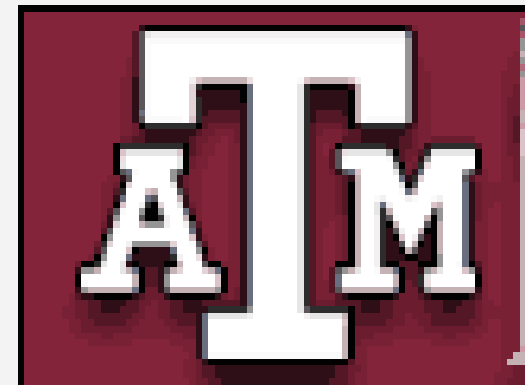


Open and Hidden

Heavy-Flavor Transport in Medium

Ralf Rapp

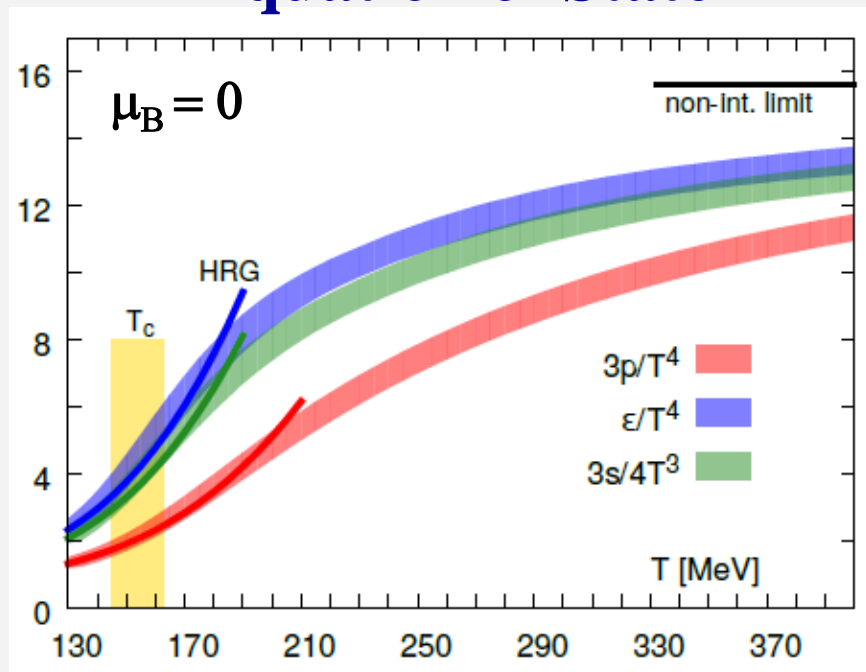
**Cyclotron Institute +
Dept of Phys & Astro
Texas A&M University
College Station, TX
USA**



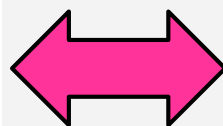
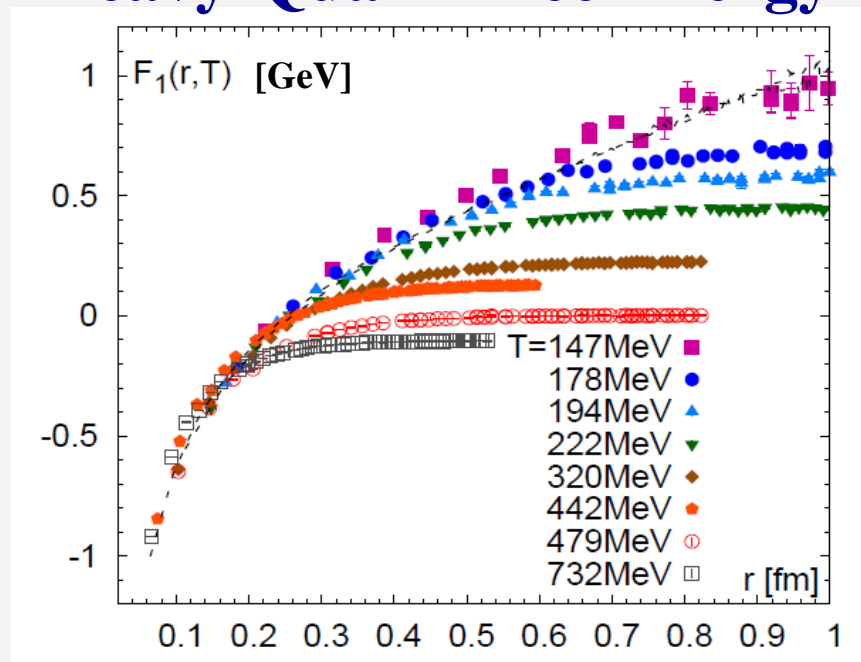
**8th International Workshop on
Heavy-Flavour Production in Nuclear Collisions
Torino (Italy), July 14-16, 2022**

1.) QCD Matter and In-Medium Force

Equation of State



Heavy-Quark Free Energy



[Bazavov et al '13,'14]

- Change in dofs above $T \sim 170\text{MeV}$

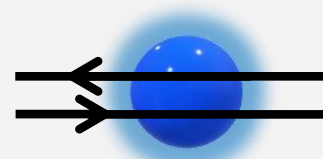
$$\bullet F_{Q\bar{Q}} = U_{Q\bar{Q}} - T S_{Q\bar{Q}}$$

- Non-perturbative above T_c

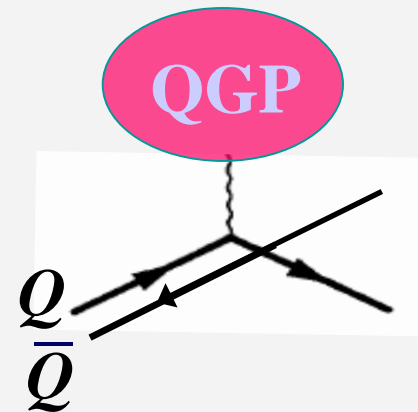
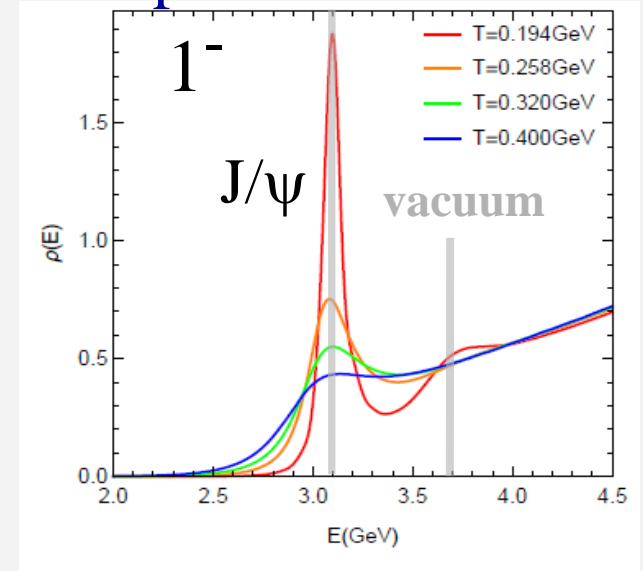
- Heavy-quark diffusion and quarkonia sensitive to long-range forces in medium, probing non-perturbative QGP structure

1.2 Quarkonia in Medium

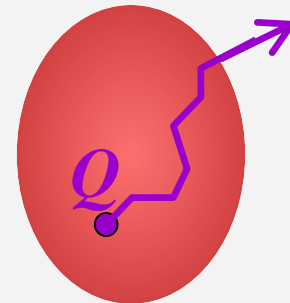
- In-medium spectral functions:
 - Mass / binding energy $E_B(p, T)$
 - Inelastic reaction rate $\Gamma(p, T; E_B)$
- $\Upsilon(1S)$: color-Coulomb force
 J/ψ , $\Upsilon(2S)$, ...: confining force
- Connection to open heavy-flavor transport



Spectral Function



1.3 Heavy-Flavor Diffusion



- **Brownian motion via elastic interactions**
(radiation suppressed: $q_0 \sim q^2/2m_Q \ll q \sim T \ll m_Q$)
- Thermalization **delayed** by $m_Q/T \Rightarrow$ measure of **interaction strength**
- **Scattering rates:** collisional width \leftrightarrow quantum effects
 \Rightarrow **Implications for QGP structure**
- **Non-perturbative interactions**
- **Hadronization:** $c \rightarrow D, D_s, \Lambda_c, \dots$ ($m_Q \gg T_c$)



Outline

1.) Introduction

2.) Open Heavy-Flavor Transport

- **Transport Equations + Heavy-Light Interaction**
- **Implications for Bulk Medium**

3.) Quarkonia in Medium

- **Production Models + Transport Coefficients**
- **Connection to Open Heavy Flavor**

4.) Conclusions

2.1 Transport Equations

- Boltzmann equation for HQ phase-space distribution f_Q

$$\left[\frac{\partial}{\partial t} + \frac{\mathbf{p}}{\omega_p} \frac{\partial}{\partial \mathbf{x}} + \mathbf{F} \frac{\partial}{\partial \mathbf{p}} \right] f_Q(t, \mathbf{x}, \mathbf{p}) = C[f_Q]$$

- Collision term: semiclassical simulation of medium + HQ quasiparticles

- Fokker-Planck equation

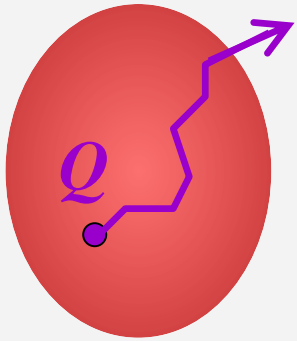
$$\frac{\partial}{\partial t} f_Q(t, \mathbf{p}) = \frac{\partial}{\partial p_i} \left\{ A_i(\mathbf{p}) f_Q(t, \mathbf{p}) + \frac{\partial}{\partial p_j} [B_{ij}(\mathbf{p}) f_Q(t, \mathbf{p})] \right\}$$

- follows from Boltzmann with $\mathbf{p}^2 \sim m_Q T \gg \mathbf{q}^2 \sim T^2$ (ok for $m_Q/T \geq 5$)

- does not require quasi-particle medium

→ well suited for **strongly coupled medium** where $E_{\text{th}} \leq \Gamma_{q,Q} < m_Q$

2.2 Transport Coefficients



$$\frac{\partial}{\partial t} f_Q(t, p) = \gamma \frac{\partial}{\partial p_i} [p_i f_Q(t, p)] + D_p \Delta_{\vec{p}} f_Q(t, p)$$

Fokker-Planck

thermalization rate

$$\begin{aligned} \gamma p &= \int d^3 q w_Q(q, p) q \\ &\sim \int |T_{Qi}|^2 (1 - \cos \theta) f^i \rho_i \rho_i' \end{aligned}$$

momentum diffusion coefficient

$$D_p = \int d^3 q w_Q(q, p) q^2$$

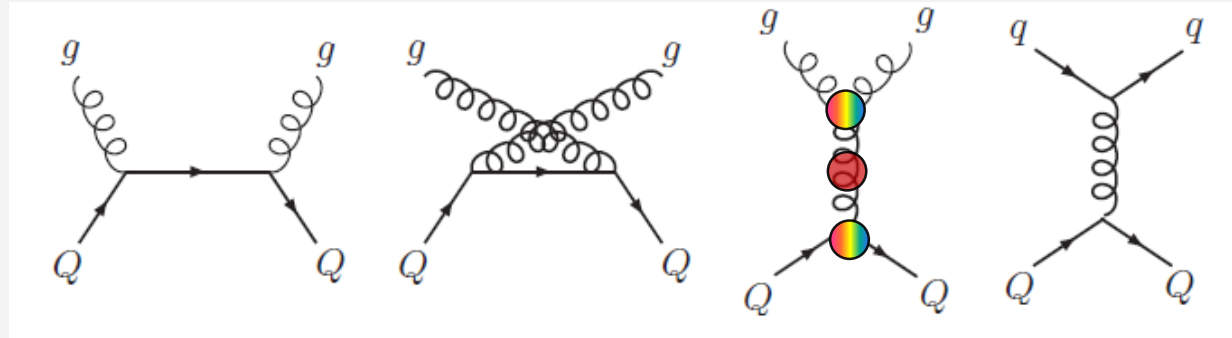
- Relaxation time $\tau_Q = 1/\gamma$
- **Spatial diffusion constant:** $\mathcal{D}_s = T / \gamma(p=0) m_Q$
- Quantum effects via thermal-parton **spectral functions** $\rho_i(k_0, k)$
- Key ingredients: - **heavy-light scattering amplitude** T_{Qi}
- **ambient medium properties**

2.3.1 Heavy-Light Interaction: Perturbative QCD

• Leading Order

- screened 1-gluon exchange

$$T_{Qi} \sim \frac{\alpha_s}{t - m_D^2}$$



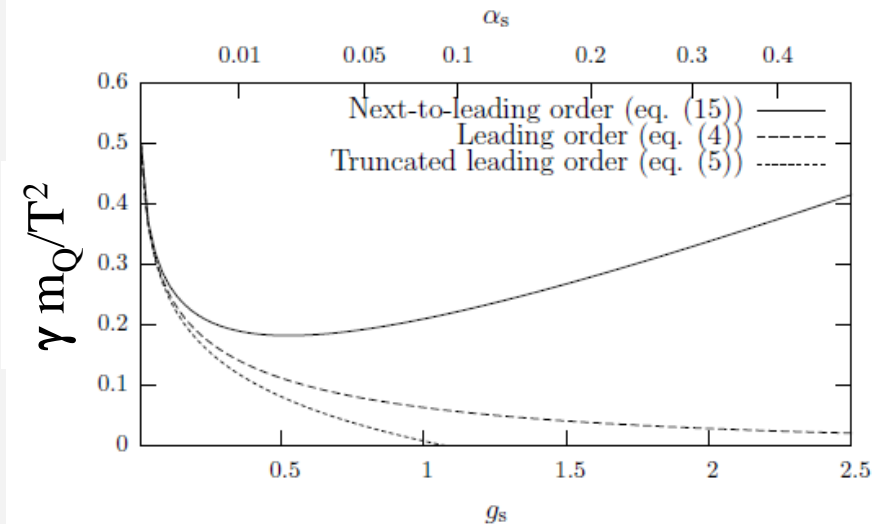
- $\tau_c \geq 20 \text{ fm/c}$ ($T \leq 300 \text{ MeV}$, $\alpha_s = 0.4$)

[Svetitsky '88, Mustafa et al '98, Molnar et al '04, AMPT, Teaney+Moore'04, Torino, Catania, PHSD, BAMPS, LBT,...]

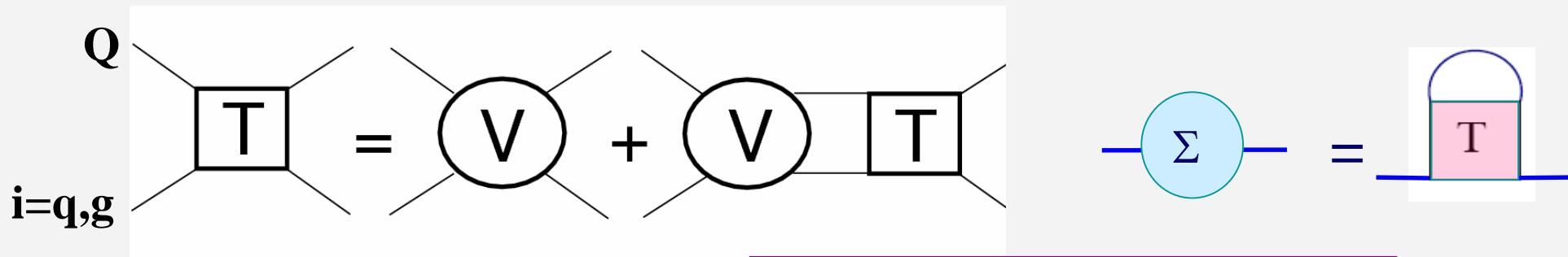
- pQCD*: running coupling with reduced Debye mass [Nantes '08]

• Next-to-Leading Order

- bad convergence (even for $\alpha_s \sim 0.1$)
- uncontrolled result [Caron-Huot+Moore '08]

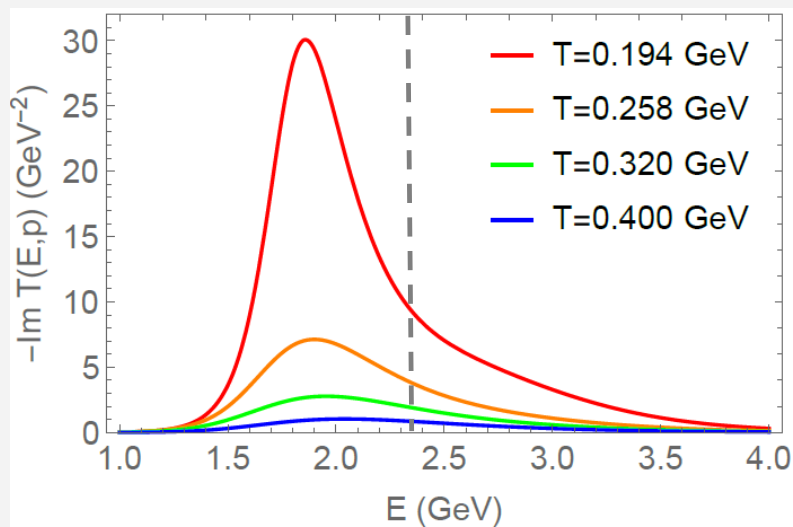
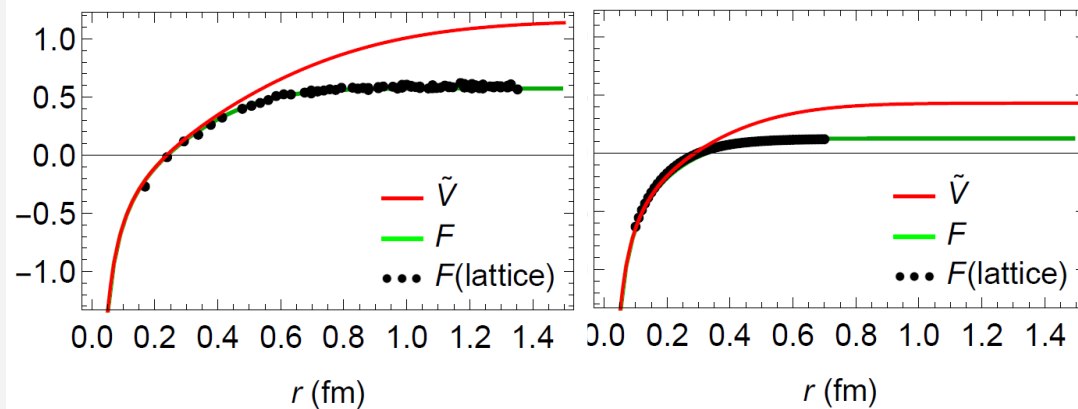


2.3.2 Heavy-Light Interactions: **T**-Matrix



- **Strong coupling** \rightarrow resummation:
- **Key input:** in-med potential V_{qi} , constrained by lattice QCD
 \rightarrow “string” force above T_c

$$T_{qi} = V_{qi} + \int V_{qi} D_Q D_i T_{qi}$$



- **Non-perturbative effects:**
 - thermalization time $\tau_c \approx 3-4 \text{ fm}/c$
 - **hadronization**

2.4 QGP Medium Models: Light Partons

→ essential for quantitative evaluation of transport coefficients

- Massless quarks + gluons (EoS too large)
- Perturbative quasiparticles ($m_{q,g} \sim gT$, EoS for $T > 1.5T_c$)
- Effective quasiparticles ($m_{q,g} \sim gT$ with g large near T_c)
- Dynamical quasiparticles ($\Gamma_{g,q} \sim 0.1-0.2 \text{ GeV}$)
- **T**-matrix + string potential ($\Gamma_{g,q} \sim 0.5 \text{ GeV}$)

- Consistency of heavy-light interaction and EoS?

2.5 Heavy-Quark Recombination

- **Instantaneous Coalescence Model (ICM)**

[Hwa '80, Likhoded et al '83, ...
Greco et al + Fries et al '03,...]

$$f_h(\mathbf{p}'_h) = \int \left[\prod_i d\mathbf{p}_i f_i(\mathbf{p}_i) \right] W(\{\mathbf{p}_i\}) \delta(\mathbf{p}'_h - \sum_i \mathbf{p}_i)$$

$$W_s = g_h \frac{(2\sqrt{\pi}\sigma)^3}{V} e^{-\sigma^2 \mathbf{k}^2} \quad \sigma \sim \text{radius of hadron } \mathbf{h}$$

- energy not conserved → challenge for chemical + thermal equilibrium

- **Resonance Recombination Model (RRM)**

[Ravagli et al '07, He et al '12]

- derived from Boltzmann equation

$$f_M(\vec{x}, \vec{p}) = \frac{\gamma_M(p)}{\Gamma_M} \int \frac{d^3\vec{p}_1 d^3\vec{p}_2}{(2\pi)^3} f_q(\vec{x}, \vec{p}_1) f_{\bar{q}}(\vec{x}, \vec{p}_2) \sigma_M(s) v_{\text{rel}}(\vec{p}_1, \vec{p}_2) \delta^3(\vec{p} - \vec{p}_1 - \vec{p}_2)$$

- $\sigma_M(s) v_{\text{rel}} \sim |T_{qi}|^2$: resonant heavy-light amplitude

→ compatible with **equilibrium limits**

Outline

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4.) Conclusions

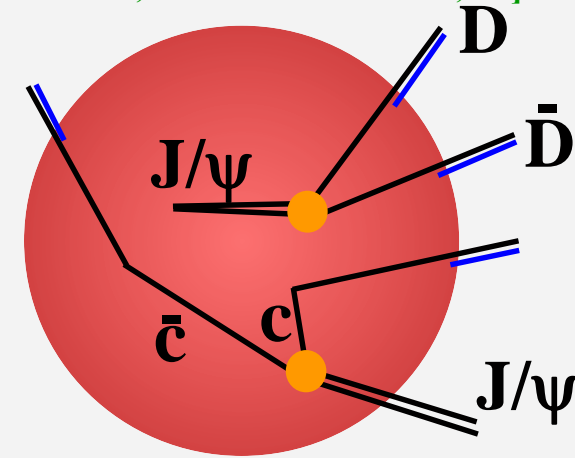
3.1 Quarkonia Transport in Heavy-Ion Collisions

[PBM+Stachel '00, Thews et al '01, Grandchamp+RR '01, Gorenstein et al '02, Ko et al '02, Gossiaux et al '05, Zhuang et al '05, Ferreiro et al '11, Strickland et al '11, ...]

• Semi-classical

Boltzmann eq. $\hat{p}^\mu \partial_\mu f^\Psi = -\Gamma_\Psi f^\Psi + \beta$

→ Rate eq.
$$\frac{dN_\Psi}{dt} = -\Gamma_\Psi (N_\Psi - N_\Psi^{eq})$$



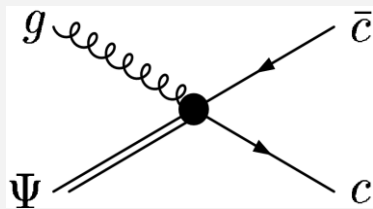
• Transport coefficients

- Equilibrium limit: $N_\Psi^{eq}(E_B, T; N_{cc}) = \gamma_c^2 n_\Psi(T) V_{FB}$

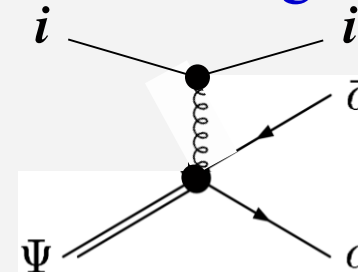
depends on **cc** Xsec + (thermal) environment: $\sigma_{cc} \sim N_{cc} \sim \gamma_c n_c(m_c, T) V_{FB}$

- Reaction rate: $\Gamma_\Psi[p, E_B, T]$

“Large” binding $E_B \geq T$



“Small” binding $E_B \leq m_D$



• gluo-dissociation (“singlet-to-octet”)

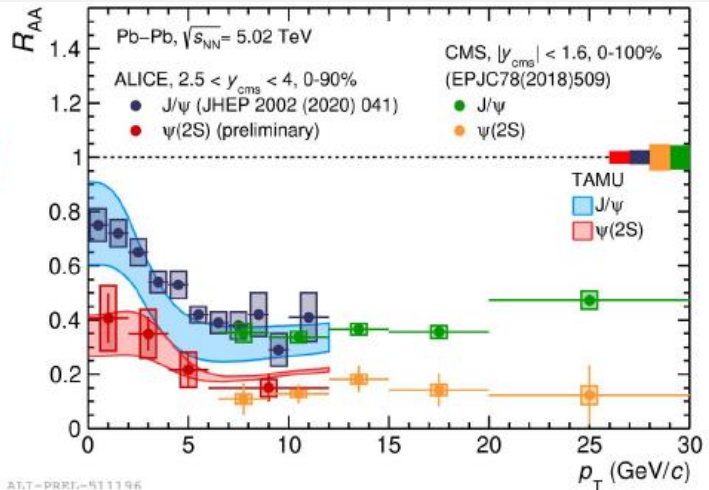
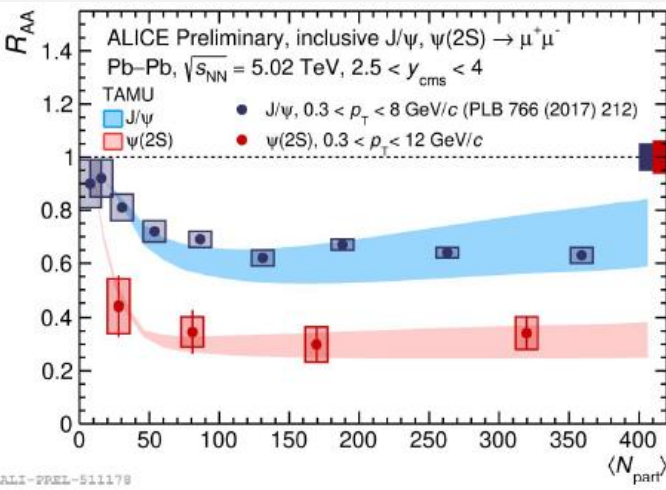
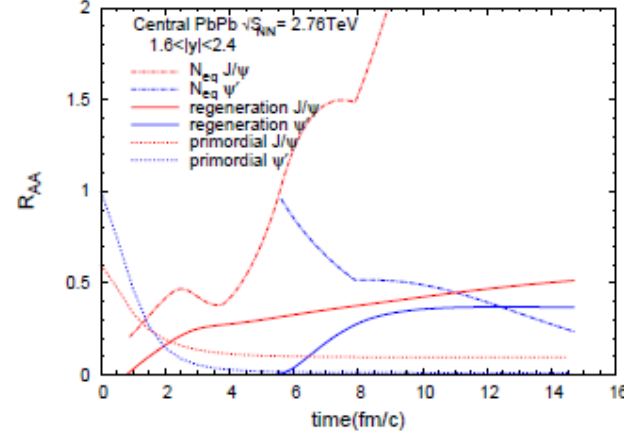
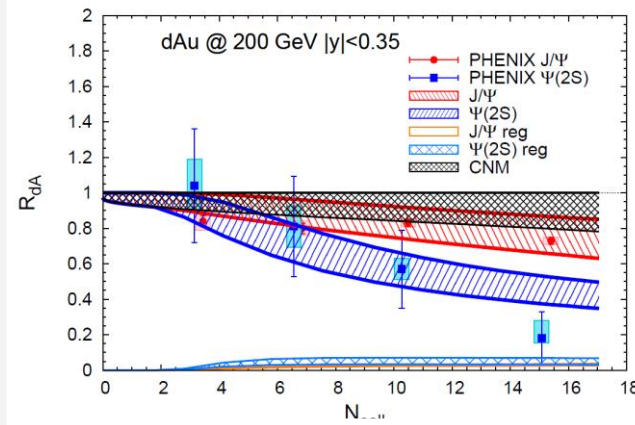
• “quasi-free”/ Landau damping

3.2 Interplay of Transport Coefficients: Charmonia

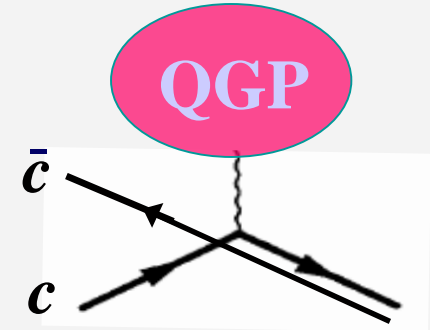
- Phenomenological constraints on Γ_ψ from **d/p-A**:
J/ψ suppressed in QGP, **ψ(2S)** in hadronic matter

- Transport in **AA**:
J/ψ / ψ(2S) near equilibrium in late QGP / hadronic phase

→ **“Sequential Regeneration”** [Du+RR '15]



3.3 Coupling of Quarkonia + Open HF Transport



• Boltzmann:

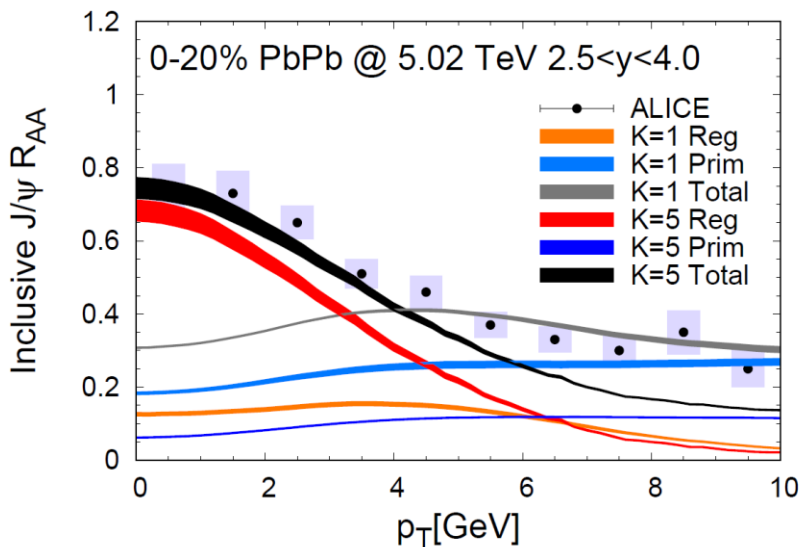
$$\frac{\partial f_{\Psi}}{\partial t} + \vec{v} \cdot \nabla f_{\Psi} = -\alpha f_{\Psi} + \beta$$

• Gain term:

$$\beta(\vec{p}, T) = \sum_i \int d\Pi_{\Psi} d_{\tilde{i}} d_c d_{\bar{c}} \overline{|M_{\tilde{i}c\bar{c} \rightarrow i\Psi}|^2} [1 \pm f_i(p_i)] f_i(\tilde{p}_i) \gamma_c f_c(p_c) \gamma_c f_{\bar{c}}(p_{\bar{c}})$$

• Quasifree:

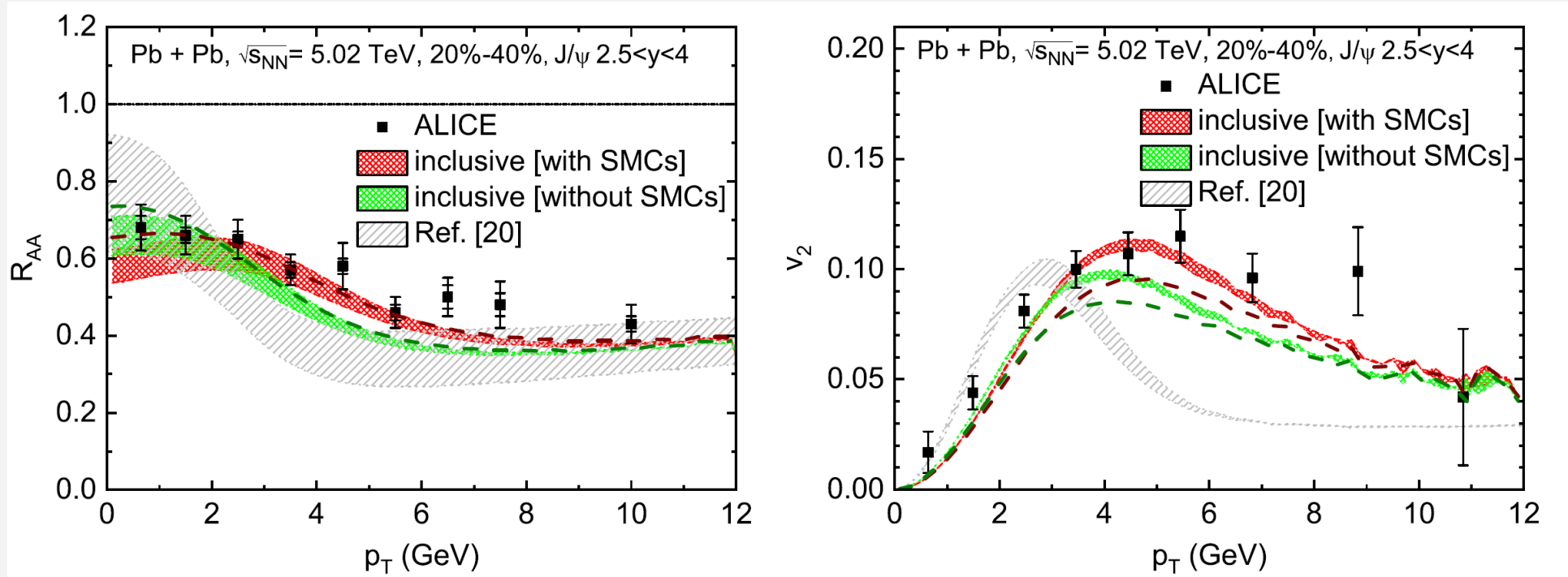
$$\overline{|M_{i\Psi \rightarrow \tilde{i}c\bar{c}}(p_i, p_{\Psi}, \tilde{p}_i, p_c, p_{\bar{c}})|^2} \approx 2 \left| M_{ic \rightarrow \tilde{i}c} \left(p_i, \frac{m_{\Psi} - m_c}{m_{\Psi}} p_{\Psi}, \tilde{p}_i, p_c \right) \right|^2 (2\pi)^3 2E_{\bar{c}} \delta^{(3)} \left(\vec{p}_{\bar{c}} - \frac{m_c}{m_{\Psi}} \vec{p}_{\Psi} \right)$$



- Simultaneous **c**-quark diffusion + charmonium kinetics (same $M_{pQCD}^2 * K$)
- **J/ψ** yield + spectral shape require $K \geq 5$
- Strong coupling for open + hidden HF

[Du+RR '22]

3.4 Charmonia - Open-HF Coupling II: Elliptic Flow

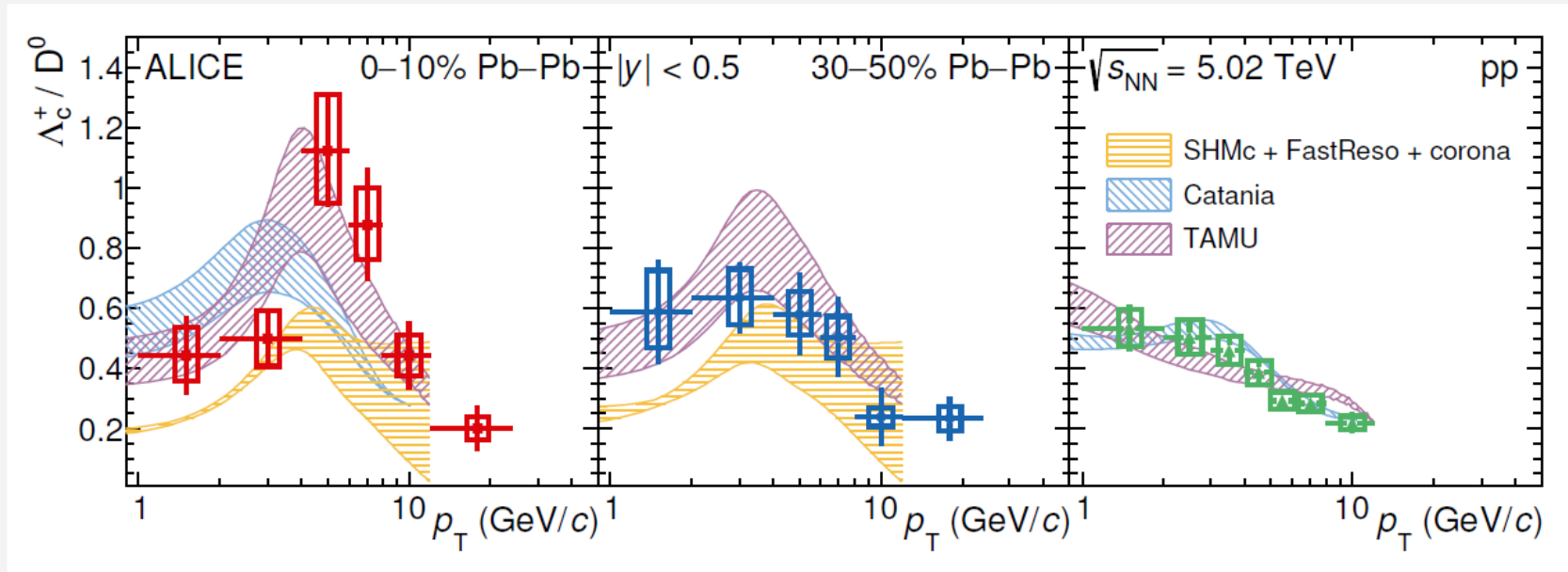


- QGP charm-quark diffusion + RRM
- Recombination reaches out to $p_T \sim 8$ GeV [He,Wu+RR '22]
- See also [Villar et al '22]

4.) Summary

- HF particles can probe non-perturbative QGP structure, by connecting lattice QCD to phenomenology (HQ free energy, correlators, susceptibilities; EoS)
- Production models need to satisfy equilibrium constraints
- Coupling of open + hidden HF sectors provides strong constraints
- Large collision rates require careful assessment of quantum effects; relations to semi-classical approaches to be understood

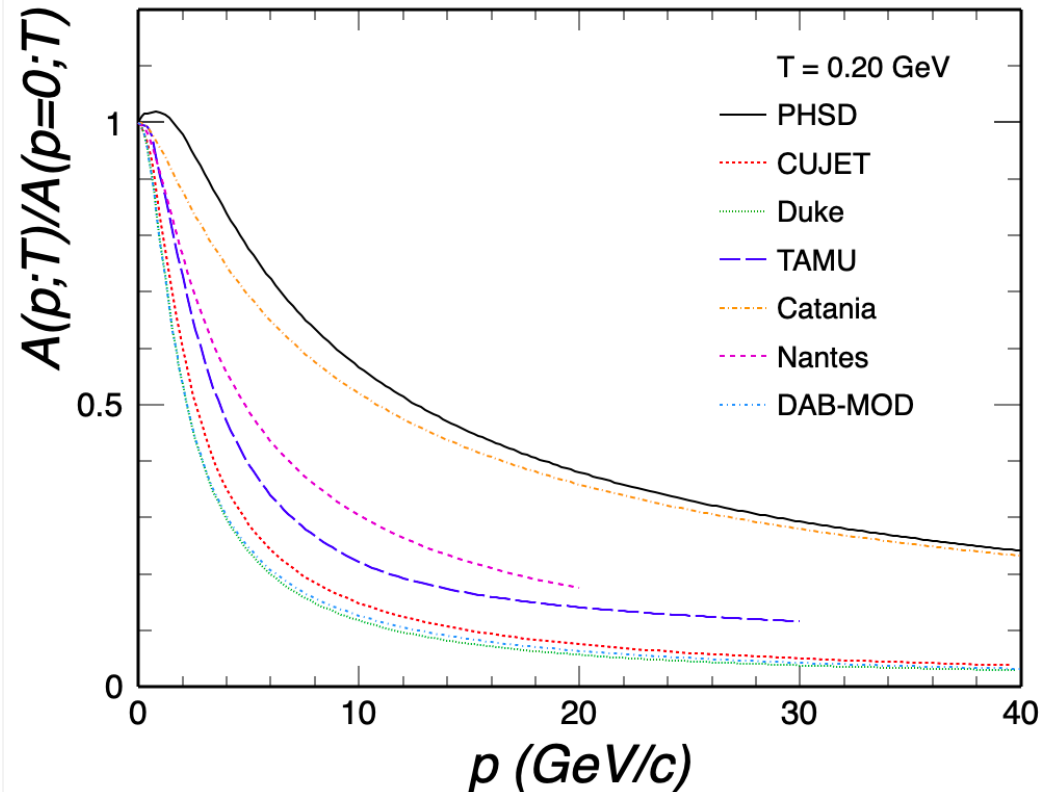
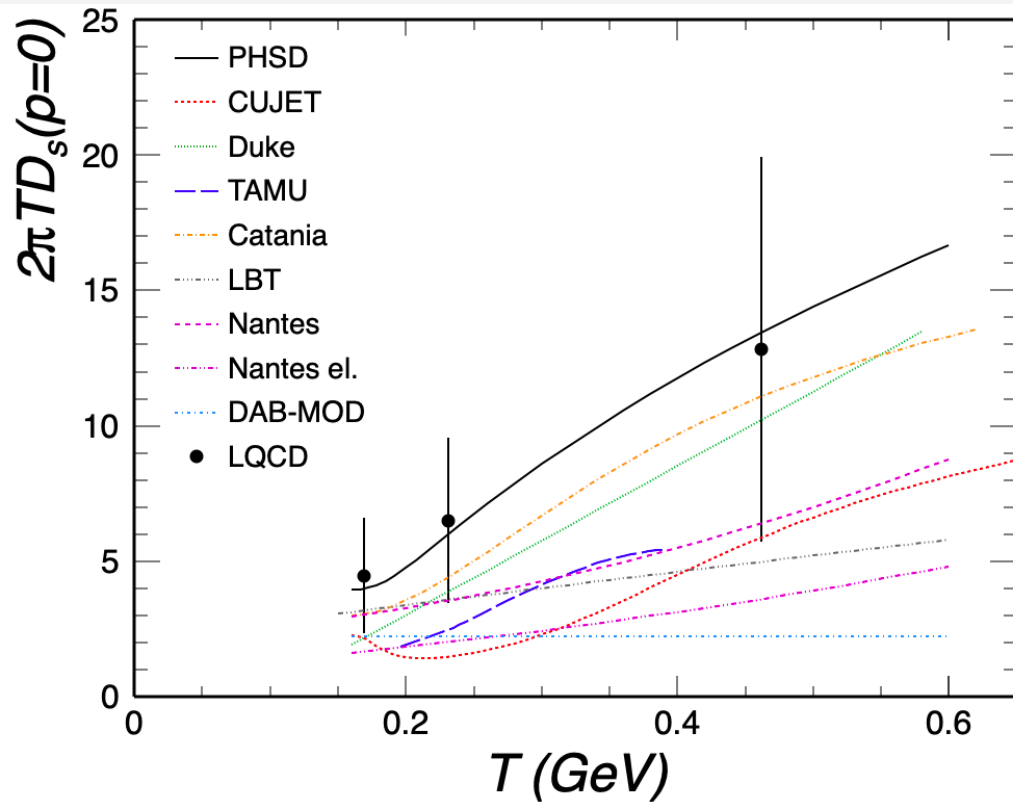
2.5 Charm-Hadron Spectra from pp to AA



- “Flow bump” vs. yield enhancement:
control over equilibrium limit essential for interpretation

2.4 Momentum Dependence of Relaxation Rate

[ECT* HF workshop '21]

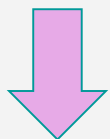


- correlation: large $\mathcal{D}_s \leftrightarrow$ weak \mathbf{p} -dependence
- need handle on momentum dependence

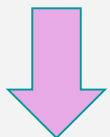
2.2 Many-Body Theory of Heavy Flavor in QGP

Microscopic Theory

Lat-QCD HQ
free energy



Heavy-quark
potential



Quarkonium
binding E_B



Transport in HI collisions

HQ interactions
in QGP (\mathcal{D}_s)



Quarkonium
reaction rate Γ_Ψ ,
equil. limit N_Ψ^{eq}



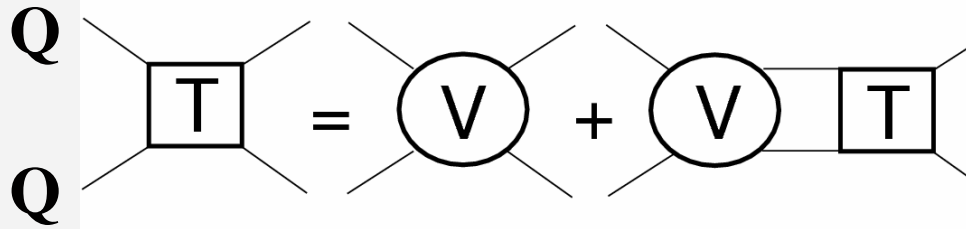
Test in Experiment

Open HF
observables

Quarkonium
observables

2.3 Thermodynamic T-Matrix Approach

- T-matrix equation



In-medium potential?

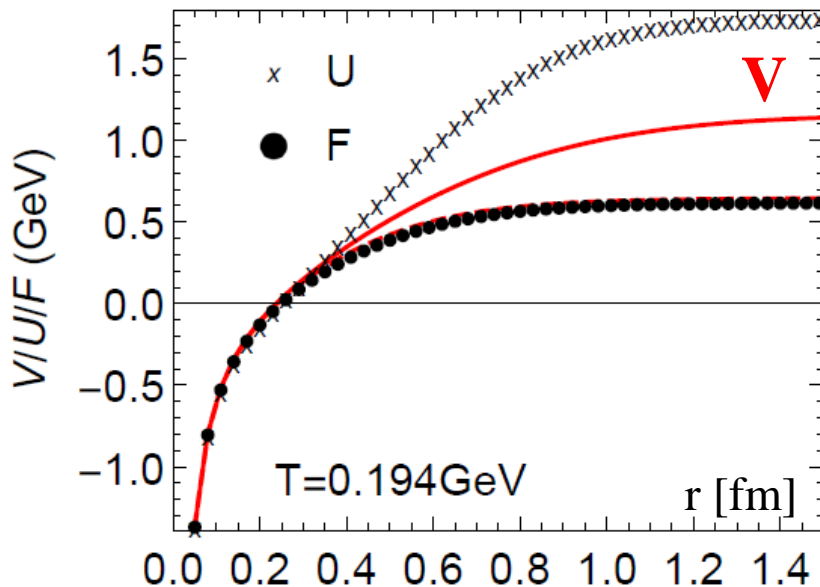
- Free energy:

$$F_{Q\bar{Q}}(r_1 - r_2) = -\frac{1}{\beta} \ln \left(\int_{-\infty}^{\infty} d\omega \sigma(\omega, r_1 - r_2) e^{-\beta\omega} \right)$$

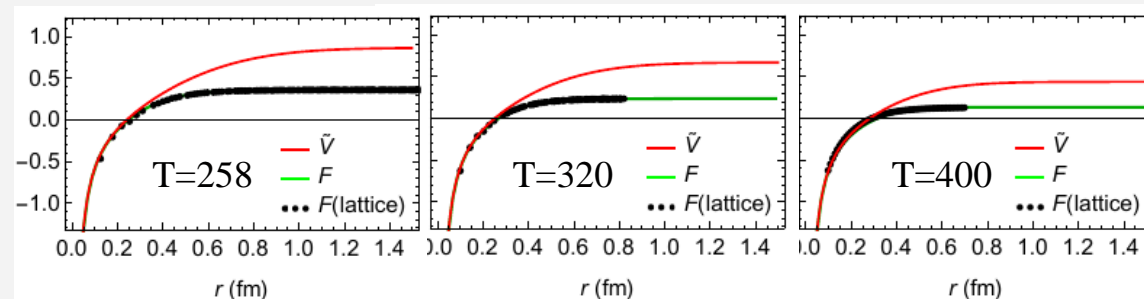
Spectral function:

$$\sigma(\omega, r) = \frac{1}{\pi} \frac{(V + \Sigma)_I(\omega)}{(\omega - (V + \Sigma)_R)^2 + (V + \Sigma)_I^2(\omega)}$$

⇒ infer in-medium potential



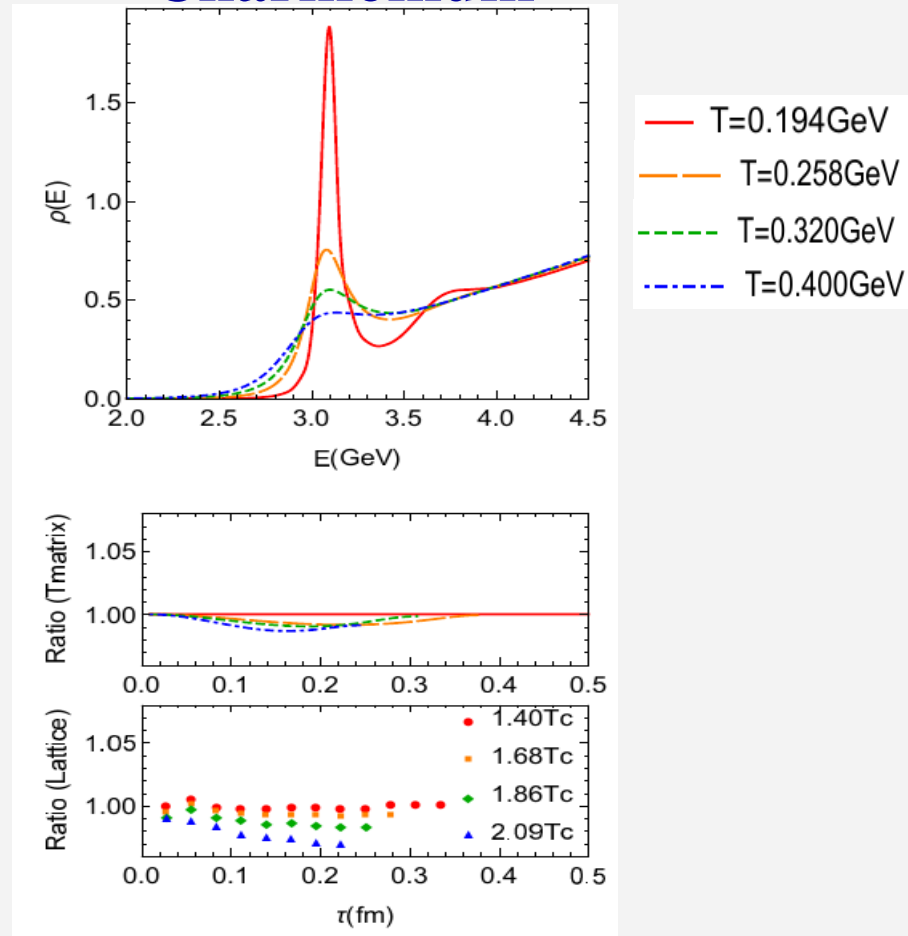
- Large imaginary parts $\Sigma_I \Rightarrow V > F$
- **Remnants of confining force above T_c !**



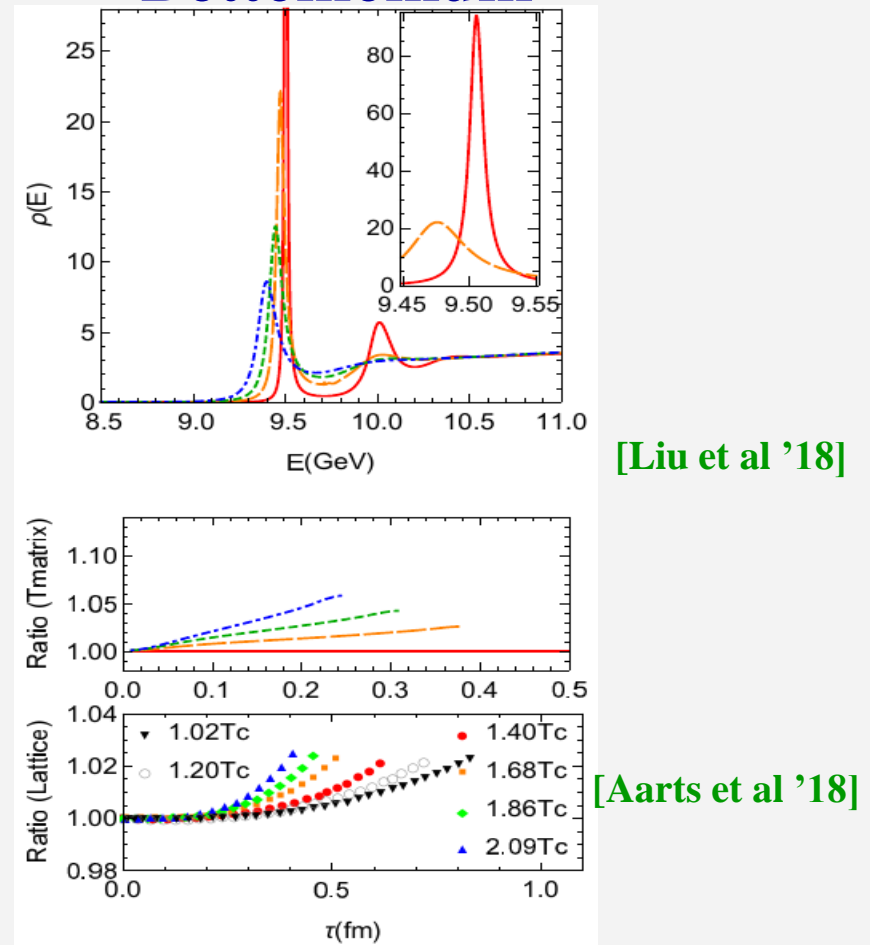
2.4 More IQCD Constraints: Euclidean Correlators

$$G_\alpha(\tau, T) = \int \frac{dE}{2\pi} \rho_\alpha(E, T) \frac{\cosh [E (\tau - 1/2T)]}{\sinh [E/2T]}$$

Charmonium



Bottomonium



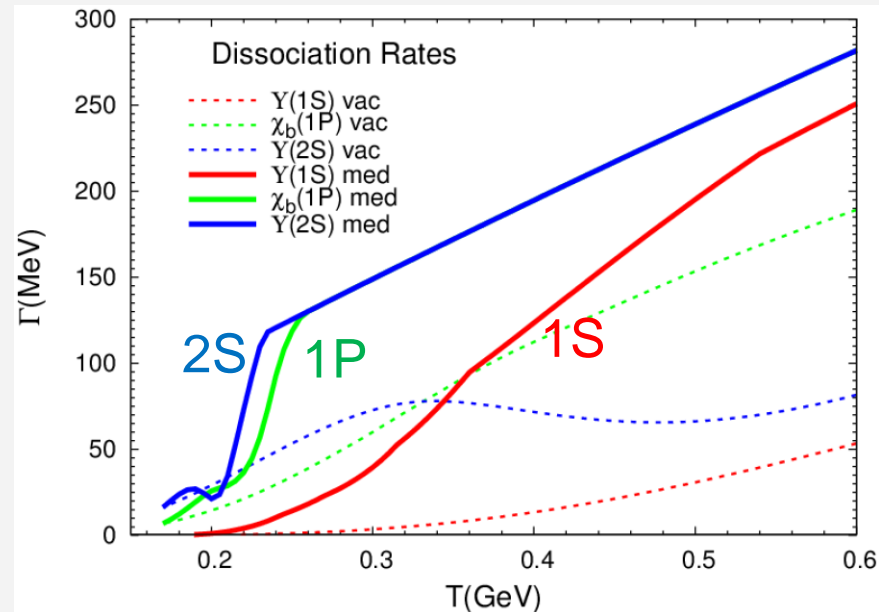
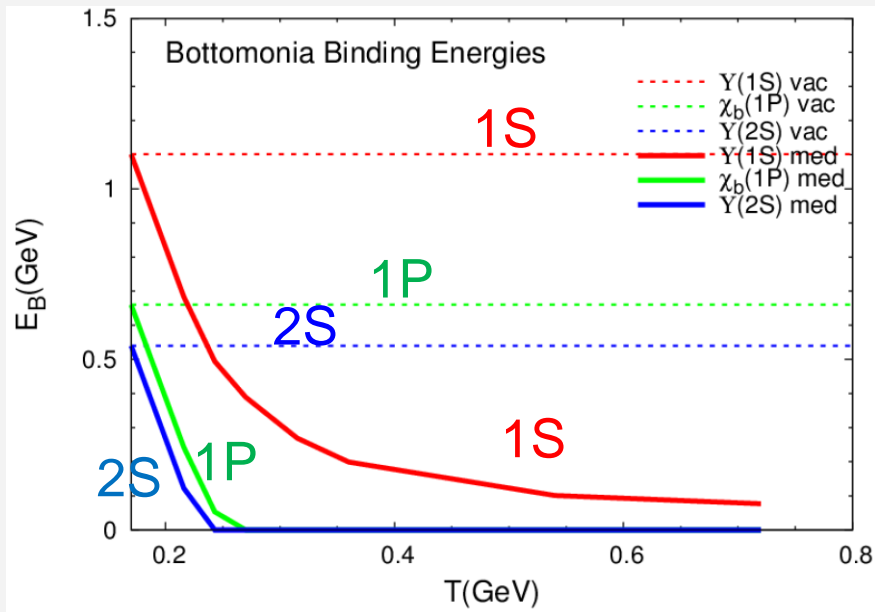
[Liu et al '18]

[Aarts et al '18]

• J/ψ “melts” near $T \sim 250\text{-}300$ MeV

• $\Upsilon(1S)$ survives until $T > 400$ MeV

2.5 Binding Energies + Reaction Rates: Y States



[Du et al '18]

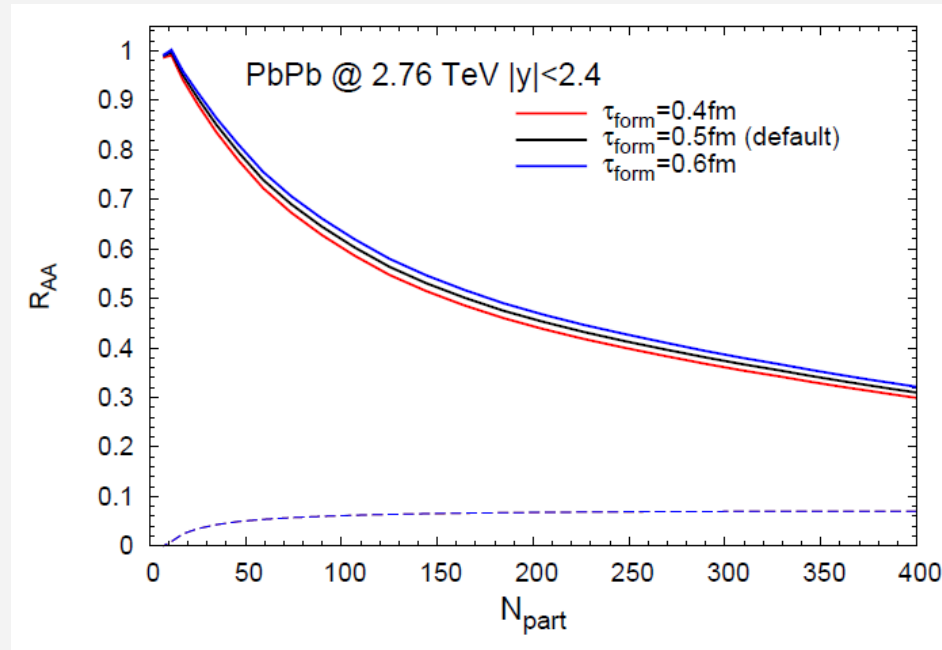
$$\Gamma_Y^{\text{qf}}(p, T) = \sum_p \int \frac{d^3 p_p}{(2\pi)^3} d_p f_p(\omega_p, T) v_{\text{rel}} \sigma_{Y p \rightarrow b \bar{b} p}(s)$$

- Reduced binding “accelerates” dissociation (opens phase space)
- Dissociation rate depends on “medium” particle distribution

2.6 Quarkonium Formation Times

- Pair production time $\tau_{QQ} \leq 0.1 \text{ fm}/c$
- Bound-state formation time $\tau_{\text{form}} \sim 1/E_B \sim 0.2 - 2 \text{ fm}/c$
- Build-up of wave function reduces dissociation rate

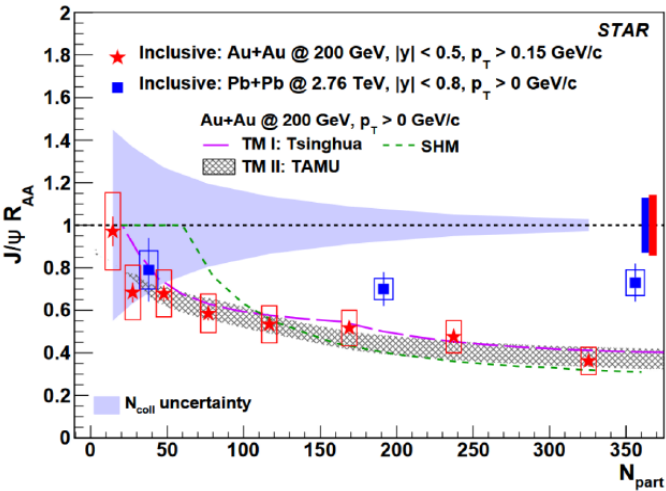
$$\alpha_Y(\vec{p}, T(\tau)) \equiv \Gamma_Y(\vec{p}, T(\tau)) \frac{\tau}{\tau_{\text{form}}} \frac{m_Y}{\sqrt{p^2 + m_Y^2}}$$



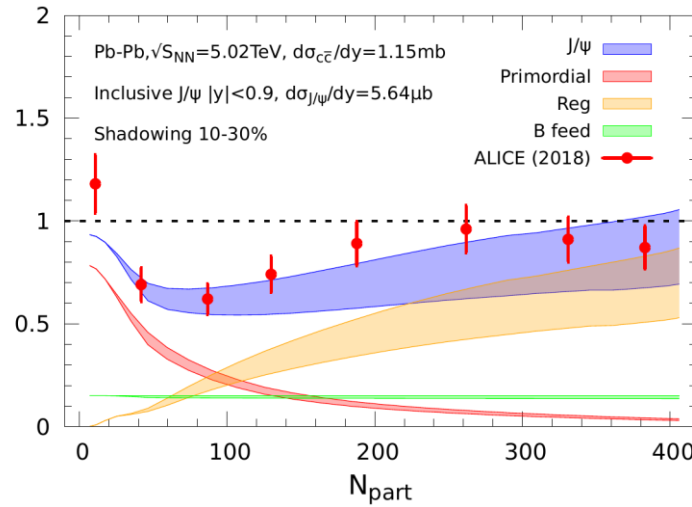
- Relatively small effect in heavy-ion collisions
- Microscopically: quantum evolution of wave package

3.1 Snapshot of Quarkonium Transport Results

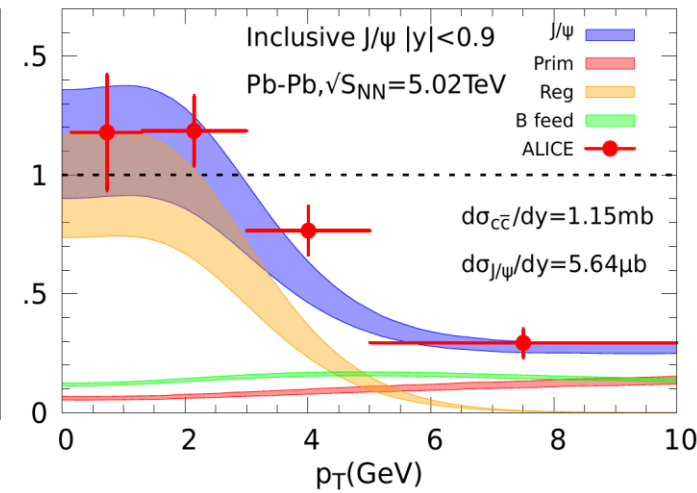
RHIC



LHC: centrality

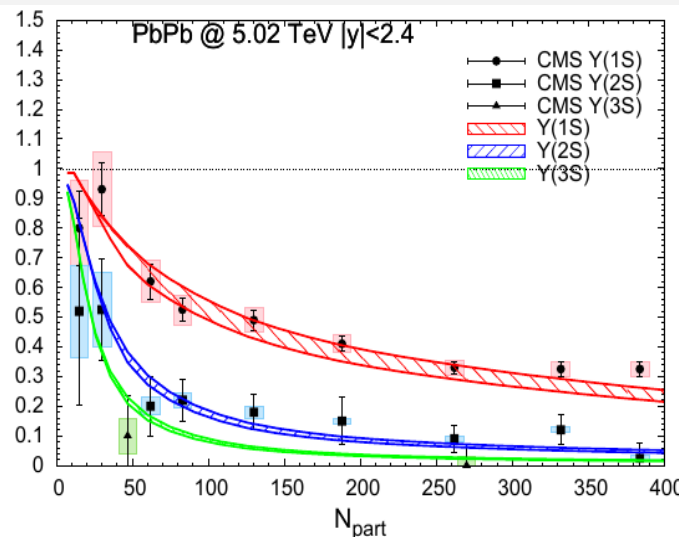
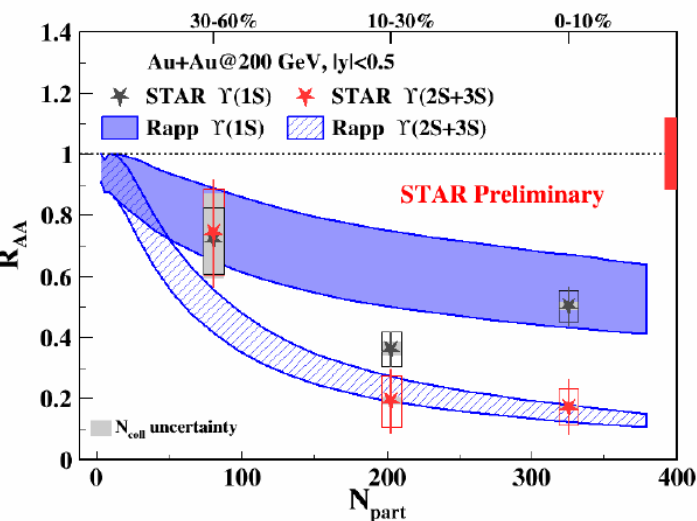


LHC: momentum



- mostly suppression

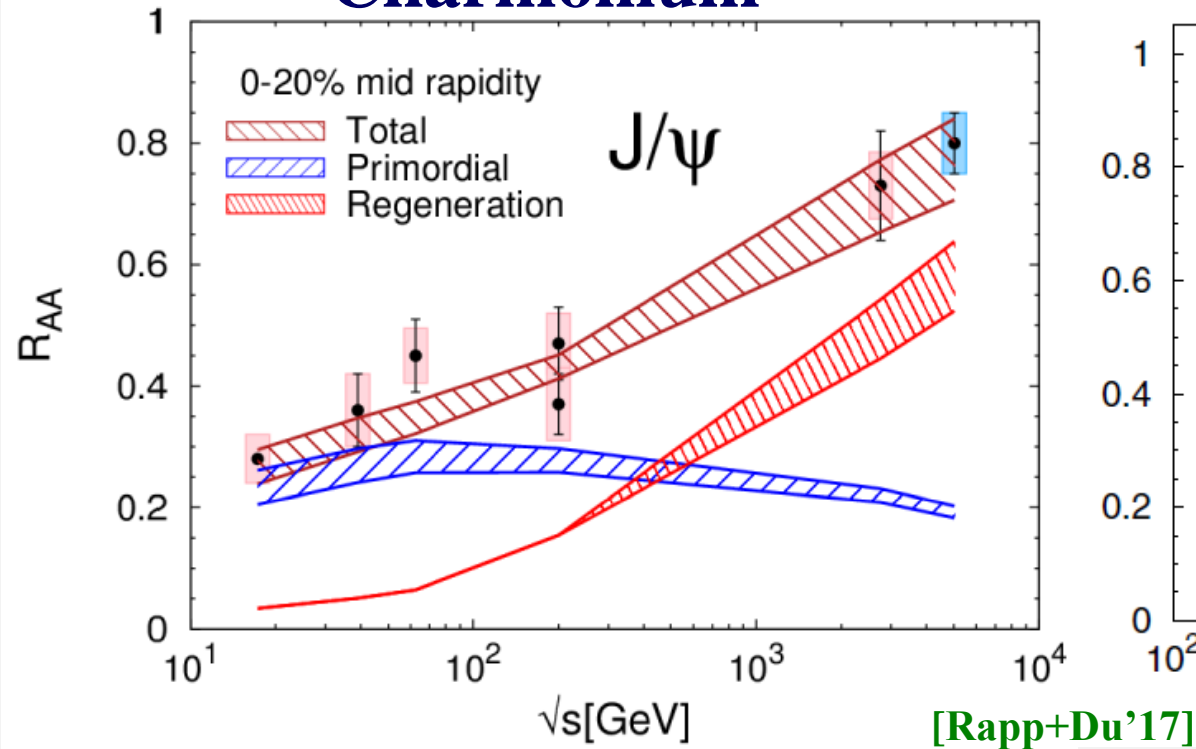
- substantial regeneration, concentrated at low p_T



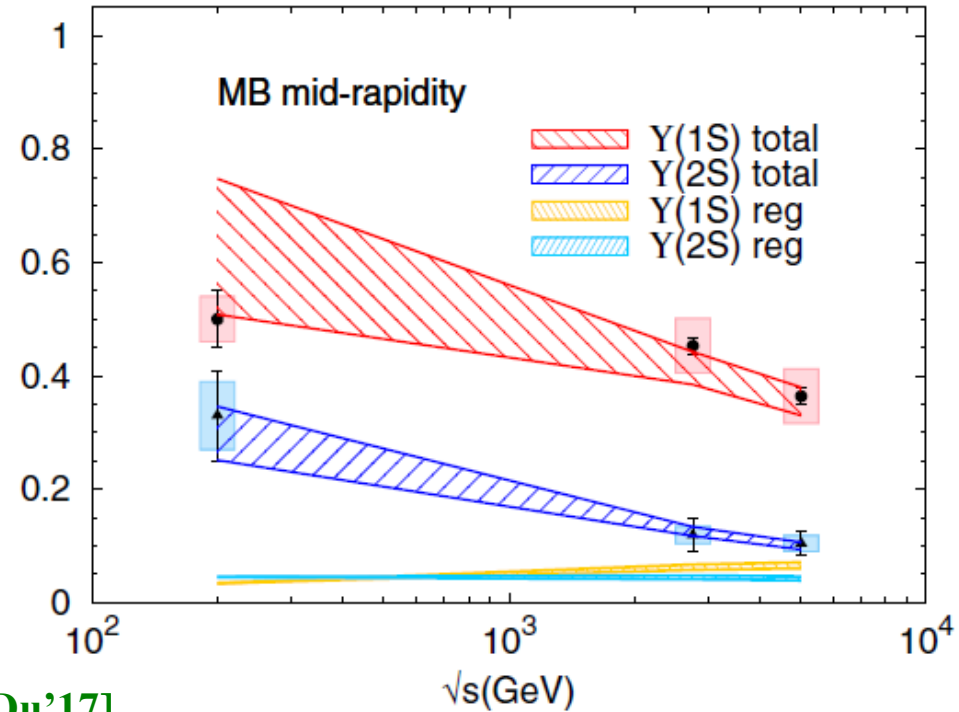
- Bottomonia mostly driven by suppression, regeneration small even at the LHC

3.2 Excitation Functions: **SPS - RHIC - LHC**

Charmonium



Bottomonium

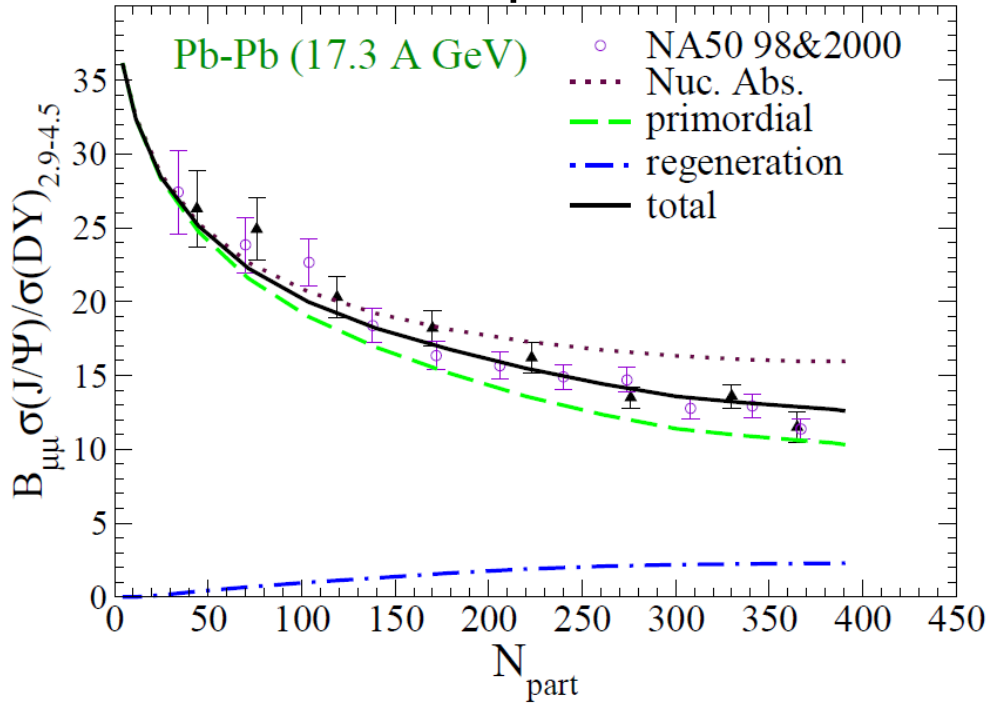


- Gradual **increase** of total J/ψ R_{AA}
- Regeneration **and** suppression increase
- Regeneration concentrated at low p_T
- Gradual **suppression**
- Regeneration (N_Y^{eq}) small
- Qualitative difference from J/ψ

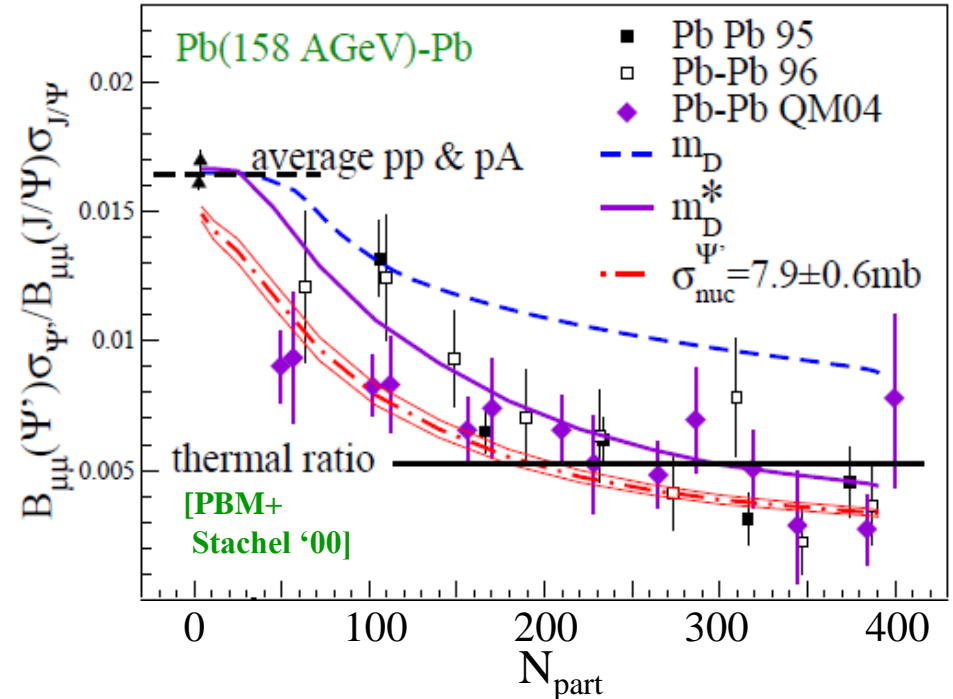
Data: NA50, PHENIX, STAR, ALICE, CMS

3.4 Charmonia at SPS Energy ($\sqrt{s}=17$ GeV)

J/ψ



ψ'(3686)



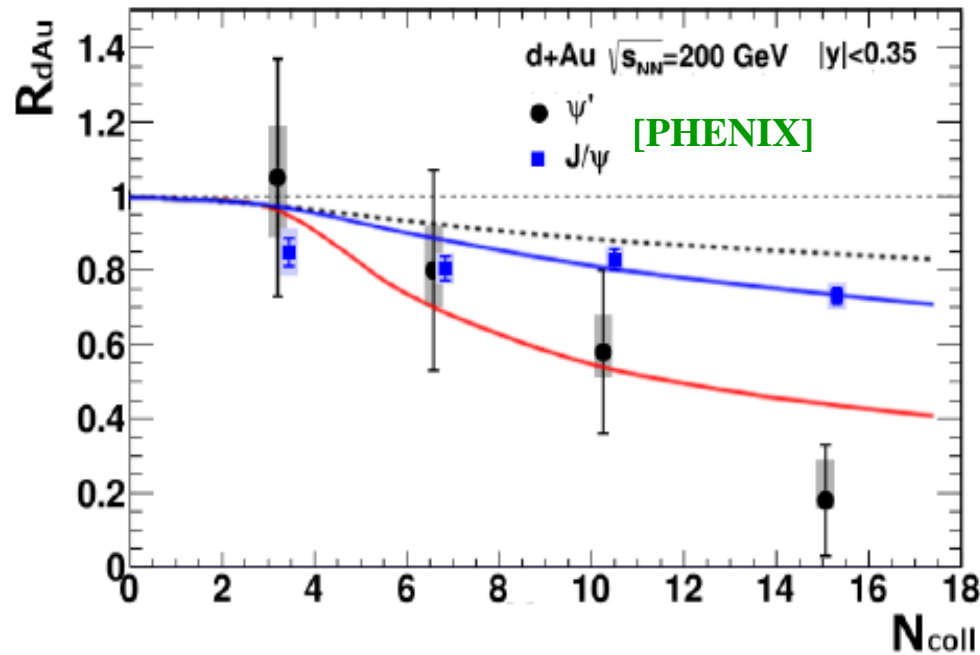
- Large contributions from initial “nuclear absorption” at low \sqrt{s}
- requires high-energy $N-\Psi$ Xsec

- ψ' dissociation in hadronic matter important in transport models

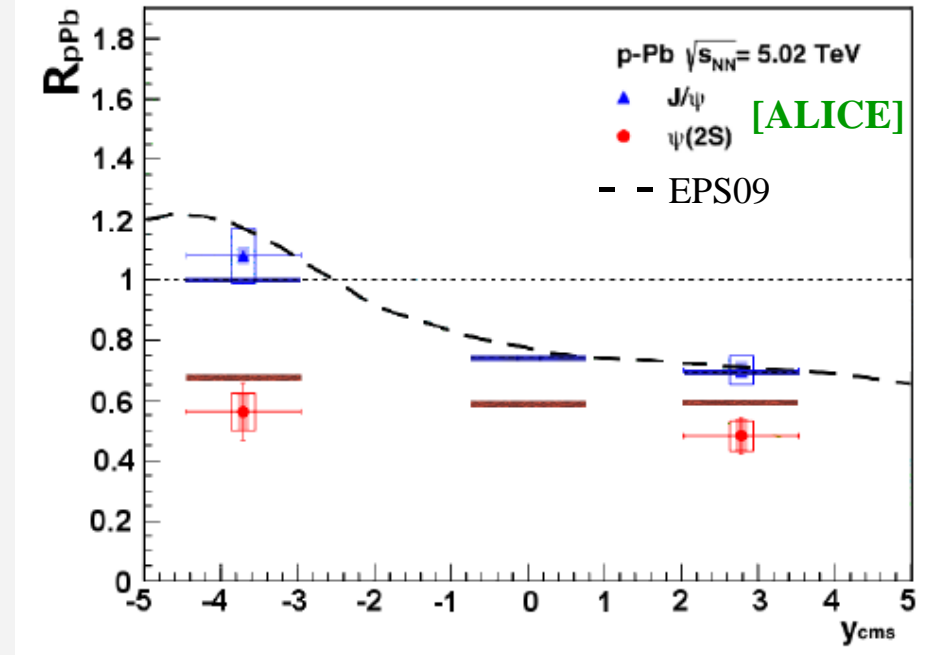
[Sorge et al '97, Grandchamp et al '03, ...]

3.5 $\psi(2S)$ in d/p-A Collisions

d-Au (0.2TeV)



p-Pb (5.02TeV)



- noticeable ψ' and little J/ψ suppression, consistent with “comovers”

[Ferreiro '15]

- supports fireball formation with:

$$\tau_{\text{FB}} \Gamma(\psi') \sim 1 \Rightarrow \Gamma_{\text{avg}}(\psi') \sim 50\text{-}100 \text{ MeV}$$

similar to thermal

$$\tau_{\text{FB}} \Gamma(J/\psi) \ll 1 \Rightarrow \Gamma_{\text{avg}}(J/\psi) < 20 \text{ MeV}$$

widths at $T \approx 200 \text{ MeV}$

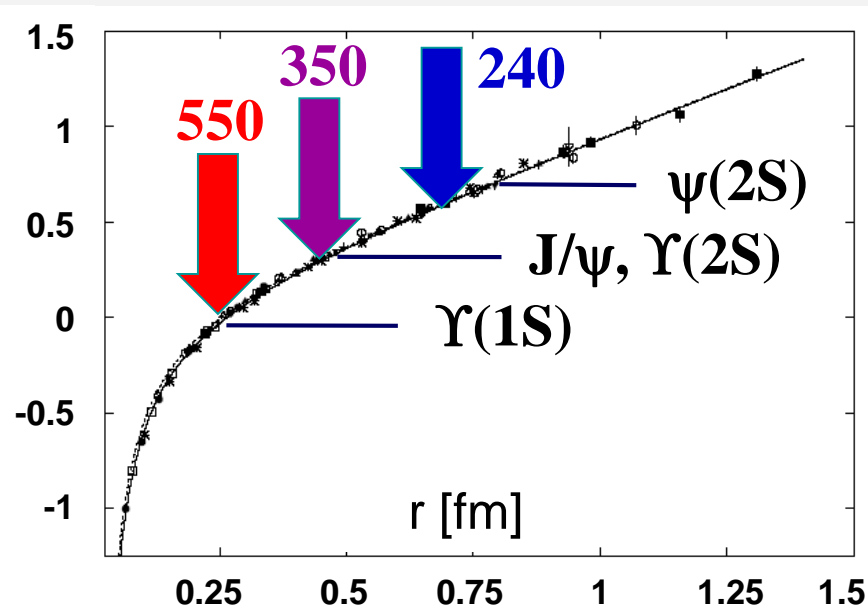
[Du et al '15]

2.3 Upshot of Quarkonium Phenomenology

Use temperature estimates from hydro/photons/dileptons to infer:

$$T_0^{\text{SPS}} (\sim 240) < T_{\text{melt}}(\text{J}/\psi, \Upsilon') \leq T_0^{\text{RHIC}} (\sim 350) < T_{\text{melt}}(\Upsilon) \leq T_0^{\text{LHC}} (\sim 550)$$

- Remnants of confining force survive at **SPS** [hold J/ψ together]
- Confining force screened at **RHIC+LHC** [“melts” $\text{J}/\psi + \Upsilon(2\text{S})$]
- Color-Coulomb screening at **LHC** [$\Upsilon(1\text{S})$ suppression]
- Thermalizing charm quarks recombine at **LHC** [large J/ψ yield]



2.4 Heavy-Quark Potential Extraction from Υ Data

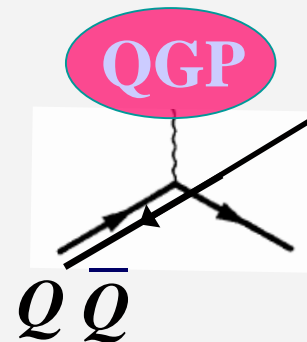
- Ansatz for in-med. potential

$$V_{Q\bar{Q}}(r) = -\frac{4}{3}\alpha_s \frac{e^{-m_D r}}{r} + \sigma (r - R_{SB}) \Theta (R_{SB} - r)$$

- Parameterize T -dependence of $m_D \sim \alpha T$, $1/R_{SB} \sim m_S(T; \beta, \gamma, \delta)$
 \Rightarrow determine in-med. binding energies $E_Y(m_D, m_S)$

- Compute $\Gamma_Y [E_Y]$

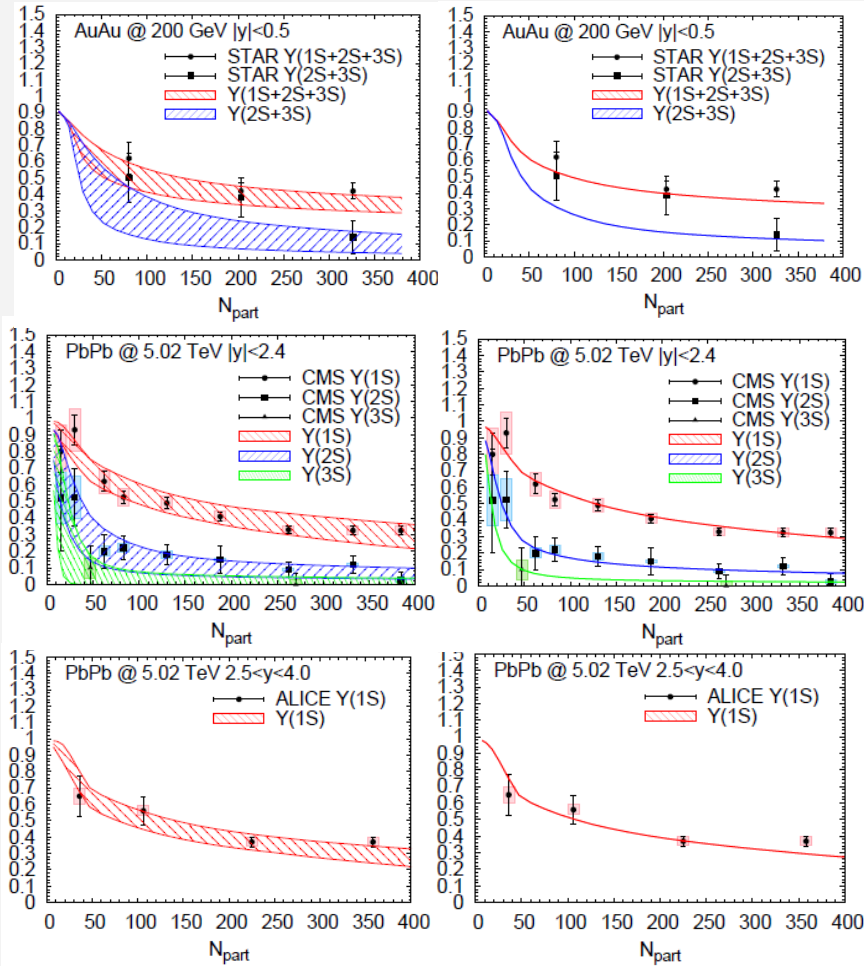
- also requires HQ-medium coupling
- perturbative: $\alpha_s \sim 0.3$
- non-perturbative: $K \geq 5$



- Deploy transport approach to fit $(\alpha, \beta, \gamma, \delta; K)$ to Υ data at **RHIC + LHC**

3.4 Statistical Extraction of Heavy-Quark Potential

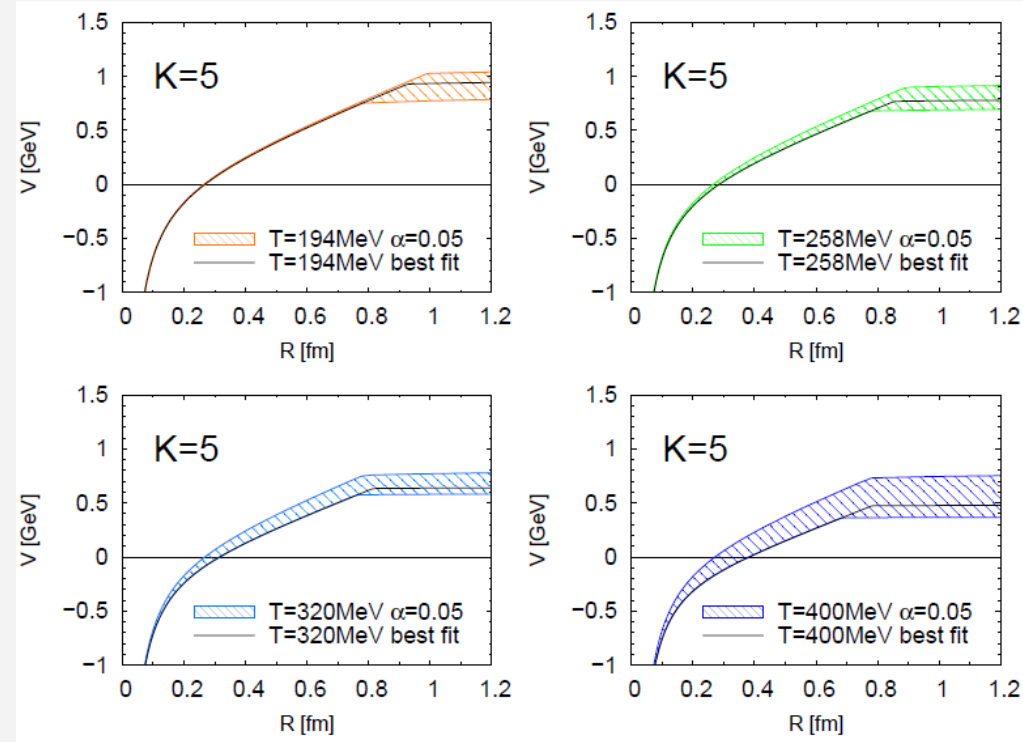
Fit Results



[Du et al '19]

• $\chi^2/\text{dof} \leq 1$

In-Medium Potential

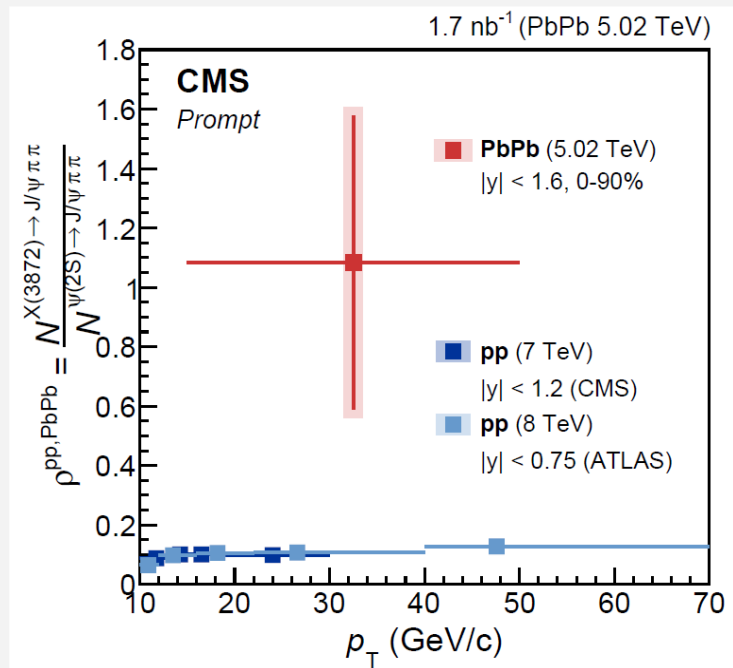


- “Strongly-coupled” solution: remnants of confining force survive well above T_c
- Not unique ...

3.5 X(3872)

Does **X(3872)** production in heavy-ion collisions reveal **internal structure** information?

Larger size (molecule vs. tetraquark) \Rightarrow larger yield -- or else?



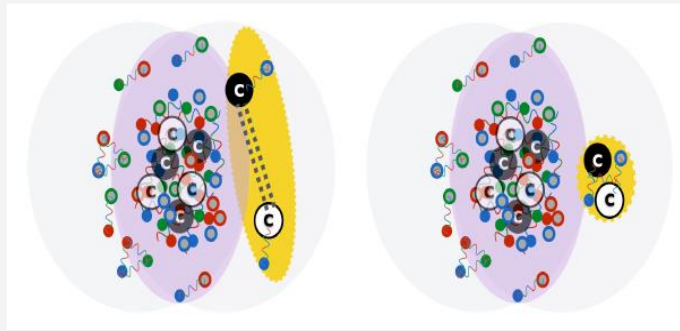
- With $R_{AA}^{\text{exp}}(\psi(2S)) \sim 0.14$, $R_{AA}(X)$ compatible with unity
- $p_T > 15\text{GeV}$ might be outside coalescence/regeneration regime

3.5.2 In-Medium X(3872) Production Models

- **Coalescence Models:** partons (f_i) projected on hadron wave fct. (W)

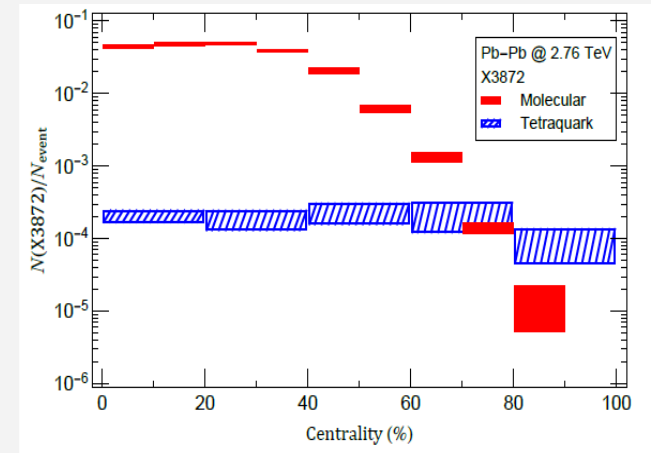
[Cho et al '11, Fontoura et al. '19, ...]

$$f_h(\mathbf{p}'_h) = \int \left[\prod_i dp_i f_i(\mathbf{p}_i) \right] W(\{\mathbf{p}_i\}) \delta(\mathbf{p}'_h - \sum_i \mathbf{p}_i) \quad W = W(\sigma) \text{ with hadron radius } \sigma$$



$$N_X \sim \sigma^3 \Rightarrow N_{\text{mol}} \gg N_{\text{tetra}}$$

[H. Zhang et al '20]



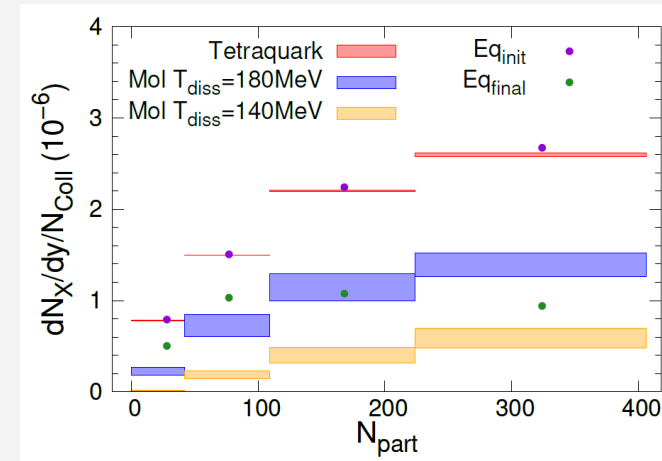
- **Transport Approach:**

[B. Wu et al '20]

$$\frac{dN_X}{dt} = -\Gamma_X (N_X - N_X^{eq})$$

- Reaction rate larger for weakly bound molecule \Rightarrow molecule freeze-out later than tetraquark

- Equilibrium limit drops with $T \Rightarrow N_{\text{mol}} < N_{\text{tetra}}$

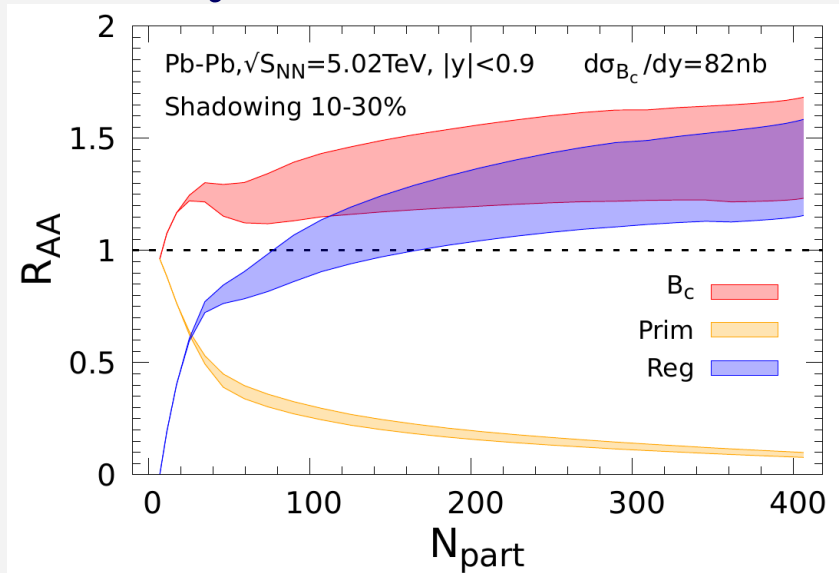


3.6 $B_c(6275)$

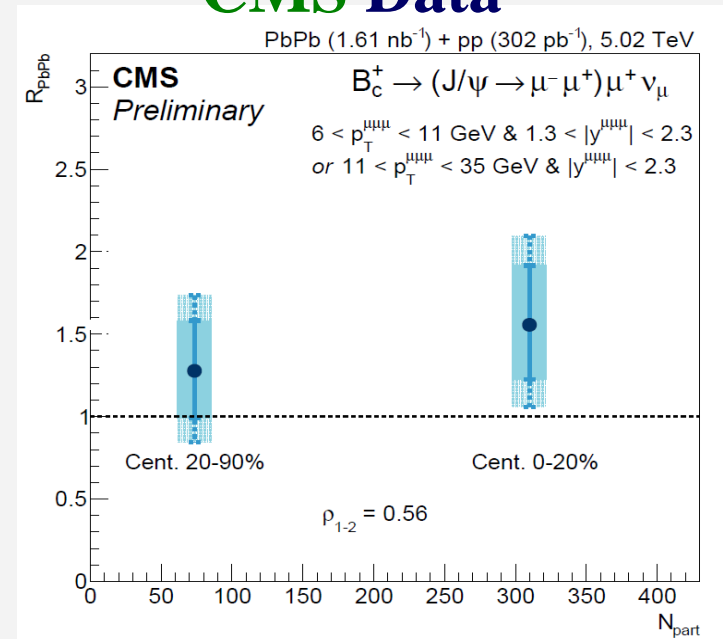
- Critical “universality test” for transport models
- Expect large yields in **AA** relative to **pp** [$dN_c/dy \sim 30$ in **PbPb(5TeV)**]
- Need to control **pp** production cross section (denominator):

$$R_{AA}(p_T) \equiv (dN/dp_T)_{AA} / N_{coll} (dN/dp_T)_{pp}$$

Theory Prediction [Wu et al. '21]



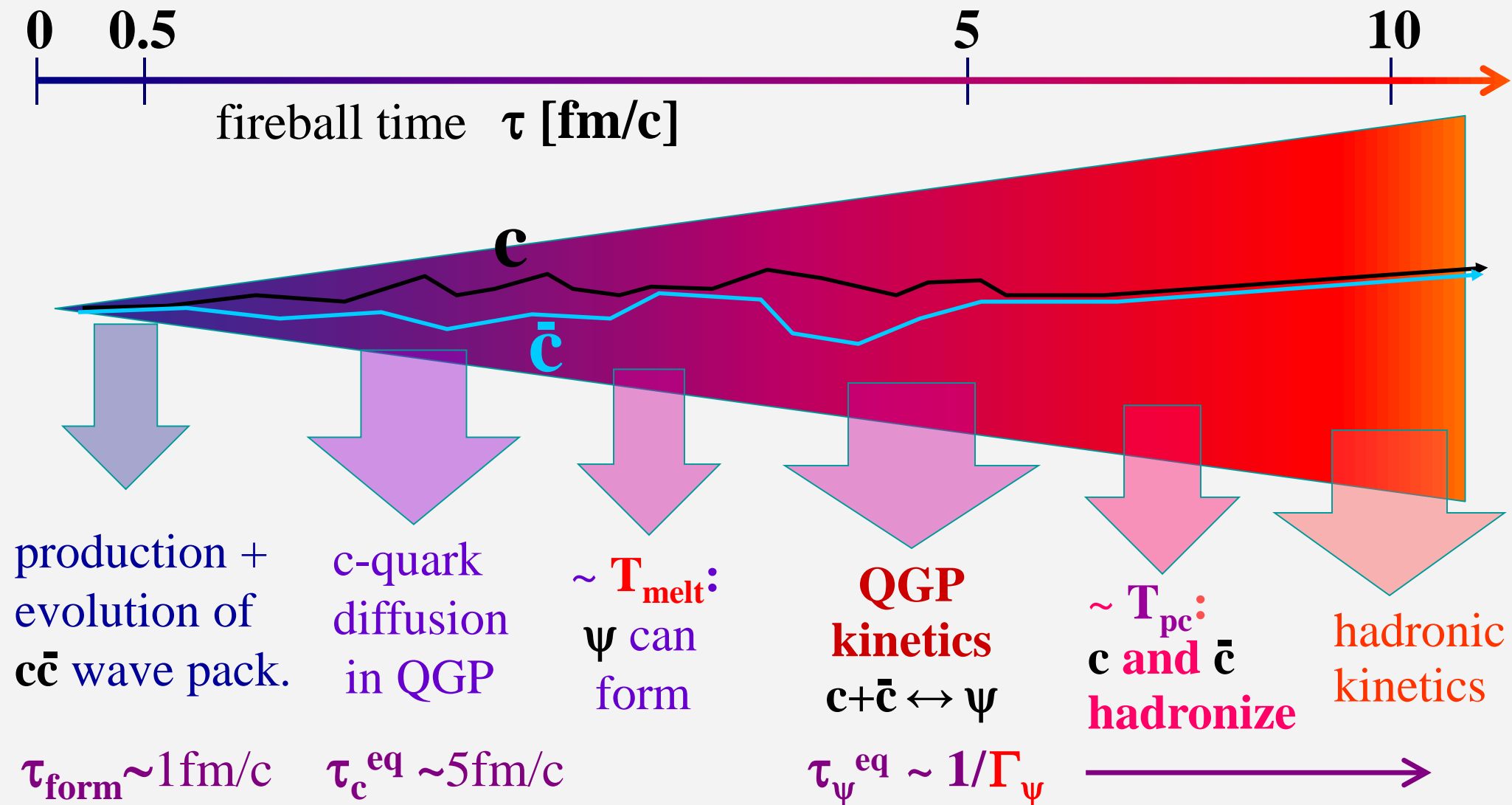
CMS Data



- Predictions appear to agree with data...

caveats: **pp** cross section, p_T range (theory: $p_T > 0$, data: $p_T^{3\mu} > 6 \text{ GeV}$)

2.1 Quarkonium Transport in URHICs



[Satz et al, Capella et al, Spieles et al, PBM et al, Thews et al, Grandchamp et al, Ko et al, Zhuang et al, Zhao et al, Chaudhuri, Gossiaux et al, Young et al, Ferreiro et al, Strickland et al, Brambilla et al, ...]