

# QUARK MASS FUNCTION FROM A OGE-TYPE INTERACTION IN MINKOWSKI SPACE

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## COLLABORATORS AND WORKS

**Teresa Peña** (LIP and IST)

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### PAPERS:

**Phys. Rev. D 98, 114033 (2018)**

Phys. Lett. B764, 38 (2017)

Phys. Rev. D 96, 074007 (2017)

Phys. Rev. D 90, 096008 (2014)

Phys. Rev. D 89, 016005 (2014)

# MESON PHENOMENOLOGY — MOTIVATION

- upcoming experiments  
**COMPASS++/AMBER (CERN)**  
GlueX (JLab), GlueX (JLab)  
PANDA (FAIR-GSI):  
**meson** properties  
origin of **hadronic mass**
- need better theoretical understanding of  **$q\bar{q}$  mesons**
- study **spectrum** and **structure**

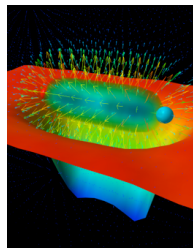


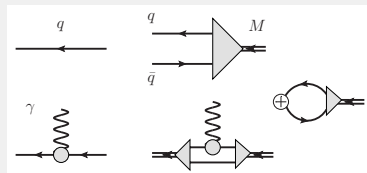
Image credit: D. LEINWEBER

## OBJECTIVES

- unified description for **all**  $q\bar{q}$  mesons from pion (0.14 GeV) to  $b\bar{b}$  ( $> 10$  GeV)
  - ✓ already achieved for **heavy** and **heavy-light** mesons
  - LEITÃO, STADLER, PEÑA, EB, PLB (2017), PRD (2017)
  - ✗ to do: **light** mesons
- spectrum and decay properties  $\Rightarrow$  information about the Lorentz structure of the **confining** interaction
- describe mass-generation mechanism of **dynamical chiral-symmetry breaking** (Talk by C. ROBERTS)

### Calculation of

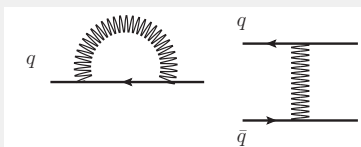
- **dynamical quark mass function**
- vertex functions
- quark-photon vertex
- meson form factors
- meson decay properties



# COVARIANT SPECTATOR THEORY (CST)

## Guiding principles

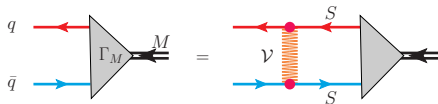
- manifest Lorentz **covariance**
- solved in **Minkowski** space
- **confinement**: covariant interaction kernel that reduces in nonrelativistic limit to 'linear-confining+Coulomb' potential
- **dynamical chiral-symmetry breaking**: axial-vector Ward Takahashi identity  
massless pion in chiral limit  
quark masses dynamically generated through self-interaction with  $q\bar{q}$  interaction



# GROSS (CST) EQUATION

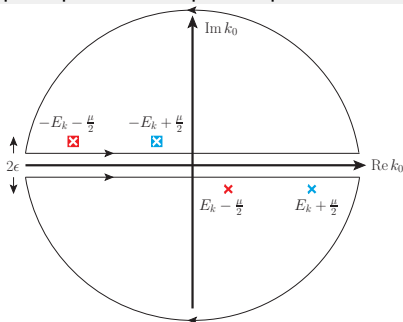
GROSS, PR 186 (1969)

## Bethe-Salpeter equation (BSE)



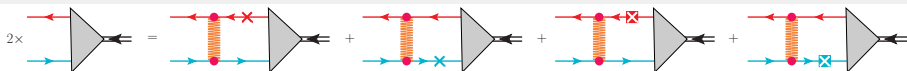
- assume:  $\exists$  **real** quark-mass poles
- keep only **quark pole** contributions
- **3D covariant Minkowski** integrations
- beyond RL: correct **1-body** Dirac limit

## quark poles in complex $k_0$ plane

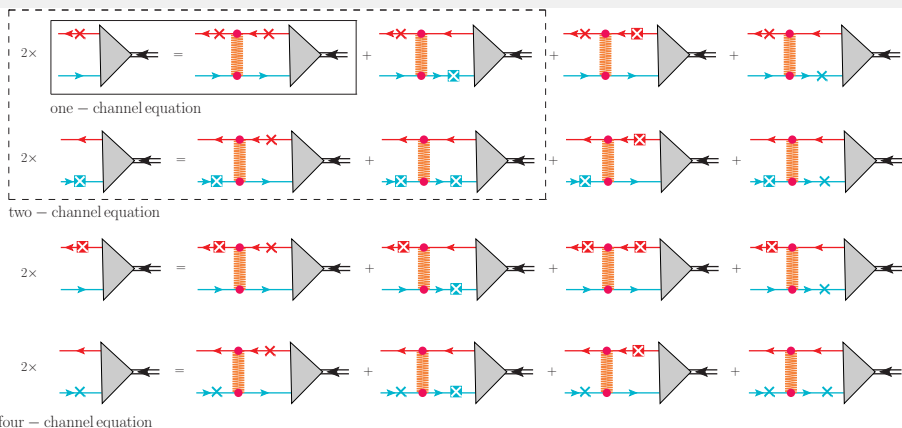


## charge-conjugation symmetric Gross equation

SAVKLI, GROSS PRC (2001), EB, GROSS, PEÑA, STADLER PRD (2014)



## 4-CHANNEL EQUATIONS



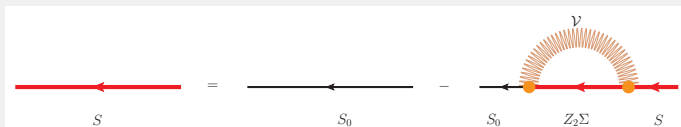
- 1-channel equation: ✓ unequal quark masses, ✗  $\mathcal{C}$ -symmetry, ✗ light mesons
- 2-channel equation: ✓  $\mathcal{C}$ -symmetry, ✗ light mesons
- 4-channel equation: ✓ all mesons, also light mesons (pion)

All have smooth 1-body (Dirac) and nonrelativistic (Schrödinger) limits

# CST DYSON EQUATION

EB, GROSS, PEÑA, STADLER, LEITÃO, PRD (2018)

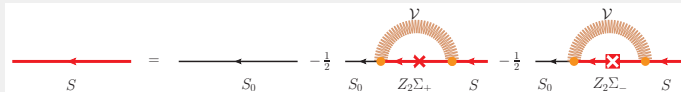
Dyson equation for **dressed** quark propagator (self-energy  $\Sigma$ )



$$S(p) = S_0(p) \sum_n [-Z_2 \Sigma(p) S_0(p)]^n = \frac{1}{m_0 + Z_2 \Sigma(p) - \not{p} - i\epsilon} \equiv \frac{Z(p^2)[M(p^2) + \not{p}]}{M^2(p^2) - p^2 - i\epsilon}$$

dressed quark **mass function**  $M(p^2)$

CST:  $M(m^2) = m$  pole equation for **constituent quark mass**  $m$



$$Z_2 \Sigma(p) = \frac{Z_2^2}{2} \sum_{\sigma} \int_{\mathbf{k}} (m + \hat{k}_{\sigma}) \mathcal{V}(p, \hat{k}_{\sigma})$$

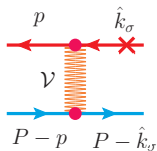
with renormalization  $Z_2(m) \equiv \frac{Z(m^2)}{1 - 2mM'(m^2)}$



## COVARIANT INTERACTION KERNEL

$$\mathcal{V}(p, \hat{k}_\sigma) = \underbrace{\mathcal{V}_\ell(p, \hat{k}_\sigma)}_{\text{'linear confining'}} + \underbrace{\mathcal{V}_g(p, \hat{k}_\sigma)}_{\text{OGE-type}} + \underbrace{\mathcal{V}_c(p, \hat{k}_\sigma)}_{\text{covariant 'constant'}}$$

reduces to  $V^{\text{nr}}(r) = \sigma r - \frac{4}{3} \frac{\alpha_s}{r} + C$  in nonrelativistic limit



$$\mathcal{V}_\ell(p, \hat{k}_\sigma) = \frac{1}{4} \underbrace{\sum_a \lambda_a \otimes \lambda_a}_{\frac{4}{3} \text{ (color singlets)}} \left[ (1 - \lambda)(\mathbf{1} \otimes \mathbf{1} + \gamma^5 \otimes \gamma^5) - \lambda \gamma^\mu \otimes \gamma_\mu \right] V_\ell(p, \hat{k}_\sigma)$$

covariant off-shell generalization of 'linear confinement'

- $q^2 \rightarrow -q_\sigma^2 = -(p - \hat{k}_\sigma)^2 \Rightarrow \int_{\mathbf{k}} V_\ell(p, \hat{k}_\sigma) \psi(\hat{k}_\sigma) = -\sigma \int_{\mathbf{k}} \frac{\psi(\hat{k}_\sigma) - \psi(\hat{k}_{R\sigma})}{q_\sigma^4}$

*confinement: meson vertex function vanishes if both quarks are on-shell!* ✓

SAVKLI, GROSS PRC (2001)

- $\mathcal{V}_\ell$  consistent with  $\mathbf{D}\chi\mathbf{SB}$  (AVWTI satisfied) ✓

EB, PEÑA, RIBEIRO, STADLER, GROSS PRD (2014)

- $\lambda = 0 \Rightarrow \mathcal{V}_\ell$  does not contribute to self-energy, consistent with meson spectrum ✓

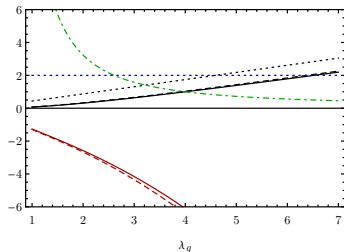
LEITÃO, STADLER, PEÑA, EB, PLB (2017), PRD (2017)

## OGE-TYPE KERNEL

EB, GROSS, PEÑA, STADLER, LEITÃO, PRD (2018)

$$\mathcal{V}_g(p, \hat{k}_\sigma) = 4\pi\alpha_s g(y) \gamma_\mu \otimes \gamma_\nu \frac{1}{M_g^2 + |q_\sigma^2|} \left[ g^{\mu\nu} - (1 - \xi) \frac{q_\sigma^\mu q_\sigma^\nu}{q^2} \right] \frac{1}{4} \sum_a \lambda_a \otimes \lambda_a$$

- finite **gluon mass**  $M_g = 0.6$  GeV and **prescription**  $q^2 \rightarrow -|q^2|$   
 $\Rightarrow$  **removes** singularity in gluon propagator ✓
- form factor  $g(y) = \frac{\lambda_g^{4n}}{\lambda_g^{4n} + (y^2 - 1)^n}$  (with  $y^2 = \frac{E_k^2}{m^2}$ ,  $\lambda_g = \frac{\Lambda_g}{m}$ )  
**regularizes**  $\int_{\mathbf{k}}$  (also at  $p^2 = 0$ ) ✓
- self-energy  $Z_2 \Sigma_g(\not{p}) = \frac{1}{4}(3 + \xi)Z_2 A_g(p^2) + \not{p} \frac{1}{2}(3 - \xi)Z_2 B_g(p^2) - \not{p}(1 - \xi)Z_2 R_g(p^2)$   
 where  $A_g, B_g, R_g \propto Z_2 \alpha_s \int_{\mathbf{k}} \dots \Rightarrow$  **renormalized** coupling  $\alpha_s^r(m) \equiv Z_2^2(m) \alpha_s$
- mass pole equation for  $m$  is gauge **independent** ✓



$$m_0 = 0: \quad \alpha_s^r T_g = 1 \quad [Z_2 \alpha_s T_g \equiv \frac{A_g(m^2)}{m} + B_g(m^2)]$$

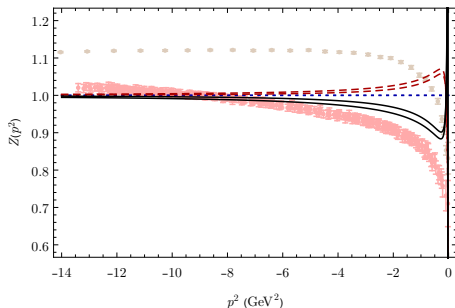
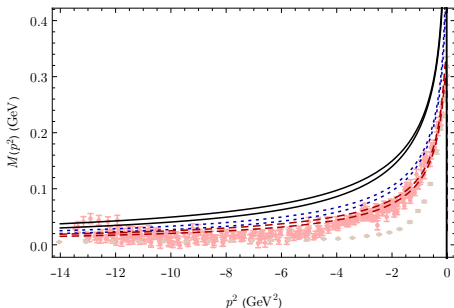
$T_g$  with/without  $|q^2|$ -prescription

(solid:  $n = 4$ , dashed:  $n = 3$ ,  
 dotted:  $M_g = 0$ ), dotdashed:  $\alpha_s^r$

$\Rightarrow$  only **with**  $|q^2|$ -prescription:

$$T_g \simeq 2, \alpha_s^r \simeq 0.5 \Rightarrow \lambda_g \simeq 7$$

## $M(p^2)$ AND $Z(p^2)$ FOR SPACELIKE $p^2$ AND LATTICE DATA



lattice data: **Bowman et al PRD (2005)**, **Oliveira et al, PRD (2019)**

**Landau gauge** ( $\xi = 0$ )

**Feynman-'t Hooft gauge** ( $\xi = 1$ )

**Yennie gauge** ( $\xi = 3$ )

$m = 0.3 \text{ GeV}$ ,  $m_0 = 0$ :

$\lambda_g = 5$  ( $\alpha_s^r = 0.722$ )

$\lambda_g = 3$  ( $\alpha_s^r = 1.577$ )

$\Rightarrow$  **reasonable** behaviour similar to LQCD (no fitting)

for  $\alpha_s^r \simeq 0.5 \Rightarrow$  need constant kernel

# COVARIANT OFF-SHELL CONSTANT KERNEL IN GENERAL GAUGE

Covariant off-shell generalization of **constant** potential:

$$\mathcal{V}_c(p, \hat{k}_\sigma) = \frac{CE_k}{2m} (2\pi)^3 \delta^3\left(\mathbf{k} - \frac{m}{\sqrt{p^2}} \mathbf{p}\right) h(p^2) h(m^2) \gamma_\mu \otimes \gamma_\nu \left[ g^{\mu\nu} - (1 - \xi) \frac{q_\sigma^\mu q_\sigma^\nu}{q_\sigma^2} \right] \frac{1}{4} \sum_a \lambda_a \otimes \lambda_a$$

- nonrelativistic limit:  $\mathcal{V}_c \rightarrow \mathcal{V}_c^{\text{nr}} \propto C \delta^3(\mathbf{k} - \mathbf{p})$  ( $\equiv$  constant potential)
- **correction** to OGE part: strong quark form factor  $h(p^2) \equiv \frac{A_g(p^2)}{A_g(m^2)}$

self-energy

- satisfy same mass pole eq. as  $\mathcal{V}_g$ :

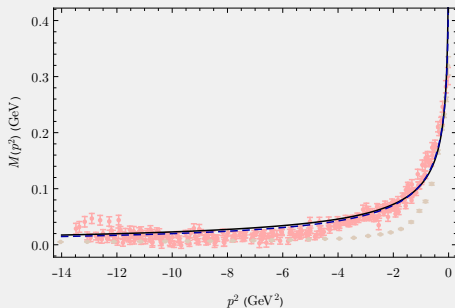
$$C \rightarrow \frac{3m \alpha_g T_g}{3 + \xi}$$

- contributes only to scalar part of self-energy:  $Z_2 \Sigma_c(p) = m \frac{A_g(p^2)}{A_g(m^2)}$

$$Z(p^2) = 1, M(p^2) = m \frac{A_g(p^2)}{A_g(m^2)}$$

(gauge independent)

- curves:  $\lambda_g = 5, \lambda_g = 3$



## CONSTANT AND OGE SELF-ENERGY

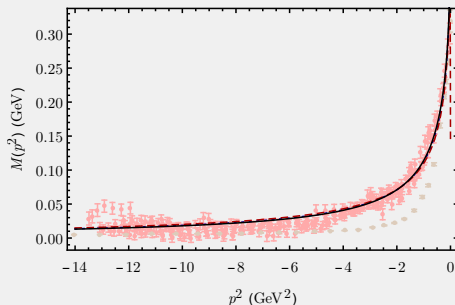
**IDEA:** each contribution to self-energy satisfies gap equation  $\Rightarrow$  also linear combination

$$Z_2 \Sigma(\not{p}) = \underbrace{\eta Z_2 \Sigma_g(\not{p})}_{\text{OGE}} + \underbrace{(1 - \eta) m \frac{A_g(p^2)}{A_g(m^2)}}_{\text{'constant'}}$$

fix parameters in chiral limit where  $m_0 = 0$

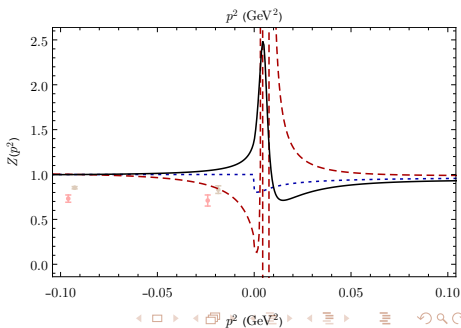
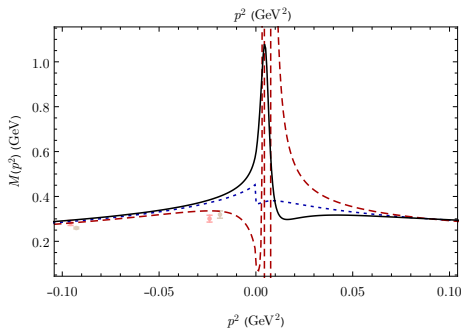
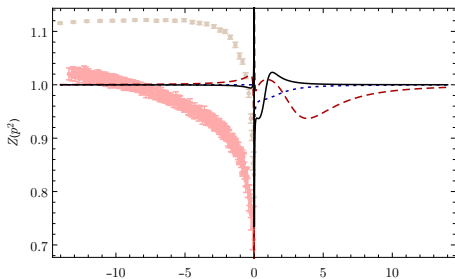
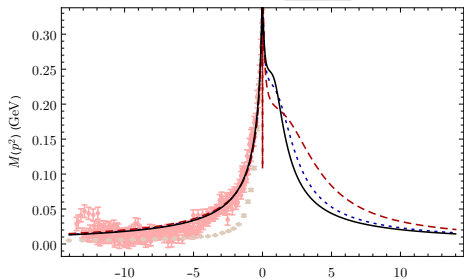
- held **fixed**:  $m \rightarrow m_\chi = 0.3 \text{ GeV}$ ,  
 $M_g = 0.6 \text{ GeV}$ ,  $n = 4$ ,  
 $\alpha_s^r(m_\chi) \rightarrow \eta \alpha_s^r(m_\chi) \equiv \alpha_s^p(m_\chi) = 0.5$   
 $\Rightarrow$  'constant' **decreases** strength of OGE
- roughly **adjust**  $\lambda_g$  to agree with LQCD

$\xi$	0	1	3
$\lambda_g$	3	2	1.5
$\eta(\lambda_g)$	0.317	0.155	0.087
$(1 - \eta) \frac{Z_2^2 C}{m}$	0.911	0.845	0.608



# RESULTS

Landau, Feynman-'t Hooft, Yennie ; lattice data: Bowman et al PRD (2005), Oliveira et al, PRD (2019)



## SUMMARY AND OUTLOOK

- Covariant Spectator Theory: dynamical quark model in Minkowski space with **confinement** and **dynamical chiral-symmetry breaking**
- $q\bar{q}$  interaction kernel
  - 2-body equation: very good description of **heavy and heavy-light** meson spectrum
  - 1-body equation: reasonable dressed **quark mass function**

### Outlook and work in progress:

- 1 mass function for finite bare quark masses
- 2 running quark-gluon coupling
- 3 running gluon mass
- 4  $\lambda \neq 0$ : vector structures for  $V_\ell$  in mass function calculation
- 5 quark mass function into bound-state calculations

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# IST-ID

Associação do Instituto Superior Técnico  
para a Investigação e Desenvolvimento

# FCT

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**THANK YOU!**