

The Nucleon Axial Form Factor for Neutrino Oscillation from First Principles

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Neutrino-Nucleus Interactions
in the Few GeV Region

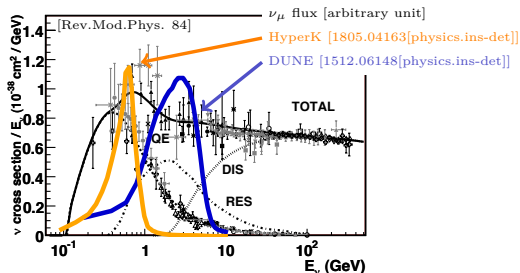
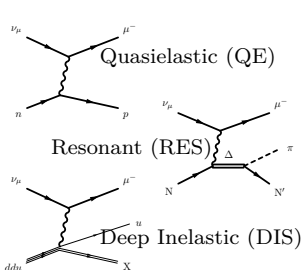
Outline

- ▶ Introduction
- ▶ Quasielastic Scattering from Experiment
 - Constraints from Deuterium Scattering
- ▶ LQCD Survey of $F_A(Q^2)$
 - Summary of $F_A(Q^2)$ Calculations
 - T2K/DUNE Implications
- ▶ Concluding Remarks

Note: all references in online slides are hyperlinked

Introduction

Neutrino Cross Sections



Energy range spans several *nucleon* interaction topologies

Nucleon amplitudes used to build *nuclear* cross sections

⇒ inputs to Monte Carlo simulations, E_ν reconstruction

Goal: isolate, quantify, improve *nucleon* amplitudes

⇒ Precise, theoretically robust *nucleon* inputs

→ definitive statements about *nuclear* uncertainties

Neutrino Event Topologies

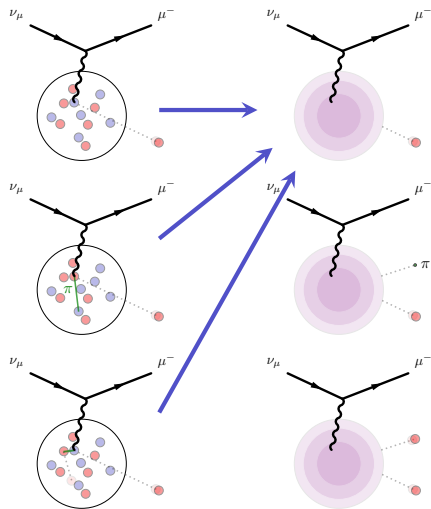
Use large nucleus for more nucleons to interact with

Nuclear environment complicates measurements:

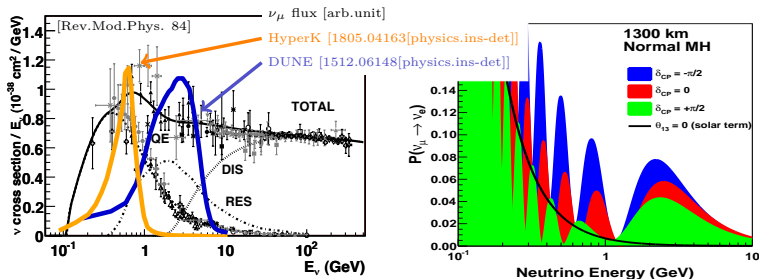
- ▶ Many allowed kinematic channels
- ▶ Reinteractions within nucleus
- ▶ Only final state particles are observable

Precise cross sections need precise nucleon amplitudes

Nucleon amplitudes assumed to be precisely known



Neutrino Cross Sections from Elementary Targets



Quasielastic is lowest E_ν , simplest \implies most important

Question:

How well do we know free nucleon quasielastic cross section
from elementary target sources?

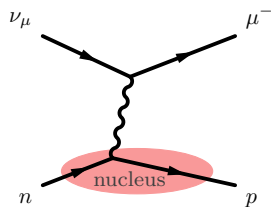
Two main sources:

- ▶ D₂ scattering
- ▶ Lattice QCD

Quasielastic Constraints

Quasielastic Form Factors

Quasielastic (QE) scattering assumes quasi-free nucleon inside nucleus



$$\mathcal{M}_{\text{nucleon}} = \langle \ell | \mathcal{J}^\mu | \nu_\ell \rangle \langle N' | \mathcal{J}_\mu | N \rangle$$

$$\langle N'(p') | (V - A)_\mu(q) | N(p) \rangle$$

$$= \bar{u}(p') \left[\begin{array}{l} \gamma_\mu F_1(q^2) + \frac{i}{2M_N} \sigma_{\mu\nu} q^\nu F_2(q^2) \\ + \gamma_\mu \gamma_5 F_A(q^2) + \frac{1}{2M_N} q_\mu \gamma_5 F_P(q^2) \end{array} \right] u(p)$$

- ▶ F_1, F_2 : constrained by eN scattering
- ▶ F_P : subleading in cross section,
 $\propto F_A$ from pion pole dominance constraint

Axial form factor F_A is leading contribution to nucleon cross section uncertainty

Form Factor Parameterizations

Most common in experimental literature: dipole ansatz —

$$F_A(Q^2) = g_A \left(1 + \frac{Q^2}{m_A^2}\right)^{-2}$$

- ▶ Overconstrained by both experimental and LQCD data (revisit later)
- ▶ Inconsistent with QCD, requirements from unitarity bounds
- ▶ Motivated by $Q^2 \rightarrow \infty$ limit, data restricted to low Q^2

Model independent alternative: z expansion [Phys.Rev.D 84 (2011)] —

$$F_A(z) = \sum_{k=0}^{\infty} a_k z^k \quad z(Q^2; t_0, t_{\text{cut}}) = \frac{\sqrt{t_{\text{cut}} + Q^2} - \sqrt{t_{\text{cut}} - t_0}}{\sqrt{t_{\text{cut}} + Q^2} + \sqrt{t_{\text{cut}} - t_0}} \quad t_{\text{cut}} \leq (3M_\pi)^2$$

- ▶ Rapidly converging expansion
- ▶ Controlled procedure for introducing new parameters

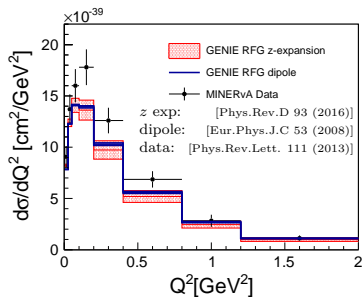
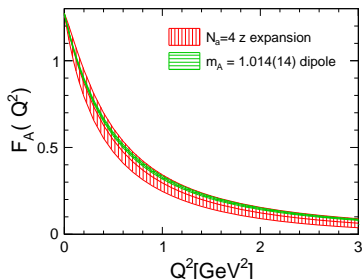
Deuterium Constraints on F_A

~ free nucleon constraints on F_A

- ▶ Outdated bubble chamber experiments:
 - Total $O(10^3)$ ν_μ QE events
 - Original data lost
 - Unknown corrections to data
 - **Deficient deuterium correction**
- ▶ Dipole overconstrained by data
underestimated uncertainty $\times O(10)$
- ▶ **Prediction discrepancies could be from nucleon and/or nuclear origins**

Coming soon:

Joint fit including MINER ν A $\bar{\nu}_\mu p \rightarrow \mu^+ n$
See Cai [NuInt22 talk] [thesis (2021)]



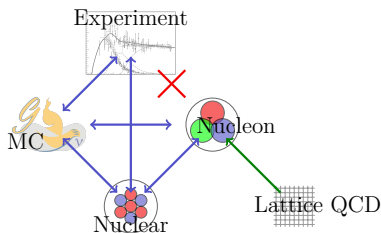
LQCD as Disruptive Technology

How can we improve precision?

Ideal: Modern high stats ν -D₂ scattering bubble chamber experiment

⇒ LQCD as a alternative/complement to experiment,
especially with experimentally inaccessible quantities

- ✓ No nuclear effects
- ✓ Realistic uncertainty estimates
- ✓ Systematically improvable
- ✓ Computers are (relatively) inexpensive



Build from the ground up:

Nucleon amplitudes from first principles

Robust uncertainty quantification

Well motivated theory inputs to nuclear models/EFTs

Lattice QCD & Implications

Lattice QCD Formalism

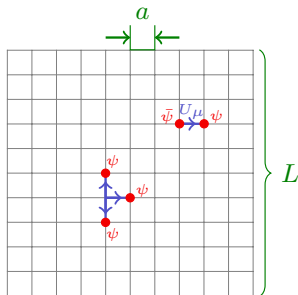
Numerical evaluation of path integral

Quark, gluon DOFs —

$$\langle \mathcal{O} \rangle = \frac{1}{Z} \int \mathcal{D}\psi \mathcal{D}\bar{\psi} \mathcal{D}U \exp(-S) \mathcal{O}_\psi [U]$$

Inputs: $am_{(u,d),\text{bare}}$
 $am_{s,\text{bare}}$
 $\beta = 6/g_{\text{bare}}^2$

Matching: e.g. $\frac{M_\pi}{M_\Omega}$, $\frac{M_K}{M_\Omega}$, M_Ω
Only 1 per computational input



Results — first principles predictions from QCD,
gluons to all orders

“Complete” error budget \implies extrapolation in a , L , M_π guided by EFT, FV χ PT

- ▶ $a \rightarrow 0$ (continuum limit)
- ▶ $L \rightarrow \infty$ (infinite volume limit)
- ▶ $M_\pi \rightarrow M_\pi^{\text{phys}}$ (chiral limit)

Now published!



Status of Lattice QCD Determination of Nucleon Form Factors and Their Relevance for the Few-GeV Neutrino Program

Annual Review of Nuclear and Particle Science

Vol. 72 - (Volume publication date September 2022)

Review in Advance first posted online on July 8, 2022. (Changes may still occur before final publication.)

<https://doi.org/10.1146/annurev-nucl-010622-120608>

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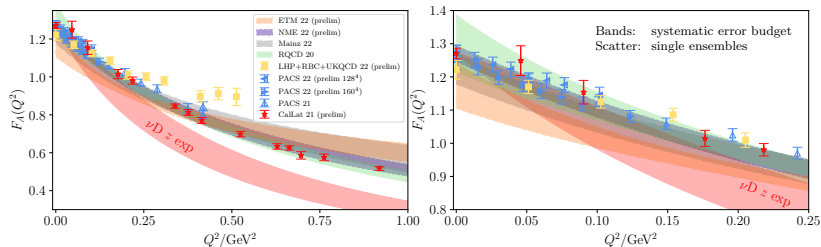
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Abstract

Calculations of neutrino–nucleus cross sections begin with the neutrino–nucleon interaction, making the latter critically important to flagship neutrino oscillation experiments despite limited measurements with poor statistics. Alternatively, lattice quantum chromodynamics (LQCD) can be used to determine these interactions from the Standard Model with quantifiable theoretical uncertainties. Recent LQCD results of g_A are in excellent agreement with data, and results for the (quasi-)elastic nucleon form factors with full uncertainty budgets are expected within a few years. We review the status of the field and LQCD results for the nucleon axial form factor, $F_A(Q^2)$, a major source of uncertainty in modeling sub-GeV neutrino–nucleon interactions. Results from different LQCD calculations are consistent but collectively disagree with existing models, with potential implications for current and future neutrino oscillation experiments. We describe a road map to solidify confidence in the LQCD results and discuss future calculations of more complicated processes, which are important to few-GeV neutrino oscillation experiments.

Expected final online publication date for the *Annual Review of Nuclear and Particle Science*, Volume 72 is September 2022. Please see <http://www.annualreviews.org/page/journal/pubdates> for revised estimates.

Nucleon Axial Form Factor

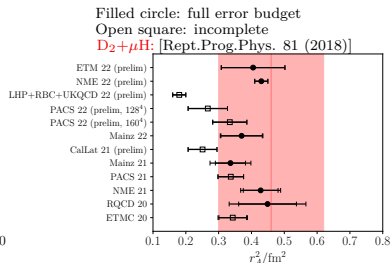
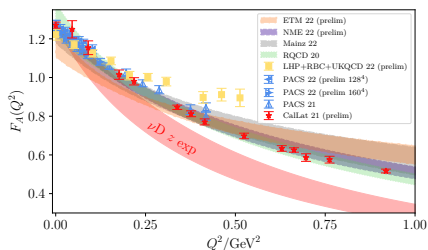


LQCD results maturing:

- ▶ Many results, all physical M_π :
independent data & different methods
- ▶ Agreement w/ single ensemble
 \implies Small systematic effects observed
(expectation: largest at $Q^2 \rightarrow 0$)
- ▶ Extrapolated results (bands) satisfy nontrivial PCAC checks
Lots of recent effort to understand

Evidence of slow Q^2 falloff, **situation unlikely to change drastically**

Axial Radius (r_A^2)



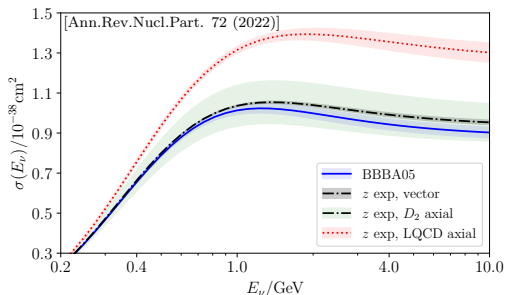
Radius related to slope: $r_A^2 = -\frac{6}{g_A} \frac{dF_A}{dQ^2} \Big|_{Q^2=0}$

Good agreement with r_A^2 from experiment, poor agreement with large Q^2

Fixing radius to agree at large Q^2 would bring radius down to $r_A^2 \sim 0.25 \text{ fm}^2$

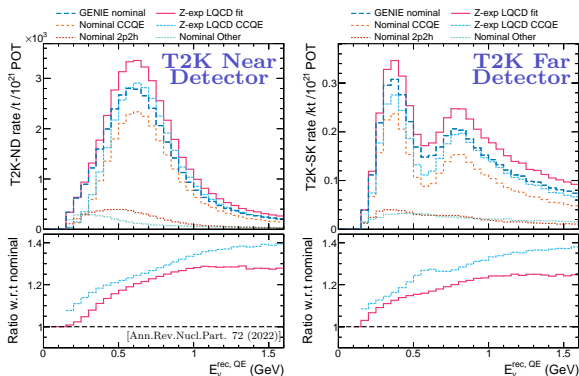
\Rightarrow Incompatible with dipole ansatz

Free Nucleon Cross Section



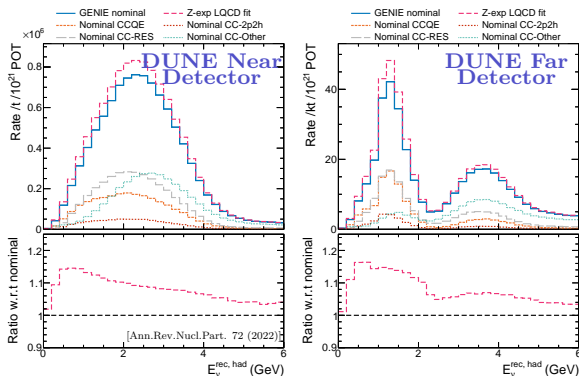
- ▶ Representative LQCD dataset (Callat 21)
- ▶ Integral over $Q^2 \implies$ enhancement of discrepancy
- ▶ LQCD prefers 30-40% enhancement of ν_μ CCQE cross section
- ▶ recent Monte Carlo tunes require 20% enhancement of QE
[Phys.Rev.D 105 (2022)] [2206.11050 [hep-ph]]
similar trend with continuum Schwinger function methods
[Phys.Rev.D 105 (2022)] [2206.12518 [hep-ph]]
- ▶ With improved precision, sensitive to vector FF tensor (black vs blue)
[Phys.Rev.D 102 (2020)] VS [Nucl.Phys.B Proc.Suppl. 159 (2006)]

T2K Implications



- ▶ Dashed dark blue (GENIE nominal) vs solid magenta (z exp LQCD fit)
- ▶ QE enhancements produce 10-20% event rate enhancement, E_{ν} -dependent
- ▶ Monte Carlo tuning makes more detailed comparisons complicated
 - ⇒ All channels are adjusted to compensate for QE changes
- ▶ cross section changes at ND \neq effective cross section changes at FD:
 - insufficient CCQE model freedom → bias in FD prediction

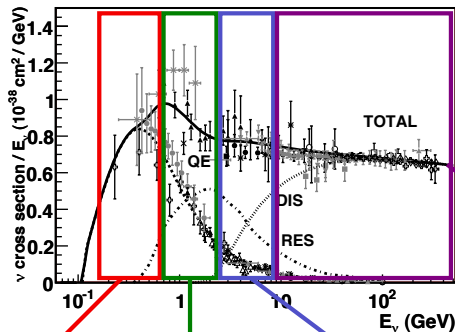
DUNE Implications



- ▶ Solid dark blue (GENIE nominal) vs dashed magenta (z exp LQCD fit)
- ▶ QE enhancements produce 10-20% event rate enhancement, E_{ν} -dependent
- ▶ Monte Carlo tuning makes more detailed comparisons complicated
 \implies All channels are adjusted to compensate for QE changes
- ▶ cross section changes at ND \neq effective cross section changes at FD:
 insufficient CCQE model freedom \rightarrow bias in FD prediction

Concluding Remarks

Future Directions for LQCD



Deep Inelastic Scattering
-Axial quasi/pseudo PDF

Quasielastic

- Nucleon Form Factors
- Full Error Budgets
- Detailed Systematics

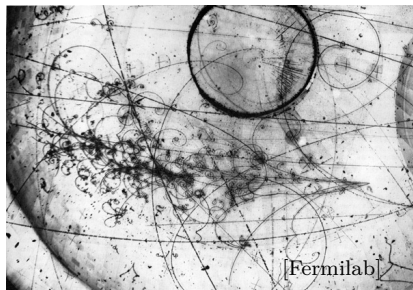
$N \rightarrow \Delta, N \rightarrow N^*$

- Transition Matrix Elements
- Multiparticle Operators
- Initial Calculations

Shallow Inelastic Scattering

- Hadronic Tensor
- Four Point Functions
- Exploration

Outlook

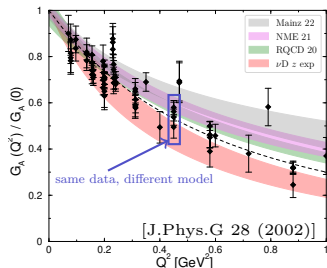


- ▶ *Nucleon* form factor uncertainty significantly underestimated
- ▶ LQCD is a proxy for missing experimental data
- ▶ Mounting evidence that ν QE cross section underestimated
 \implies Attention needed to avoid biased results
- ▶ Lots of opportunities for LQCD to provide key missing inputs for ν oscillation

Thank you for your attention!

Backup

Electro Pion Production



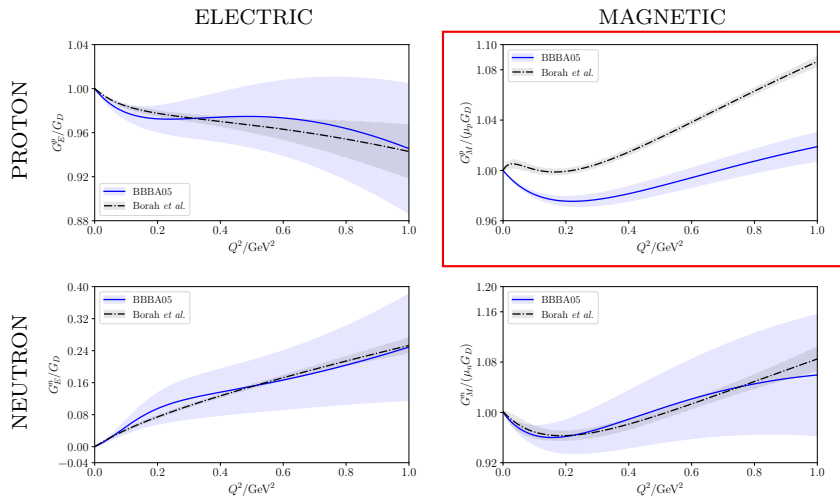
- ▶ Large model uncertainty, not included in world averages
- ▶ Valid only in $M_\pi \rightarrow 0, q \rightarrow 0$ limits
- ▶ Expansion to $O(M_\pi^2, Q^2)$:
 - restricted Q^2 validity
 - lacks shape freedom in Q^2
- ▶ Predates Heavy Baryon χ PT, no systematic power counting

Modern experiments do not report $F_A(Q^2) \implies$ averages out of date

Possible argument for comparing to r_A^2 from low Q^2 ; high Q^2 untrustworthy

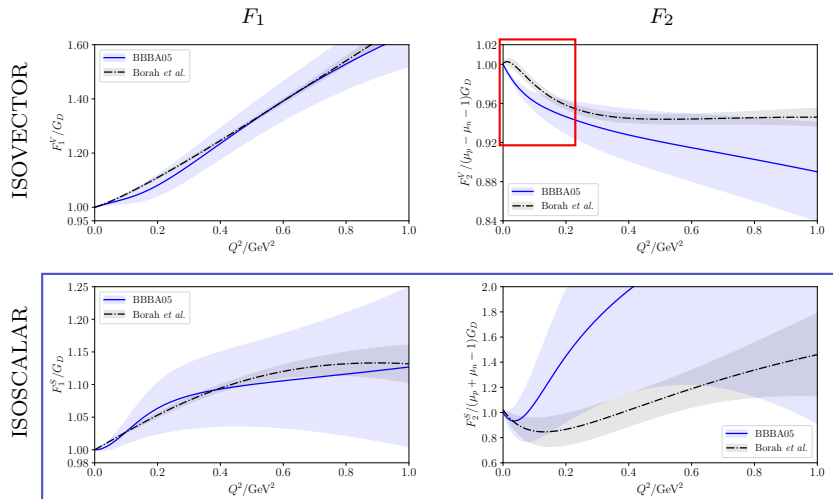
Effort needed to update prediction from photo/electro pion production

Vector Form Factors - Proton/Neutron



Large tension in proton magnetic form factor

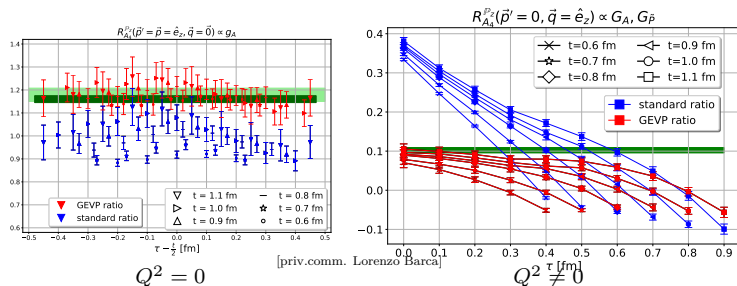
Vector Form Factors - Isospin Symmetric



Uncertain slope of F_2^V

Large uncertainty on isoscalar form factors

Axial FF - $N\pi$ Interpolating Operators



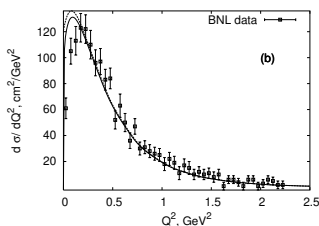
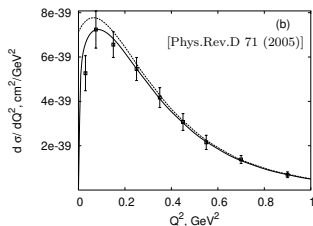
Address primary source of excited state contamination: $N\pi$

2×2 operator basis, explicit 3- and 5-quark interpolating operators

Significantly flatter ratios, simplified analysis

Will analysis with only 3-quark operators be consistent?

Resonance Production - $N \rightarrow \Delta$



$N \rightarrow \Delta$ transition form factors are poorly known, but needed

1π production cross section known to 30% [Phys.Rev.C 88 (2013)]

DUNE error budget anticipates $\lesssim 10\%$ precision [2002.03005 [hep-ex]]

Completely unconstrained axial form factors in other $J^P = 3/2^-$ channels
 \implies 100% uncertainties from $V - A$, $A - A$ interference terms

[Phys.Rev.D 74 (2006)]

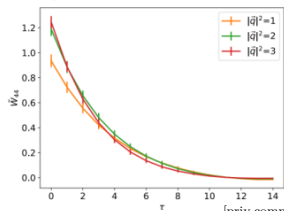
Previous work by ETM: [Phys.Rev.D 83 (2011)] [Phys.Rev.Lett. 98 (2007)]

Formal developments:

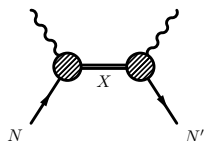
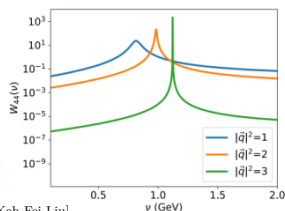
$$1 + \mathcal{J} \rightarrow 2 \quad (N\gamma^* \rightarrow N\pi) \quad [\text{Phys.Rev.D 103 (2021)}]$$

$$1 + \mathcal{J} \rightarrow 2 + \mathcal{J} \quad (N\gamma^* \rightarrow \Delta \rightarrow N\pi\gamma^*) \quad [\text{Phys.Rev.D 105 (2022)}]$$

Resonance Production - $N \rightarrow N^*$



[priv.comm. Keh-Fei Liu]



See also: [Phys.Rev.D 101 (2020)]

Four point function with $\langle \mathcal{O}(0) \mathcal{J}_4(-q) \mathcal{J}_4(q) \bar{\mathcal{O}}(0) \rangle$, $M_\pi \sim 370$ MeV

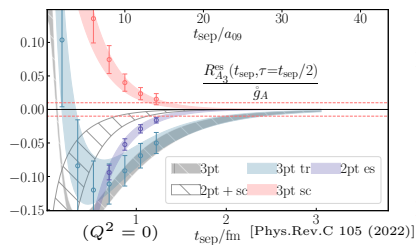
Removed elastic contribution \implies resonant response (strong overlap with Roper)

Hadronic tensor methods for addressing SIS ($1.4 \text{ GeV} \leq W \leq 2.0 \text{ GeV}$)

Large $N\pi$, $N\pi\pi$ contributions; strong interferences between resonant/nonresonant

Currently no practical $Q^2 \neq 0$ data in this region [S.Nakamura - NuSTEC S&DIS]

Excited States in g_A

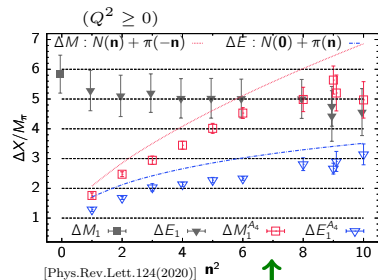


Compare fit to correlator data ratio

Remnant contamination dominated by “transition” states ($m \rightarrow n$, violet)

Statistically significant until 2 fm typical data $\lesssim 1$ fm

Excited states still present in practically achievable large time limit



NME collab:

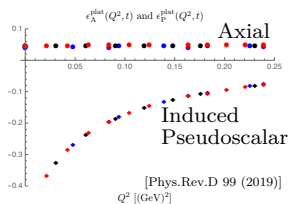
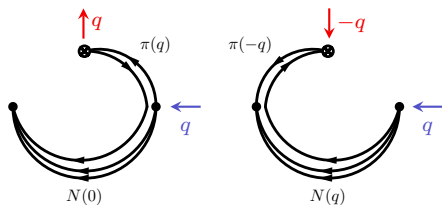
Q^2 contamination from $N \rightarrow N\pi$

Dominant contribution agrees with χ Pt expectation

$N\pi$ is important for $g_A(Q^2)$

NOTE: expect only approx agreement between data/curves

Excited States - χ PT and $N\pi$



Contamination primarily from enhanced $N\pi$, mostly from induced pseudoscalar

Intermediate fits not sensitive to $N\pi$

\implies need simultaneous fits including intermediate correlators

[Phys.Rev.C 105 (2022)] [Phys.Rev.D 105 (2022)]

Alternate fit strategies to remove $N\pi$ (are they comparable?):

- ▶ Kinematic constraints ($F_P = 0$)
- ▶ include \mathcal{A}_4 (strong $N\pi$ coupling)
- ▶ explicit $N\pi$ operators

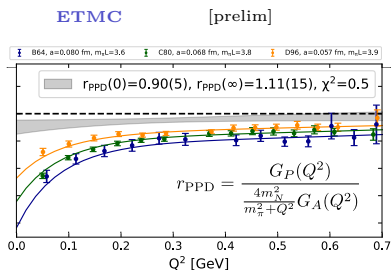
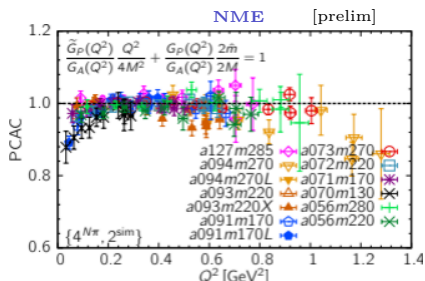
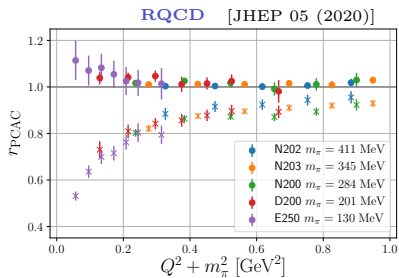
Prediction from χ PT: [Phys.Rev.D 99 (2019)]

First demonstration by NME: [Phys.Rev.Lett. 124 (2020)]

χ PT-inspired fit methods for fitting form factor data

[Phys.Rev.D 105 (2022)] [JHEP 05 (2020) 126]

PCAC Checks



- ▶ Relation btw F_A , F_P , \tilde{F}_P via PCAC
- ▶ Contamination in F_A and \tilde{F}_P , F_P very different [Phys.Rev.D 99 (2019)]
 \Rightarrow nontrivial consistency check