

# Resummation in artemide

LHC EW precision sub-group meeting (pT W/Z benchmarking)

Alexey Vladimirov

Regensburg University



Universität Regensburg

## ARTEMIDE

is made to explore **TMD factorization** and **TMD distributions**.

### Specific points

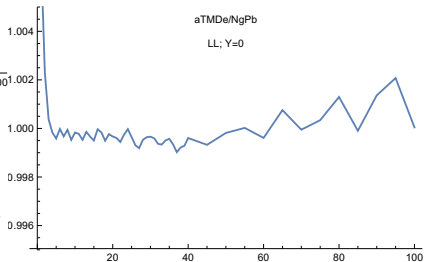
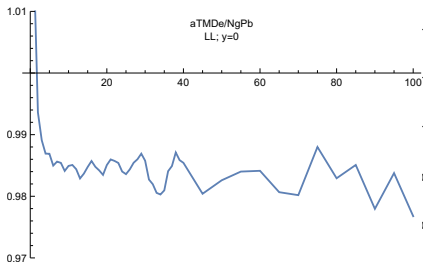
- ▶ Variety of principle processes: DY, SIDIS, SIDIS→jet (*ee*-annihilation to be added)
- ▶ Variety of sub-subprocesses: targets, polarizations, etc. (constantly updating list)
- ▶ Special emphasis of non-perturbative structure → full flexibility to define a model ( $b^*$  different kind of locality,  $f_{NP}$ 's of arbitrary form, RAD of arbitrary form, etc.)
- ▶ Modular structure
- ▶ Perturbative parts up to best known
- ▶  $\zeta$ -prescription as a main universality concept

### artemide is also exploration code (all switchable)

- ▶ Different resummation versions
- ▶ Power corrections: target mass corrections, kinematic corrections
- ▶ Large-x resummation (in progress)
- ▶ PDF variations
- ▶ ...

## On earlier discrepancies

- ▶ Mistake: wrong type of resummation has been used (see next slides)
- ▶ Bug: There was a bug in the code for  $\alpha_{\text{QED}}$  (incorrect determination of boundary condition  $\rightarrow -1.5\%$  of common normalization) **FIXED!**



## Resummation formalisms

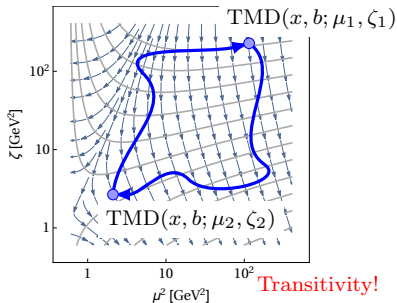
$$\begin{array}{c}
 q_T\text{-res.} \\
 \propto \\
 \text{PB} \\
 e^{2S} [f_1 \otimes \mathcal{H} \otimes f_2] \\
 \\
 \left( \frac{d\sigma}{dq_T} \right)_{\text{res.}} \quad \left( \overset{\text{TMD}}{\propto} H \times F_1 \times F_2 \right) + \mathcal{O} \left[ \left( \frac{q_T}{Q} \right)^m \right] \\
 \\
 \overset{\text{SCET}}{\propto} H \times B_1 \times B_2 \times S \\
 \\
 \left( H \times F_1 \times R \times F_2 \right) \\
 \uparrow \quad \uparrow \quad \uparrow \\
 \text{three non-perturbative functions}
 \end{array}$$

- ▶ Each non-perturbative function has its own definition in terms of operators
- ▶ Each non-perturbative function responsible for different degree of freedom

$$\mu^2 \frac{dF(x, b; \mu, \zeta)}{d\mu^2} = \frac{\gamma_F(\mu, \zeta)}{2} F(x, b; \mu, \zeta),$$

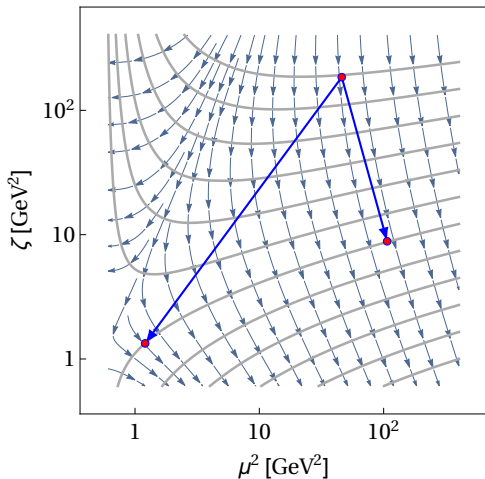
$$\zeta \frac{dF(x, b; \mu, \zeta)}{d\zeta} = -\mathcal{D}(\mu, b) F(x, b; \mu, \zeta)$$

## 2D evolution



$$R[(\mu_1, \zeta_1) \rightarrow (\mu_2, \zeta_2)] = \exp \int_P \left( \frac{d\mu}{\mu} \gamma_F - \frac{d\zeta}{\zeta} \mathcal{D} \right)$$

TMD distribution is not defined by a scale  $(\mu, \zeta)$   
It is defined by an equipotential line.

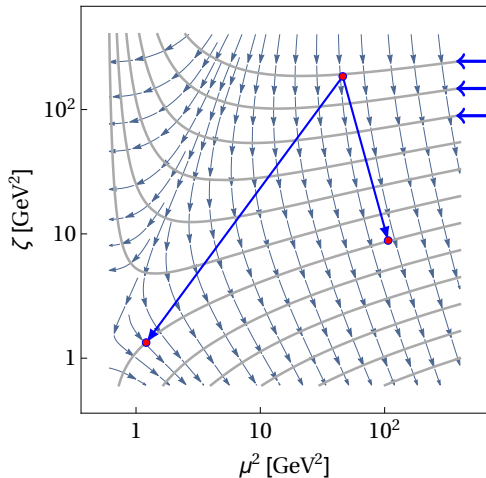


The scaling is defined by  
~~a difference between scales~~  
a difference between potentials

Evolution factor to both points  
is the same  
although the scales are  
different by  $10^2 \text{GeV}^2$



TMD distributions on the same equipotential line are equivalent.



TMD( $x, b, 1$ )

TMD( $x, b, 2$ )

TMD( $x, b, 3$ )

We can enumerate them by a lines  
not by  $(\mu, \zeta)$

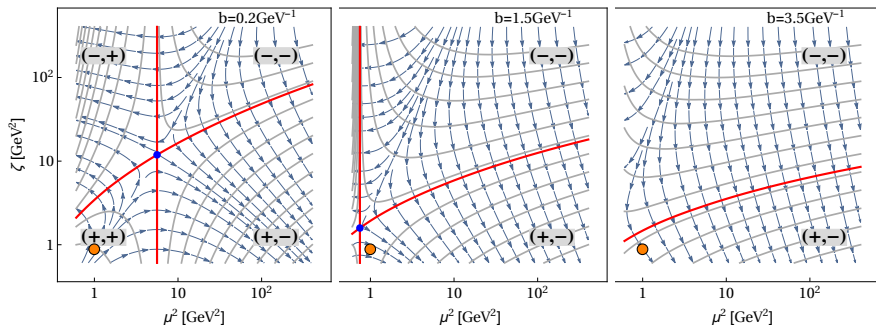
$$F(x, b; \mu, \zeta) \rightarrow F(z, b; \text{line})$$



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## Boundary condition problem

How to select same equipotential line for each values of  $b$ ?



### Standard solution: **fix scale**

If one set  $(\mu^2, \zeta) = (1, 1)\text{GeV}^2$  (f.i.);

$\Rightarrow$  this point belongs to different equipotentials at different values of  $b$

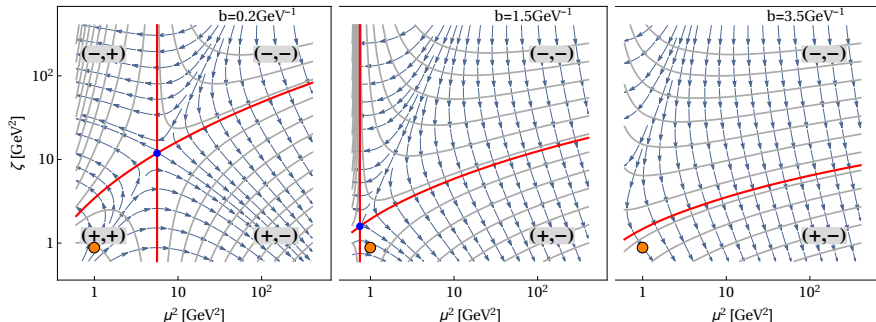
$\Rightarrow F(x, b; 1, 1)$  contains admixture of  $\mathcal{D}$  at different  $b$

another example CS-scheme  $\mu = \zeta = 1/b^*(b)$



## Boundary condition problem

How to select same equipotential line for each values of  $b$ ?



Better solution: **fix line**

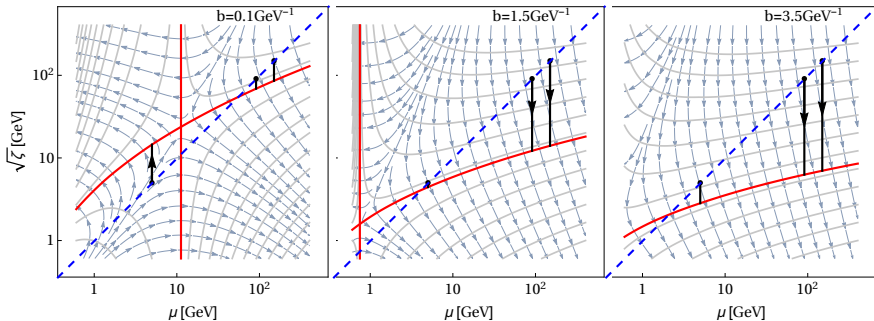
Select an equipotential line and follow it at all values of  $b$

**Problem:** Lines are discontinuous and difficult to define

**Solution:** Take the line that passes through saddle point (unique, continues, simple)

If the equipotential line is know evolution is simple

$$F(x, b; \mu, \zeta) = \left( \frac{\zeta}{\zeta_{\mu}(b)} \right)^{-\mathcal{D}(b, \mu)} F(x, b)$$



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## How to find equipotential line?

- ▶ Let  $\zeta_\mu(b) = \zeta(b, \mu)$  be equipotential line

$$\Gamma_{\text{cusp}}(\mu) \ln \left( \frac{\mu^2}{\zeta_\mu(b)} \right) - \gamma_V(\mu) = 2\mathcal{D}(\mu, b) \frac{d \ln \zeta_\mu}{d \ln \mu^2}$$

- ▶ Special boundary condition

$$\zeta_{\mu_0}(b) = \zeta_0, \quad \text{where} \quad \mathcal{D}(b, \mu_0) = 0, \quad \gamma_F(\mu_0, \zeta_0) = 0.$$

Does it help? **Not really...**

- ▶ The saddle point goes below  $\Lambda_{QCD}$  are large  $b$
- ▶  $\mathcal{D}$  is generically unknown  $\Rightarrow$  for each version of  $\mathcal{D}$  solve SP, and solve equation  $\Rightarrow$  numerically expensive



Exact solution for arbitrary  $\mathcal{D}$   
 $(\mu, \zeta) \rightarrow (a_s, \mathcal{D})$

- ▶ Let  $\zeta_\mu(b) = \zeta(b, \mu)$  be equipotential line

$$\zeta_\mu(b) = \mu^2 \exp\left(-\frac{g(a_s, \mathcal{D})}{\mathcal{D}}\right)$$

$$2\mathcal{D} + 2\beta(a_s) \frac{\partial g(a_s, \mathcal{D})}{\partial a_s} - \Gamma_{\text{cusp}}(a_s) \frac{\partial g(a_s, \mathcal{D})}{\partial \mathcal{D}} + \gamma_V(a_s) = 0$$

- ▶ Special boundary condition

$$g(a_s, 0) = 0$$

Does it help? **Extremely!!**

- ▶ Has **solution for arbitrary values of  $\mathcal{D}$**
- ▶ Does not care about Landau pole (automatic analytic continuation)
- ▶ Can be solved exactly (if needed)
- ▶ Has **very good perturbative convergence** since  $(\mu = Q)$  at arbitrary  $b$
- ▶ Very fast numerically (just algebraic function)

Exact solution for arbitrary  $\mathcal{D}$   
 $(\mu, \zeta) \rightarrow (a_s, \mathcal{D})$

- ▶ Let  $\zeta_\mu(b) = \zeta(b, \mu)$  be equipotential line

$$\zeta_\mu(b) = \mu^2 \exp\left(-\frac{g(a_s, \mathcal{D})}{\mathcal{D}}\right)$$

- ▶ LO+NLO ( $p = 2\beta_0\mathcal{D}/\Gamma_0$ )

$$g(a_s, \mathcal{D}) = \frac{2\beta_0^2}{a_s\Gamma_0} \left[ (e^{-p} - 1 + p) + a_s \frac{\beta_1}{\beta_0} \left( e^{-p} - 1 + p - \frac{p^2}{2} + \dots \right) \right]$$

see 1912.06532 appendix C, for all formulas up to N<sup>3</sup>LO

Does it help? **Extremely!!**

- ▶ Has **solution for arbitrary values of  $\mathcal{D}$**
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- ▶ Has **very good perturbative convergence** since ( $\mu = Q$ ) at arbitrary  $b$
- ▶ Very fast numerically (just algebraic function)

In perturbative regime reproduces resummation (up to a given order in  $a_s$ )

$$\begin{aligned} F(x, b) &= C(x, b; \mu_{OPE}) \otimes f(x, \mu_{OPE}) \\ &= [\delta(\bar{x}) + a_s(P(x)\mathbf{L}_\mu + 2\bar{x} - \zeta_2\delta(\bar{x})) + \dots] \otimes f(x, \mu_{OPE}) \end{aligned}$$

- ▶ Different expression for perturbative matching (**no double logs!**) (effect  $x$ -dependence at NNLO)
- ▶ Free parameter  $\mu_{OPE}$ , independent on the scales of process

$$\mathcal{D}(b, \mu) = a_s \frac{\Gamma_0}{2} \mathbf{L}_\mu + a_s^2(\dots) + \dots$$

$$\zeta_\mu(b) = C_0 \frac{\mu}{b} \exp\left(\frac{\gamma_1}{\Gamma_0} + a_s^2(\dots) + \dots\right)$$

**Collect together reexpand get the same expression.**



In *artemide* a TMD is defined on special equipotential line

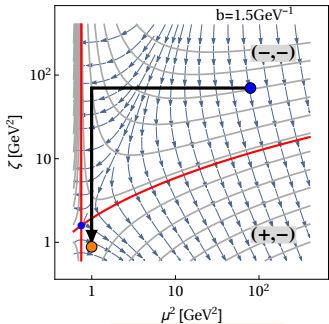
$$d\sigma \simeq H(Q) \int db e^{ibq_T} \left( \frac{Q^2}{\zeta_Q(b)[\mathcal{D}]} \right)^{-2\mathcal{D}(Q,b)} F_1(x_1, b) F_2(x_2, b)$$

- ▶ Three non-perturbative functions; each responsible for single variable (and  $b$ )
- ▶ In comparison to “standard CSS” matching some terms are exponentiated (whose that are related to  $\mathcal{D}$ )
- ▶ It gives different numerics at large- $b$ /small- $Q$  (even if evolution is taken the same)



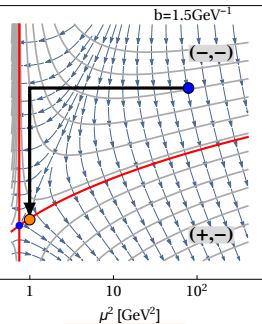
$$\mu_i^2 = \zeta_i = Q^2$$

canonic CSS  
NangaParbat



$$\mu_f^2 = \zeta_f \simeq \frac{C_0}{b^*(b)}$$

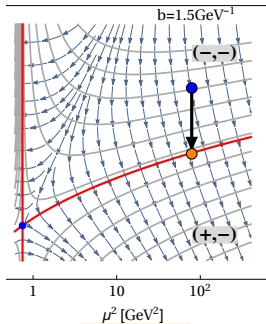
artemide  
CSS-like



$$\mu_f^2 \simeq \frac{C_0}{b^*(b)}$$

$$\zeta = \zeta_{\mu_f}(b)$$

artemide  
default



$$\mu_f^2 = Q$$

$$\zeta = \zeta_Q(b)$$

$$R = \exp \left[ \int \frac{d\mu}{\mu} \left( \Gamma \ln \left( \frac{\mu^2}{\zeta_f} \right) - \gamma_V \right) \right] \left( \frac{\zeta_f}{\zeta_i} \right)^{-D}$$

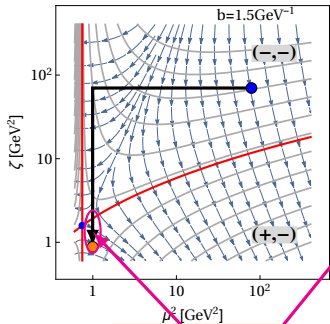
$$R = \left( \frac{\zeta_f}{\zeta_{\mu}(b)} \right)^{-D}$$





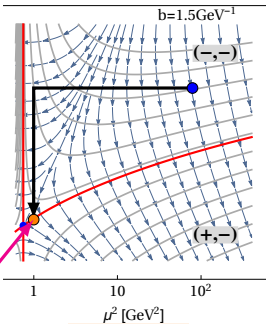
$$\mu_i^2 = \zeta_i = Q^2$$

canonic CSS  
NangaParbat



$$\mu_f^2 = \zeta_f \simeq \frac{C_0}{b^*(b)}$$

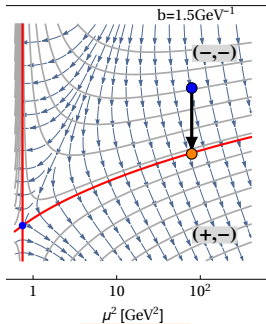
artemide  
CSS-like



$$\mu_f^2 \simeq \frac{C_0}{b^*(b)}$$

$$\zeta = \zeta_{\mu_e}(b)$$

artemide  
default



$$\mu_f^2 = Q$$

$$\zeta = \zeta_Q(b)$$

$$\Delta = \left( \frac{\zeta_{\mu_i}(b)b^*}{C_0} \right)^{-D}$$

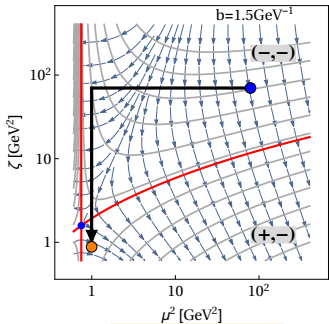
Part of matching coefficient  
leads to different definition of “unprimed” orders  
( $\Delta \equiv 1$  at LL)

$$R = \left( \frac{\zeta_f}{\zeta_{\mu}(b)} \right)^{-D}$$



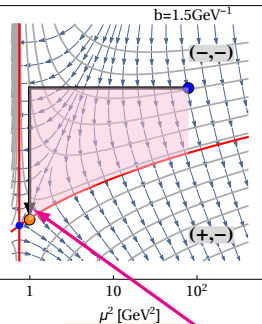
$$\mu_i^2 = \zeta_i = Q^2$$

canonic CSS  
NangaParbat



$$\mu_f^2 = \zeta_f \simeq \frac{C_0}{b^*(b)}$$

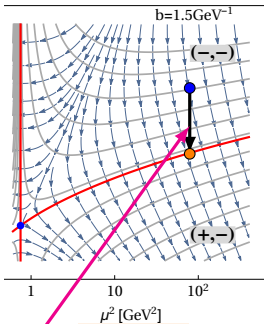
artemide  
CSS-like



$$\mu_f^2 \simeq \frac{C_0}{b^*(b)}$$

$$\zeta = \zeta_{u,c}(b)$$

artemide  
default



$$\mu_f^2 = Q$$

$$\zeta = \zeta_Q(b)$$

$$R = \exp \left[ \int \frac{d\mu}{\mu} \left( \Gamma \ln \left( \frac{\mu^2}{\zeta_f} \right) - \gamma_V \right) \right] \left( \frac{\zeta_f}{\zeta_i} \right)$$

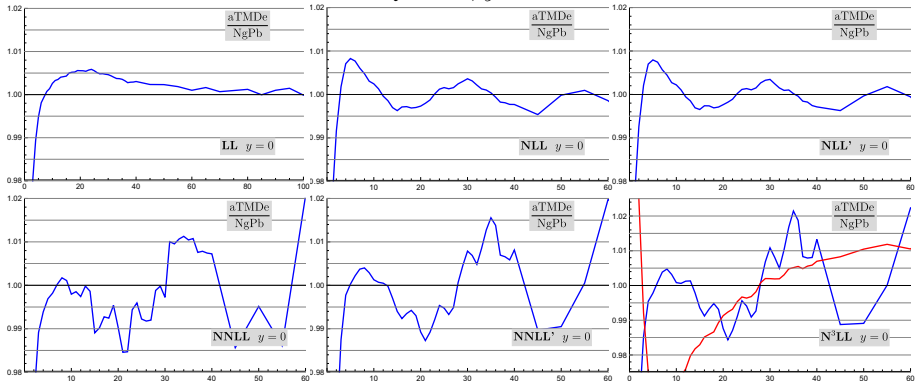
$$\Delta = \exp \int_C \left( \frac{d\mu}{\mu} \gamma_F - \frac{d\zeta}{\zeta} \mathcal{D} \right)$$

$$\sim a_s^{n+1} \times \text{area}$$

decrease with the increase of order



$$Q = M_Z, y = 0$$



- ▶  $\alpha_s$  and  $N_f$ -scheme is different (artemide=(LHA,variable), NangaParbat=(self, $N_f = 5$ ))
- ▶ CSS-like code in artemide is not well-optimized (noisy)
- ▶  $f_{NP} = 1$  is badly convergent (noise at large- $q_T$ )



$$Q = M_Z, y = 0$$

