

# Hadron Spectroscopy

## Part II

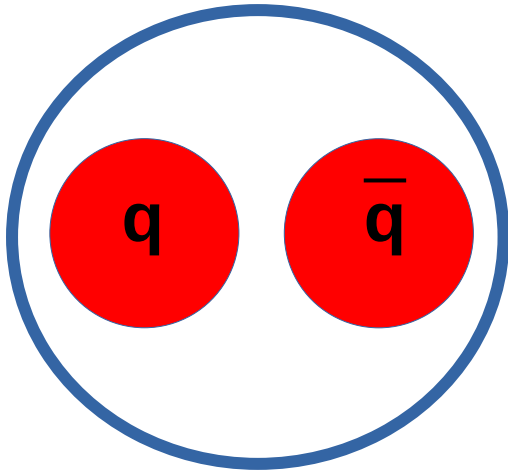
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Joint Institute for Nuclear Research  
Dubna, Russia

*e-mail: [zhemchugov@jinr.ru](mailto:zhemchugov@jinr.ru)*

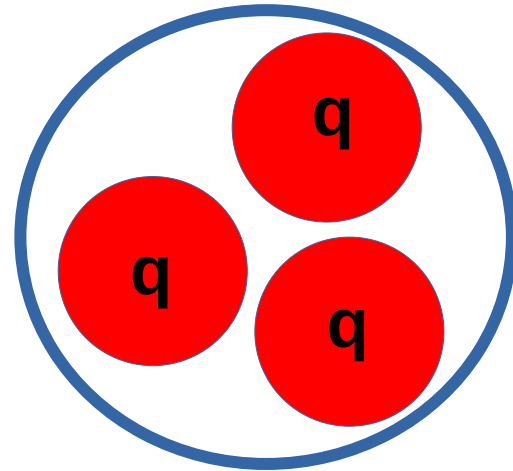
# Outline

- Why do we call it spectroscopy?
- What are we really measuring?
- Conventional hadron spectra
- Light hadrons: hunt for glueballs, search for diquarks
- Heavy hadrons: multiquarks and other exotics

# Quark model



*mesons*



*baryons*

# More quarks per particle?

Volume 8, number 3

PHYSICS LETTERS

1 February 1964

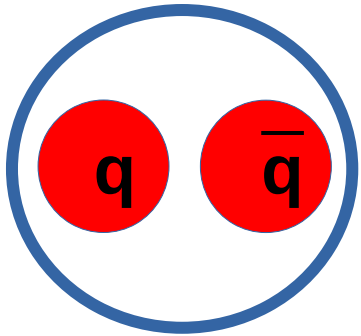
## A SCHEMATIC MODEL OF BARYONS AND MESONS \*

M. GELL-MANN

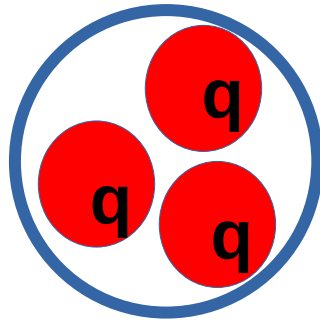
*California Institute of Technology, Pasadena, California*

A simpler and more elegant scheme can be constructed if we allow non-integral values for the charges. We can dispense entirely with the basic baryon  $\Lambda$  if we assign to the triplet  $t$  the following properties: spin  $\frac{1}{2}$ ,  $z = -\frac{1}{3}$ , and baryon number  $\frac{1}{3}$ . We then refer to the members  $u^{\frac{2}{3}}$ ,  $d^{-\frac{1}{3}}$ , and  $s^{-\frac{1}{3}}$  of the triplet as "quarks"  $q$  and the members of the anti-triplet as anti-quarks  $\bar{q}$ . Baryons can now be constructed from quarks by using the combinations  $(qqq)$ ,  $(qqq\bar{q})$ , etc., while mesons are made out of  $(q\bar{q})$ ,  $(q\bar{q}\bar{q})$ , etc. It is assumed that the lowest baryon configuration  $(qqq)$  gives just the representations **1**, **8**, and **10** that have been observed, while the lowest meson configuration  $(q\bar{q})$  similarly gives just **1** and **8**.

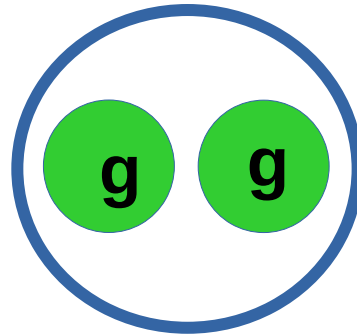
# Quark model



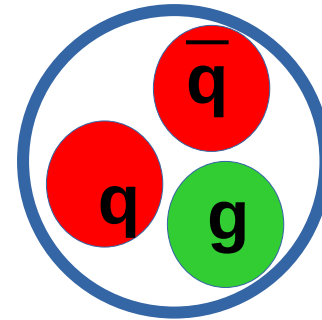
*mesons*



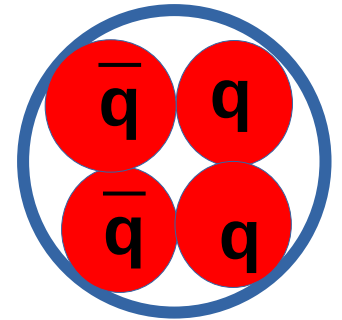
*baryons*



*glueballs*



*hybrids*



*multiquarks*

# Question

Can we call a deuteron a multiquark state?

# How to identify exotic hadrons?

- Exotic quantum numbers:
  - e.g.  $J^{PC} = 0^{--}, 0^{+-}, 1^{-+}, 2^{+-}, 3^{-+}, \dots$
- Surplus particles
- Unexpected decay patterns
  - e.g.  $X^+ \rightarrow J/\psi \pi^+$

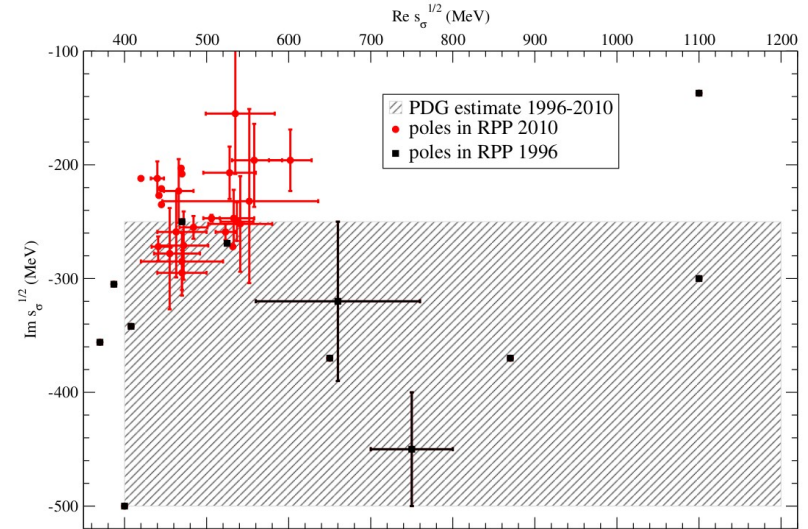




# Light hadrons

# $f_0(500)$ or $\sigma$

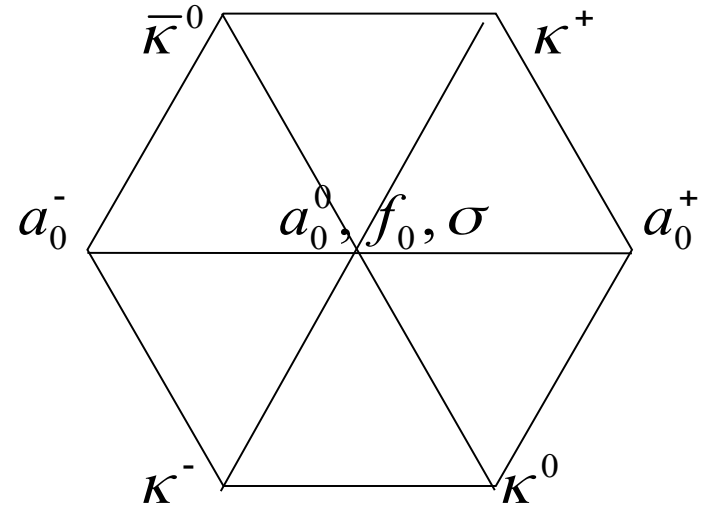
- $\pi\pi$  S wave
- Suggested in 1955 to explain short-range nucleon-nucleon interactions by two-pion exchange
- Controversial experimental status for about 40 years
- General agreement on the meson properties reached in mid-2000s after extensive theoretical and experimental efforts
- It is well established that it cannot be interpreted as predominantly made of a quark and an antiquark
  - Glueball?
  - Tetraquark?



J. Pelaez, arXiv:1510.00653v2

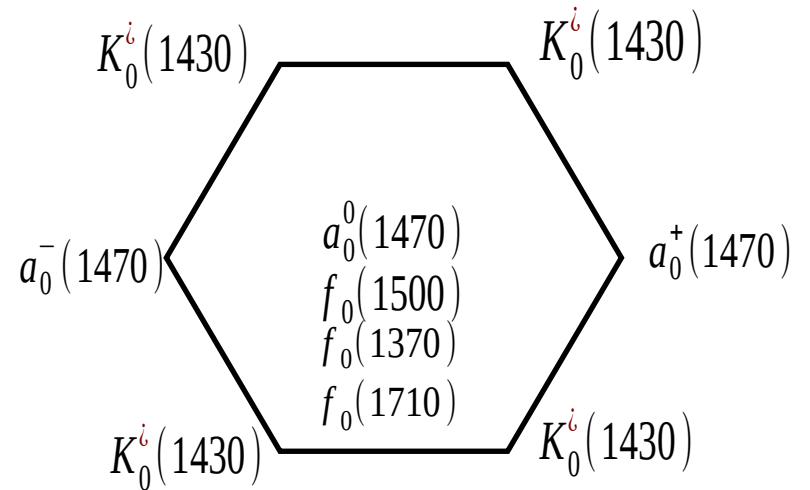
# Other scalar mesons below 1 GeV

- $K^*(700)$  or  $\kappa$ :  $K\pi$  S wave similar to  $f_0(500)$
- $a_0(980)$  and  $f_0(980)$ 
  - Very close to  $KK$  threshold
  - Same mass
- Nonet:  $I=0$ :  $f_0(500)$ ,  $f_0(980)$ ,  $I=1/2$ :  $K^*(700)$ ,  $I=1$ :  $a_0(980)$  ?
- What is it? Meson-meson molecule, or diquark-antidiquark, or compact tetraquark, or what?



# Scalar mesons above 1 GeV

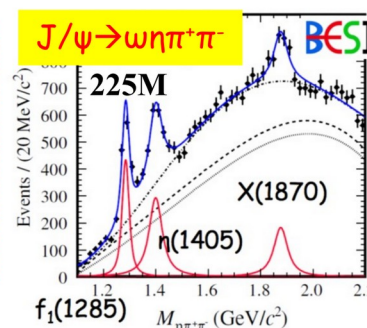
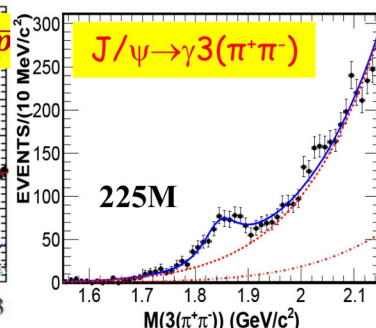
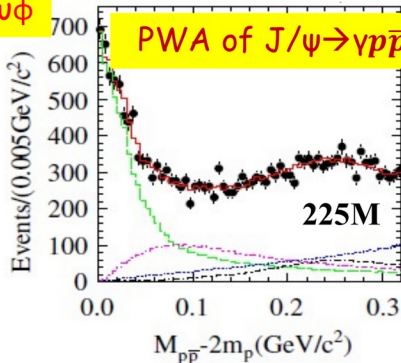
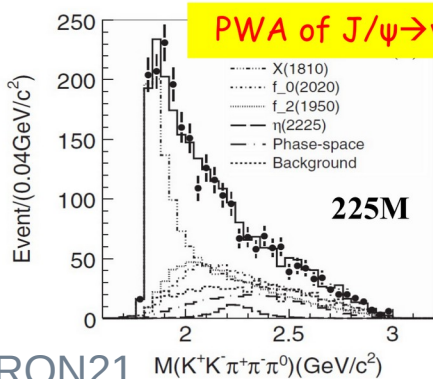
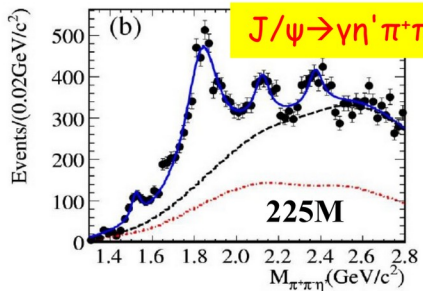
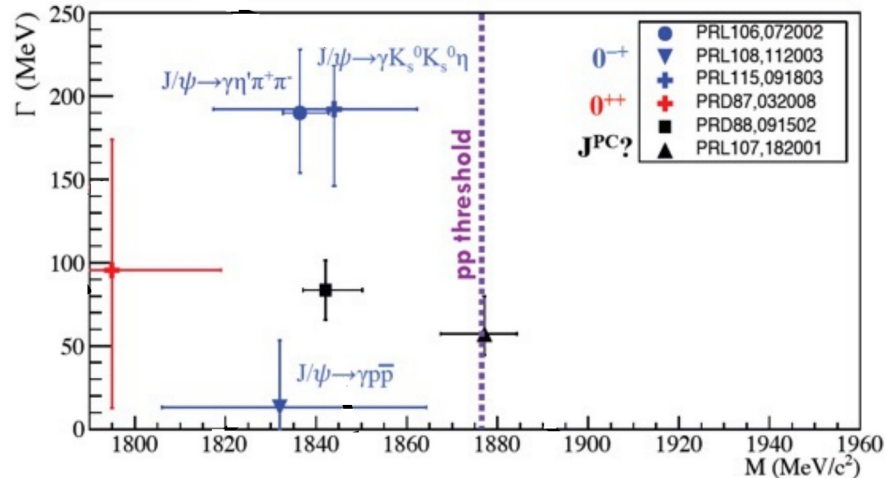
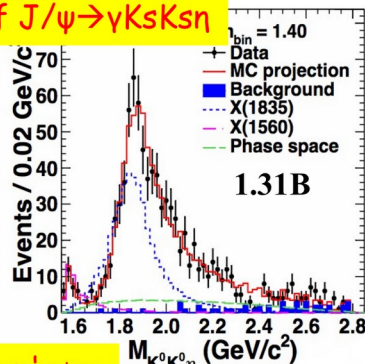
- Nonet:  $K_0^*(1430)$ ,  $a_0(1470)$ ,  $f_0(1370)$ ,  
 $f_0(1500)$ ,  $f_0(1710)$
- Which of  $f_0(1370)$ ,  $f_0(1500)$ ,  
 $f_0(1710)$  is an extra state?



# X(1835?) near the $pp$ mass threshold

BES III

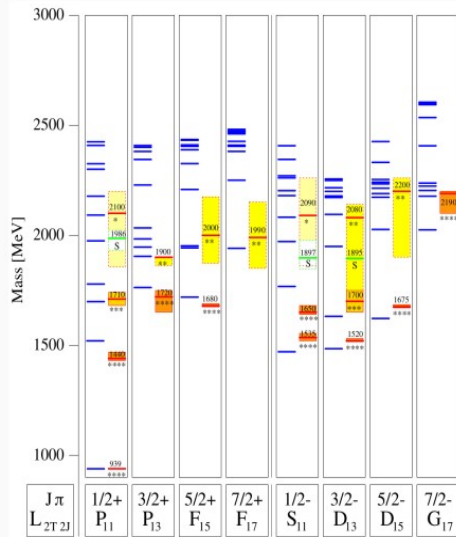
PWA of  $J/\psi \rightarrow \gamma K_s^0 K_s^0 \eta$



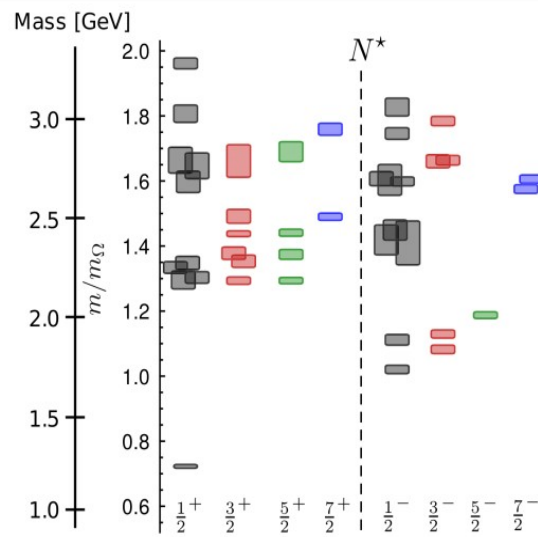
# Missing light baryons

A. Thiel, HADRON21

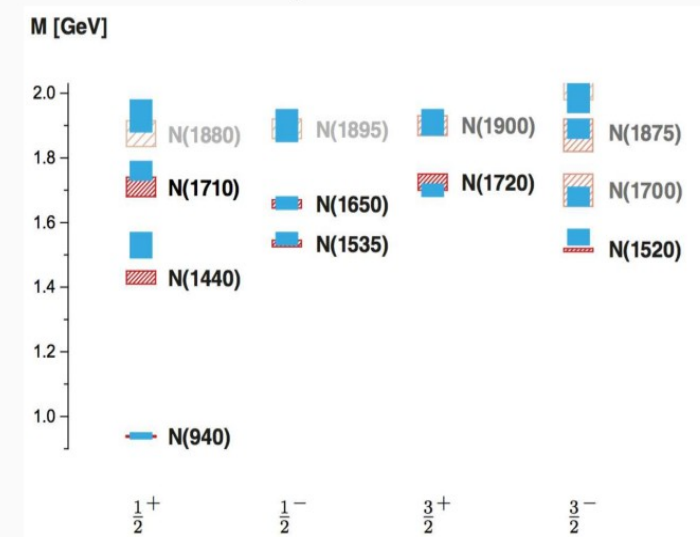
Quark Model



Lattice QCD Calculations



Dyson–Schwinger/Bethe–Salpeter approach



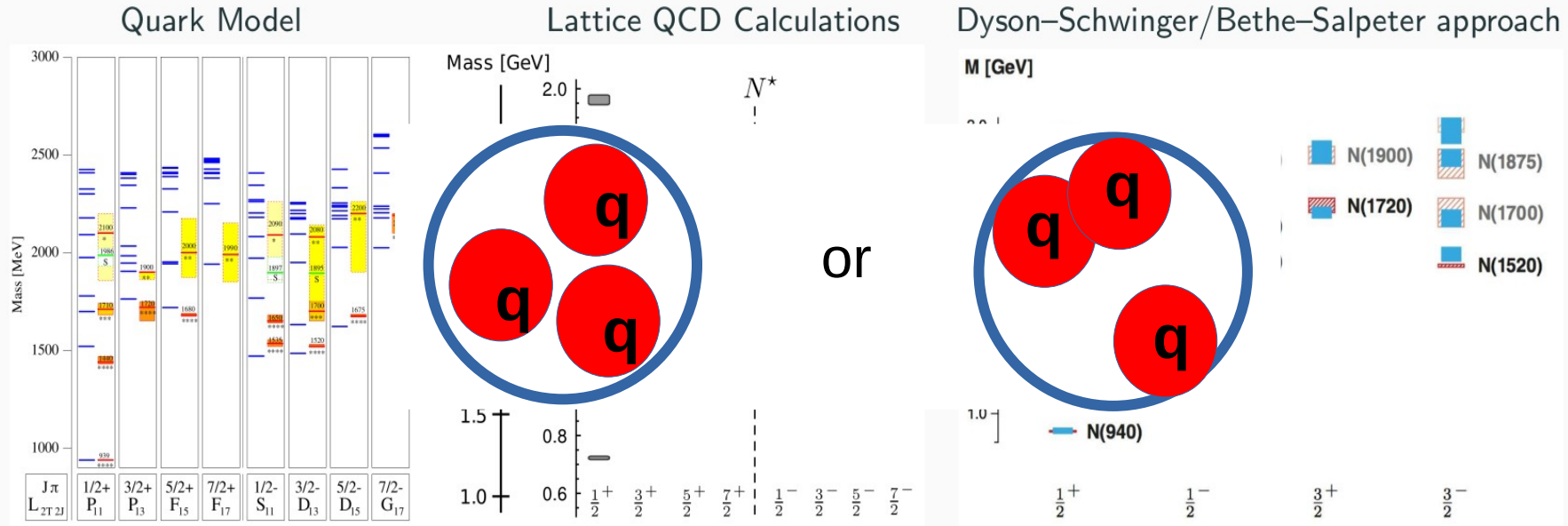
[U. Loering, et al., Eur.Phys.J.A10:395 (2001)]

[R. Edwards et al., Phys.Rev.D 84 (2011) 07450]

[Eichmann, Fischer, Few Body Syst. 60 (2019) 1,2]

- Large inconsistency between calculation and experiment
- Lack of experimental data or wrong theoretical assumptions?

# Missing light baryons



[U. Loering, et al., Eur.Phys.J.A10:395 (2001)]

[R. Edwards et al., Phys.Rev.D 84 (2011) 07450]

[Eichmann, Fischer, Few Body Syst. 60 (2019) 1,2]

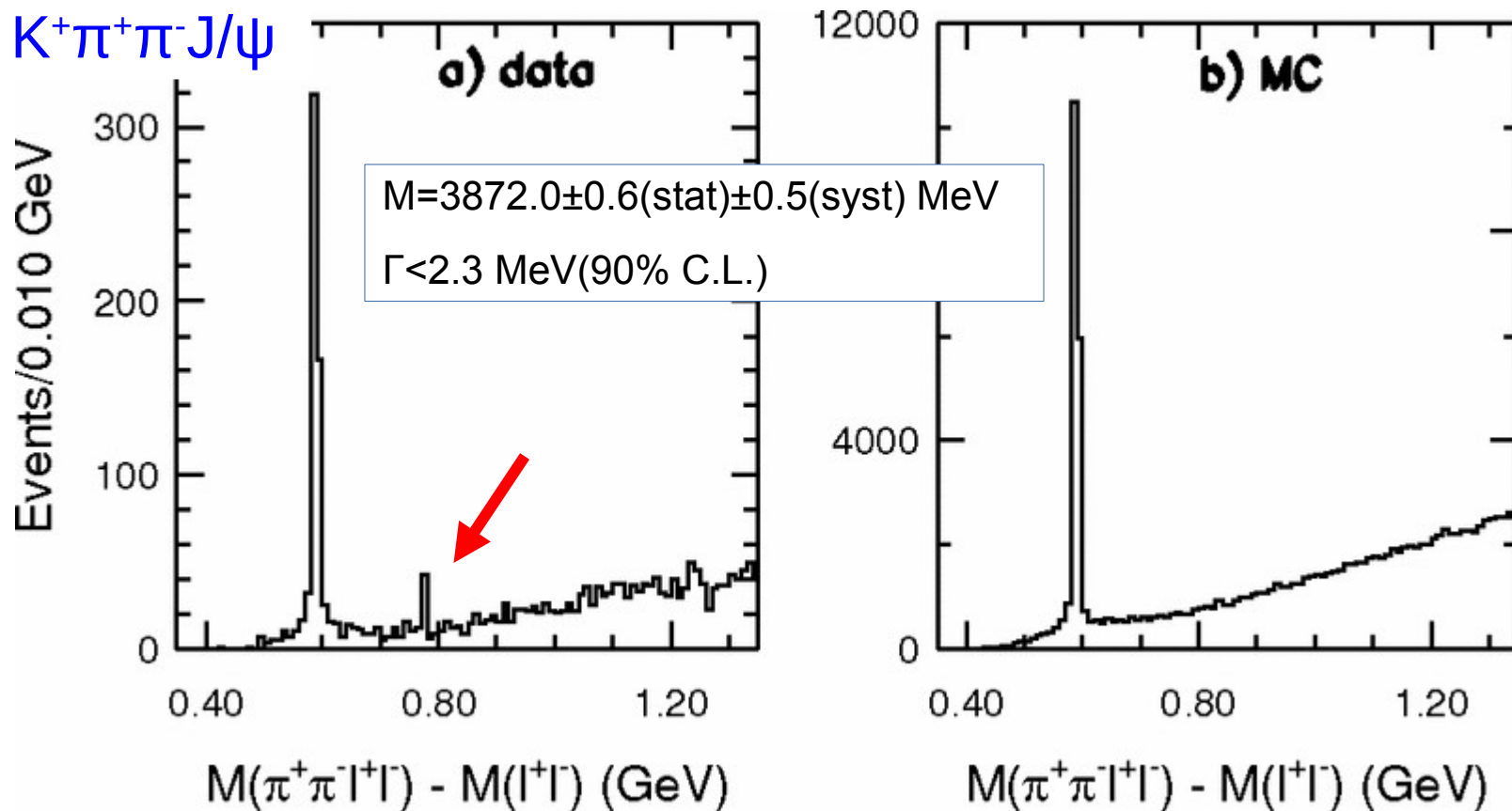
Since diquarks are not color singlets, they cannot exist as free particles. How to identify a diquark?

# Heavy hadrons



# First exotic heavy hadron - X(3872)

$B^+ \rightarrow K^+ \pi^+ \pi^- J/\psi$



S.-K. Choi et al. (Belle Collaboration)  
Phys. Rev. Lett. 91 (2003) 262001

## X(3872) two decades later

- $J^{PC} = 1^{++}$  (LHCb PRL 110, 222001 (2013))
- Mass is within 1 MeV with the mass of  $D^*D$
- Decay rate to  $\omega J/\psi$  and  $\rho J/\psi$  is approximately the same
- Comparable decay rate to  $D^*D$  and  $\gamma\psi(3686)$ : not a  $D^*D$  bound state
- Charged partner is not found

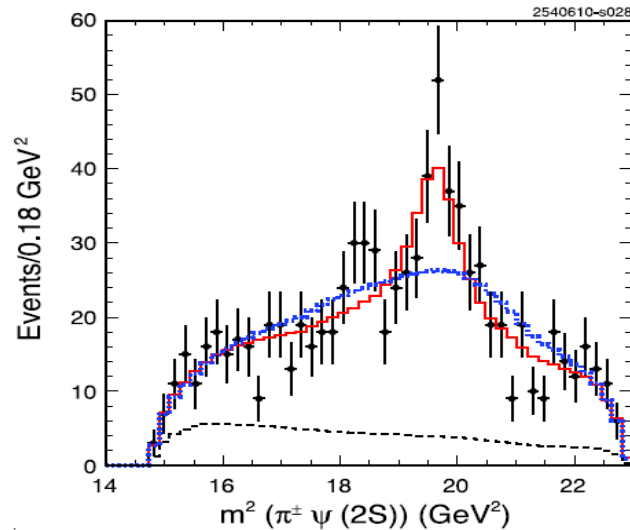
# Charged charmonium-like states

$Z_c^\pm(4430)$ ,  $Z_c^\pm(4050)$ ,  $Z_c^\pm(4250)$

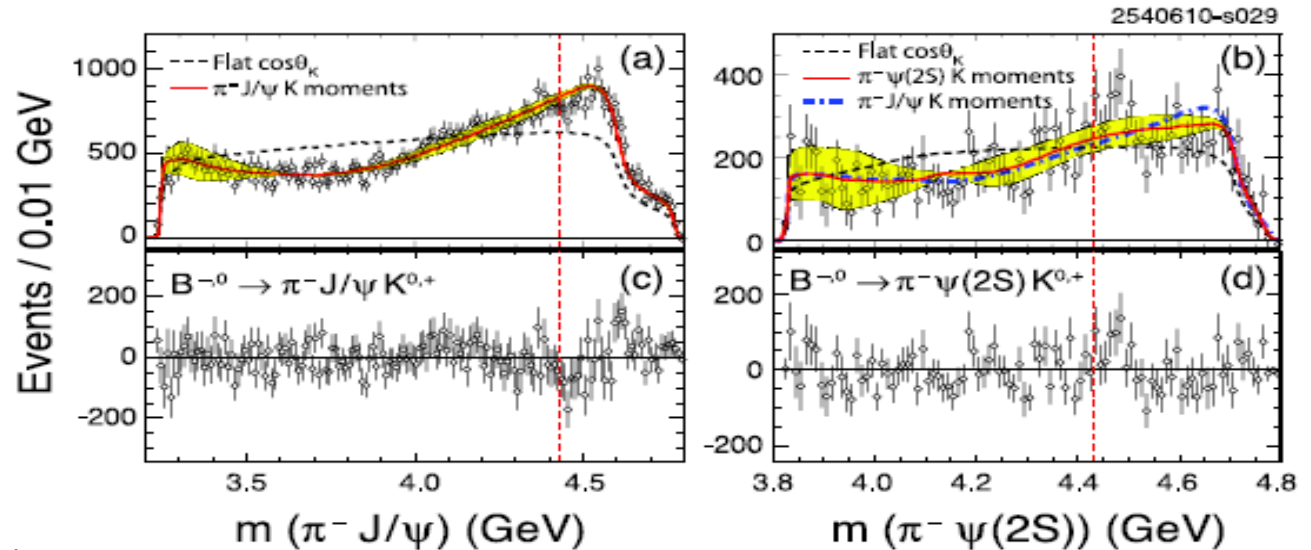
Observed by BELLE in  $B \rightarrow K\pi\psi(2S)$ ,  $B \rightarrow K\pi\chi_{c1}$

Not observed by BaBar in very similar conditions

## BELLE



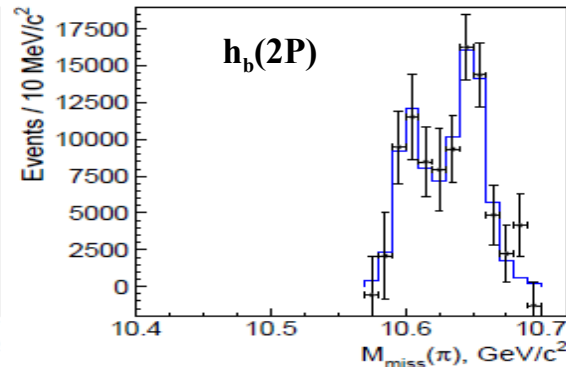
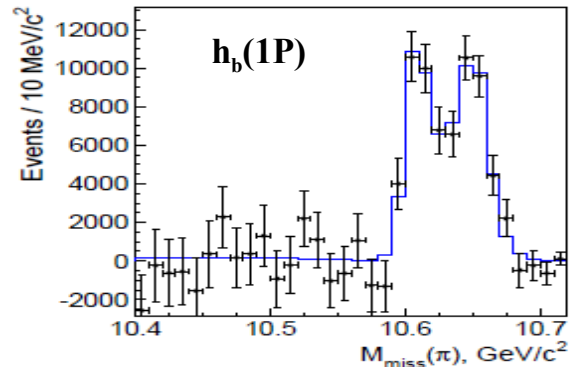
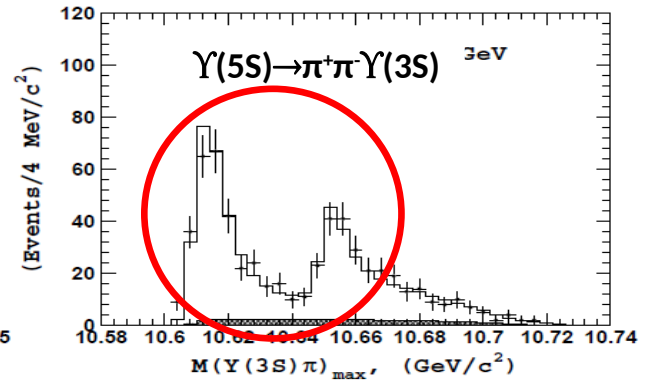
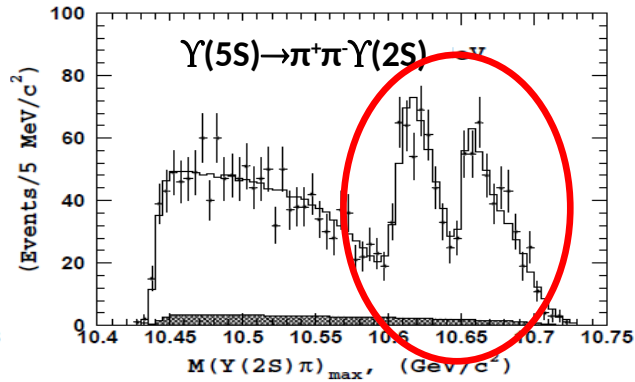
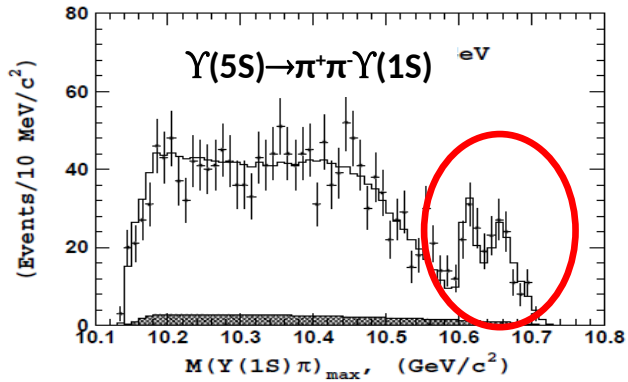
## BaBar



# Observation of charged bottomonium-like state $Z_b(10610)$ and $Z_b(10650)$

$\Upsilon(5S) \rightarrow \pi^+\pi^-\Upsilon(nS), \pi^+\pi^-h_b(nP)$

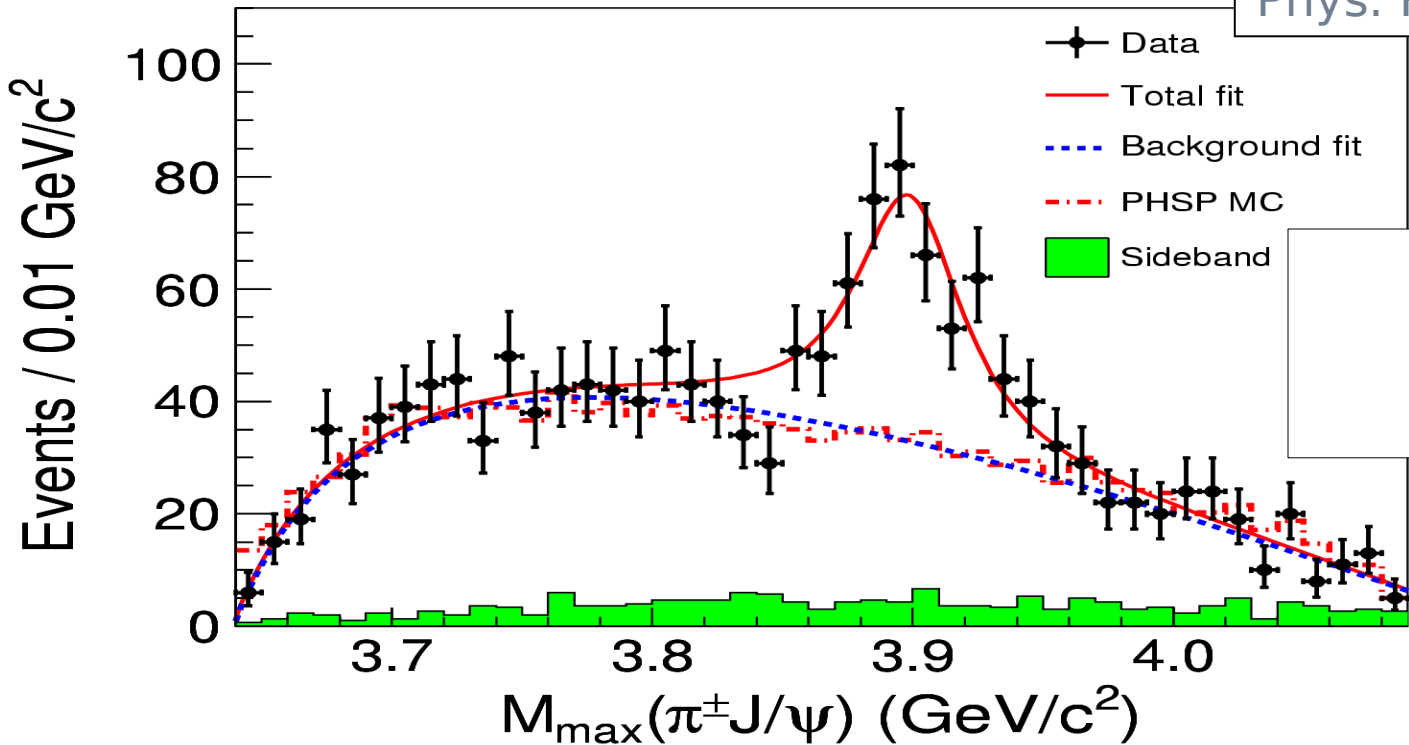
BELLE, PRL 108, 122001 (2012)



# The $Z_c^\pm(3900)$ observaton @ BESII.

$e^+e^- \rightarrow \pi^+\pi^-J/\psi$  at  $\sqrt{s}=4260$  MeV

BESIII: arXiv:1303.5949  
 Phys. Rev. Lett (2013) 252001



Mass =  $(3899.0 \pm 3.6 \pm 4.9)$  MeV  
 Width =  $(46 \pm 10 \pm 20)$  MeV  
 Fraction =  $(21.5 \pm 3.3 \pm 7.5)\%$

# $Z_c(3900)^0$

Observation of neutral  $Z(3900)^0$

Phys. Rev. Lett. 115, 112003 (2015)

**Process:  $e^+e^- \rightarrow \pi^0\pi^0 J/\psi$**

@4.23 GeV ( $1 \text{ fb}^{-1}$ )

@4.26 GeV ( $0.8 \text{ fb}^{-1}$ )

@4.36 GeV ( $0.5 \text{ fb}^{-1}$ )

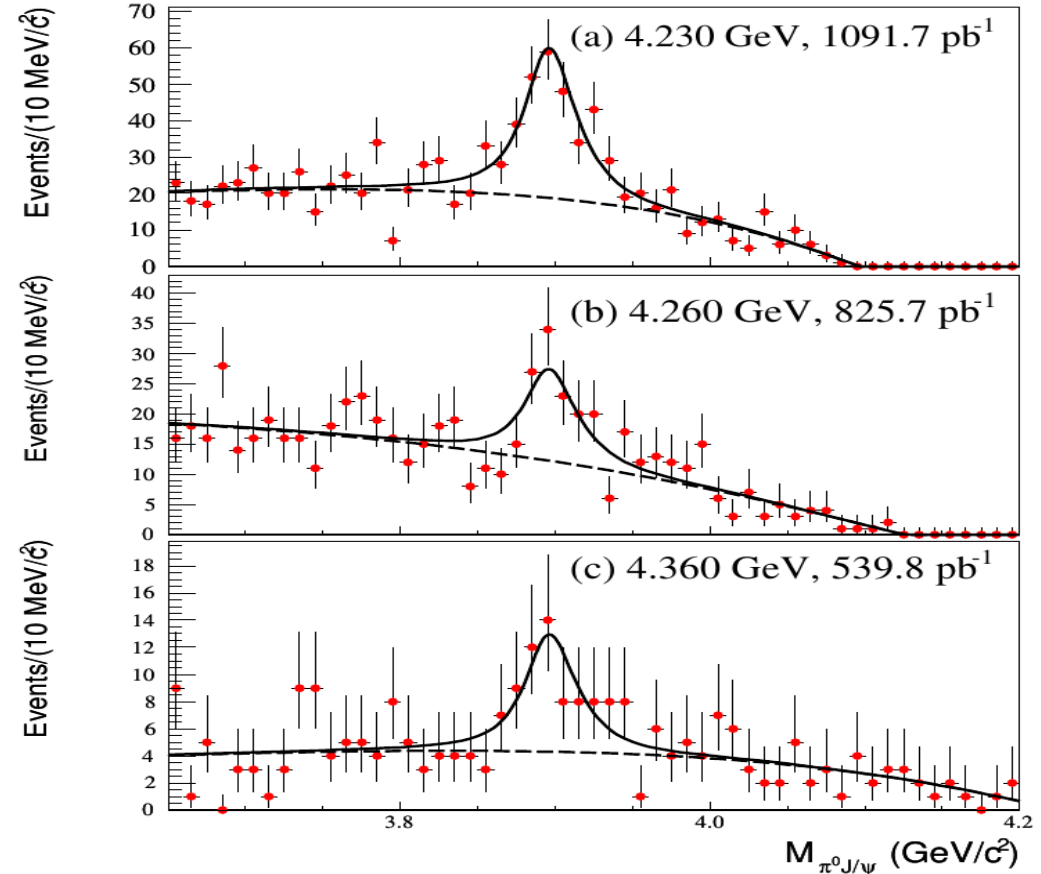
+ 7 small samples 4.19 – 4.42 GeV

$M = 3894.8 \pm 2.3 \pm 3.2 \text{ MeV}/c^2$

$\Gamma = 29.6 \pm 8.2 \pm 8.2 \text{ MeV}$

Statistical significance:  $10.4\sigma$

Interpreted as isospin partner of  $Z(3900)^\pm$



# $Z_c(3885)^\pm$

Phys. Rev. Lett. 112, 022001 (2014)

$Z_c(3900)$  lies  $\sim 20$  MeV above  $D\bar{D}^*$  mass threshold.

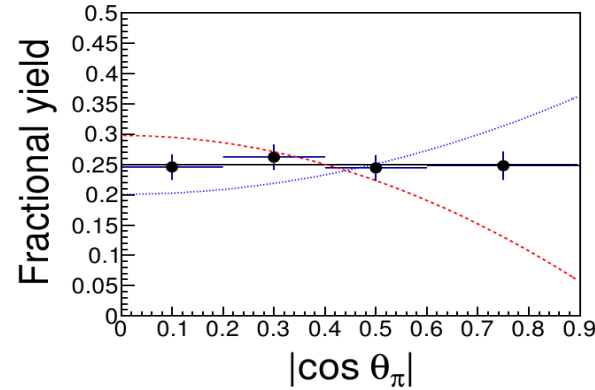
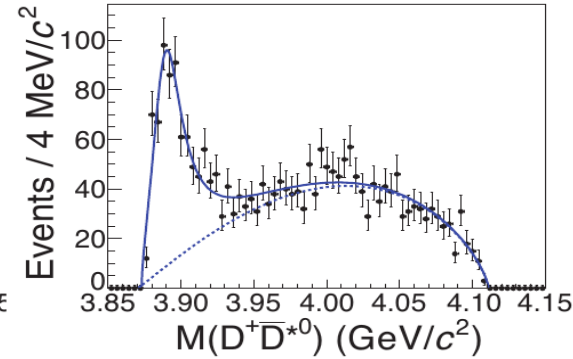
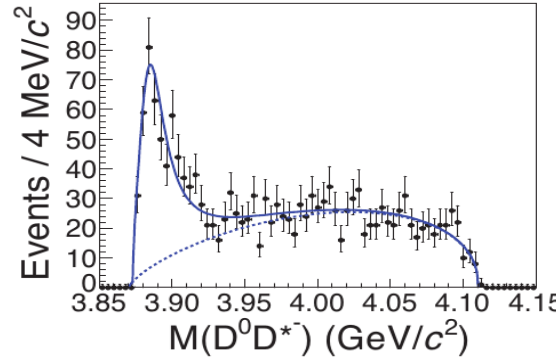
Process:  $e^+e^- \rightarrow (D^0D^{*-})\pi^+ + (D^+\bar{D}^{*0})\pi^-$   
@4.26 ( $0.5 \text{ fb}^{-1}$ )

	$Z_c(3885) \rightarrow D\bar{D}^*$	$Z_c(3900) \rightarrow \pi J/\psi$
Mass ( $\text{MeV}/c^2$ )	$3883.9 \pm 1.5 \pm 4.2$	$3899 \pm 3.6 \pm 4.9$
$\Gamma$ (MeV)	$24.8 \pm 3.3 \pm 11.0$	$46 \pm 10 \pm 20$
$\sigma \times \mathcal{B}$ (pb)	$83.5 \pm 6.6 \pm 22.0$	$13.5 \pm 2.1 \pm 4.8$

Assuming  $Z_c(3885)$  and  $Z_c(3900)$  are one state

$$\frac{\Gamma(Z_c(3885) \rightarrow D\bar{D}^*)}{\Gamma(Z_c(3900) \rightarrow \pi J/\psi)} = 6.2 \pm 1.1 \pm 2.7$$

$$\frac{\mathcal{B}(\psi(3770) \rightarrow D\bar{D})}{\mathcal{B}(\psi(3770) \rightarrow J/\psi \pi^+ \pi^-)} = 482 \pm 84$$



Black:  $1^+$   
Red:  $0^-$   
Blue:  $1^-$

# $Z_c(3885)^0$

Phys. Rev. Lett. 115, 222002 (2015)

Process:  $e^+e^- \rightarrow (D^+D^{*-})^0\pi^0 + (D^0\bar{D}^{*0})^0\pi^0$

@4.23 ( $1.1 \text{ fb}^{-1}$ )

@4.26 ( $0.8 \text{ fb}^{-1}$ )

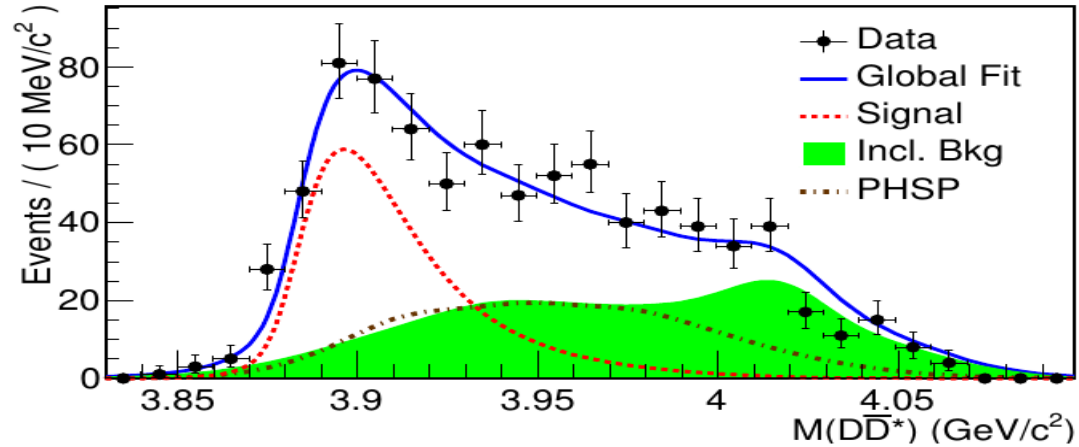
$M = 3885.7_{-5.7}^{+5.7} \pm 8.4 \text{ MeV}/c^2$

$\Gamma = 47 \pm 9 \pm 10 \text{ MeV}$

Born cross section consistent for  
 $e^+e^- \rightarrow Z_c \pi^0 \rightarrow (D\bar{D}^*)^0\pi^0 + \text{c.c.}$

is consistent with half of

$e^+e^- \rightarrow Z_c^+\pi^- \rightarrow (D\bar{D}^*)^+\pi^- + \text{c.c.}$



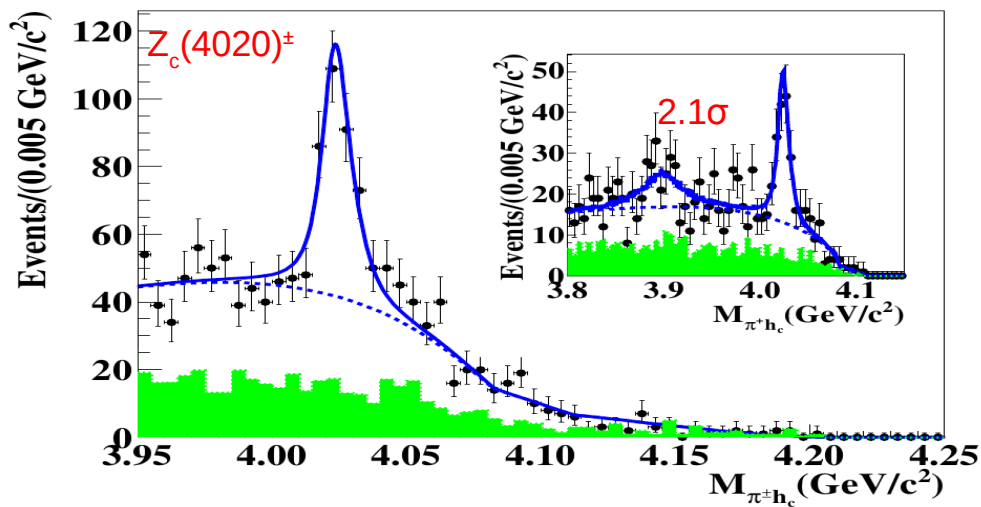


# $Z_c(4020)^\pm$ and $Z_c(4020)^0$

$e^+e^- \rightarrow \pi^+\pi^-h_c$  and  $\pi^0\pi^0h_c$

$h_c$  reconstructed through E1 transition  $h_c \rightarrow \gamma\eta_c$ , reconstructed from 16 exclusive hadronic modes.

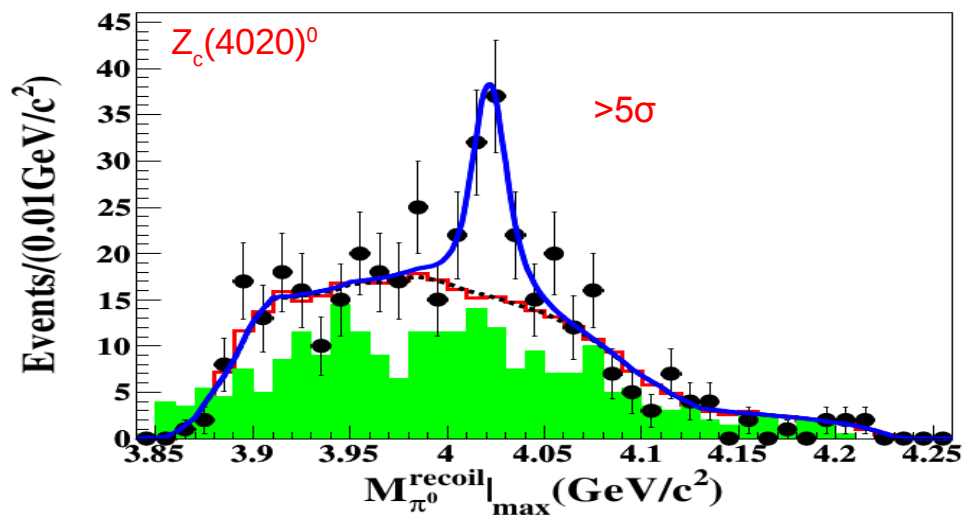
$\sqrt{s} = 4.23, 4.26, \text{ and } 4.36 \text{ GeV}$



Phys.Rev.Lett.111, 242001 (2013)

$M=4022.9 \pm 0.8 \pm 2.7 \text{ MeV}/c^2$   
 $\Gamma = 7.9 \pm 2.7 \pm 2.6 \text{ MeV}$

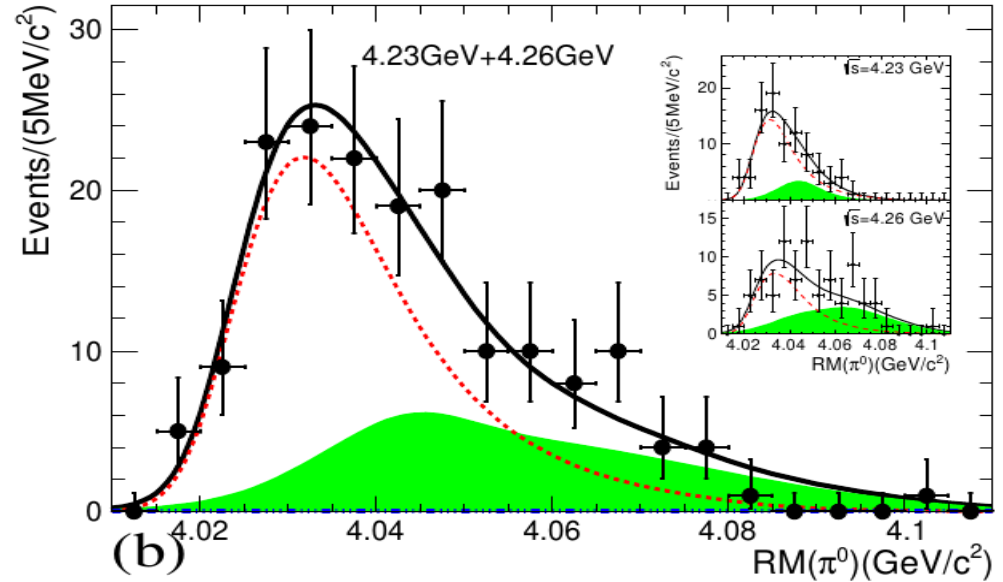
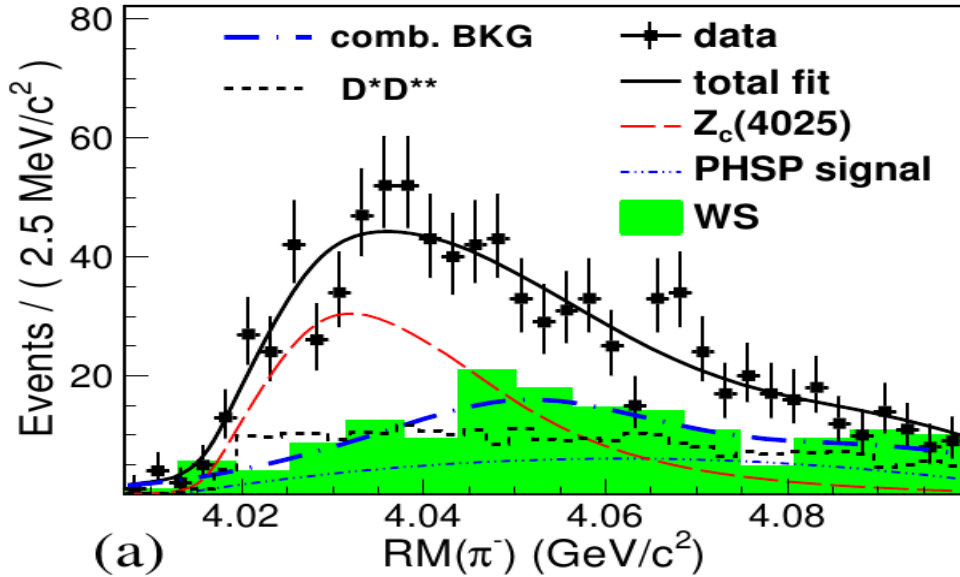
Close to  $D^*\bar{D}^*$  threshold



Phys.Rev.Lett.113.212002(2014)

$M=4023.9 \pm 2.2 \pm 3.8 \text{ MeV}/c^2$   
 Fixed  $\Gamma$

# $Z_c(4025)^\pm$ and $Z_c(4025)^0$



Phys. Rev. Lett. 112, 132001 (2014)

$e^+e^- \rightarrow (D^*\bar{D}^*)^+\pi^- + \text{c.c.}$

$0.8 \text{ fb}^{-1}$  @ 4.26 GeV

Partial reconstruction technique

$M = 4026.3 \pm 2.6 \pm 3.7 \text{ MeV}/c^2$

$\Gamma = 24.8 \pm 5.6 \pm 7.7 \text{ MeV}$

Phys. Rev. Lett. 115, 182002 (2015)

$e^+e^- \rightarrow (D^*\bar{D}^*)^0\pi^0$

$1.1 \text{ fb}^{-1}$  @ 4.23 and  $0.8 \text{ fb}^{-1}$  @ 4.26 GeV

$M = (4025.5^{+2.0}_{-4.7} \pm 3.1) \text{ MeV}/c^2$

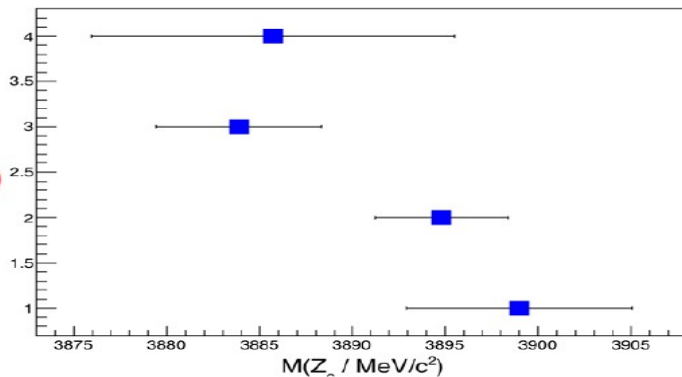
$\Gamma = (23.0 \pm 6.0 \pm 1.0) \text{ MeV}$

$\Gamma =$

# Summary on $Z_c$

### Mass

$Z_c(3900)$



Graph

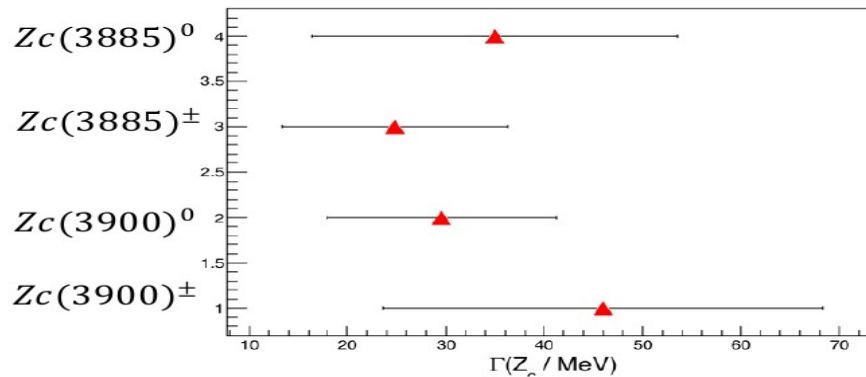
### Width

$Z_c(3885)^0$

$Z_c(3885)^\pm$

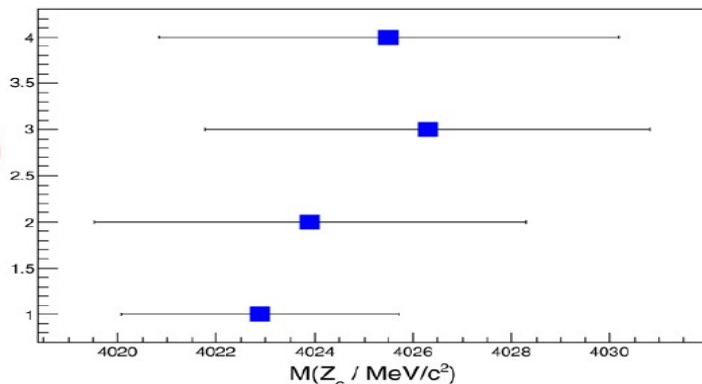
$Z_c(3900)^0$

$Z_c(3900)^\pm$



Graph

$Z_c(4020)$

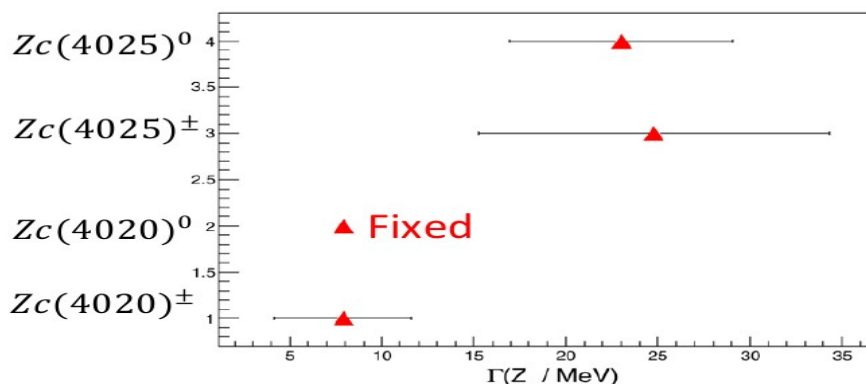


$Z_c(4025)^0$

$Z_c(4025)^\pm$

$Z_c(4020)^0$

$Z_c(4020)^\pm$

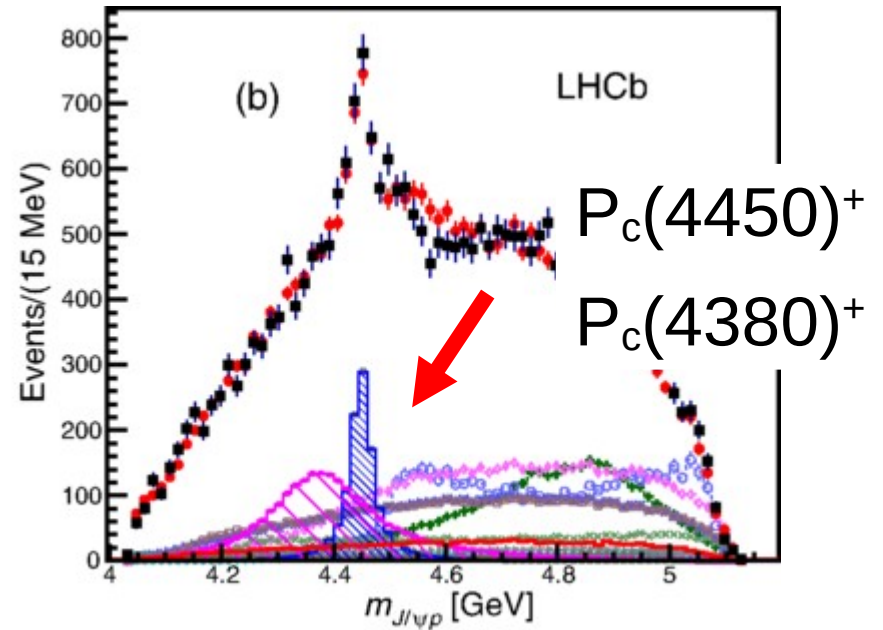
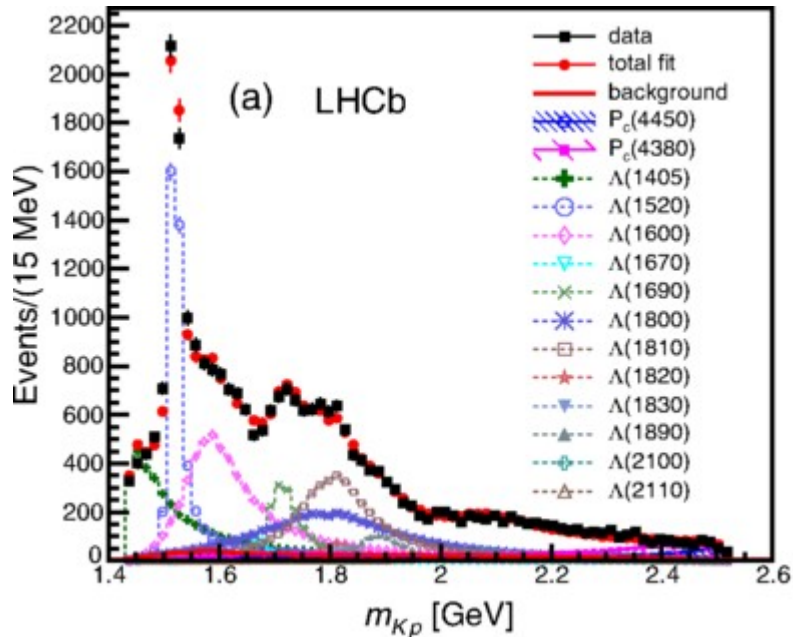
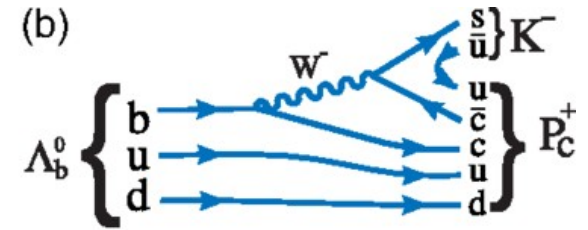
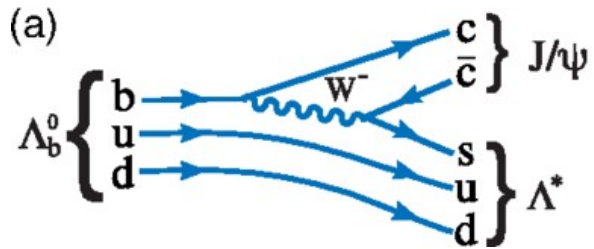


# Question

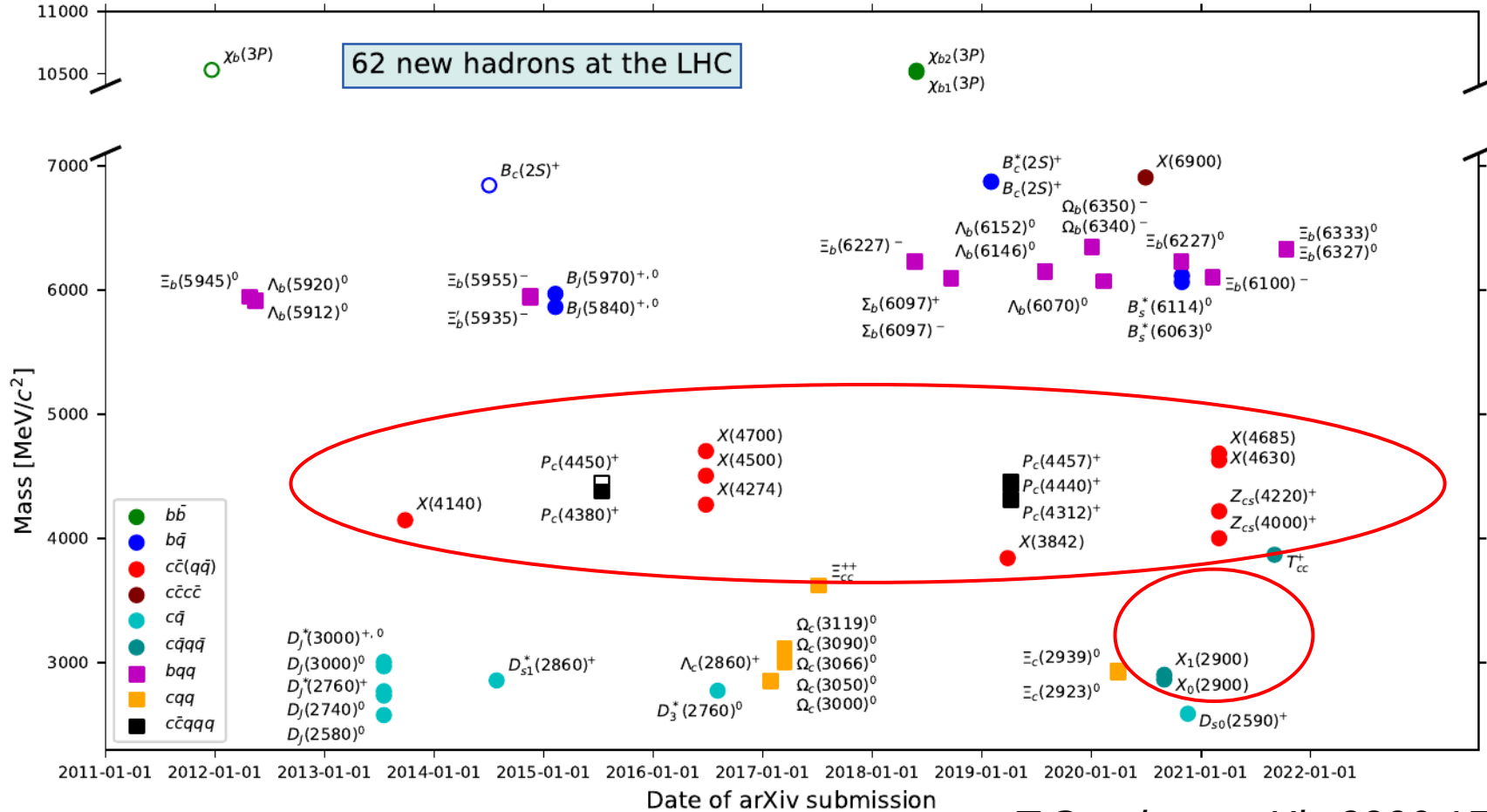
Is it possible to look for a similar state with a strange quarkonium?

# Pentaquarks at LHC

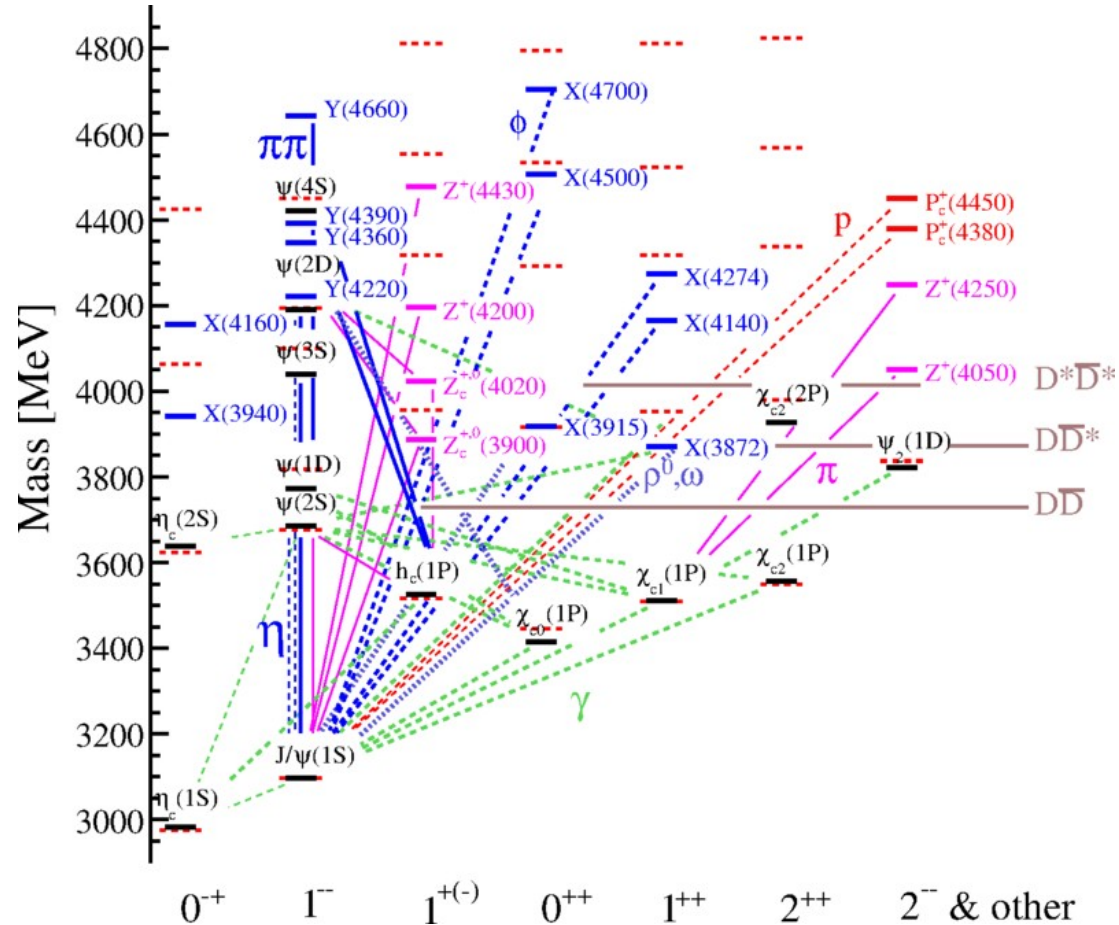
R. Aaij et al. (LHCb Coll.) Phys. Rev. Lett. 115 (2015) 072001



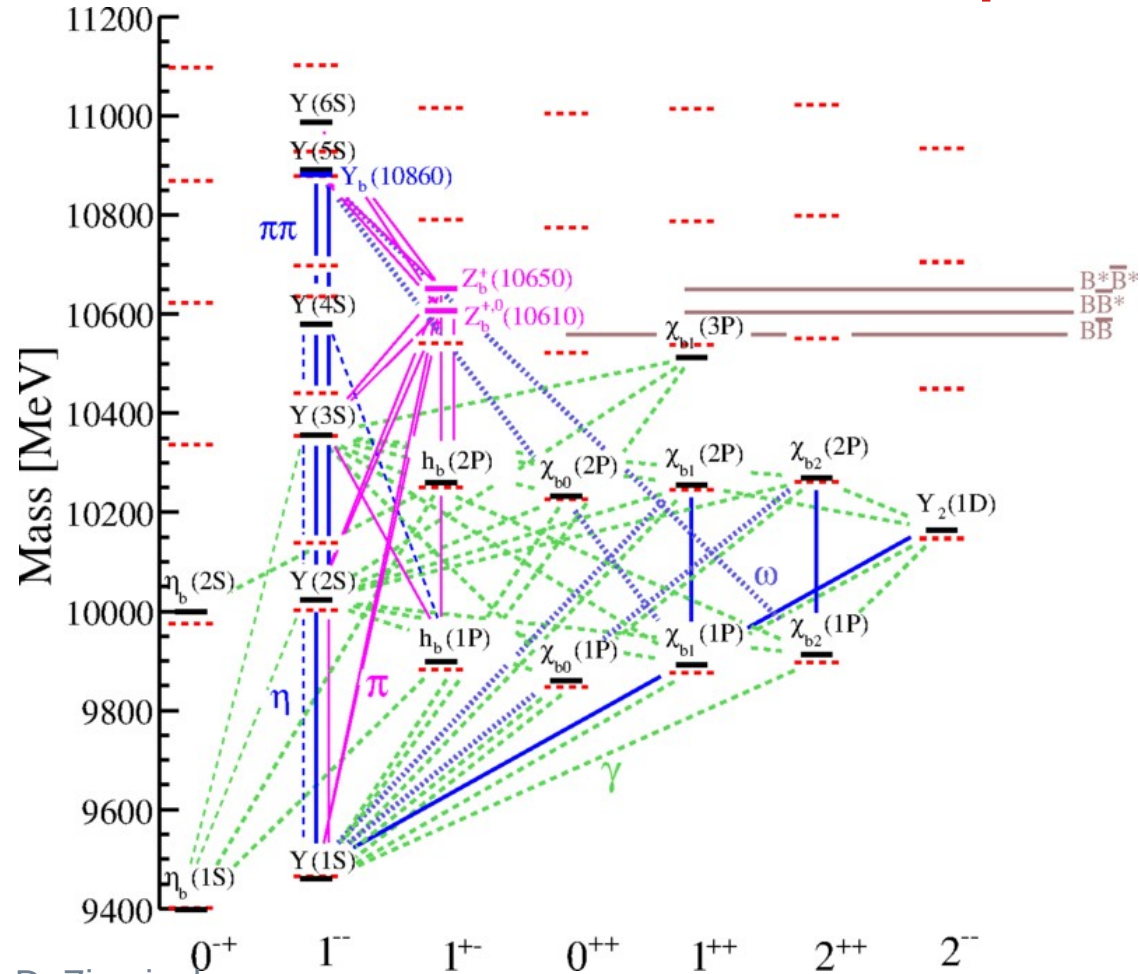
# New hadrons at LHC



# Modern charmonium-like spectrum



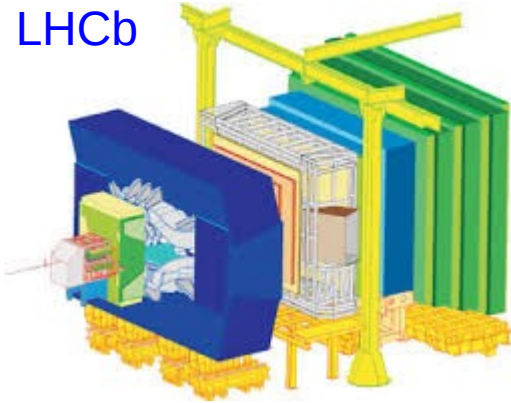
# Modern bottomonium-like spectrum



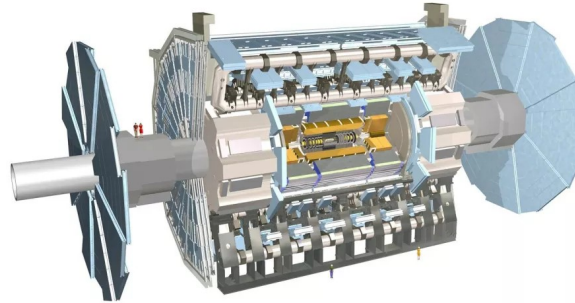


# Experimental landscape

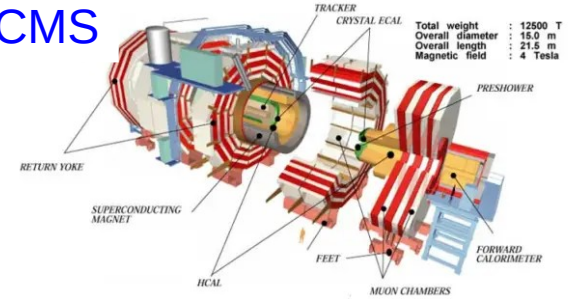
LHCb



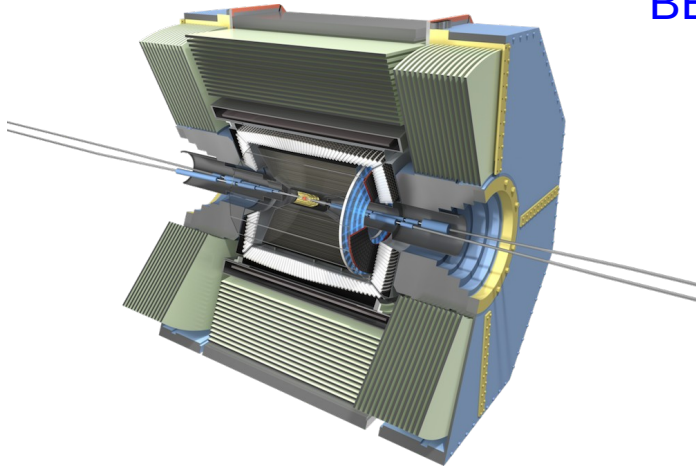
ATLAS



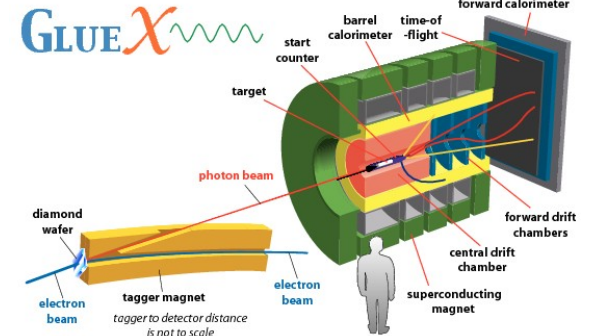
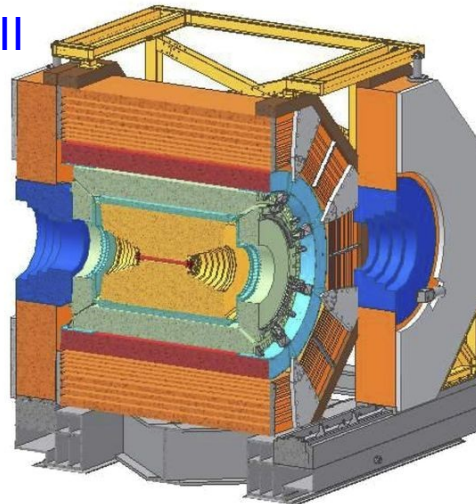
CMS



BELLEII

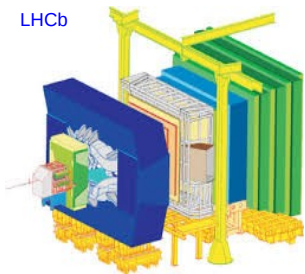


BESIII

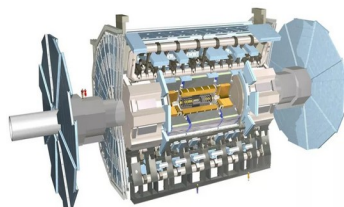


# Experimental landscape (near future)

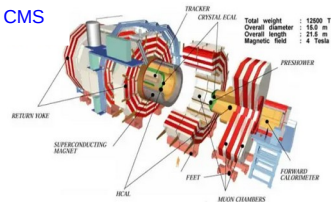
LHCb



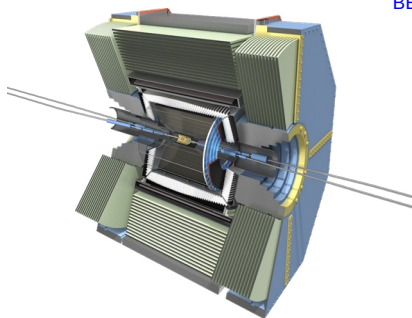
ATLAS



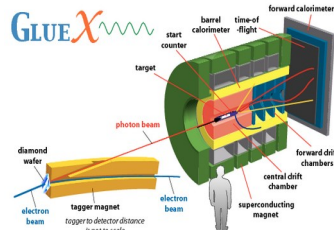
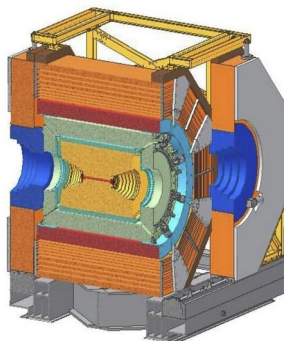
CMS



BELLEII



BESIII

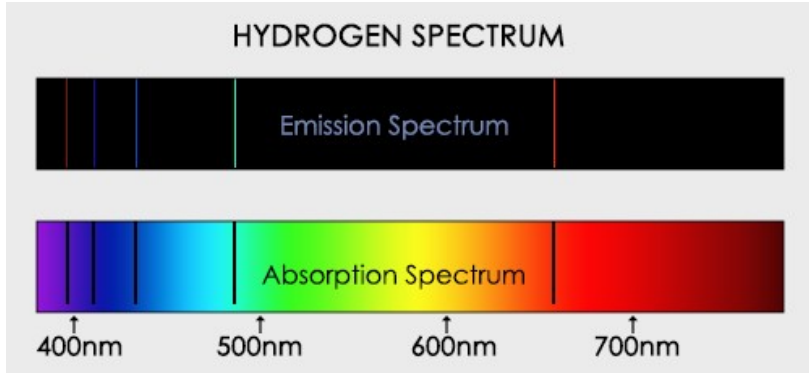


+

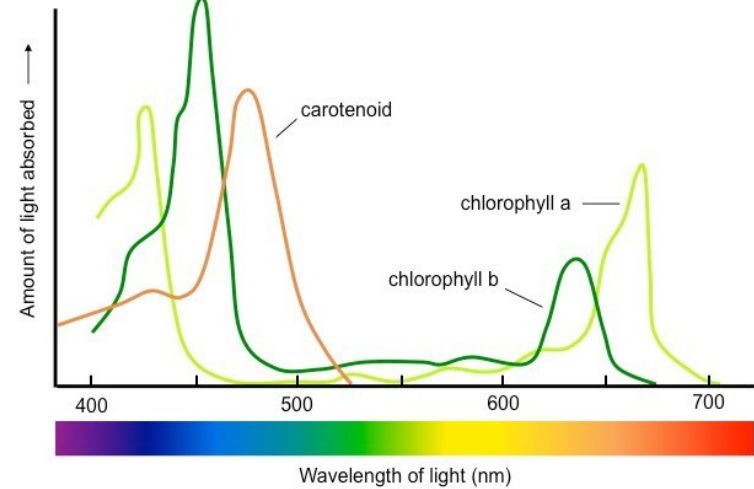
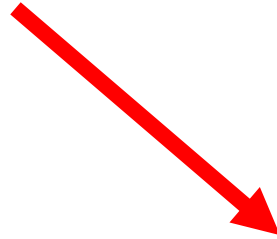
PANDA  
AMBER  
EIC  
SCTF

# Summary

- Hadron spectroscopy is a unique tool to gain knowledge about the intrinsic properties and the composition of hadrons
- Discovery of a large number of multiquark states is the most interesting event since  $J/\psi$  observation
- More precise experimental data are needed to understand XYZ states
- Surplus light mesons and missing light baryons are still there
- Lattice QCD is making a good progress, but theory describing all experimental findings is not yet built



A lot of exciting things ahead of us!



Thank you for your attention!