:
 - is there a TeV-scale solution to the hierarchy problem?
 - is DM a thermal WIMP?
 - could the cosmological EW phase transition have been 1st order?
 - could baryogenesis have taken place during the EW phase transition?
 - could neutrino masses have their origin at the TeV scale?

• ...

Event rates: examples

FCC-ee	н	Z	W	t	т(←Z)	b(←Z)	c(←Z)
	10 ⁶	5 10 ¹²	10 ⁸	10 ⁶	3 10 ¹¹	1.5 10 ¹²	10 ¹²
FCC-hh		н	b	t	W(*	←t) т (←W←t)
	2.5	10 ¹⁰	10 ¹⁷	10 ¹²	10	12	10 ¹¹
FCC-e	h		н			t	
			2.5 10 ⁶			2 10 ⁷	

(1) guaranteed deliverables: Higgs properties

https://arxiv.org/pdf/1708.08912.pdf

	Model	$b\overline{b}$	$c\overline{c}$	gg	WW	au au	ZZ	$\gamma\gamma$	$\mu\mu$
1	MSSM [40]	+4.8	-0.8	- 0.8	-0.2	+0.4	-0.5	+0.1	+0.3
2	Type II 2HD $[42]$	+10.1	-0.2	-0.2	0.0	+9.8	0.0	+0.1	+9.8
3	Type X 2HD $[42]$	-0.2	-0.2	-0.2	0.0	+7.8	0.0	0.0	+7.8
4	Type Y 2HD $[42]$	+10.1	-0.2	-0.2	0.0	-0.2	0.0	0.1	-0.2
5	Composite Higgs [44]	-6.4	-6.4	-6.4	-2.1	-6.4	-2.1	-2.1	-6.4
6	Little Higgs w. T-parity [45]	0.0	0.0	-6.1	-2.5	0.0	-2.5	-1.5	0.0
$\overline{7}$	Little Higgs w. T-parity $[46]$	-7.8	-4.6	-3.5	-1.5	-7.8	-1.5	-1.0	-7.8
8	Higgs-Radion $[47]$	-1.5	- 1.5	+10.	-1.5	-1.5	-1.5	-1.0	-1.5
9	Higgs Singlet [48]	-3.5	-3.5	-3.5	-3.5	-3.5	-3.5	-3.5	-3.5

5 – 10 %

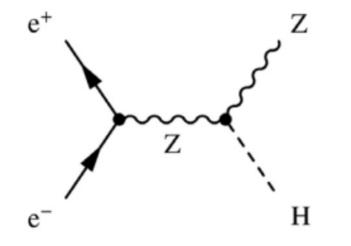
> 10%

NB: when the b coupling is modified, BR deviations are smaller than the square of the coupling deviation. Eg in model 5, the BR to b, c, tau, mu are practically SM-like

(sub)-% precision must be the goal to ensure 3-5σ evidence of deviations, and to cross-correlate coupling deviations across different channels

<u>The absolutely unique power of $e^+e^- \rightarrow ZH$ (circular or linear):</u>

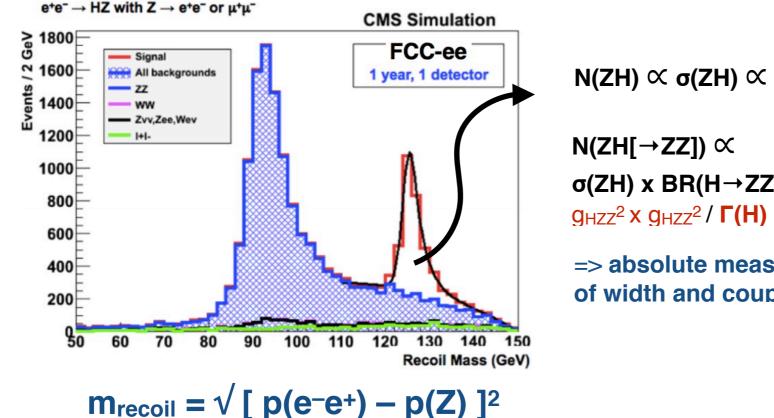
- the model independent absolute measurement of HZZ coupling, which allows the subsequent:
 - sub-% measurement of couplings to W, Z, b, T
 - % measurement of couplings to gluon and charm



 $p(H) = p(e^{-}e^{+}) - p(Z)$

 $= [p(e^-e^+) - p(Z)]^2$ peaks at m²(H)

reconstruct Higgs events independently of the Higgs decay mode!



 $N(ZH) \propto \sigma(ZH) \propto g_{HZZ}^2$

 σ (ZH) x BR(H \rightarrow ZZ) \propto

=> absolute measurement of width and couplings

The absolutely unique power of pp \rightarrow H+X:

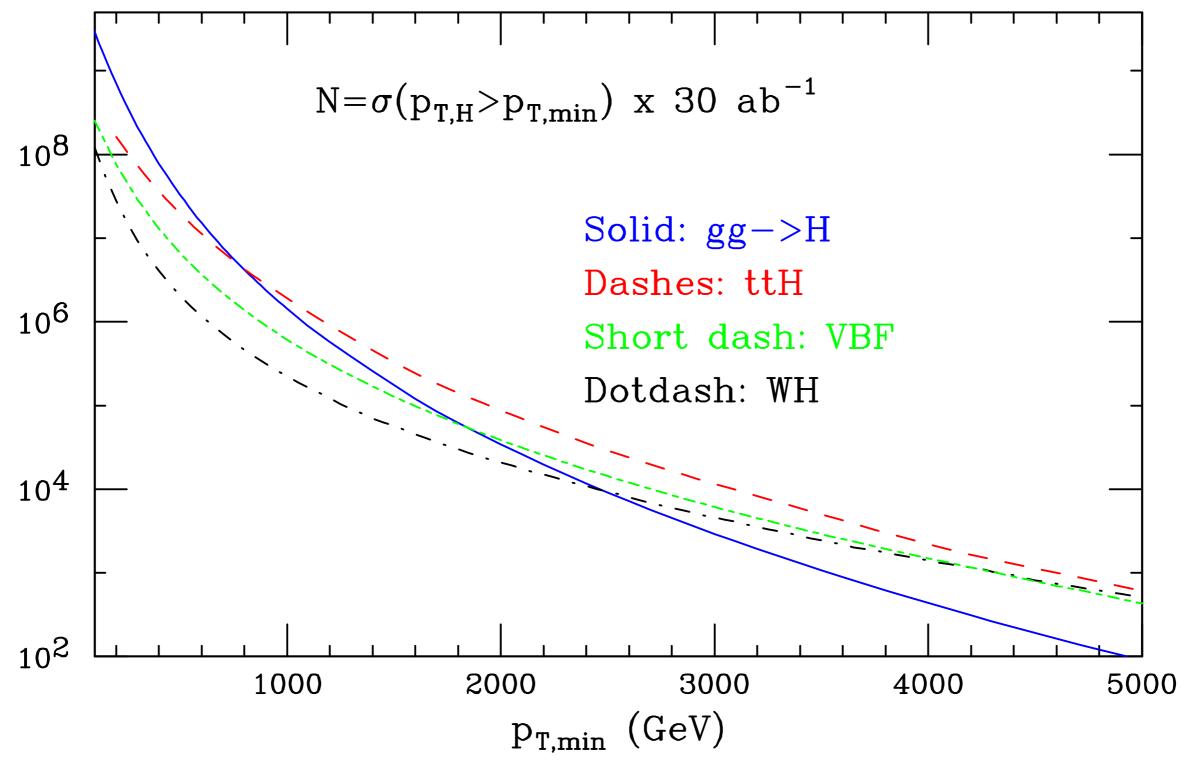
- the extraordinary statistics that, complemented by the per-mille e^+e^- measurement of eg BR(H \rightarrow ZZ*), allows
 - the sub-% measurement of rarer decay modes
 - the $\leq 5\%$ measurement of the Higgs trilinear selfcoupling
- the huge dynamic range (eg pt(H) up to several TeV), which allows to
 probe d>4 EFT operators up to scales of several TeV
 - search for multi-TeV resonances decaying to H, or extensions of the Higgs sector

	gg→H	VBF	WH	ZH	ttH	HH
N ₁₀₀	24 x 10 ⁹	2.1 x 10 ⁹	4.6 x 10 ⁸	3.3 x 10 ⁸	9.6 x 10 ⁸	3.6 x 10 ⁷
N100/N14	180	170	100	110	530	390

 $N_{100} = \sigma_{100 \text{ TeV}} \times 30 \text{ ab}^{-1}$

 $N_{14} = \sigma_{14 \text{ TeV}} \times 3 \text{ ab}^{-1}$

H at large рт



- Hierarchy of production channels changes at large p_T(H):
 - $\sigma(ttH) > \sigma(gg \rightarrow H)$ above 800 GeV
 - $\sigma(VBF) > \sigma(gg \rightarrow H)$ above 1800 GeV

Three kinematic regimes

- Inclusive production, $p_T > 0$:
 - largest overall rates
 - most challenging experimentally:
 - triggers, backgrounds, pile-up \Rightarrow low efficiency, large systematics
 - \blacksquare det simulations challenging, likely unreliable \Rightarrow regime not studied so far

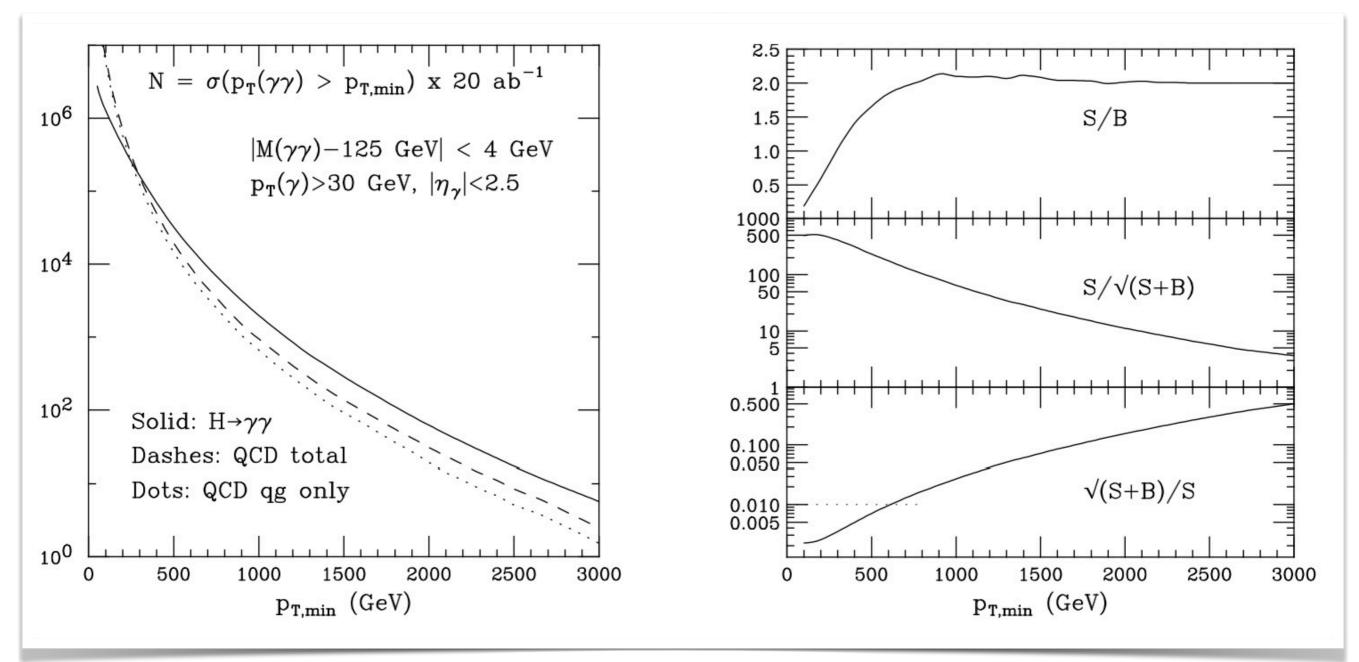
• <u>p⊤ ≳ 100 GeV :</u>

- stat uncertainty ~few × 10⁻³ for $H \rightarrow 4I, \gamma\gamma, ...$
- improved S/B, realistic trigger thresholds, reduced pile-up effects ?
- current det sim and HL-LHC extrapolations more robust
- ➡ focus of FCC CDR Higgs studies so far
- sweet-spot for precision measurements at the sub-% level

• <u>p⊤ ≳ TeV :</u>

- stat uncertainty O(10%) up to 1.5 TeV (3 TeV) for $H \rightarrow 4I$, $\gamma\gamma$ ($H \rightarrow bb$)
- new opportunities for reduction of syst uncertainties (TH and EXP)
- different hierarchy of production processes
- indirect sensitivity to BSM effects at large Q² , complementary to that emerging from precision studies (eg decay BRs) at Q~m_H

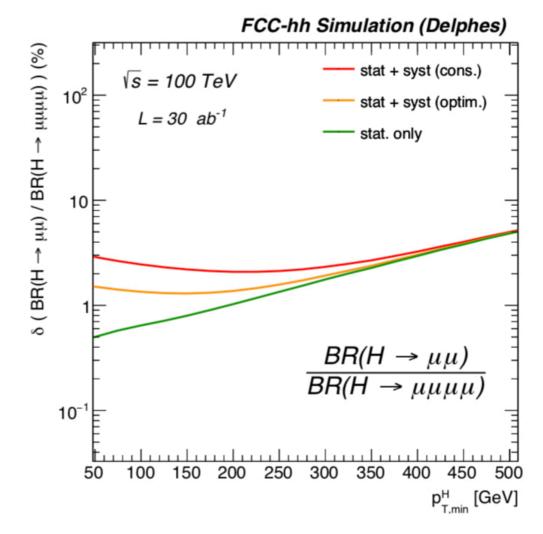
$gg \rightarrow H \rightarrow \gamma \gamma$ at large p_T



lacksquare	At LHC, S/B in the $H \rightarrow \gamma \gamma$ channel is O(few %)

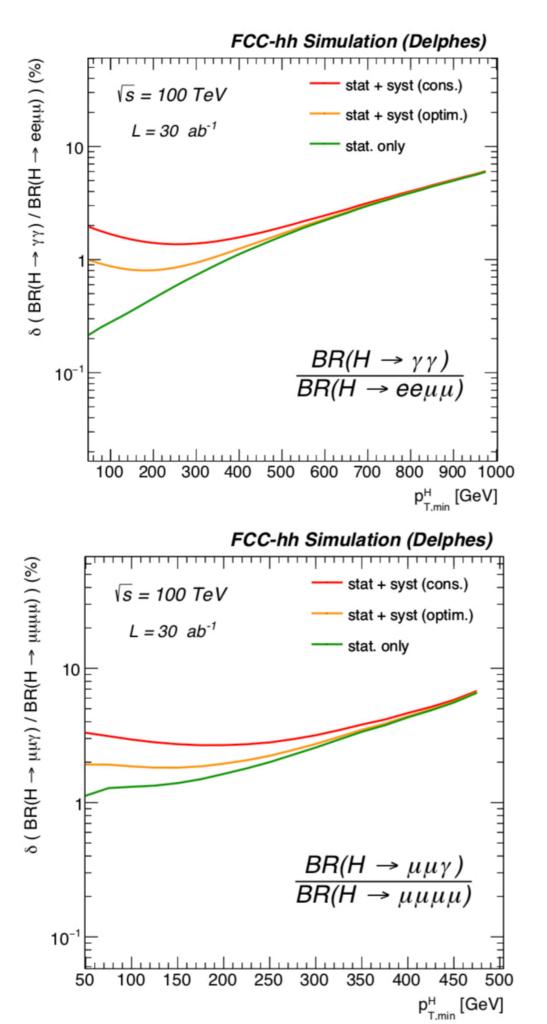
- At FCC, for $p_T(H)>300$ GeV, S/B~I
- Potentially accurate probe of the H pt spectrum up to large pt

δ _{stat}	р _{т,min} (GeV)
0.2%	100
0.5%	400
1%	600
10%	1600



Normalize to BR(4I) from ee => sub-% precision for absolute couplings

Future work: explore in more depth data-based techniques, to <u>validate and</u> <u>then reduce</u> the systematics in these ratio measurements, possibly moving to lower pt's and higher stat

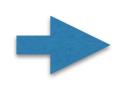


Higgs couplings after FCC-ee / hh

	HL-LHC	FCC-ee	FCC-hh
δΓΗ / ΓΗ (%)	SM	1.3	tbd
δg _{HZZ} / g _{HZZ} (%)	1.5	0.17	tbd
δg _{Hww} / g _{Hww} (%)	1.7	0.43	tbd
δд _{ньь} / д _{ньь} (%)	3.7	0.61	tbd
$\delta g_{Hcc} / g_{Hcc} (\%)$	~70	1.21	tbd
δg _{Hgg} / g _{Hgg} (%)	2.5 (gg->H)	1.01	tbd
δg _{Hττ} / g _{Hττ} (%)	1.9	0.74	tbd
δg _{Hµµ} / g _{Hµµ} (%)	4.3	9.0	0.65 (*)
δg _{Hγγ} / g _{Hγγ} (%)	1.8	3.9	0.4 (*)
δg _{Htt} / g _{Htt} (%)	3.4	~10 (indirect)	0.95 (**)
δg _{HZγ} / g _{HZγ} (%)	9.8	—	0.9 (*)
δдннн / дннн (%)	50	~44 (indirect)	5
BR _{exo} (95%CL)	$BR_{inv} < 2.5\%$	< 1%	BR _{inv} < 0.025%

NB

BR(H→ZY,YY) ~O(10⁻³) ⇒ O(10⁷) evts for Δ_{stat} ~% BR(H→µµ) ~O(10⁻⁴) ⇒ O(10⁸) evts for Δ_{stat} ~%

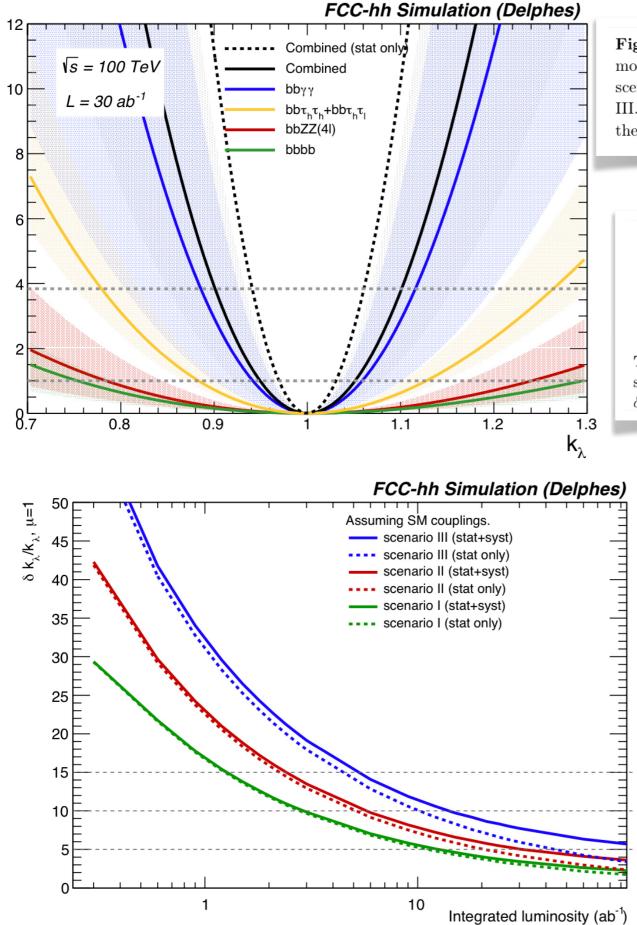


pp collider is essential to beat the % target, since no proposed ee collider can produce more than O(10⁶) H's

* From BR ratios wrt B(H \rightarrow ZZ*) @ FCC-ee

** From pp \rightarrow ttH / pp \rightarrow ttZ, using B(H \rightarrow bb) and ttZ EW coupling @ FCC-ee

The Higgs self-coupling at FCC-hh https://arxiv.org/abs/2004.03505



-2∆ In L

Figure 13. Expected negative log-Likelihood scan as a function of the trilinear self-coupling modifier $\kappa_{\lambda} = \lambda_3 / \lambda_3^{\text{SM}}$ in all channels, and their combination. The solid line corresponds to the scenario II for systematic uncertainties. The band boundaries represent respectively scenario I and III. The dashed line represents the sensitivity obtained including statistical uncertainties only, under the assumptions of scenario I.

Syst scenarios

	@68% CL	scenario I	scenario II	scenario III
S	stat only	2.2	2.8	3.7
δ_{μ}	stat + syst	2.4	3.5	5.1
s	stat only	3.0	4.1	5.6
$\delta_{\kappa_{\lambda}}$	stat + syst	3.4	5.1	7.8

Table 7. Combined expected precision at 68% CL on the di-Higgs production cross- and Higgs self coupling using all channels at the FCC-hh with $\mathcal{L}_{int} = 30 \text{ ab}^{-1}$. The symmetrized value $\delta = (\delta^+ + \delta^-)/2$ is given in %.

- I. Target det performance: LHC Run 2 conditions
- II. Intermediate performance
- III. Conservative: extrapolated HL-LHC performance, with today's algo's (eg no timing, etc)

Expected precision on the Higgs self-coupling as a function of the integrated luminosity.

3-5 ab⁻¹ are sufficient to get below the 10% level

=> within the reach of the first 5yrs of FCC-hh running,

in the "low" luminosity / low pileup phase

=> compatible with the timescale for a similar precision measurement by CLIC @ 3 TeV

Higgs as a BSM probe: precision vs dynamic reach

$$L = L_{SM} + \frac{1}{\Lambda^2} \sum_k \mathcal{O}_k + \cdots$$

$$O = |\langle f|L|i\rangle|^2 = O_{SM} \left[1 + O(\mu^2/\Lambda^2) + \cdots\right]$$

For H decays, or inclusive production, $\mu \sim O(v, m_H)$

$$\delta O \sim \left(\frac{v}{\Lambda}\right)^2 \sim 6\% \left(\frac{\text{TeV}}{\Lambda}\right)^2 \implies \text{precision probes large } \Lambda$$

e.g. $\delta O = 1\% \Rightarrow \Lambda \sim 2.5 \text{ TeV}$

For H production off-shell or with large momentum transfer Q, $\mu \sim O(Q)$

 $\delta O \sim \left(\frac{Q}{\Lambda}\right)^2$ \Rightarrow kinematic reach probes large Λ even if precision is low e.g. $\delta O=15\%$ at Q=1 TeV $\Rightarrow \Lambda\sim2.5$ TeV

Precision and extensive kinematic reach provide unique complementarity and redundancy, crucial to interpret possible SM deviations manifest in either of these observabes

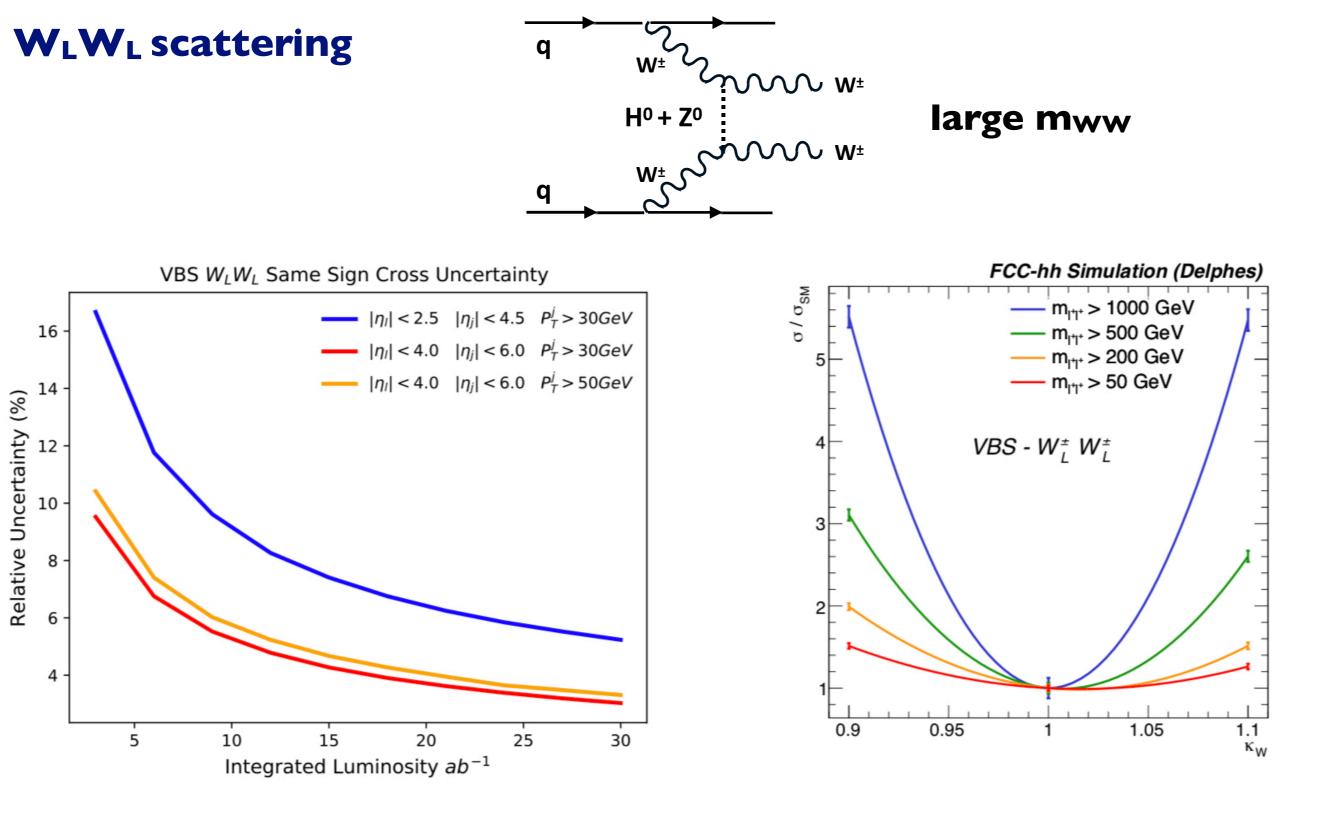


Table 4.5: Constraints on the HWW coupling modifier κ_W at 68% CL, obtained for various cuts on the di-lepton pair invariant mass in the $W_L W_L \rightarrow HH$ process.

$m_{l^+l^+}$ cut	> 50 GeV	$> 200 { m ~GeV}$	$> 500~{ m GeV}$	$> 1000 { m ~GeV}$	$\kappa_{-} - \frac{g_{HWW}}{g_{HWW}}$
$\kappa_W \in$	[0.98,1.05]	[0.99,1.04]	[0.99,1.03]	[0.98,1.02]	$\kappa_W - g_{HWW}^{SM}$

(I) guaranteed deliverables: EW observables

The absolutely unique power of **Circular** e⁺e⁻:

e+e- → Z	e+e- → WW	т(←Z)	b(←Z)	c(←Z)
5 10 ¹²	10 ⁸	3 10 ¹¹	1.5 10 ¹²	10 ¹²

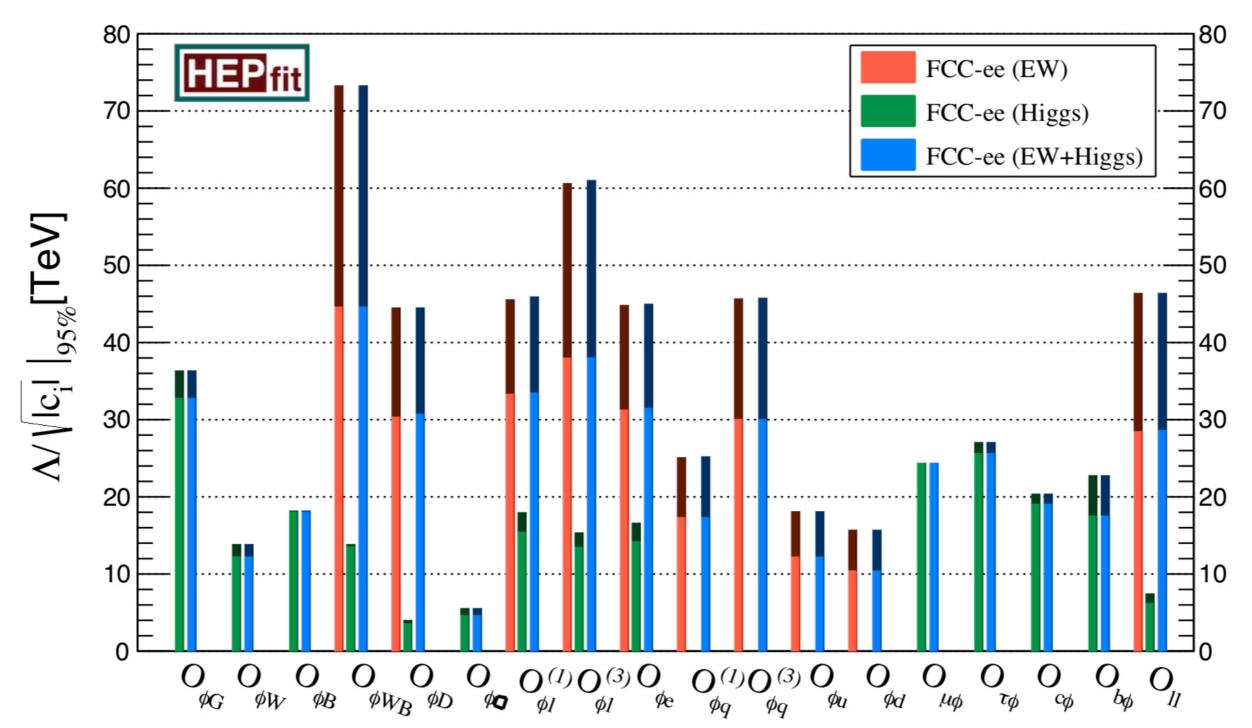
=> O(10⁵) larger statistics than LEP at the Z peak and WW threshold

	Observable	present value ± error	FCC-ee stat.	FCC-ee syst.
	m _Z (keV)	91186700±2200	5	100
	$\Gamma_{\rm Z}$ (keV)	2495200±2300	8	100
	R_l^Z (×10 ³)	20767 ± 25	0.06	0.2-1.0
	$\alpha_{s} \ (m_{Z}) \ (\times 10^{4})$	1196±30	0.1	0.4-1.6
	R_{b} (×10 ⁶)	216290±660	0.3	<60
	$\sigma_{\rm had}^{0}$ (×10 ³) (nb)	41541±37	0.1	4
	N_{ν} (×10 ³)	2991±7	0.005	1
	$\sin^2 \theta_W^{eff}$ (×10 ⁶)	231480±160	3	2-5
eters	$1/\alpha_{QED}(m_Z)$ (×10 ³)	128952±14	4	Small
e	$A_{\rm FB}^{b,0}$ (×10 ⁴)	992±16	0.02	1-3
	$A_{\rm FB}^{{\rm pol}, \tau}$ (×104)	1498±49	0.15	<2
	m _W (MeV)	80350±15	0.6	0.3
	$\Gamma_{\rm W}$ (MeV)	2085±42	1.5	0.3
	α_s (m _W) (×10 ⁴)	1170 ± 420	3	Small
	$N_{\nu}(\times 10^3)$	2920±50	0.8	Small
	m _{top} (MeV)	172740 ± 500	20	Small
	$\Gamma_{\rm top}$ (MeV)	1410±190	40	Small
	$\lambda_{\rm top}/\lambda_{\rm top}^{\rm SM}$	1.2±0.3	0.08	Small
	ttZ couplings	±30%	0.5 - 1.5%	Small

EW parameters @ FCC-ee

(2) Direct discovery reach at high mass: the power of 100 TeV

Global EFT fits to EW and H observables at FCC-ee



Constraints on the coefficients of various EFT op's from a global fit of (i) EW observables, (ii) Higgs couplings and (iii) EW+Higgs combined. Darker shades of each color indicate the results neglecting all SM theory uncertainties.



100 TeV is the appropriate CoM energy to directly search for new physics appearing indirectly through precision EW and H measurements at the future ee collider

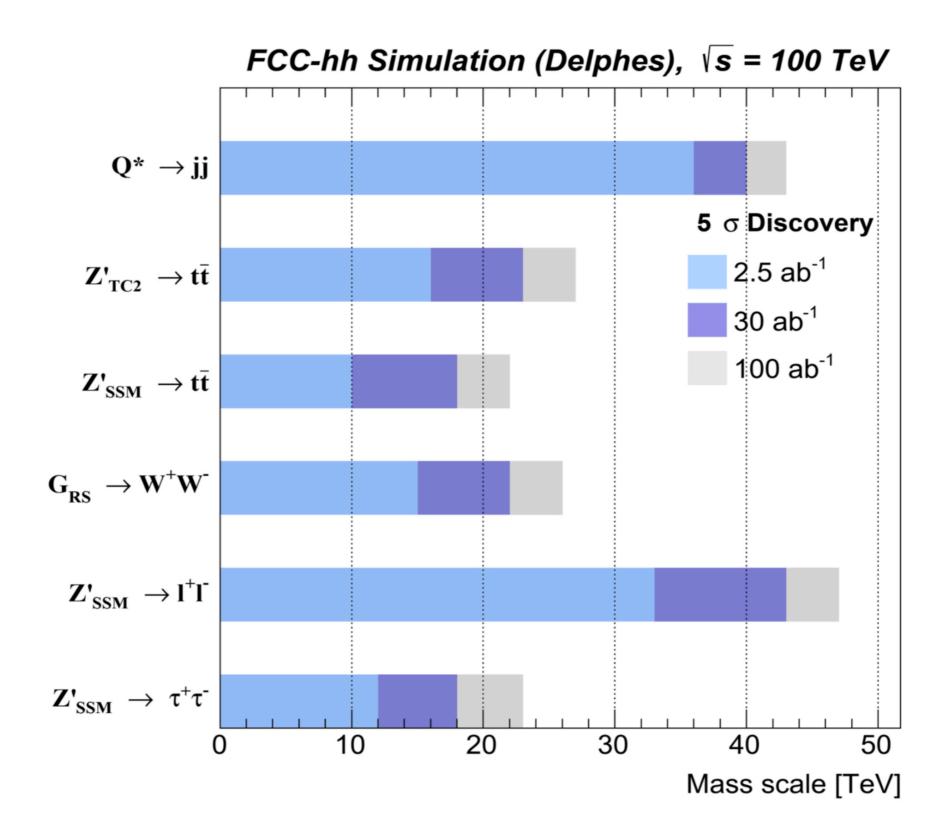
ATLAS Preliminary

ATLAS SUSY Searches* - 95% CL Lower Limits

	ATLAS SUSY Sear Narch 2019	'ches'	· - 95%	% CI	L LO	wer Limits			ATLAS Preliminary $\sqrt{s} = 13 \text{ TeV}$	
	Model	5	Signatur	e j	f£ di [16-	Masslimit			Reference	
	$\hat{q}\hat{q}, \hat{q} \rightarrow \hat{q}\hat{t}_{1}^{0}$	0 a, p mono-jet	2-6 jets 1-3 jets	$E_{\gamma}^{\rm triss}$	36.1 36.1	ð [2x, 8x Digin] 0,9 ð [1x, 8x Digin] 0,43 0,71	1.55	m(∂ ⁰)<100 GeV m(∂)-m(∂ ⁰)=5 CeV	1712.00302 1711.00301	
Inclusive Searches	$33. S \rightarrow q \overline{q} \overline{r}_1^O$	0 e, p	2-6 jets	$E_7^{\rm mins}$	\$6.1	R Forbidden	0.95-1.6	m(ℓ ⁰)<200 Ce/V m(ℓ ⁰)=200 Ce/V	1712.02352 1712.02352	
e See	<i>13. 2→490</i> 00 ² 1	3 г., р сс, рр	4 jats 2 jets	E_7^{min}	36.1 36.1	2 2	1.85	m(ℓ) ⁴ <800GeV m(ℓ)-m(ℓ) ⁴)=50GeV	1706.00701 1805.11381	
Arishe	£ġ, g→qqWZŽ ⁰	0 e, p 3 e, p	7-11 jets 4 jeta	E_7^{min}	36.1 36.1	8 2 0.96	1.8	m(8) m(3)-m(7) (=200 GeV	1706.02794 1706.08731	
pul	λ3. <i>δ→u</i> ? ⁰ 1	0-1 e, µ 3 e, µ	3 <i>b</i> 4 jets	$E_7^{\rm mix}$	79.8 36.1	2	2.25	m(i ⁰)<200 GeV m(i)-200 GeV	ATLAS-CONF-2018-041 1706.03731	
	$\hat{h}_1 \hat{h}_1, \hat{h}_1 {\rightarrow} \lambda \hat{t}_1^2 / d\hat{t}_1^+$		Multiple Multiple Multiple		36.1 36.1 36.1	δ1 Forbidden 0.9 δ1 Forbidden 0.58-0.82 δ1 Forbidden 0.7	m(\vec{t}_1^2)=300 m(\vec{t}_2^2)=300 GeV/	$\Pi(\hat{t}_{1}^{0})=300 \text{ GeV}, BR(\partial \hat{t}_{1}^{0})=1$ $GeV, BR(\partial \hat{t}_{1}^{0})=BR(\partial \hat{t}_{1}^{0})=0.5$ $\Pi(\hat{t}_{1}^{0})=300 \text{ GeV}, BR(\partial \hat{t}_{1}^{0})=1$	1708.09266, 1711.03301 1708.08286 1706.08731	
arks	$\hat{b}_1 \hat{b}_1, \hat{b}_1 \rightarrow b \hat{\ell}_2^0 \rightarrow b h \hat{\ell}_1^0$	$0 \ e_s \mu$	6.0	E_7^{\min}	130	λ ₁ Forbidden δ ₁ 0.23-0.46		$(1 = 130 \text{ GeV}, m(\tilde{t}^0) = 100 \text{ GeV})$ $(\tilde{t}^0_1) = 130 \text{ GeV}, m(\tilde{t}^0_1) = 0 \text{ GeV})$	SUSY-2018-31 SUSY-2018-01	
n, squ	$\tilde{i}_1 \tilde{i}_1, \tilde{i}_1 \rightarrow Wh \tilde{\xi}_1^0 \text{ at } \tilde{i}_1^0$ $\tilde{i}_1 \tilde{i}_1, Well-Tempered LSP$	$0\text{-}2e,\mu$	0-2 jets/1-2 Multiple	$b E_7^{miss}$	86.1 86.1	λ ₁ 1./ λ ₁ 0.49-0.84		$m(\hat{\xi}_1^n)=1$ GeV $\chi_m(\hat{\xi}_1^n)=m(\hat{\xi}_1^n)=5$ GeV, $\hat{z}_1 \approx \hat{z}_2$	1506.00010, 1706.04100, 1711.11520 1709.04150, 1711.11520	
^d ge	$\vec{s}_1 \vec{s}_1, \vec{s}_1 \rightarrow \vec{\tau}_1 h v, \vec{\tau}_1 \rightarrow \pi \vec{G}$,τ 2 jets'i λ		36.1	II.	1.16	m(71)=600 GeV	1809.10178	
80	$\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow c \tilde{t}_1^0 / \tilde{c} \tilde{v}, \tilde{c} \rightarrow c \tilde{t}_1^0$	$0 < \mu$	2.0	$E_7^{\rm mino}$	36.1	2 0.85 Ž ₁ 0.46		m(?))=0GeV m(2)=60GeV m(2,2)-m(2))=5GeV	1805.01649 1805.01649	
	$\tilde{s}_2 \tilde{s}_2, \tilde{s}_2 \rightarrow \tilde{s}_1 + h$	0 г.,р 1-2 г.,µ	4 b	Enites Enites	36.1 36.1	7, 0.43 7, 0.32-0.68	$m(\tilde{t}_{i}^{2})=0$	$m(\tilde{t}_1, \tilde{t}) \cdot m(\tilde{t}_1^2) = 5 \text{ GeV}$ $0 \text{ GeV}, m(\tilde{t}_1) \cdot m(\tilde{t}_1^2) = 180 \text{ GeV}$	1711.00001 1700.00906	
	$\tilde{x}_1^* \tilde{x}_2^0$ via $\overline{w}Z$	2-3 e.µ se.µµ	21	Etiis Etiis	56.1 36.1	并示形 ³ 式示形 ³ 0.17		$m(\tilde{\xi}_1^3)=0$ $m(\tilde{\xi}_1^2)=10$ CeV	1403.5294, 1505.02293 1712.001 (9	
	$\tilde{E}_{1}^{+}\tilde{E}_{1}^{+}$ via W W	Z e, µ	21	$E_7^{\rm min}$	139	3 [±] 042		$m(\hat{x}_1) - m(x_1) = 10$ GeV $m(\hat{x}_1^2) = 0$	ATLAS-CONF-2019-008	
	λ ¹ ₁ λ ¹ ₂ via We	0-1 <i>4.</i> ,µ	2 b	Erra R	86.1	$\tilde{k}_{1}^{a}/\tilde{k}_{2}^{b}$ 0.68		$m(\tilde{\ell}_1^0)=0$	1812.09432	
EW Meet	$\mathcal{X}_{1}^{*}\mathcal{X}_{1}^{*}$ via \tilde{b}_{1}/\tilde{r} $\mathcal{X}_{1}^{*}\mathcal{X}_{1}^{*}/\mathcal{X}_{2}^{*}, \mathcal{X}_{1}^{*} \rightarrow \tilde{r}_{1}v(\tau\tilde{r}), \mathcal{X}_{2}^{0} \rightarrow \tilde{r}_{1}v(\tau\tilde{r})$	2 <, µ 2 +		$E_{\gamma}^{\rm miss}$	139 56.1	$\frac{x_1^4}{\hat{x}_1^4 K_2^3}$ $\frac{x_1^2 K_2^3}{\hat{x}_1^2 K_2^3}$ 0.78	$m(\tilde{r}_{1}^{i})$ -	$m(\tilde{t}, \tilde{t}) = 0.5(m(\tilde{t}_1^+) + m(\tilde{t}_1^+))$ = 0, $m(\tilde{t}, \tilde{t}) = 0.5(m(\tilde{t}_1^+) + m(\tilde{t}_1^+))$ = 24, $m(\tilde{t}, \tilde{t}) = 0.5(m(\tilde{t}_1^+) + m(\tilde{t}_1^+))$	ATLAS-CONF-2018-008 1708.07875 1708.07875	
	$\tilde{t}_{1,\mathbf{R}}\tilde{t}_{1,\mathbf{R}}, \tilde{t} \rightarrow t \hat{x}_{\perp}^{2}$	2 e, p 2 e, p	0 jets ≥ 1	$E_7^{\rm miss}$ $E_7^{\rm miss}$	139 36.1	2 0.7 2 0.18	ng ping ping a	$m(\xi_1^2)=0$ $m(\xi_1^2)=0$ $m(\xi_1^2)=5$ GeV	ATLAS-CONF-2019-008 1712.08119	
	$\hat{H}\hat{H}, \hat{H} \rightarrow hG/ZG$	0 е. µ 4 е. µ	≥3 ē 0 jets	Eyin Eyin	36.1 36.1	R 0.13-0.23 0.29-0.88		$BR(\vec{k}_1^2 \rightarrow \Lambda \vec{G}) \approx 1$ $BR(\vec{k}_1^2 \rightarrow Z \vec{G}) \approx 1$	1604,0000	
-lived	$\operatorname{Direct} \mathcal{R}_1^+ \mathcal{R}_1^-$ proof., long-lived \mathcal{R}_1^+	Disapp. 14	k 1 jet	$E_7^{\rm miles}$	36.1	$\hat{X}_{1}^{\pm} = 0.46$ $\hat{X}_{1}^{\pm} = 0.15$		Pure Wins Pure Higgsins	1712.02118 ATL PHYS PUB 2017-019	
Long-	Stable ĝ R-hadron Metestable ĝ R-hadron, ĝ⇔ge∛1		Muttiple Muttiple		36.1 36.1	ž ∦ [r(ĝ) =10 ns. 0.2 ns]	2.0	$m(\tilde{\xi}_{a}^{2})$ =100 GeV	1902.01636.1605.04095 1710.04901.1803.04095	
	$\Box FV \; pp \!\rightarrow\! \delta_{\tau} + X_{\tau} \vartheta_{\tau} \!\rightarrow\! \epsilon \mu / \sigma \tau / \mu \tau$	querjut			3.2	7.	1.9	λ_{311}^{*} =0.11, $\lambda_{122/129,229}$ =0.07	1607.08079	
	$\tilde{\chi}_{1}^{\pm} \tilde{\chi}_{1}^{\pm} / \tilde{\chi}_{2}^{2} \rightarrow WW/ZUUUvv$ $\tilde{\chi}_{2}^{0}, \tilde{\chi}_{-}^{0} aqg \tilde{\chi}_{1}^{0}, \tilde{\chi}_{1}^{0} \rightarrow qaq$	4 c.µ	0 jets 4-6 large- <i>R</i> ji	E ^{rcito} ets	36.1 36.1	$\hat{x}_1^+/\hat{x}_2^- = [A_{13} \neq 0, A_{12} \neq 0]$ 0.02 $\hat{x}_1^- [n(\hat{x}_1^2) + 200 \text{ GeV}]$ (100 GeV)	1.33	$m(\tilde{k}_{n}^{2})=100 \text{ GeV}$ Large \mathcal{X}_{n}^{2}	1004.00502 1804.00568	
ΡР			Multiple		36.1	$\frac{1}{k} = [k_{112}^{\mu}/2a/4, 2a/5]$ 1.	5 2.0	$m(\tilde{\mathcal{R}}_1^0)$ -200 GeV, birs-like	ATLAS-CONF-2016-019	
ΩC	$\hat{x}, \hat{t} \rightarrow \hat{x}_{1}^{0}, \hat{x}_{1}^{0} \rightarrow tbx$ $\hat{t}_{1}\hat{x}_{1}, \hat{t}_{1} \rightarrow bx$		Multiple 2 jets + 2 /	in .	86.1 86.7	$\bar{x} = [\frac{1}{20}] = 20 - 5, 10 - 2[0.56 1. 0.56 1. 0.42 0.61 1. 0.4$	5	m(l ⁰)⇒200 GeV, bino-lika	ATLAS-CONF-2015-009 1710.07171	
	$\tilde{I}_1 \tilde{I}_1, \tilde{I}_1 \rightarrow g \ell$	2 e.μ 1 μ	2 # DV		36.1 135	$\frac{h}{h} = \frac{1}{20} =$	0.4-1.45 1.6	BD(7)-Hgr(=100%, eccel	1710.05544 ATLAS-CONF-2019-006	
	•									
'Only	a selection of the available mas	s limits on	new state	95 OF	1	0-1	1 . N	lass scale [TeV]	Contraction of the second second	
simi	nomena is shown. Many of the li silfied models, c.f. refs. for the au	ssumptions	ased on s made.							
							@14 TeV		0.4-1.4	15
~							The second se		1.0	1.6
	Good rule of thumb to estimate FCC discovery reach at high mass: scale up by ~6x the LHC potential									
E	olicitly verif	ied i	in m	an	v e	kamples, which helped		THE REAL PROPERTY.	7	
	etting detecto				-	• • •	-		@100 Te ^v	V
										V

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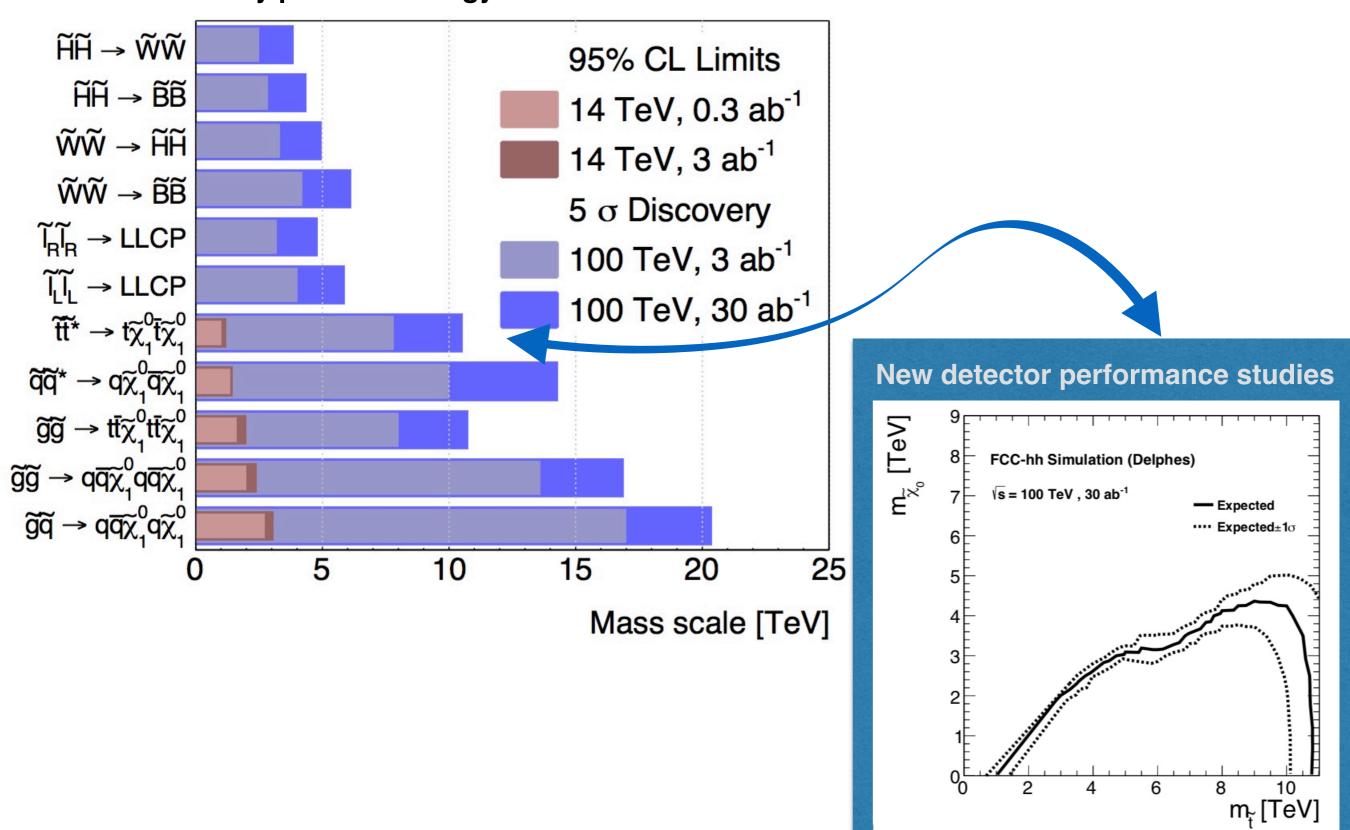
s-channel resonances



FCC-hh reach ~ 6 x HL-LHC reach

SUSY reach at 100 TeV

Early phenomenology studies



(3) The potential for yes/no answers to important questions

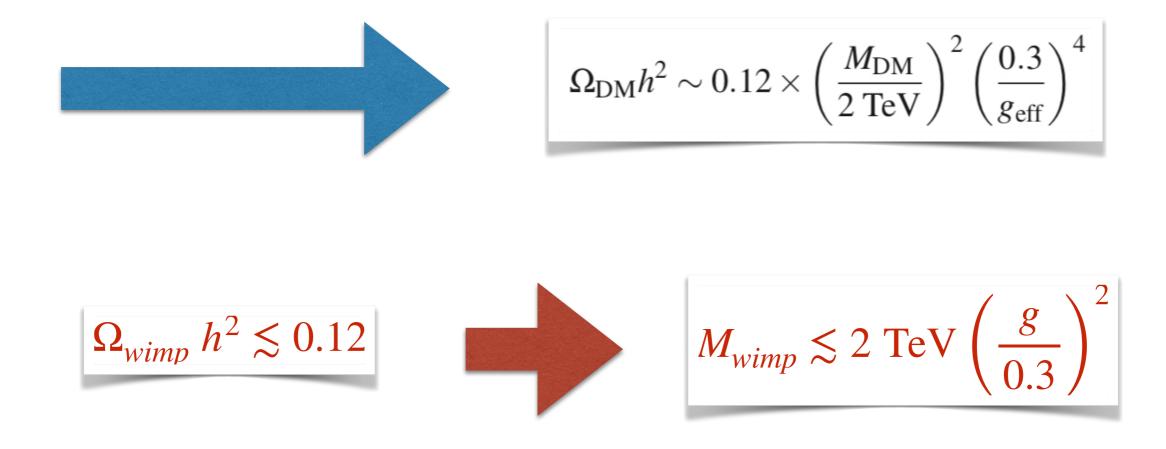
WIMP DM theoretical constraints

For particles held in equilibrium by pair creation and annihilation processes, ($\chi \ \chi \leftrightarrow SM$)

$$\Omega_{\rm DM} h^2 \sim rac{10^9 {
m GeV}^{-1}}{M_{
m pl}} rac{1}{\langle \sigma v
angle}$$

For a particle annihilating through processes which do not involve any larger mass scales:

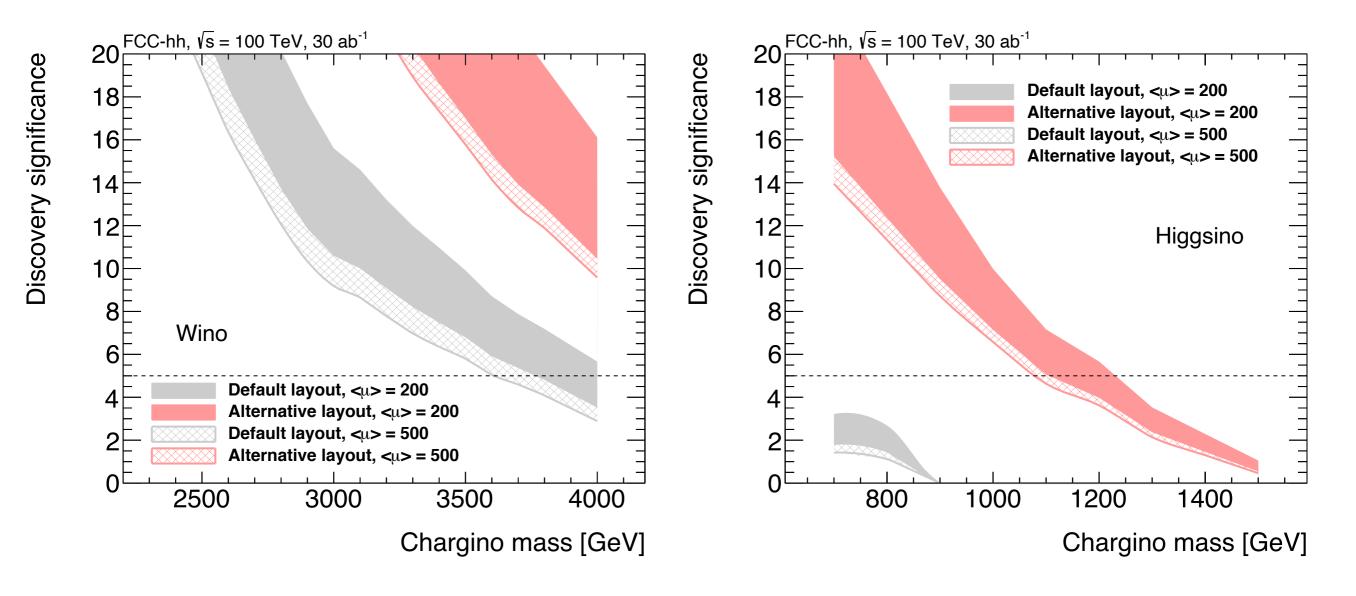
 $\langle \sigma v \rangle \sim g_{\rm eff}^4 / M_{\rm DM}^2$



K. Terashi, R. Sawada, M. Saito, and S. Asai, *Search for WIMPs with disappearing track signatures at the FCC-hh*, (Oct, 2018) . https://cds.cern.ch/record/2642474.

DM WIMP searches in the most elusive, compressed scenarios:

Disappearing charged track analyses (at ~full pileup)



=> coverage beyond the upper limit of the thermal WIMP mass range for both higgsinos and winos !!



Not covered

- Countless studies of discovery potential for multiple BSM scenarios, from SUSY to heavy neutrinos, from very low masses to very high masses, LLPs, DM, etcetcetc, at FCC-ee, FCC-hh and FCC-eh
- Sensitivity studies to SM deviations in the properties of top quarks, flavour physics in Z decays: huge event rates offer unique opportunities, that cannot be matched elsewhere

 Operations with heavy ions: new domains open up at 100 TeV in the study of high-T/high-density QCD. Broaden the targets, the deliverables, extend the base of potential users, and increase the support beyond the energy frontier community

Final remarks

- The study of the SM will not be complete until we clarify the nature of the Higgs mechanism and exhaust the exploration of phenomena at the TeV scale: many aspects are still obscure, many questions are still open.
- The exptl program possible at a future collider facility, combining a versatile high-luminosity e⁺e⁻ circular collider, with a follow-up pp collider in the 100 TeV range, offers unmatchable breadth and diversity: concrete, compelling and indispensable Higgs & SM measurements enrich a unique direct & indirect discovery potential
- The unique feature of a circular ee + pp collider is the possibility to match the indirect high-mass-scale sensitivity of precision measurements to the direct search potential at large mass
- The next 5-6 years, before the next review of the European Strategy for Particle Physics, will be critical to reach the scientific consensus and political support required to move forward

Additional material: recent reports on Future Circular Colliders

• FCC CDR:

- Vol.1: Physics Opportunities (CERN-ACC-2018-0056) <u>http://cern.ch/go/Nqx7</u>
- Vol.2: The Lepton Machine (CERN-ACC-2018-0057) http://cern.ch/go/7DH9
- Vol.3: The Hadron Machine (CERN-ACC-2018-0058), <u>http://cern.ch/go/Xrg6</u>
- Vol.4: High-Energy LHC (CERN-ACC-2018-0059) <u>http://cern.ch/go/S9Gq</u>
- "Physics at 100 TeV", CERN Yellow Report: https://arxiv.org/abs/1710.06353
- CEPC CDR: Physics and Detectors