

Update on Bodek-Yang Model

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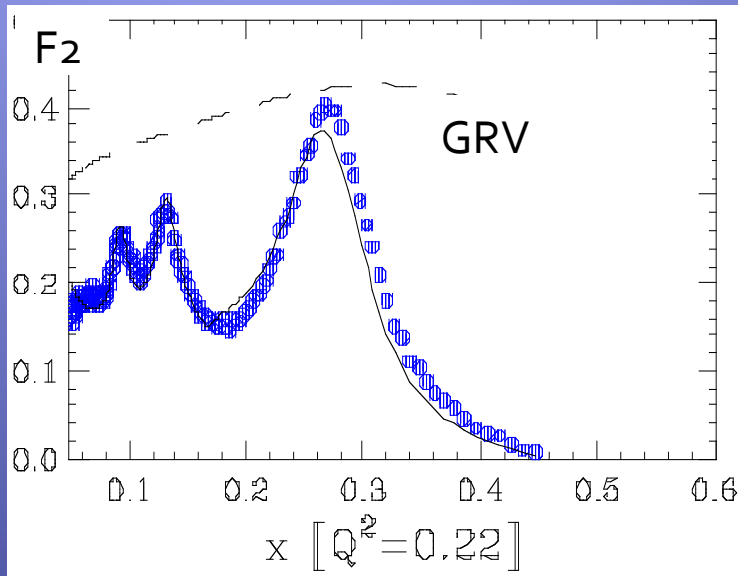
Un-ki Yang

Seoul National University

NuINT 2022 Workshop, October 24-29, 2022

Bodek-Yang Model

- Bodek-Yang model: describe DIS cross section in all Q^2 regions
- Challenges in e/μ -N DIS
 - High x PDFs at low Q^2
 - Resonance region overlapped with a DIS contribution
 - Hard to extrapolate DIS contribution to low Q^2 region from high Q^2 data due to non-perturbative QCD effects



- A model in terms of quark-parton model (easy to convert e/μ scattering to ν scattering)
- ✓ Understanding of high x PDFs at low Q^2 ? wealthy SLAC, JLAB data.
- ✓ Understanding of resonance scattering in terms of quark-parton model? (duality works, many studies by JLAB)

Lessons from previous QCD studies

- **NLO & NNLO analyses with DIS data:** PRL 82, 2467 (1999), Eur. Phys. J. C13, 241 (2000) by Bodek and Yang
 - **Kinematic higher twist (target mass) effects** are large and must be included in the form of Georgi & Politzer x scaling.
 - **Resonance region is also well described (duality works).**
 - Most of **dynamic higher twist corrections (in NLO analysis)** are similar to missing NNLO higher order terms.
- **NNLO pQCD+TM with NNLO PDFs can describe the non-perturbative QCD effects at low Q^2**
- Thus, we reverse the approach to build the model:
 - **Use LO PDFs and “effective target mass and final state masses” to account for initial target mass, final target mass, and even missing higher orders**

$$\xi = \frac{2xQ'^2}{Q^2(1 + \sqrt{1 + 4M^2x^2/Q^2})},$$

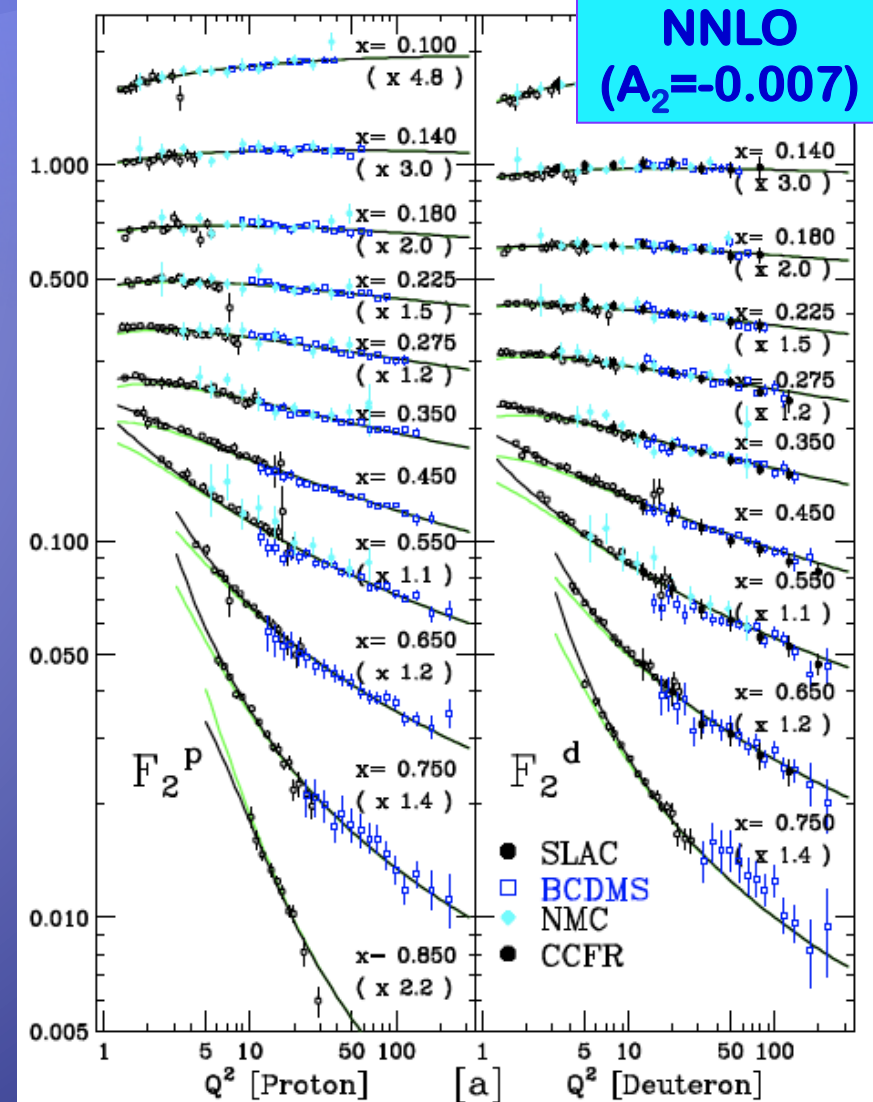
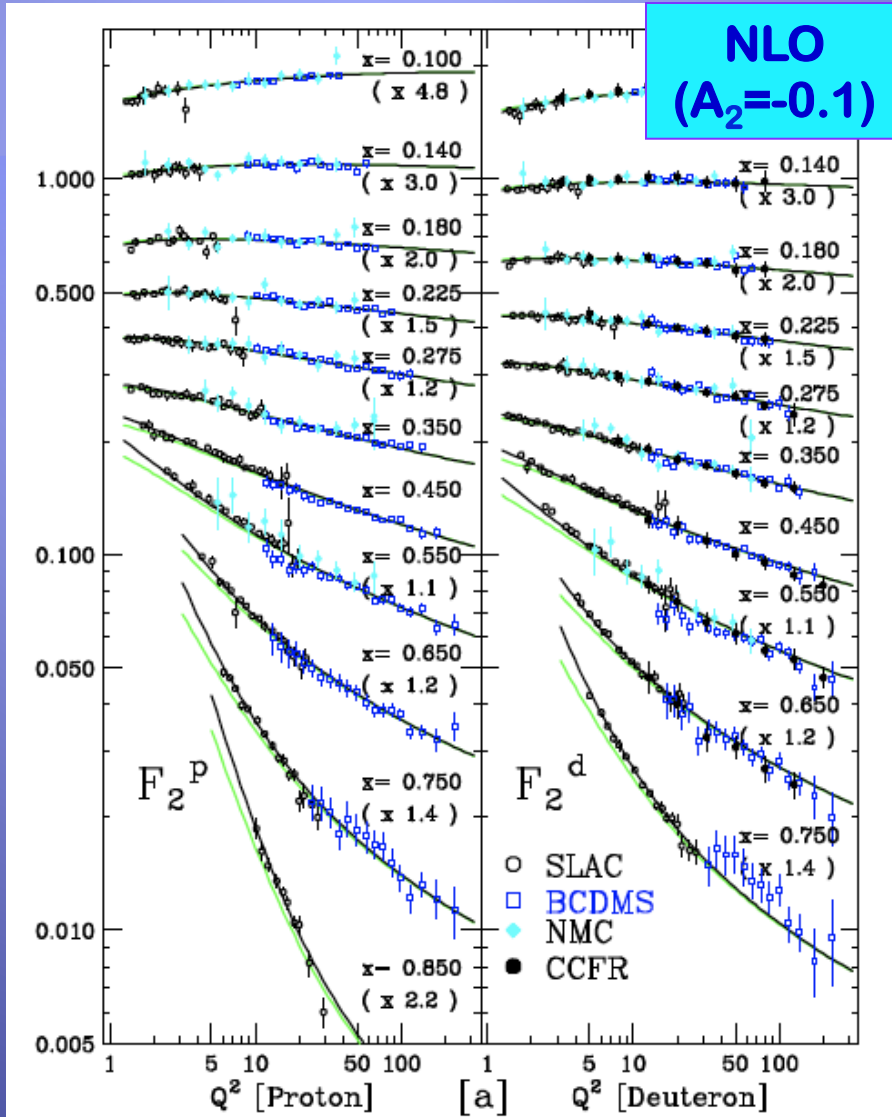
$$2Q'^2 = [Q^2 + M_f^2 - M_t^2] + \sqrt{(Q^2 + M_f^2 - M_t^2)^2 + 4Q^2(M_t^2 + P_T^2)}.$$

NLO vs NNLO

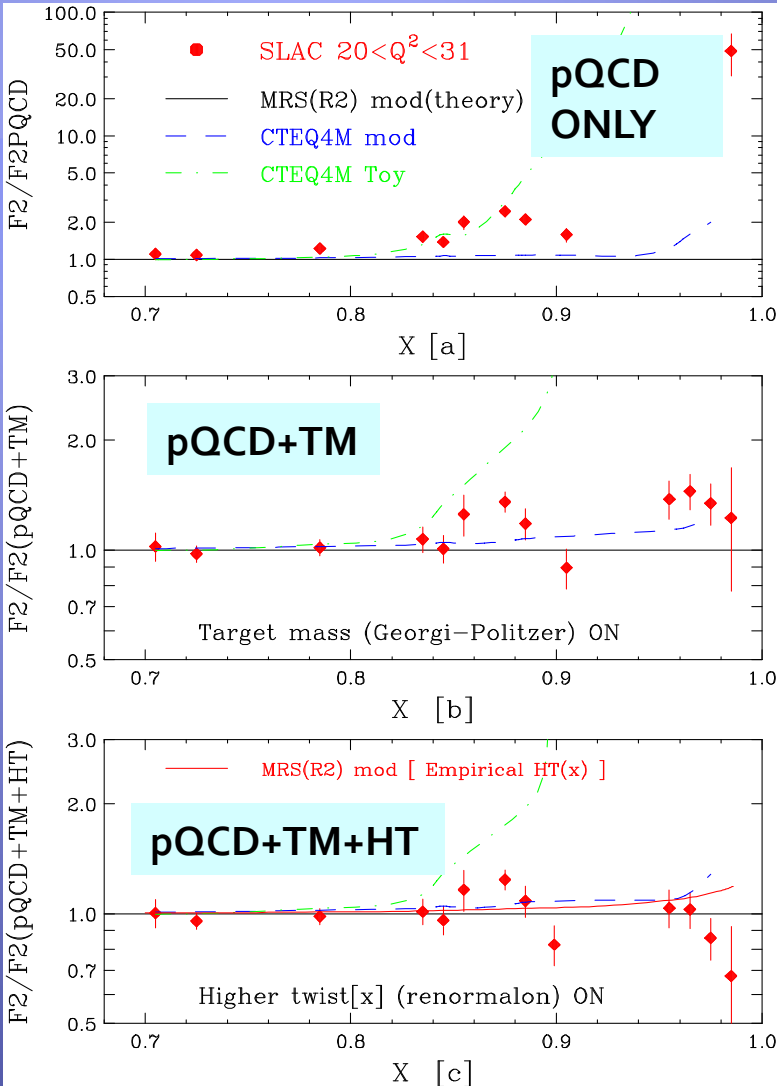
Studies of higher twist and higher order effects in NLO and NNLO QCD analysis of lepton-nucleon scattering data on F_2 and $R = \sigma_L/\sigma_T$

U.K. Yang & A. Bodek

The European Physical Journal C - Particles and Fields 13, 241–245 (2000) | [Cite this article](#)



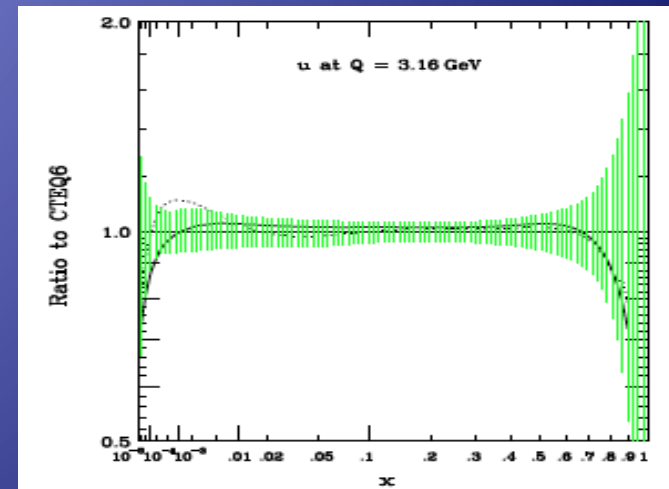
Very high x and low Q^2 data



Parton Distributions, d/u , and Higher Twist Effects at High x

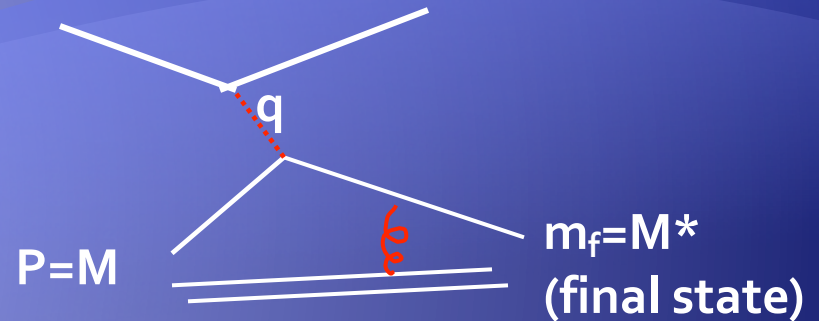
U. K. Yang and A. Bodek
Phys. Rev. Lett. **82**, 2467 – Published 22 March 1999

- Very high x and low Q^2 data is well described by the pQCD+TM+HT
- Extraction of the high x PDF is promising (1999)
- still a large uncertainty (2022)



Modeling neutrino cross sections

- NNLO pQCD +TM approach: describes the DIS region and resonance data very well



- **Bodek-Yang LO approach**: (pseudo NNLO)
 - Use effective LO PDFs with a new scaling variable, ξ_w to absorb target mass, higher twist, missing QCD higher orders

$$x_{Bj} = \frac{Q^2}{2M\nu}$$



$$\xi_w = \frac{Q^2 + B}{\{M\nu[1 + \sqrt{(1 + Q^2 / \nu^2)}] + A\}}$$

- Multiply all PDFs by K factors for photo production limit and higher twist

$$F_2(x, Q^2) \rightarrow \frac{Q^2}{Q^2 + C} F_2(\xi_w, Q^2)$$

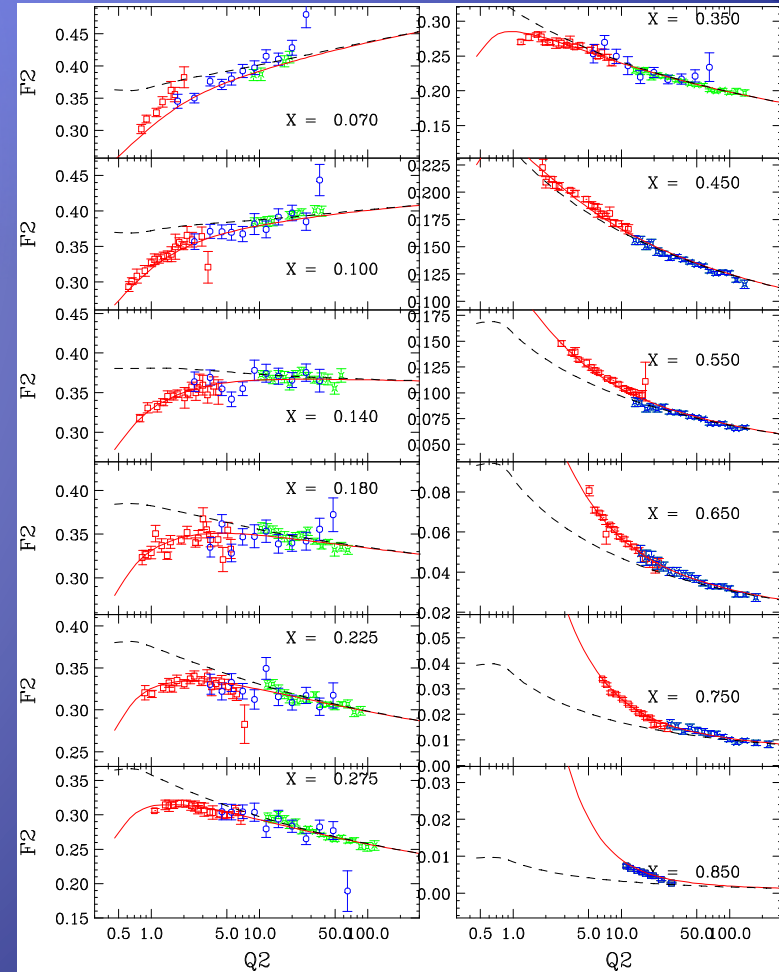
Bodek-Yang Effective LO PDFs Model

1. Start with GRV98 LO ($Q^2_{\min}=0.80$)
2. Replace x_{bj} with a new scaling, ξ_w
3. Multiply all PDFs by K factors for photo prod. limit and higher twist

$$[\sigma(\gamma) = 4\pi\alpha/Q^2 * F_2(x, Q^2)]$$

$$K_{sea} = Q^2/[Q^2+C_{sea}]$$

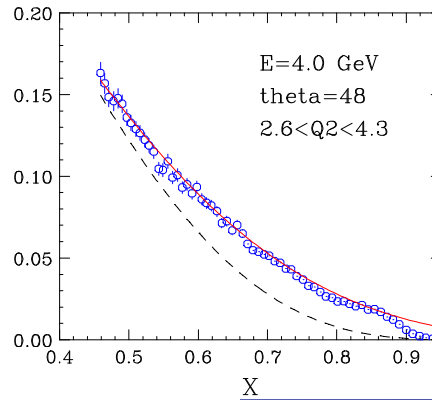
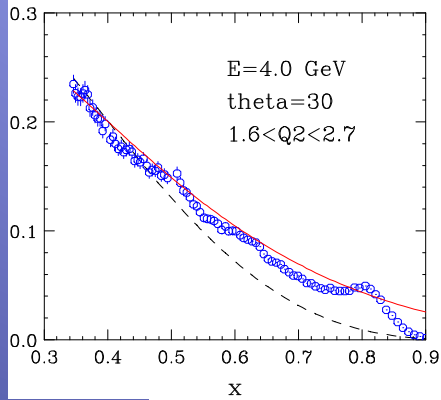
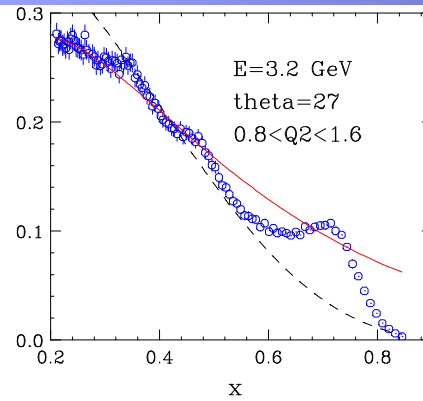
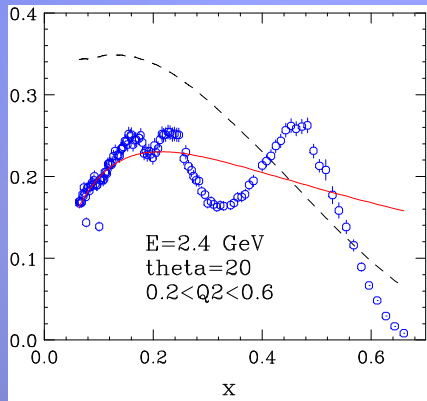
$$K_{val} = [1 - G_D^2(Q^2)] * [Q^2+C_{2V}] / [Q^2+C_{1V}]$$
 motivated by Adler Sum rule
 where $G_D^2(Q^2) = 1/[1+Q^2/0.71]^4$
4. Freeze the evolution at $Q^2 = Q^2_{\min}$
 - $F_2(x, Q^2 < 0.8) = K(Q^2) * F_2(\xi_w, Q^2=0.8)$
5. Fit all DIS $F_2(p/D)$ data: with $W > 2$ GeV
 SLAC/BCDMS/NMC/HERA data



$F_2(p)$

$$\chi^2/DOF = 1235/1200$$

Predictions for Resonance, Photo-production data



$F_2(d)$ resonance

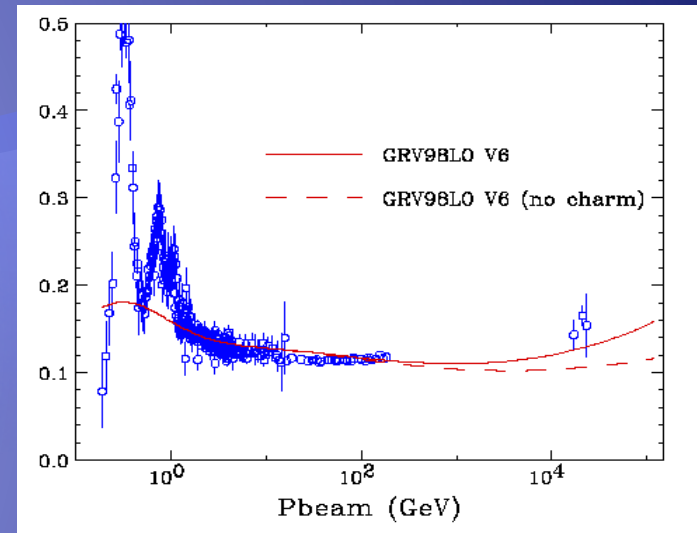


Photo-production (P)

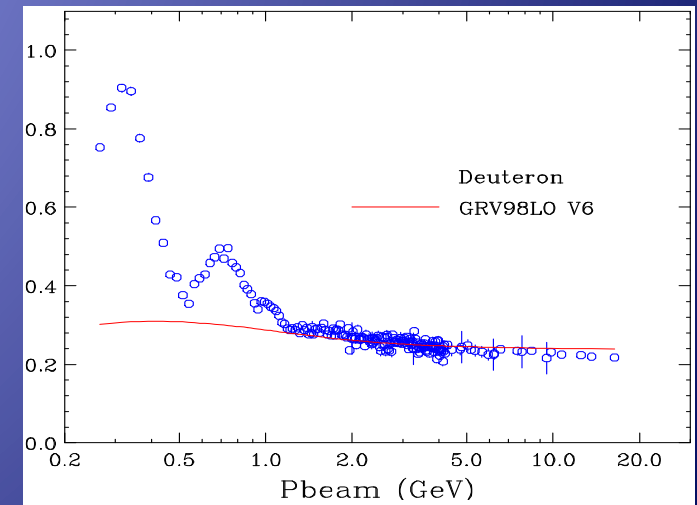


Photo-production (d)

Bodek-Yang Effective LO PDFs Model

- Include the photo-production data
- Use different K factors for up and down quark type separately

$$K_{val}(u,d) = [1 - G_D^2(Q^2)] * [Q^2 + C_{2V}] / [Q^2 + C_{1V}]$$

$$K_{sea}(u,d,s) = Q^2 / [Q^2 + C_{sea}]$$

- Additional K^{LW} factor for valence quarks:

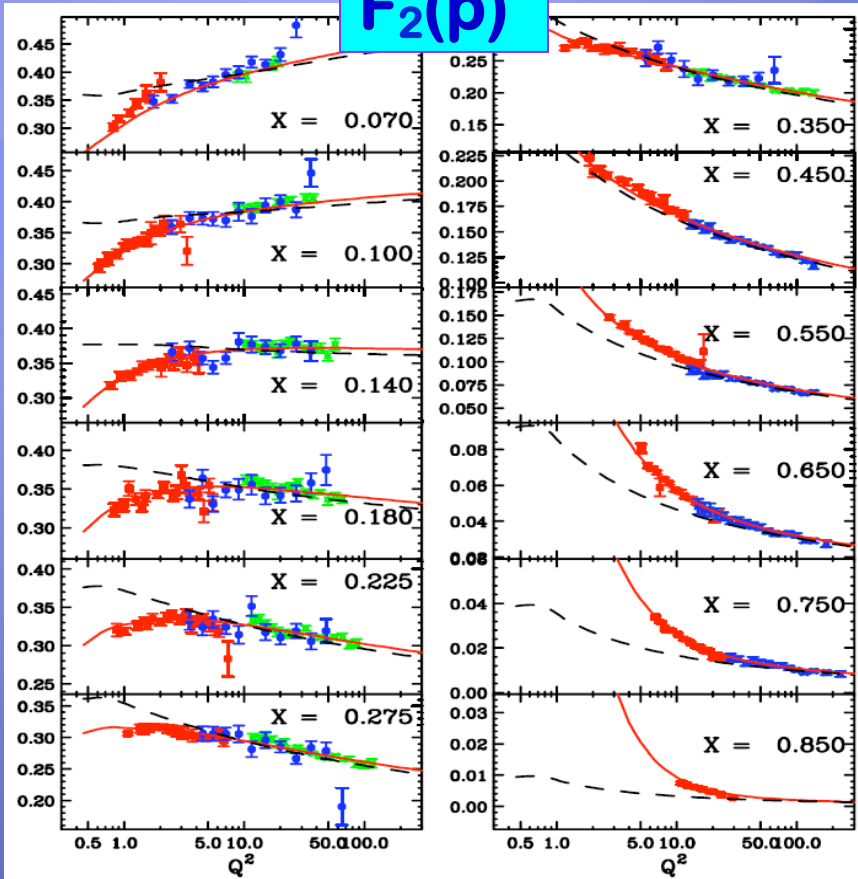
$$K_{val} = K^{LW} * [1 - G_D^2(Q^2)] * [Q^2 + C_{2V}] / [Q^2 + C_{1V}]$$

where $K^{LW} = (\nu^2 + C^V) / \nu^2$

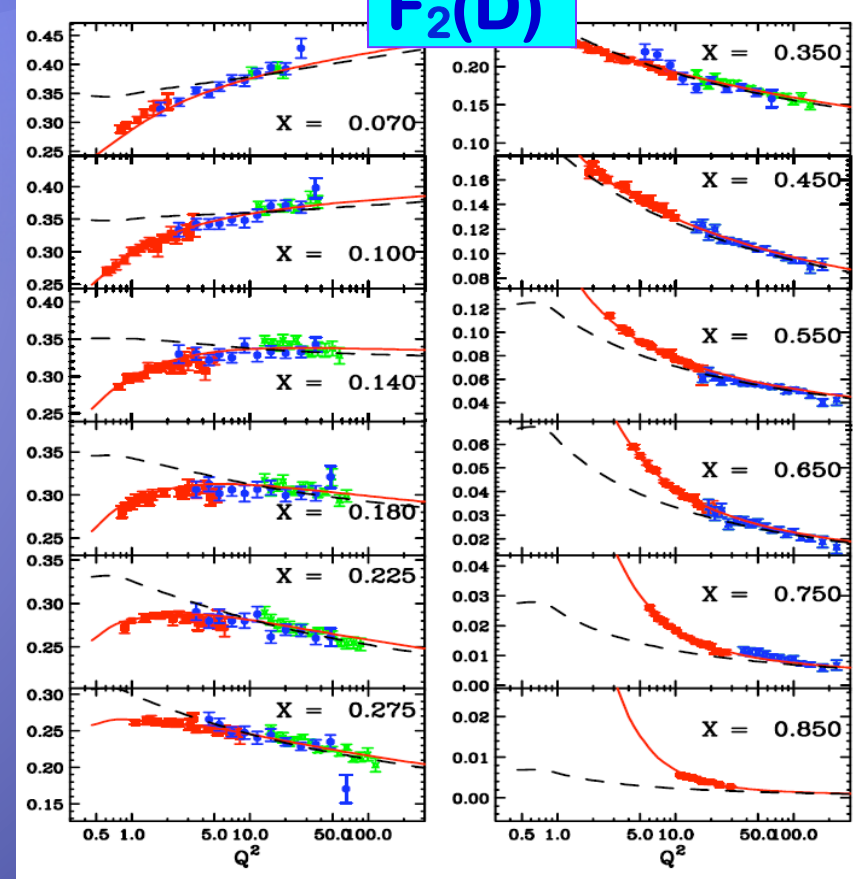
<i>A</i>	<i>B</i>	<i>C_{v2d}</i>	<i>C_{v2u}</i>
0.621	0.380	0.323	0.264
<i>C_{sea}^{down}</i>	<i>C_{sea}^{up}</i>	<i>C_{v1d}</i>	<i>C_{v1u}</i>
0.561	0.369	0.341	0.417
<i>C_{sea}^{strange}</i>	<i>C^{low-ν}</i>	<i>F_{valence}</i>	<i>N</i>
0.561	0.218	$[1 - G_D^2(Q^2)]$	1.026

Fit Results on DIS $F_2(p/D)$ data

$F_2(p)$

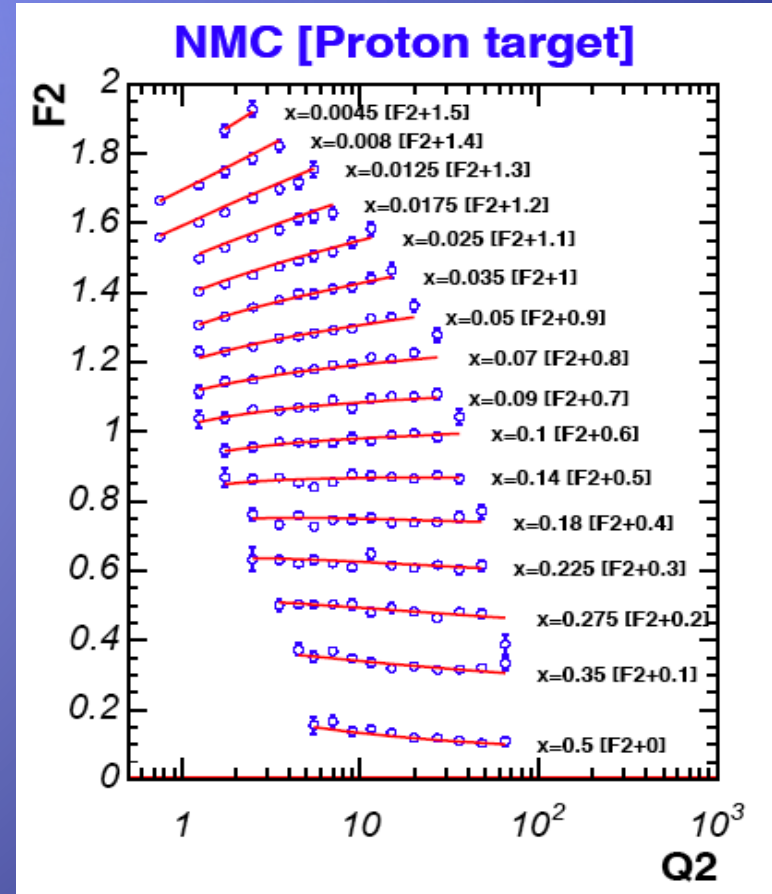
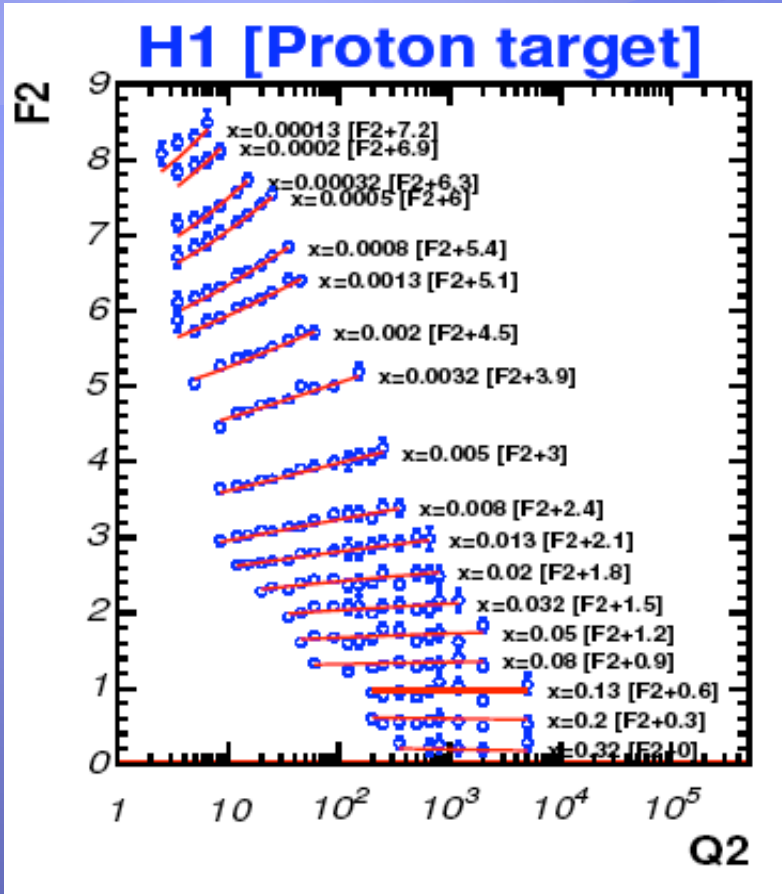


$F_2(D)$



- Excellent Fitting:
 - red solid line: effective LO using ξ_w
 - black dashed line: x_{bj}

Low x HERA and NMC data



➤ Fit works at low x

Photo-production data

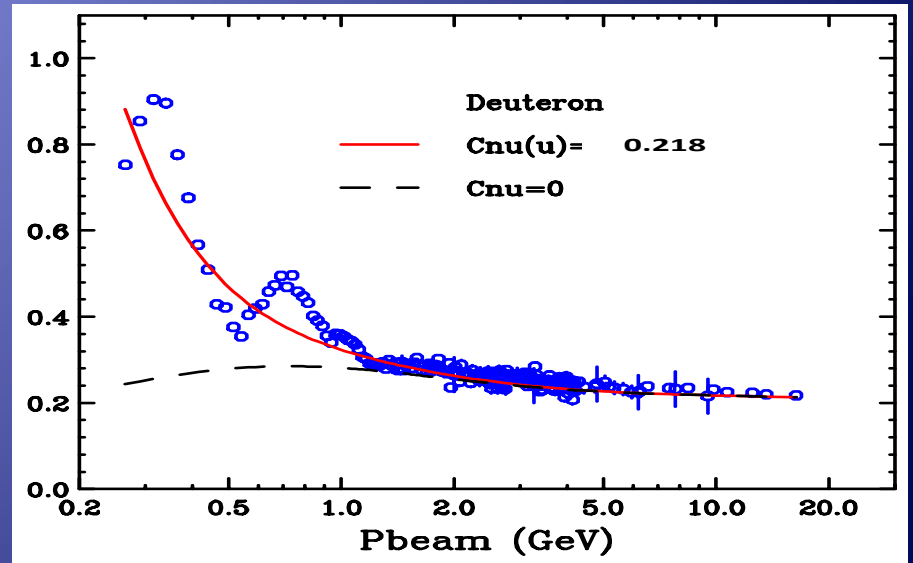
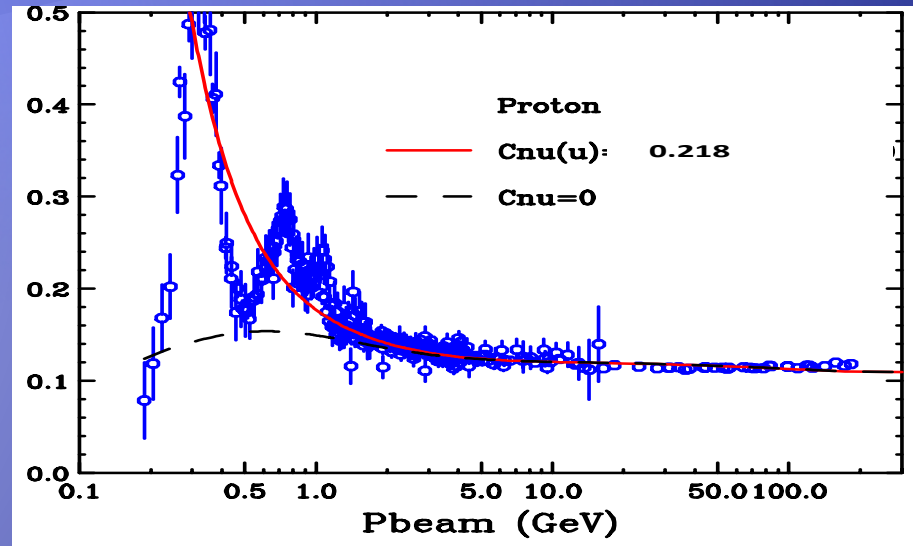
- Additional K^{LW} factor for valence quarks:

$$K_{val} = K^{LW} * [1 - G_D^2(Q^2)] * [Q^2 + C_{2V}] / [Q^2 + C_{1V}]$$

$$K^{LW} = (\nu^2 + C^V) / \nu^2$$

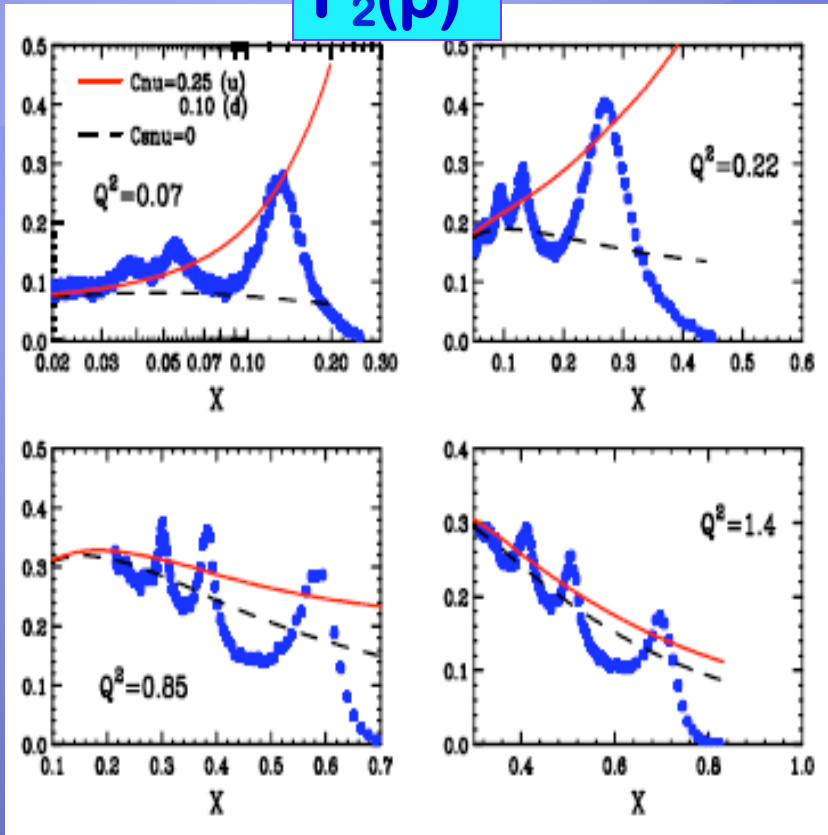
This makes a duality work all the way down to $Q^2=0$ (for charged leptons)

- Photo-production data with $\nu(P_{beam}) > 1$ GeV included in the fitting

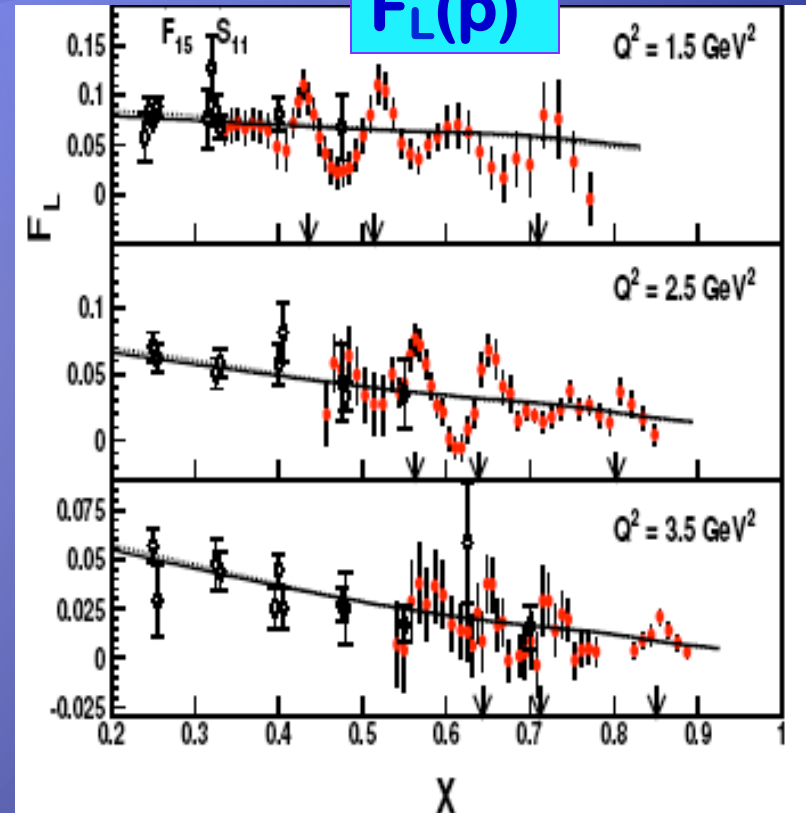


F₂ & F_L Resonance data

F₂(p)



F_L(p)



$$F_L = F_2(1 + 4M^2x^2/Q^2) \frac{R}{(1+R)}$$

- Predictions are in good agreement (not included in the fit) duality works
- F_L was calculated using F₂ and R₁₉₉₈

Neutrino cross sections

- Effective LO model with ξw describe all DIS and resonance F_2 data as well as photo-production data ($Q^2=0$ limit): vector contribution works well
- Neutrino Scattering:
 - Effective LO model works for xF_3 ?
 - Nuclear correction using e/μ scattering data
 - Axial vector contribution at low Q^2 ?
 - Use $R=R_{1998}$ to get $2xF_1$
 - Implement charm mass effect through ξw slow rescaling algorithm for F_2 , $2xF_1$, and xF_3

Submitted to EPJC (2022), arXiv:2108.09240

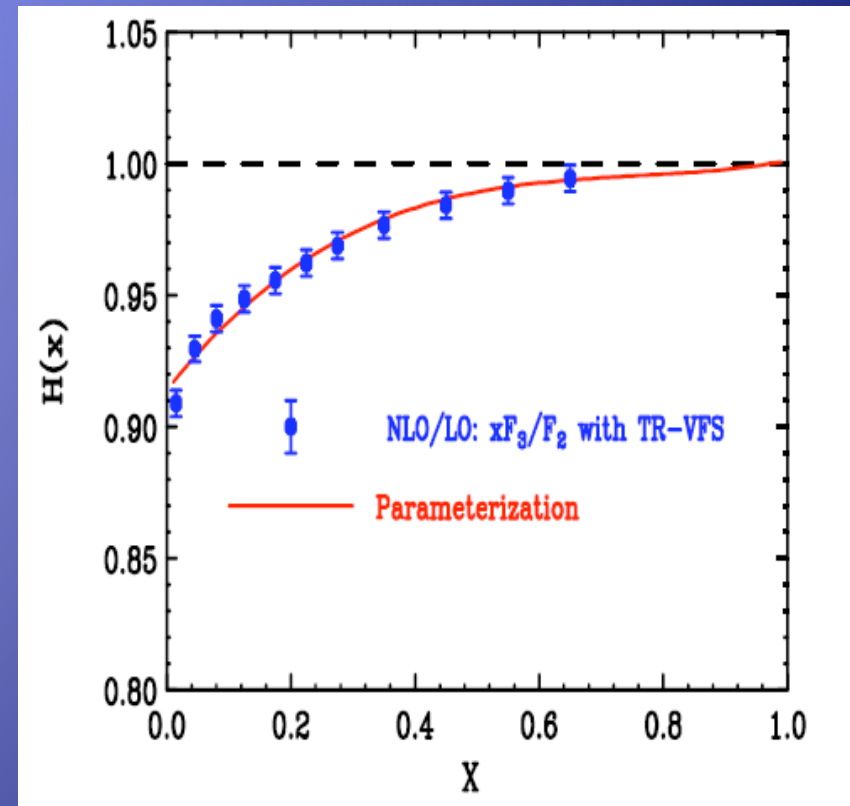
Effective LO model for xF_3 ?

- Scaling variable, ξw absorbs higher order effect for F_2 , but the higher order effects for F_2 and xF_3 are not the same
- Use NLO QCD to get double ratio

$$H(x) = \frac{xF_3(\text{NLO})}{xF_3(\text{LO})} / \frac{F_2(\text{NLO})}{F_2(\text{LO})}$$

not 1 but almost indep. of Q^2

- Enhance anti-neutrino cross section by 3%



Axial Vector Structure Functions

- At high Q^2 , vector and axial vector contribution are same, but not at low Q^2
- K factors for axial contributions: type II

$$K_{sea}^{vector} = \frac{Q^2}{Q^2 + C} \Rightarrow K_{sea}^{axial} = \frac{Q^2 + 0.55C_{sea}^{axial}}{Q^2 + C_{sea}^{axial}}$$

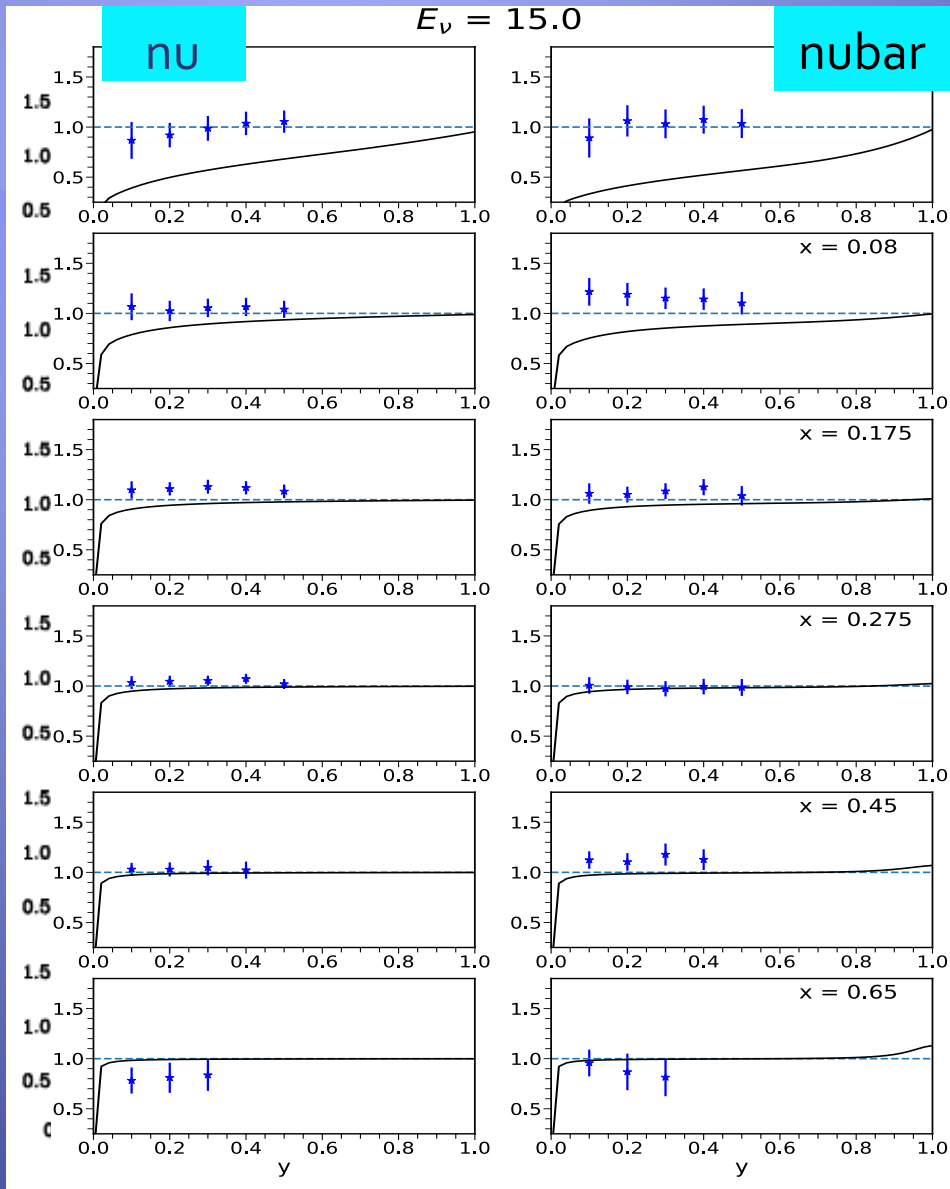
$$K_{val}^{axial} = \frac{Q^2 + 0.1C_{val}^{axial}}{Q^2 + C_{val}^{axial}}$$

where $C_{sea}^{axial} = 0.75$, $C_{val}^{axial} = 0.18$

- 0.55 was chosen to satisfy the prediction from PCAC by Kulagin, agrees with CCFR/CHROUS data for F_2 extrapolation to ($Q^2=0$)
- But, the non-zero PCAC component of F_2^{axial} at low Q^2 : mostly longitudinal

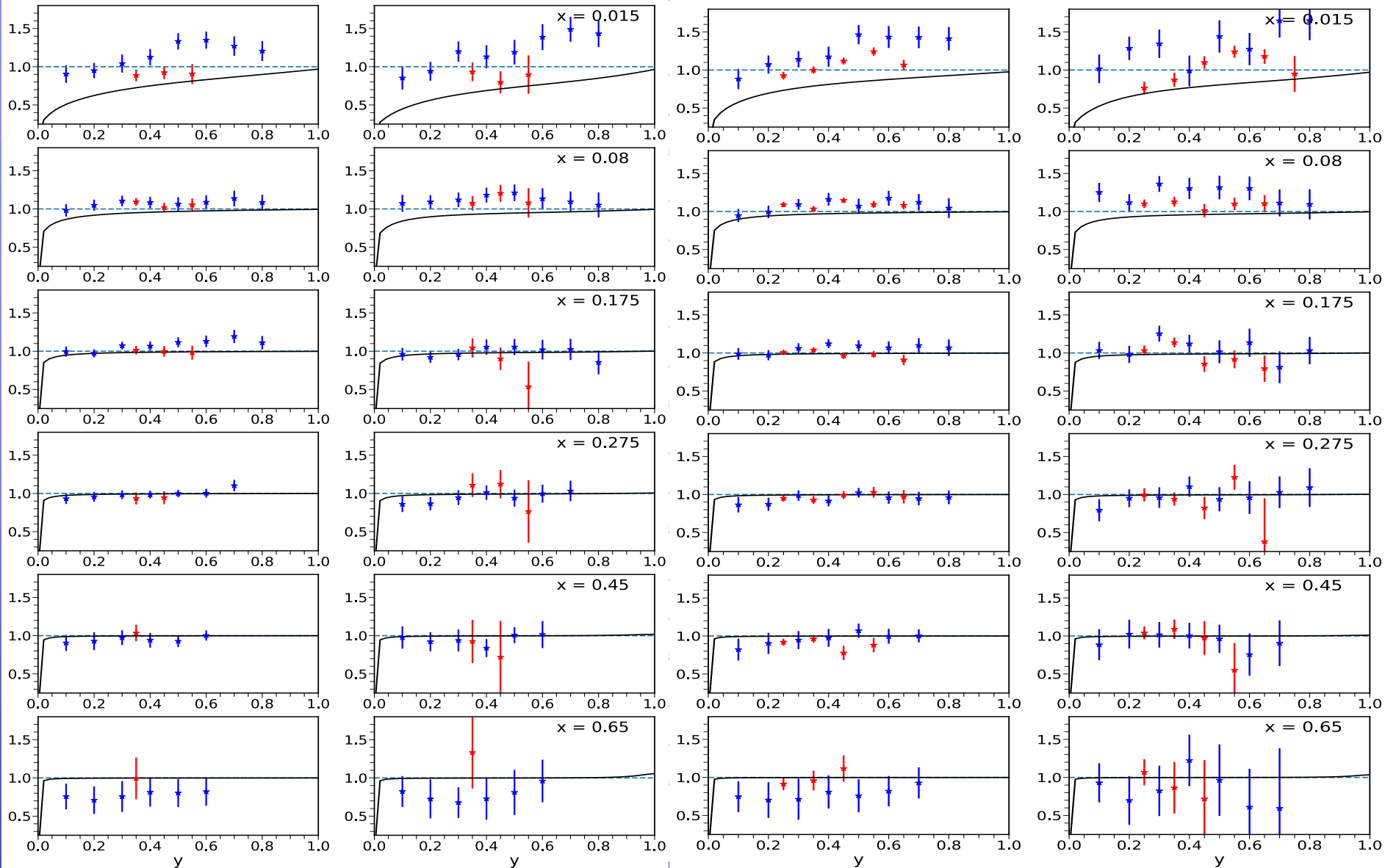
$$2xF_1^{axial} = 2xF_1^{vector}$$

Comparison with CCFR (Fe), CHORUS (Pb) data

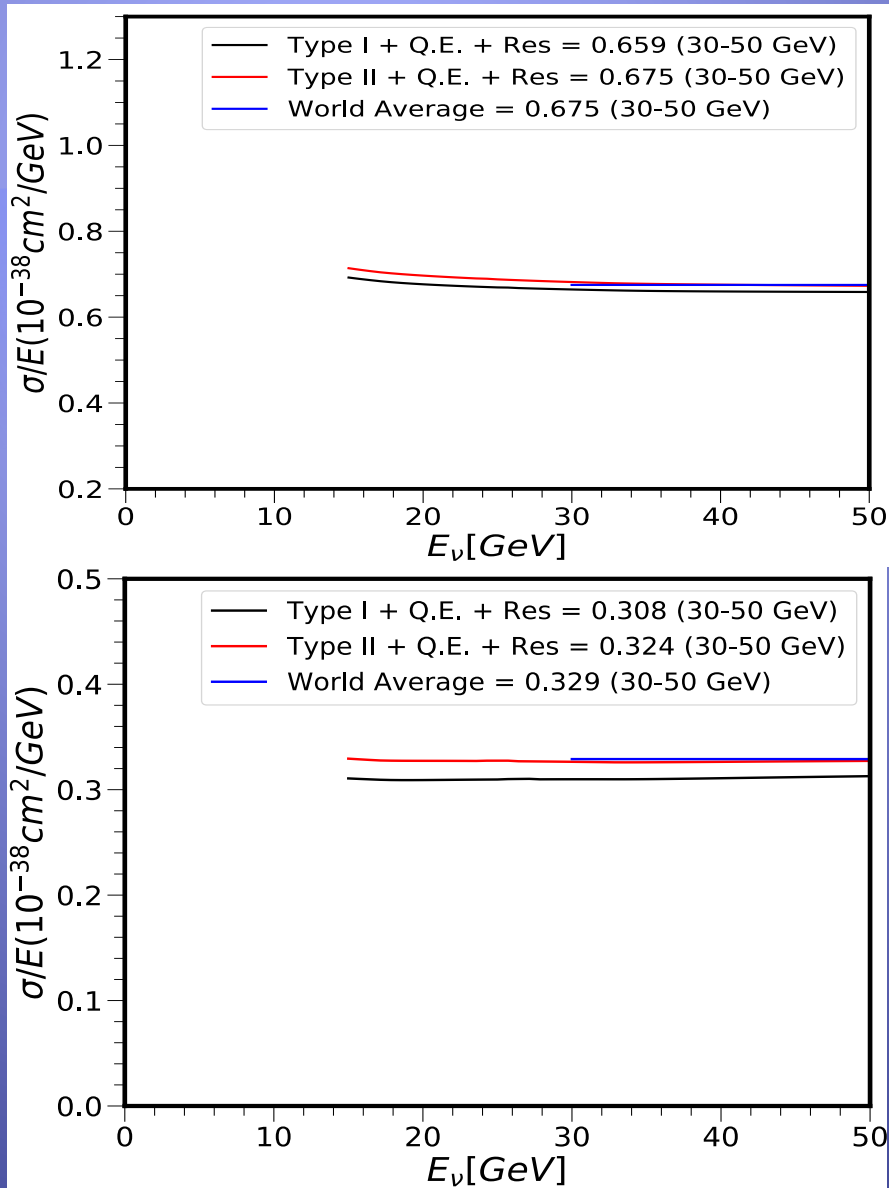


- Blue point: CHORUS/theory (type II)
- Solid line: theory (type I)/(type II)
- Type I (Vector = Axial at low Q^2)
- Type II (Vector < Axial at low Q^2)

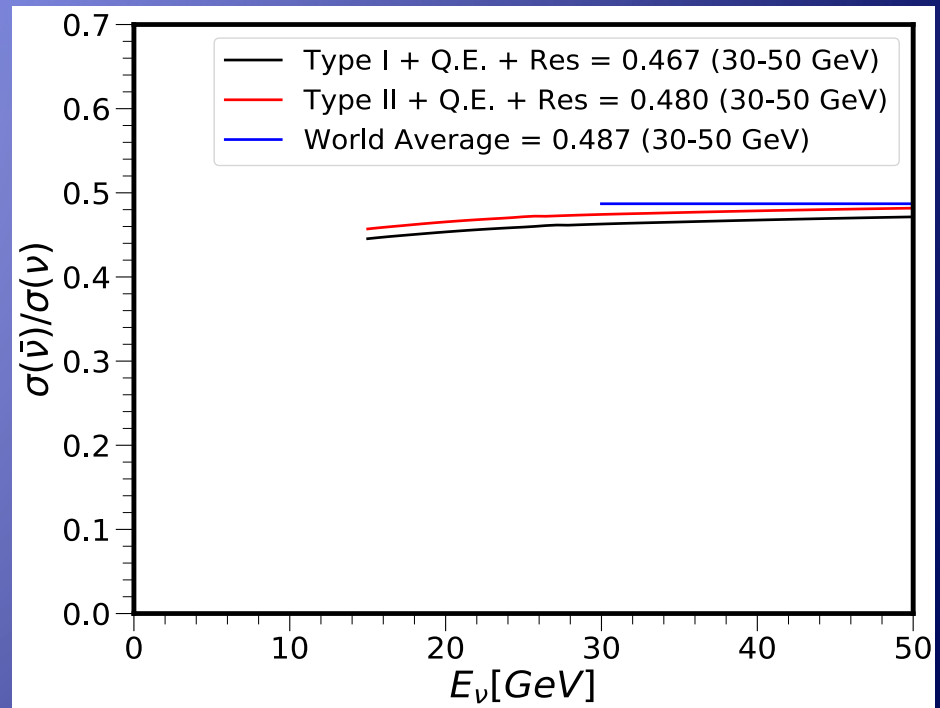
Comparison with CCFR(Fe), CHORUS (Pb) data

 $E_\nu = 35.0$ $E_\nu = 55.0$ 

Total cross sections



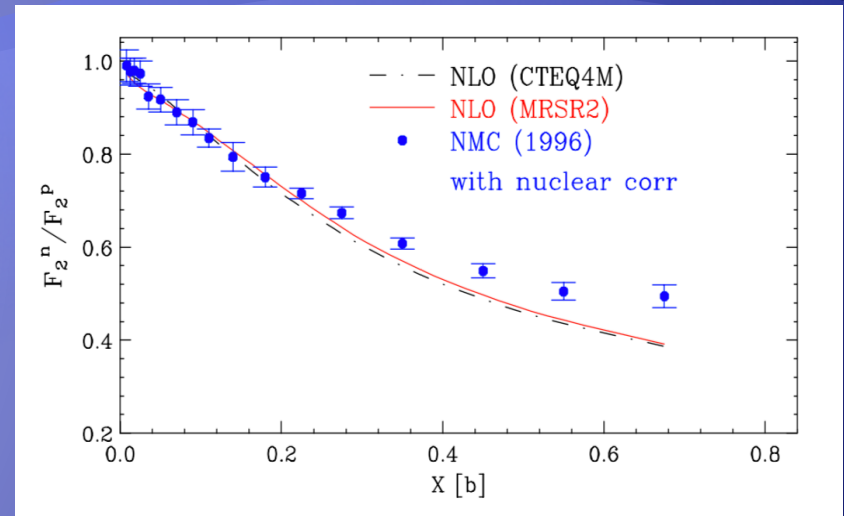
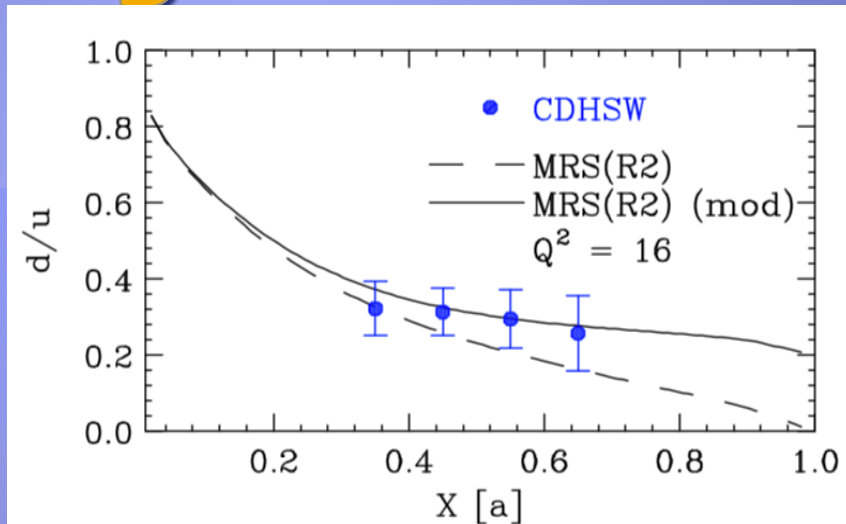
➤ BY(DIS, $W > 1.4$)
+ Q.E. + Resonance



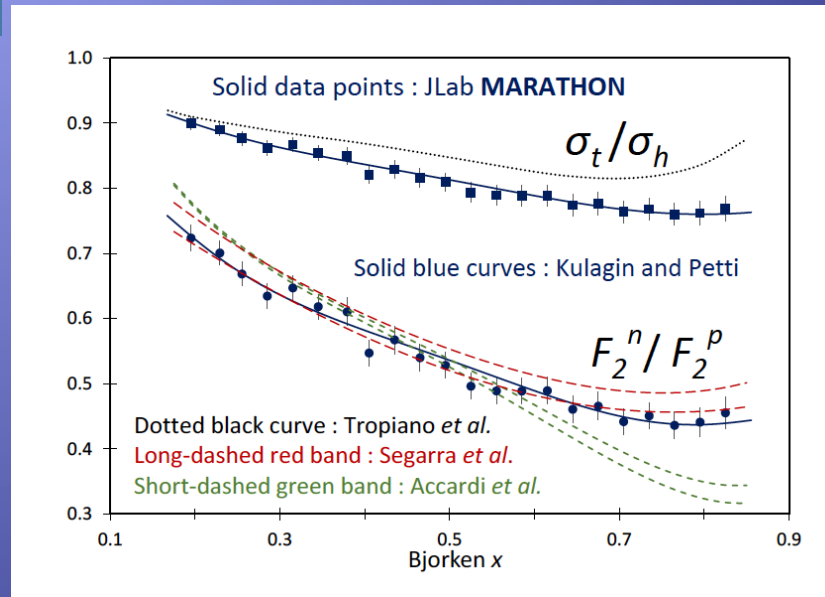
Summary & Discussions

- BY Effective LO model with ξw describe all e/ μ DIS and resonance data as well as photo-production data (down to $Q^2=0$): provide a good reference for vector SF for neutrino cross section
- $d\sigma/dx dy$ data favor updated BY(DIS) type II model
- BY(DIS) type II model (low Q^2 : axial > vector) provide a good reference for neutrino cross sections. Low energy neutrino experiments can normalize their data to our model to extract their flux
- Model also works well down to $W=1.4$ GeV, thus providing overlap with resonance models
- Future improvement: use very high-x data (nCTEQ effort)

High-x PDF



PRL 82, 2467 (1999)



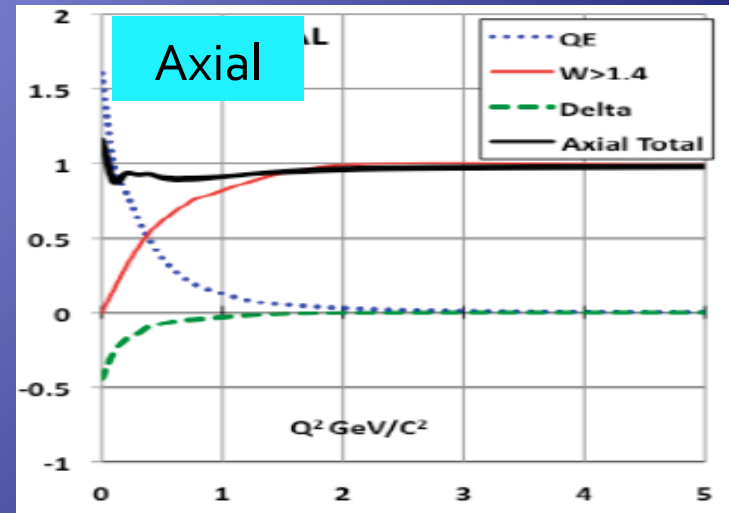
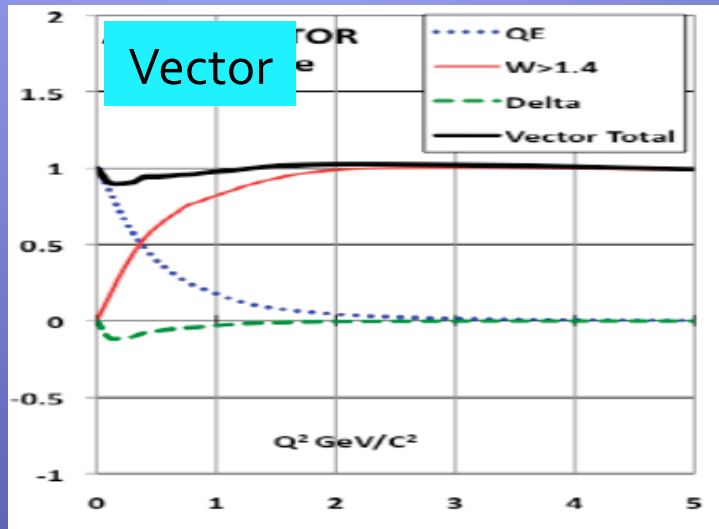
PRL 128, 132003
(2021)

Test of the Adler Sum Rule

- This sum rule should be valid at all values of Q^2

$$|F_V(Q^2)|^2 + \int_{\nu_0}^{\infty} W_{2n-sc}^{\nu-vector}(\nu, Q^2) d\nu - \int_{\nu_0}^{\infty} W_{2p-sc}^{\nu-vector}(\nu, Q^2) d\nu = 1$$

$$|F_A(Q^2)|^2 + \int_{\nu_0}^{\infty} W_{2n-sc}^{\nu-axial}(\nu, Q^2) d\nu - \int_{\nu_0}^{\infty} W_{2p-sc}^{\nu-axial}(\nu, Q^2) d\nu = 1$$



Nuclear Effects: use e/ μ data

