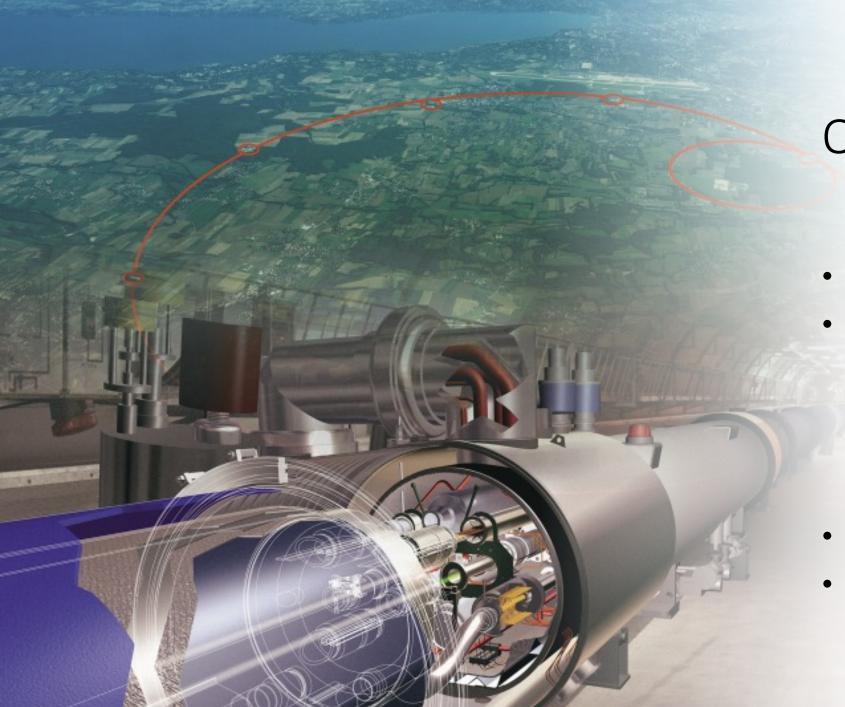
Prospects: LHC in Run 3 and the HL-LHC

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2022 ASIA EUROPE PACIFIC SCHOOL OF HIGH-ENERGY PHYSICS



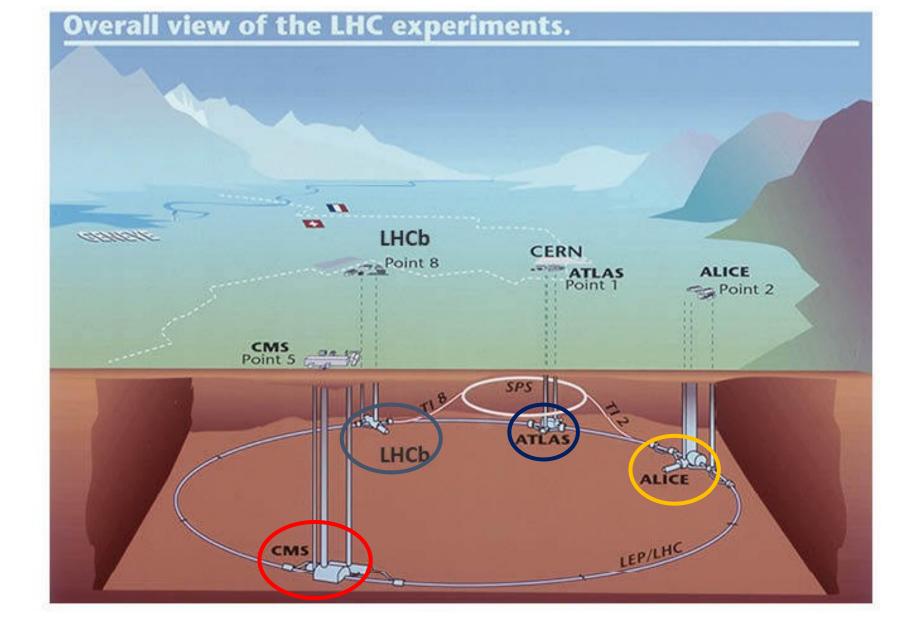


Outline

- Introduction
- Perspectives
 - Energy
 - Luminosity
 - Technology
 - Techniques
- Run 3
- High-Luminosity LHC (HL-LHC)

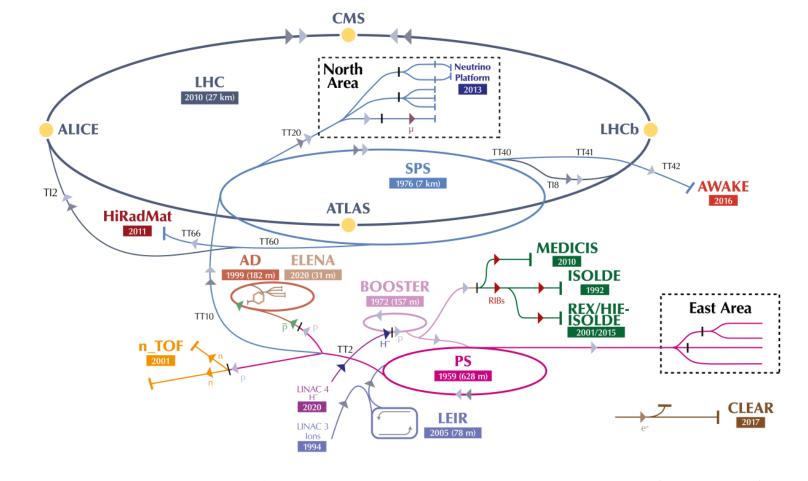
Motivation

- You've heard over the past week about the
 - Impressive performance of the standard model (SM), and what it lacks
 - Status of SM measurements
 - most are in line with predictions. We have a few several-sigma discrepancies that need to be chased down, particularly in flavor physics, but no "smoking gun"
 - The need to improve the precision of our measurements of the Higgs boson
 - Motivation for extending searches for BSM to higher-mass particles and lower cross-section processes
- We need to fully exploit the Run 3 and HL-LHC datasets, and provide direction for our field's next step



I'll focus primarily on ATLAS and CMS, but say a few words about ALICE and LHCb

The CERN accelerator complex Complexe des accélérateurs du CERN

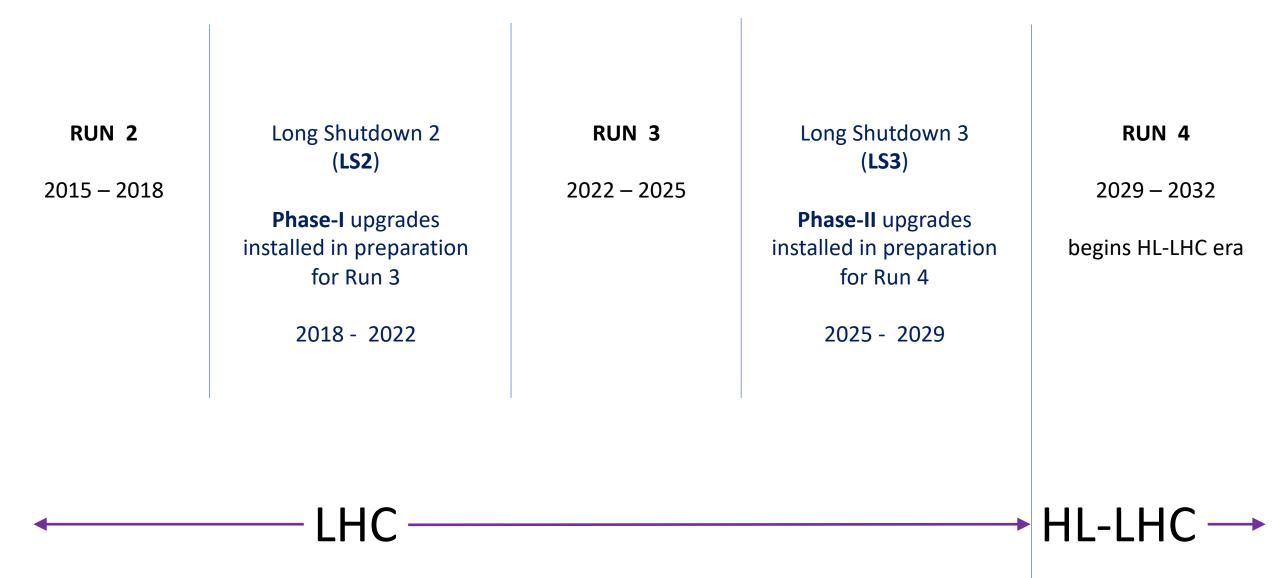


 \downarrow H⁻ (hydrogen anions) \downarrow p (protons) \downarrow ions \downarrow RIBs (Radioactive Ion Beams) \downarrow n (neutrons) \downarrow p (antiprotons) \downarrow e⁻ (electrons) \downarrow μ (muons)

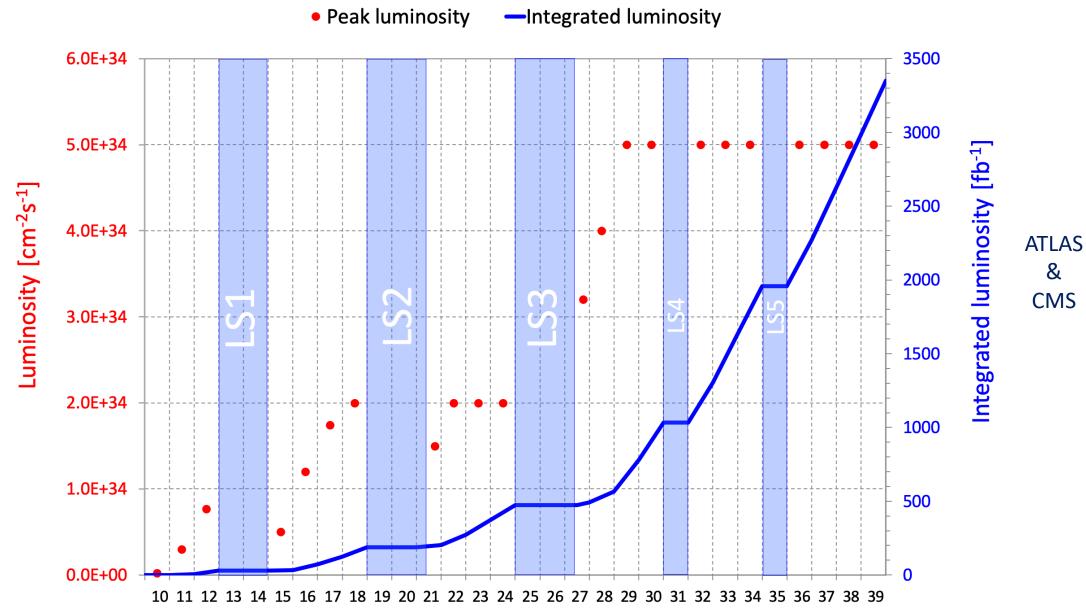
LHC - Large Hadron Collider // SPS - Super Proton Synchrotron // PS - Proton Synchrotron // AD - Antiproton Decelerator // CLEAR - CERN Linear Electron Accelerator for Research // AWAKE - Advanced WAKefield Experiment // ISOLDE - Isotope Separator OnLine // REX/HIE-ISOLDE - Radioactive EXperiment/High Intensity and Energy ISOLDE // MEDICIS // LEIR - Low Energy Ion Ring // LINAC - LINear ACcelerator // n_TOF - Neutrons Time Of Flight // HiRadMat - High-Radiation to Materials // Neutrino Platform

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A cheat sheet for recent LHC-schedule-related terminology....

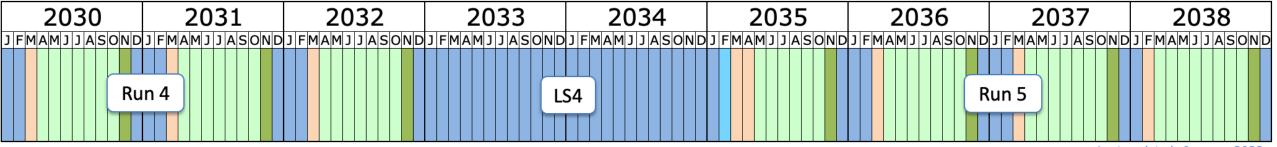


NOTE: LS3 has been shifted one year forward (see next slide)



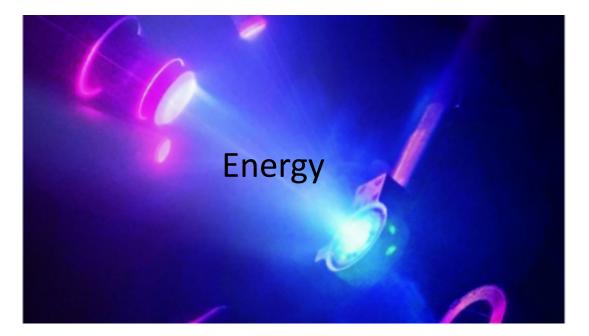
The LHC Schedule is DYNAMIC so take schedules as current drafts



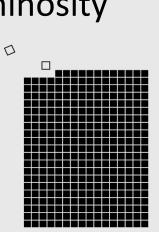


Last updated: January 2022

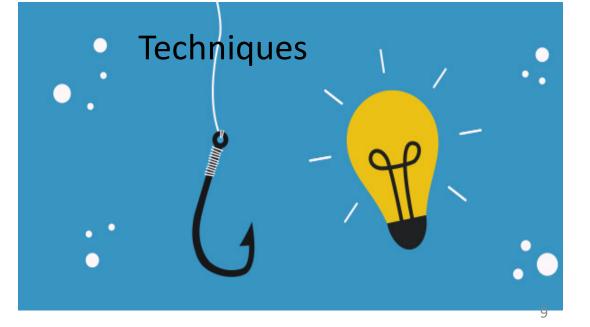
Shutdown/Technical stop Protons physics Ions Commissioning with beam Hardware commissioning/magnet training



Luminosity \diamond \Diamond D



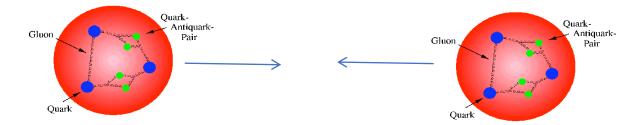




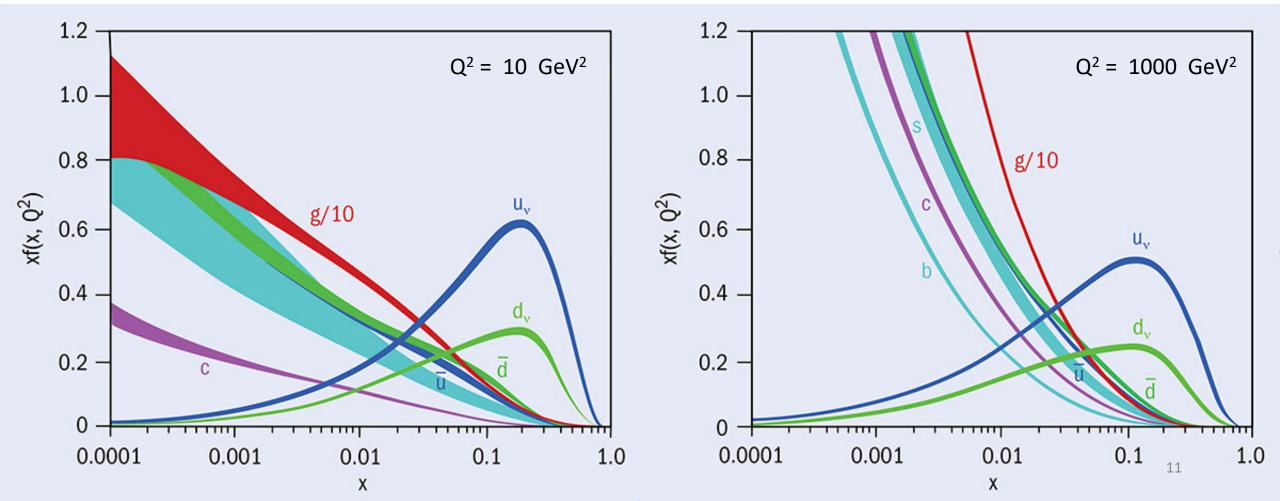


The Energy Frontier

Available Energy

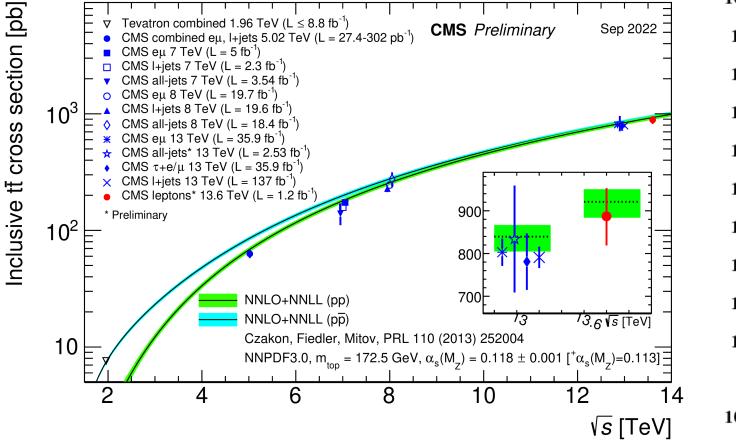


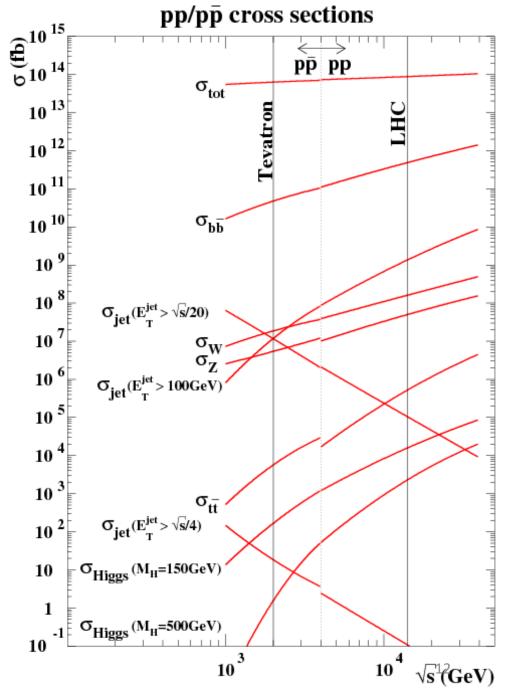
Protons are not fundamental particles. It's their constituents that collide. This explains our need for huge datasets and need to save only a small fraction of collisions we see...



Energy and Cross Section

High-mass particle cross sections rise more quickly with center-of-mass energy.



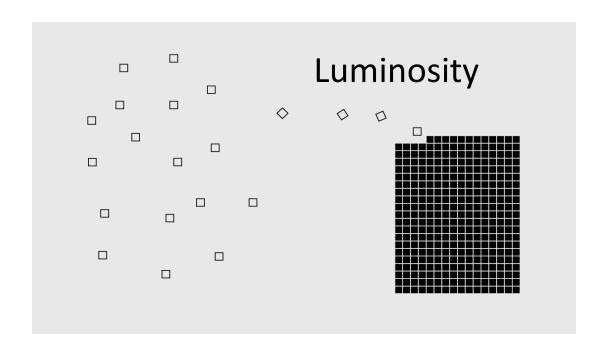


Higgs Boson Cross Sections

\sqrt{s}	σ	$\delta(ext{theory})$	$\delta(ext{PDF})$	$\delta(lpha_s)$
13 TeV	48.61 pb	+2.08 pb +4.27% -3.15 pb +4.27% -6.49%	$\left(\pm 0.89 { m pb} \left(\pm 1.85\% \right) ight)$	$^{+1.24 \mathrm{pb}}_{-1.26 \mathrm{pb}} \left(^{+2.59\%}_{-2.62\%} ight)$
			$\left({45} ight) \pm 1.00 { m pb} \left(\pm 1.85\% ight)$	

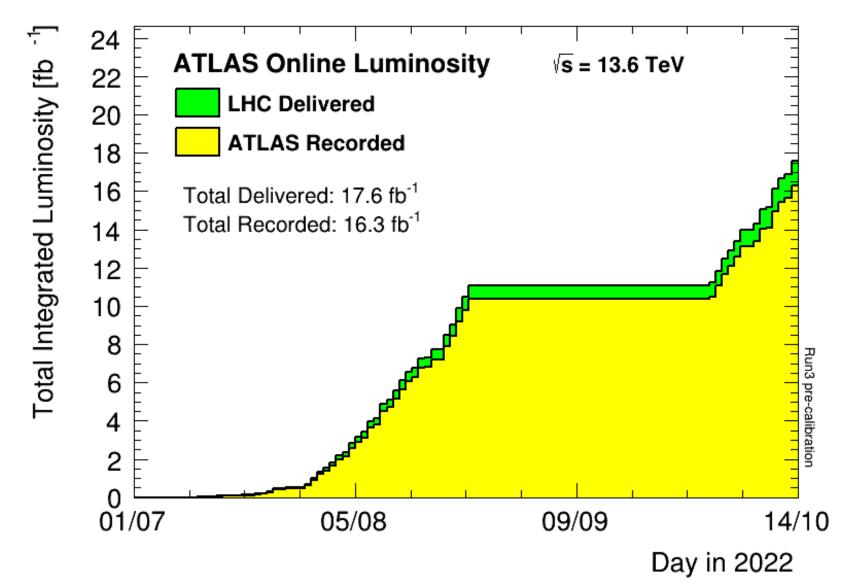
\sqrt{s} [TeV]	$\sigma^{ m VBF}$ [fb]	$\Delta_{ m scale}$ [%]	$\Delta_{\mathrm{PDF}\opluslpha_{\mathrm{s}}}$ [%]	$\sigma_{ m NNLO}^{ m DIS}$ [fb]	$\delta_{ m ELWK}$ [%]	σ_γ [fb]	$\sigma_{s\text{-ch}}$ [fb]
13	3766	$^{+0.43}_{-0.33}_{+0.45}$	± 2.1	3939	-5.3	35.3	1412
14	4260	$+0.45 \\ -0.34$	± 2.1	4460	-5.4	40.7	1555

ggH (top) and VBF (bottom) Cross Section at 13 and 14 TeV uncertainties provided as a function of the pp collider energy (From <u>CERN Yellow Report</u>)



Integrated Luminosity (data accumulated) enables

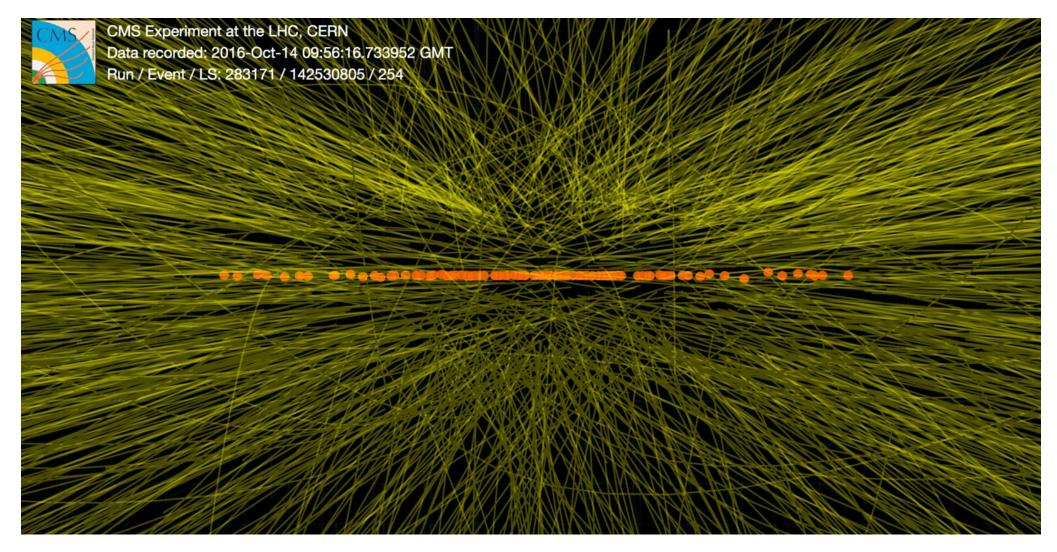
- Accessing rare decay modes
- Increasing precision of measurements
- Opening up analysis options



15

High Instantaneous Luminosity (rate) challenges

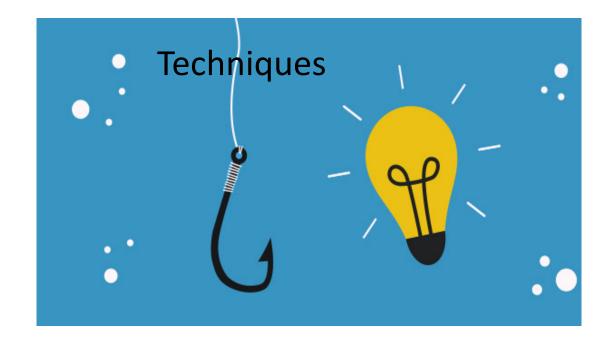
- detectors
- data acquisition systems
- analysis techniques



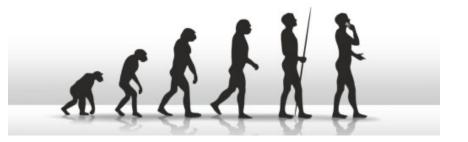


Enable or improve measurements & searches

- We need to answer the challenges of higher instantaneous luminosities and more sophisticated data analysis techniques with detector design choices.
- Improvements include
 - Higher granularity detectors to handle higher particle densities
 - Electronics that are more radiation hard
 - Lower material budgets, particularly for un-instrumented components such as support structures
 - Making additional information available, such as high-precision timing
 - Better switches and connections for detector readout
 - Making critical information, such as tracking, available earlier in trigger chain
 - Moving analysis to the trigger step

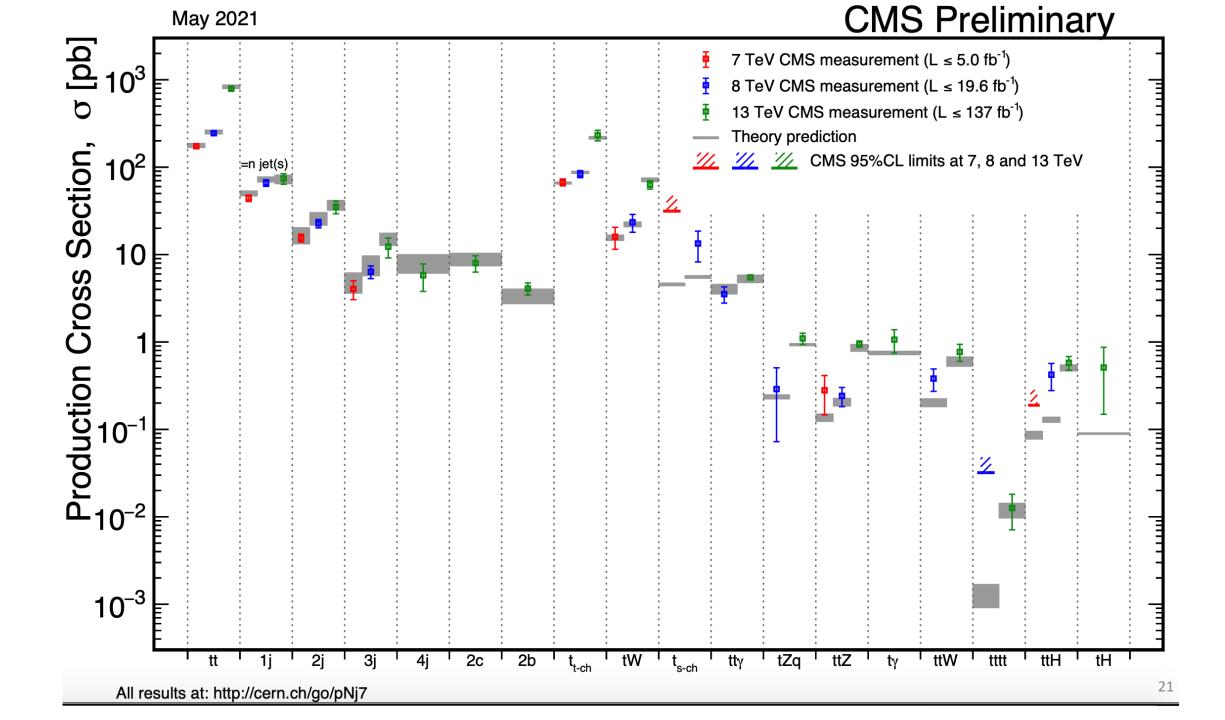


AI, AI Everywhere!



cuts -> likelihoods -> BDT -> NN -> Deep Learning

- More challenging data-taking environments require, and more sophisticated detectors enable, the use of machine-learning in many physics analyses
- Special training is required so we know how to:
 - choose an algorithm or strategy
 - optimize an algorithm
 - interpret results, including the evaluation of systematic uncertainties
- Use cases are ever increasing!
 - We've moved beyond simple object identification or classification
 - Extract sensitivity from fit to ML output rather than a kinematic variable like mass
 - Object calibration (such as energy scale determination)



ATLAS Heavy Particle Searches* - 95% CL Upper Exclusion Limits Status: July 2022

ATLAS Preliminary

 $\int \mathcal{L} dt = (3.6 - 139) \text{ fb}^{-1}$ $\sqrt{s} = 8, 13 \text{ TeV}$

	Model	ℓ, γ Jets	· E _T ^{miss} ∫	£ dt[fb	⁻¹] Limit	j - · · · · ·	.0 100/10	Reference
Extra dimensions	ADD $G_{KK} + g/q$ ADD non-resonant $\gamma\gamma$ ADD QBH ADD BH multijet RS1 $G_{KK} \rightarrow \gamma\gamma$ Bulk RS $G_{KK} \rightarrow WW/ZZ$ Bulk RS $g_{KK} \rightarrow WV \rightarrow \ell \nu qq$ Bulk RS $g_{KK} \rightarrow tt$ 2UED / RPP	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	 J Yes J/2j Yes	139 36.7 139 3.6 139 36.1 139 36.1 36.1	g _{KK} mass	9.4 TeV	$ \begin{array}{l} n=2 \\ n=3 \; \text{HLZ NLO} \\ n=6 \\ n=6, M_D=3 \; \text{TeV, rot BH} \\ k/\overline{M}_{PI}=0.1 \\ k/\overline{M}_{PI}=1.0 \\ r/m=15\% \\ \text{Tier (1,1), } \mathcal{B}(A^{(1.1)} \rightarrow tt)=1 \end{array} $	2102.10874 1707.04147 1910.08447 1512.02586 2102.13405 1808.02380 2004.14636 1804.10823 1803.09678
Gauge bosons	$\begin{array}{l} \mathrm{SSM}\; Z' \to \ell\ell \\ \mathrm{SSM}\; Z' \to \tau\tau \\ \mathrm{Leptophobic}\; Z' \to bb \\ \mathrm{Leptophobic}\; Z' \to tt \\ \mathrm{SSM}\; W' \to \ell\nu \\ \mathrm{SSM}\; W' \to \tau\nu \\ \mathrm{SSM}\; W' \to tb \\ \mathrm{HVT}\; W' \to WZ \to \ell\nu q\ell \ \ell' \ \mathrm{mode} \\ \mathrm{HVT}\; W' \to WZ \to \ell\nu \ell\ell' \ \mathrm{mode} \\ \mathrm{HVT}\; W' \to WH \to \ell\nu bb \ \mathrm{mode} \\ \mathrm{HVT}\; W' \to ZH \to \ell\ell \nu\nu bb \ \mathrm{mode} \\ \mathrm{HVT}\; W_R \to \mu N_R \end{array}$	leIC 3 <i>e</i> ,μ 2j(VBI IB 1 <i>e</i> ,μ 1-2b,1-	- Yes Yes J - J Yes -) Yes 0 j Yes	139 36.1 139 139 139 139 139 139 139 139 139 80	Z' mass Z' mass Z' mass W' mass W' mass W' mass W' mass W' mass W' mass W' mass UV mass W mass W mass UV m UV	5.1 TeV 2.42 TeV 2.1 TeV 4.1 TeV 5.0 TeV 4.4 TeV 4.3 TeV 3.3 TeV 3.2 TeV 5.0 TeV	$\Gamma/m = 1.2\%$ $g_V = 3$ $g_V c_H = 1, g_F = 0$ $g_V = 3$ $g_V = 3$ $g_V = 3$ $m(N_R) = 0.5 \text{ TeV}, g_L = g_R$	1903.06248 1709.07242 1805.09299 2005.05138 1906.05609 ATLAS-CONF-2021-025 ATLAS-CONF-2021-043 2004.14636 ATLAS-CONF-2022-005 2207.00230 2207.00230 1904.12679
C	Cl qqqq Cl llqq Cl eebs Cl µµbs Cl tttt	$\begin{array}{cccc} - & 2 j \\ 2 e, \mu & - \\ 2 e & 1 b \\ 2 \mu & 1 b \\ \geq 1 e, \mu & \geq 1 b, \geq \end{array}$	- -	37.0 139 139 139 36.1		8 TeV 2.0 TeV 2.57 TeV	$\begin{array}{c} \textbf{21.8 TeV} & \eta_{LL} \\ \textbf{35.8 TeV} & \eta_{LL} \\ \textbf{g}_* = 1 \\ \textbf{g}_* = 1 \\ \textbf{C}_{4t} = 4\pi \end{array}$	1703.09127 2006.12946 2105.13847 2105.13847 1811.02305
MQ	Axial-vector med. (Dirac DM) Pseudo-scalar med. (Dirac DM) Vector med. Z'-2HDM (Dirac D Pseudo-scalar med. 2HDM+a	M) 0 e, μ 2 b		139 139 139 139	m _{med} 376 GeV	2.1 TeV 3.1 TeV	$\begin{array}{l} g_q = 0.25, g_\chi = 1, m(\chi) = 1 \mathrm{GeV} \\ g_q = 1, g_\chi = 1, m(\chi) = 1 \mathrm{GeV} \\ \tan\beta = 1, g_Z = 0.8, m(\chi) = 100 \mathrm{GeV} \\ \tan\beta = 1, g_\chi = 1, m(\chi) = 10 \mathrm{GeV} \end{array}$	2102.10874 2102.10874 2108.13391 ATLAS-CONF-2021-036
D7	Scalar LQ 1 st gen Scalar LQ 2 nd gen Scalar LQ 3 rd gen Scalar LQ 3 rd gen Scalar LQ 3 rd gen Scalar LQ 3 rd gen Vector LQ 3 rd gen	$\begin{array}{c c} 2 e & \geq 2 j \\ 2 \mu & \geq 2 j \\ 1 \tau & 2 b \\ 0 e, \mu & \geq 2 j, \geq 2 \\ \geq 2 e, \mu, \geq 1 \tau \geq 1 j, \geq 1 \\ 0 e, \mu, \geq 1 \tau & 0 - 2 j, 2 \\ 1 \tau & 2 b \end{array}$	b –	139 139 139 139 139 139 139 139	LO mass 1.7 LO" mass 1.2 TeV LO ³ mass 1.24 TeV LO ³ mass 1.24 TeV LO ³ mass 1.24 TeV LO ³ mass 1.26 TeV		$\begin{array}{l} \beta=1\\ \beta=1\\ \mathcal{B}(\mathrm{LQ}_{3}^{\mathrm{G}}\rightarrow b\tau)=1\\ \mathcal{B}(\mathrm{LQ}_{3}^{\mathrm{G}}\rightarrow t\nu)=1\\ \mathcal{B}(\mathrm{LQ}_{3}^{\mathrm{G}}\rightarrow t\nu)=1\\ \mathcal{B}(\mathrm{LQ}_{3}^{\mathrm{G}}\rightarrow b\nu)=1\\ \mathcal{B}(\mathrm{LQ}_{3}^{\mathrm{G}}\rightarrow b\tau)=0.5, \mathrm{Y-M} \ \mathrm{coupl.} \end{array}$	2006.05872 2006.05872 2108.07665 2004.14060 2101.11582 2101.12527 2108.07665
Vector-like fermions	$ \begin{array}{l} VLQ \ TT \to Zt + X \\ VLQ \ BB \to Wt/Zb + X \\ VLQ \ T_{5/3} \ T_{5/3} \ T_{5/3} \to Wt + X \\ VLQ \ T \to Ht/Zt \\ VLQ \ T \to Ht/Zt \\ VLQ \ B \to Hb \\ VLL \ \tau' \to Z\tau/H\tau \end{array} $	$\begin{array}{l} 2e/2\mu/\geq 3e,\mu \geq 1 \ b,\geq \\ \text{multi-channel} \\ 2(SS)/\geq 3 \ e,\mu \geq 1 \ b,\geq \\ 1 \ e,\mu \qquad \geq 1 \ b,\geq \\ 1 \ e,\mu \qquad \geq 1 \ b,\geq \\ 0 \ e,\mu \qquad \geq 2b,\geq 1j, \\ \text{multi-channel} \qquad \geq 1j \end{array}$	1 j Yes 3 j Yes 1 j Yes	139 36.1 36.1 139 36.1 139 139	T mass 1. Y mass 1.8		$\begin{array}{l} & \mathrm{SU}(2) \text{ doublet} \\ & \mathrm{SU}(2) \text{ doublet} \\ & \mathcal{B}(T_{5/3} \rightarrow Wt) = 1, \ c \left(T_{5/3} Wt\right) = 1 \\ & \mathrm{SU}(2) \text{ singlet}, \ \kappa_T = 0.5 \\ & \mathcal{B}(Y \rightarrow Wb) = 1, \ c_R(Wb) = 1 \\ & \mathrm{SU}(2) \text{ doublet}, \ \kappa_B = 0.3 \\ & \mathrm{SU}(2) \text{ doublet} \end{array}$	ATLAS-CONF-2021-024 1808.02343 1807.11883 ATLAS-CONF-2021-040 1812.07343 ATLAS-CONF-2021-018 ATLAS-CONF-2022-044
Excited	Excited quark $q^* \rightarrow qg$ Excited quark $q^* \rightarrow q\gamma$ Excited quark $b^* \rightarrow bg$ Excited lepton ℓ^* Excited lepton ν^*	$\begin{array}{cccc} - & 2j \\ 1\gamma & 1j \\ - & 1b,1 \\ 3e,\mu & - \\ 3e,\mu,\tau & - \end{array}$	j – –	139 36.7 139 20.3 20.3	q* mass q q* mass b* mass b* mass μ* mass ν* mass 1.6 *	6.7 TeV 5.3 TeV 3.2 TeV 3.0 TeV TeV	only u^* and d^* , $\Lambda = m(q^*)$ only u^* and d^* , $\Lambda = m(q^*)$ $\Lambda = 3.0 \text{ TeV}$ $\Lambda = 1.6 \text{ TeV}$	1910.08447 1709.10440 1910.0447 1411.2921 1411.2921
Other			s Yes – – –	139 36.1 139 139 20.3 139 34.4	Nº mass 910 GeV N _R mass 910 GeV H ^{±±} mass 350 GeV H ^{±±} mass 1.08 TeV H ^{±±} mass 1.08 TeV multi-charged particle mass 1.59 T monopole mass 1.01 T 10 ⁻¹ 1	3.2 TeV TeV 2.37 TeV 	$\begin{split} m(W_R) &= 4.1 \text{ TeV}, g_L = g_R \\ \text{DY production} \\ \text{DY production}, \mathcal{B}(H_L^{\pm\pm} \to \ell\tau) = 1 \\ \text{DY production}, \mathcal{B}(H_L^{\pm\pm} \to \ell\tau) = 1 \\ \text{DY production}, q &= 5e \\ \text{DY production}, g &= 1g_D, \text{ spin } 1/2 \\ \hline \\ \textbf{Mass scale [TeV]} \end{split}$	2202.02039 1809.11105 2101.11961 ATLAS-CONF-2022-010 1411.2921 ATLAS-CONF-2022-034 1905.10130

*Only a selection of the available mass limits on new states or phenomena is shown.

†Small-radius (large-radius) jets are denoted by the letter j (J).

ATLAS SUSY Searches* - 95% CL Lower Limits

March 2022

ATLAS Preliminary $\sqrt{s} = 13 \text{ TeV}$

23

Model	Signature	$\int \mathcal{L} dt [\mathbf{fb}^{T}]$	Mass limit	$\sqrt{s} = 13$ Reference
$\tilde{q}\tilde{q},\tilde{q}{ ightarrow}q\tilde{\chi}_1^0$	$0 e, \mu$ 2-6 jets E_T^{miss} mono-jet 1-3 jets E_T^{miss}	139 139	<i>q</i> [1×, 8× Degen.] 1.0 1.85 m($\tilde{\chi}_1^0$)	<400 GeV 2010.14293 ⁰¹ ₁)=5 GeV 2102.10874
$\begin{split} \tilde{g}\tilde{g}, \; \tilde{g} \to q \bar{q} \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{g}, \; \tilde{g} \to q \bar{q} W \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{g}, \; \tilde{g} \to q \bar{q} (\ell \ell) \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{g}, \; \tilde{g} \to q q W Z \tilde{\chi}_{1}^{0} \end{split}$	$0 e, \mu$ 2-6 jets E_T^{miss}		ğ 2.3 m(t) ğ Forbidden 1.15-1.95 m(t/t)	(¹)=0 GeV 2010.14293 1000 GeV 2010.14293
$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\bar{q}W\tilde{\chi}_1^0$	1 e,μ 2-6 jets $ee,\mu\mu$ 2 jets E_T^{miss}	139 139	ğ 2.2 m(χ ⁰ ₁	<600 GeV 2101.01629 <700 GeV CERN-EP-2022-014
$ ilde{g} ilde{g}, ilde{g} o q ar{q}(\ell \ell) ilde{\chi}_1^0 \ ilde{g} ilde{g}, ilde{g} o q q W Z ilde{\chi}_1^0$	$\begin{array}{cccc} 0 \ e, \mu & 7-11 \ \text{jets} & E_T \\ \text{SS} \ e, \mu & 6 \ \text{jets} \end{array}$	139 139 139		<600 GeV 2008.06032
$\tilde{g}\tilde{g}, \tilde{g} \rightarrow t \tilde{t} \tilde{\chi}_1^0$	$\begin{array}{ccc} 0\text{-1} e, \mu & 3 b & E_T^{\text{miss}} \\ \text{SS} e, \mu & 6 \text{ jets} \end{array}$			<200 GeV ATLAS-CONF-2018-041
$ ilde{b}_1 ilde{b}_1$	$0 e, \mu$ $2 b E_T^{\text{miss}}$	139	\tilde{b}_1 \tilde{b}_1 \tilde{b}_1 1.255 $m(\tilde{\lambda}_1^0)$ 10 GeV < $\Delta m(\tilde{b}_1 \tilde{\lambda})$ 10 GeV < $\Delta m(\tilde{b}_1 \tilde{\lambda})$	<400 GeV 2101.12527
$\tilde{b}_1 \tilde{b}_1, \tilde{b}_1 \rightarrow b \tilde{\chi}_2^0 \rightarrow b h \tilde{\chi}_1^0$ $\tilde{\iota}_1 \tilde{\iota}_1, \tilde{\iota}_1 \rightarrow \tilde{\chi}_1^0$	$\begin{array}{ccc} 0 \ e, \mu & 6 \ b & E_{T}^{\text{miss}} \\ 2 \ \tau & 2 \ b & E_{T}^{\text{miss}} \end{array}$	139 139	b_1 Forbidden 0.23-1.35 $\Delta m(\tilde{\chi}_2^0, \tilde{\chi}_1^0) = 130 \text{ GeV}, m(\tilde{\chi}_1^0, \tilde{\chi}_1$	=100 GeV 1908.03122
$\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow t \tilde{\chi}_1^0$	0-1 $e, \mu \ge 1$ jet E_T^{miss}	139	<i>ĩ</i> 1.25 m(v ⁰ ₁)=1 GeV 2004.14060,2012.03799
$\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \to W b \tilde{\chi}_1^0$ $\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \to \tilde{t}_1 \to \tilde{t}_1 \to \tilde{t}_1$	$\begin{array}{rll} 1 \ e, \mu & 3 \ {\rm jets}/1 \ b & E_T^{\rm miss} \\ 1-2 \ \tau & 2 \ {\rm jets}/1 \ b & E_T^{\rm miss} \end{array}$	139 139		=500 GeV 2012.03799 =800 GeV 2108.07665
$ \tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow \tilde{\tau}_1 b \nu, \tilde{\tau}_1 \rightarrow \tau \tilde{G} \tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow c \tilde{\chi}_1^0 / c \tilde{c}, \tilde{c} \rightarrow c \tilde{\chi}_1^0 $	1-2 τ 2 jets/1 b E_T^{miss} 0 e, μ 2 c E_T^{miss}	36.1		=800 GeV 2108.07665
	$0 e, \mu$ mono-jet E_T^{fmiss}		<i>ī</i> ₁ 0.55 m(<i>ī</i> ₁ , <i>č</i>)-m((1)=5 GeV 2102.10874
$ \tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow t \tilde{\chi}_2^0, \tilde{\chi}_2^0 \rightarrow Z/h \tilde{\chi}_1^0 \tilde{t}_2 \tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z $	$\begin{array}{ccc} 1-2 \ e, \mu & 1-4 \ b & E_T^{\text{miss}} \\ 3 \ e, \mu & 1 \ b & E_T^{\text{miss}} \end{array}$	139 139	$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	=500 GeV 2006.05880)= 40 GeV 2006.05880
$ ilde{\chi}_1^{\pm} ilde{\chi}_2^0$ via WZ	$\begin{array}{rll} \mbox{Multiple }\ell/\mbox{jets} & E_T^{\rm miss}\\ ee, \mu\mu & \geq 1 \mbox{ jet} & E_T^{\rm miss} \end{array}$	139 139	$\vec{\chi}_{1}^{+}/\vec{\chi}_{2}^{0}$ 0.96 $m(\vec{\chi}_{1}^{0})=0$ $m(\vec{\chi}_{1}^{0})=5$ GeV	wino-bino 2106.01676, 2108.07586 wino-bino 1911.12606
$ ilde{\chi}_1^{\pm} ilde{\chi}_1^{\mp}$ via WW	$2 e, \mu \qquad E_T^{miss}$	139	$\tilde{\chi}_1^{\pm}$ 0.42 m($\tilde{\chi}_1^0$)=0.	wino-bino 1908.08215
$\tilde{\chi}_1^{\pm} \tilde{\chi}_2^0$ via <i>Wh</i>	Multiple ℓ /jets E_T^{miss}	139	$\tilde{\chi}_{1}^{\pm}/\tilde{\chi}_{2}^{0}$ Forbidden 1.06 m($\tilde{\chi}_{1}^{0}$)=70 GeV	
$\tilde{\chi}_{1}^{\pm} \tilde{\chi}_{1}^{\mp}$ via $\tilde{\ell}_{L} / \tilde{v}$ $\tilde{\tau} \tilde{\tau}, \tilde{\tau} ightarrow \tau \tilde{\chi}_{1}^{0}$	$\begin{array}{ccc} 2 \ e, \mu & E_T^{\text{miss}} \\ 2 \ \tau & E_T^{\text{miss}} \end{array}$	139 139	$\tilde{\chi}_{1}^{\pm}$ 1.0 m($\tilde{\ell}, \tilde{\nu}$)=0.5(m/ $\tilde{\ell}$	$\begin{array}{c} \overset{\pm}{}_{1}) + m(\tilde{\chi}_{1}^{0})) \\ m(\tilde{\chi}_{1}^{0}) = 0 \end{array} \qquad 1908.08215 \\ 1911.06660 \end{array}$
$ \begin{split} &\tilde{\chi}_{1}^{\pm} \tilde{\chi}_{1}^{\mp} \text{ via } \tilde{\ell}_{L} / \tilde{\nu} \\ &\tilde{\tau} \tilde{\tau}, \tilde{\tau} \rightarrow \tau \tilde{\chi}_{1}^{0} \\ &\tilde{\ell}_{L,R} \tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell \tilde{\chi}_{1}^{0} \end{split} $	$2 e, \mu$ 0 jets E_T^{miss}	139	<i>ℓ</i> 0.7	$m(\tilde{\chi}_{1}^{0})=0$ 1908.08215
)=10 GeV 1911.12606
$\tilde{H}\tilde{H}, \tilde{H} \rightarrow h\tilde{G}/Z\tilde{G}$	$\begin{array}{llllllllllllllllllllllllllllllllllll$	36.1 139	μ 0.13-0.23 0.29-0.88 BR(ℓ μ 0.55 BR(ℓ	$\hat{D} \to h\tilde{G}$)=1 1806.04030 $\to Z\tilde{G}$)=1 2103.11684
	$e, \mu = 0$ jets E_T 0 $e, \mu \geq 2$ large jets E_T	139	й 0.45-0.93 Вн(й)	$\rightarrow Z\widetilde{G}$)=1 2108.07586
Direct $\tilde{\chi}_1^+ \tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^\pm$	Disapp. trk 1 jet E_T^{miss}			Pure Wino 2201.02472 e higgsino 2201.02472
Stable \tilde{g} R-hadron Metastable \tilde{g} R-hadron, $\tilde{g} \rightarrow qq \tilde{\chi}_1^0$ $\tilde{\ell}\tilde{\ell}, \tilde{\ell} \rightarrow \ell \tilde{G}$	pixel dE/dx E_T^{miss} pixel dE/dx E_T^{miss}	139	<i>ğ</i> 2.05	CERN-EP-2022-029
Metastable \tilde{g} R-hadron, $\tilde{g} \rightarrow qq \tilde{\chi}_1^0$	pixel dE/dx E_T^{miss} Displ. lep E_T^{miss}	139		=100 GeV CERN-EP-2022-029
$\tilde{\ell}\tilde{\ell},\tilde{\ell}{\rightarrow}\ell\tilde{G}$			$\tilde{\tau}$ 0.34 τ	() = 0.1 ns 2011.07812
	pixel dE/dx E_T^{miss}	139	τ 0.36 τι	ℓ̃) = 10 ns CERN-EP-2022-029
$\tilde{\chi}_{1}^{\pm}\tilde{\chi}_{1}^{\mp}/\tilde{\chi}_{1}^{0}, \tilde{\chi}_{1}^{\pm} \rightarrow Z\ell \rightarrow \ell\ell\ell$	3 e,µ	139		Pure Wino 2011.10543
$\tilde{\chi}_{1}^{\pm} \tilde{\chi}_{1}^{\mp} / \tilde{\chi}_{2}^{0} \xrightarrow{\sim} WW/Z\ell\ell\ell\ell\nu\nu$	4 e, μ 0 jets E_T^{miss} 4-5 large jets			=200 GeV 2103.11684
$\widetilde{g}\widetilde{g}, \ \widetilde{g} \to qq\widetilde{\chi}_1^0, \ \widetilde{\chi}_1^0 \to qqq$ $\widetilde{t}, \ \widetilde{t} \to t\widetilde{\chi}_1^0, \ \widetilde{\chi}_1^0 \to tbs$	4-5 large jets Multiple	36.1 36.1	$\begin{array}{c c c c c c c c c c c c c c c c c c c $	Large λ_{112}'' 1804.03568 // bino-like ATLAS-CONF-2018-003
$\widetilde{t}\widetilde{t}, \widetilde{t} \rightarrow b\widetilde{\chi}_{1}^{\pm}, \widetilde{\chi}_{1}^{\pm} \rightarrow bbs$	$\geq 4b$	139	323	=500 GeV 2010.01015
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow bs$	2 jets + 2 b	36.7	$\tilde{t}_1 [qq, bs]$ 0.42 0.61	1710.07171
$\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow q\ell$	$\begin{array}{ccc} 2 \ e, \mu & 2 \ b \\ 1 \ \mu & DV \end{array}$	36.1 136	$\begin{array}{c c} \tilde{t}_1 & 0.4-1.45 & BR(\tilde{t}_1 \rightarrow be \\ \tilde{t}_1 & [1e-10 < \lambda'_{224} < 1e-8, 3e-10 < \lambda'_{224} < 3e-9] & 1.0 & 1.6 & BR(\tilde{t}_1 \rightarrow \mu) = 100^{\circ} \end{array}$	/ <i>bμ</i>)>20% 1710.05544 %, cosθ _r =1 2003.11956
$\tilde{\chi}_1^{\pm}/\tilde{\chi}_2^0/\tilde{\chi}_1^0, \tilde{\chi}_{1,2}^0 \rightarrow tbs, \tilde{\chi}_1^+ \rightarrow bbs$	$1-2 \ e, \mu \ge 6 \text{ jets}$	139	25K 25K	e higgsino 2106.09609
	nass limits on new states or e limits are based on		0 ⁻¹ Mass scale	



Vininun

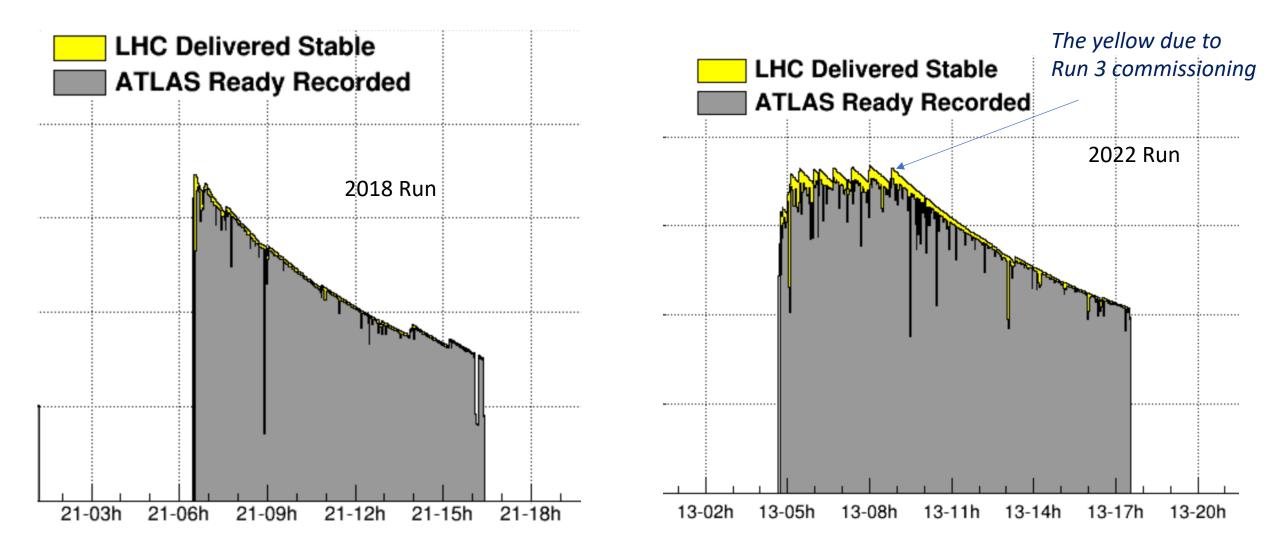
RUN 3

Run: 420624

What is new in Run 3?

- World-breaking energy achieved on July 5, 2022: $\sqrt{s} = 13.6$ TeV
- More than double the Run 2 dataset
- LHC providing luminosity leveling

Luminosity Leveling Example: Run 2 vs. Run 3







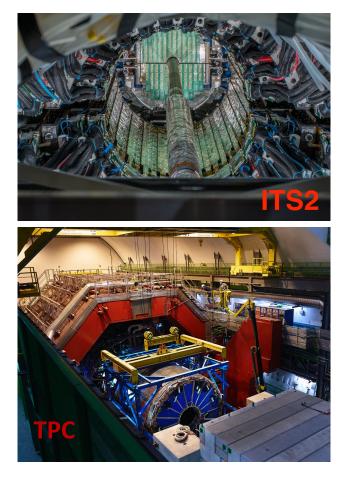


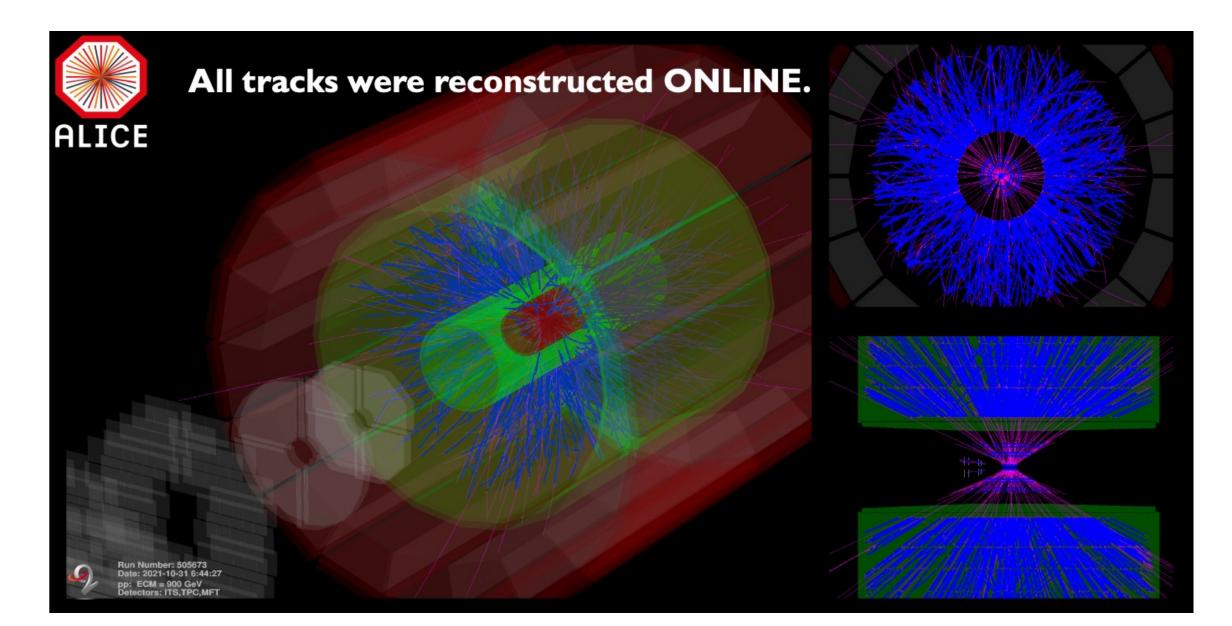




A Large Heavy Ion Experiment (ALICE)

- Significant upgrades for Run 3 to handle
 50 kHz Pb-Pb collisions & high occupancies
- Many new systems
 - Inner Tracking System (ITS2)
 - Muon Forward Tracker
 - Fast Interaction Trigger
 - Time Projection Chamber (TPC) with GEMS
 - Online-offline software framework
- 3 TB/s readout with specialized hardware and GPU-based data reduction

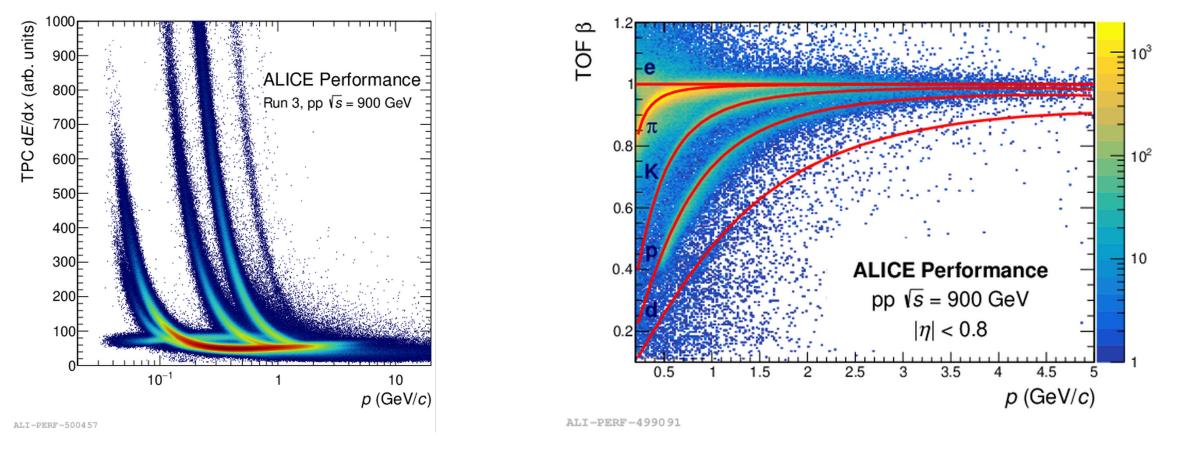






ALICE Performance: from November pilot beam

Beautiful particle ID from dE/dx and time of flight measurements







ATLAS Upgrades for Run 3

MUON NEW SMALL WHEELS (NSW)

Installed new muon detectors with precision tracking and muon selection capabilities. Key preparation for the HL-LHC.

NEW READOUT SYSTEM FOR THE NSWs

The NSW system includes two million micromega readout channels and 350 000 small strip thin-gap chambers (sTGC) electronic readout channels.

LIQUID ARGON CALORIMETER

New electronics boards installed, increasing the granularity of signals used in event selection and improving trigger performance at higher luminosity.



TRIGGER AND DATA ACQUISITION SYSTEM (TDAQ)

Upgraded hardware and software allowing the trigger to spot a wider range of collision events while maintaining the same acceptance rate.

NEW MUON CHAMBERS IN THE CENTRE OF ATLAS

Installed small monitored drift tube (sMDT) detectors alongside a new generation of resistive plate chamber (RPC) detectors, extending the trigger coverage in preparation for the HL-LHC.

ATLAS FORWARD PROTON (AFP)

Re-designed AFP time-of-flight detector, allowing insertion into the LHC beamline with a new "out-ofvacuum" solution.

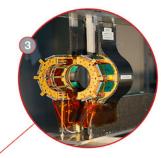
CMS Upgrades for Run 3

BEAM PIPE

Replaced with an entirely new one compatible with the future tracker upgrade for HL-LHC, improving the vacuum and reducing activation.

PIXEL TRACKER

All-new innermost barrel pixel layer, in addition to maintenance and repair work and other upgrades.



BRIL

New generation of detectors for monitoring LHC beam conditions and luminosity.

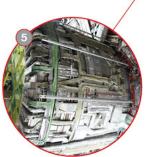


CATHODE STRIP CHAMBERS (CSC)

Read-out electronics upgraded on all the 180 CSC muon chambers allowing performance to be maintained in HL-LHC conditions.

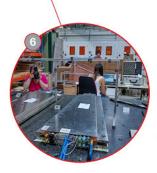
HADRON CALORIMETER

New on-detector electronics installed to reduce noise and improve energy measurement in the calorimeter.



SOLENOID MAGNET

New powering system to prevent full power cycles in the event of powering problems, saving valuable time for physics during collisions and extending the magnet lifetime.



GAS ELECTRON MULTIPLIER (GEM) DETECTORS

An entire new station of detectors installed in the endcap-muon system to provide precise muon tracking despite higher particle rates of HL-LHC.

Run 3 ATLAS and CMS Physics Program

- Take advantage of **trigger improvements** to push on searches for unconventional signatures, such as long-lived particles.
- Make **Run 2 + Run 3 combinations** to improve precision of Higgs boson production/decay cross sections and other SM measurements
 - aim to constrain $H \rightarrow bb$ and $H \rightarrow \mu\mu$ couplings at 20% level
- Benefit from increased data push mass reach on searches for supersymmetric or exotic particles
 - up to a 1 TeV increase can be achieved with respect to Run 2 in some cases
- Pushing on **techniques** such as "parking" data, "scouting" or "triggerlevel-analysis" to access challenging phase space for b-physics or lowmomentum hadronic physics.

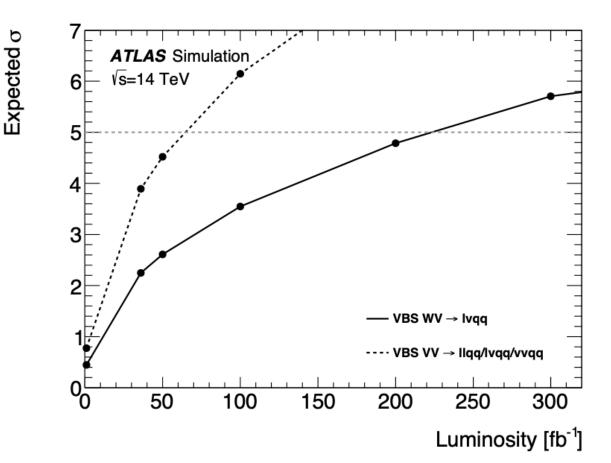
Inclusive vector boson scattering (VBS)

Important for the story of EWK symmetry breaking:

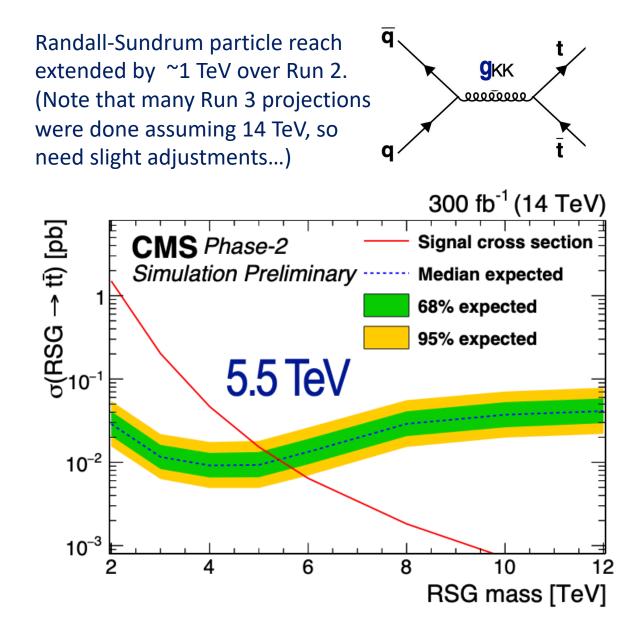
Without the SM Higgs boson, the longitudinal VV scattering cross section increases with \sqrt{s} and would violate unitarity at high energies!

Additionally, VBS provides indirect searches for BSM by giving access to anomalous quartic gauge couplings.

The study of **inclusive** VV scattering was done assuming $\sqrt{s} = 14 TeV$ so is a bit optimistic for Run 3, but gives some confidence that we will have 5 sigma significance with Run 3 data. Picking out the **longitudinal component** requires the HL-LHC for sensitivity.



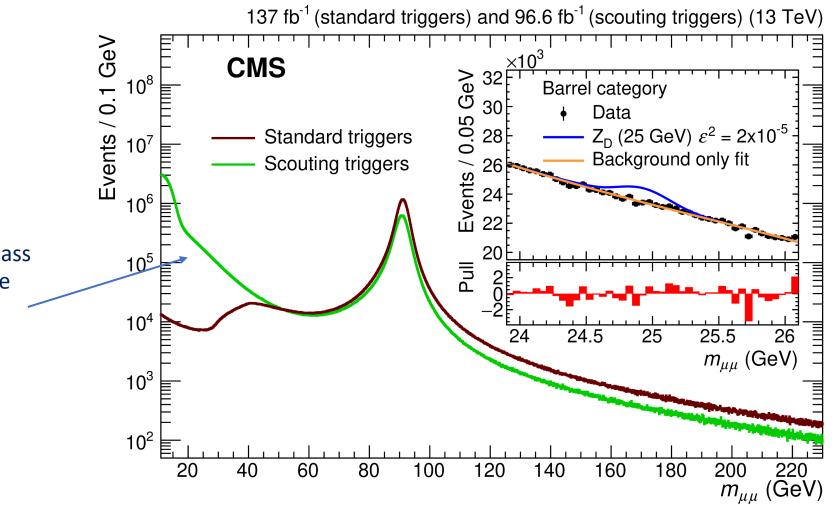
Can extend search reach with additional data



```
H/A \rightarrow \tau \tau Search exclusions
```

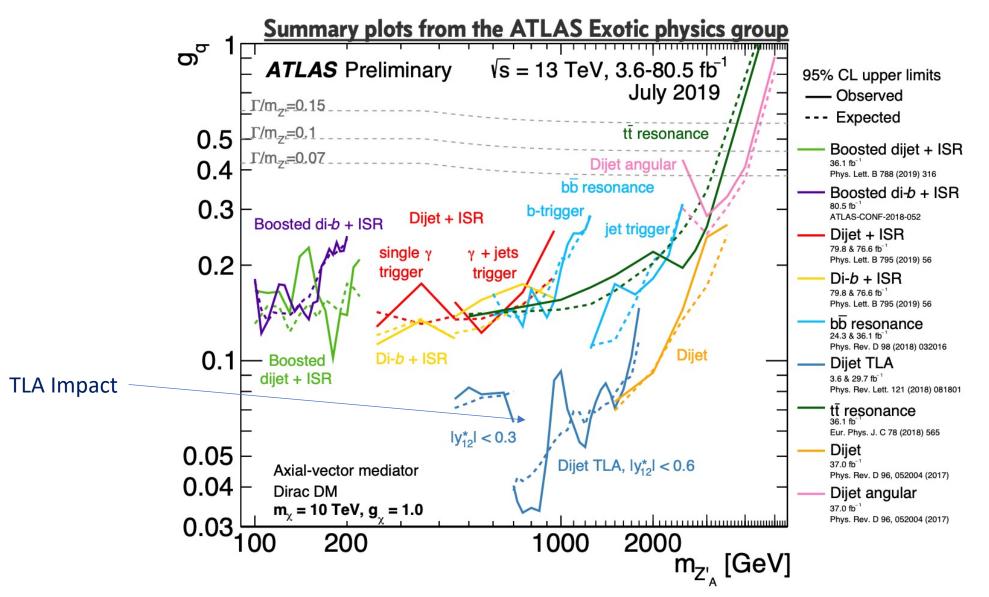
CMS Projection 95% CL expected exclusion: YR18 syst. uncert. JHEP 09(2018)007 ±1σ ±2σ ----- 3000 fb⁻¹ ----- 300 fb⁻¹ ----- 6000 fb⁻¹ $\tan \beta$ mmod+ 50 m^{MSSM} ≠ 125 ± 3 GeV 40 30 20 10 500 1000 1500 2000 m_A (GeV)

We've seen that trigger scouting (CMS) provides access to expanded phase space

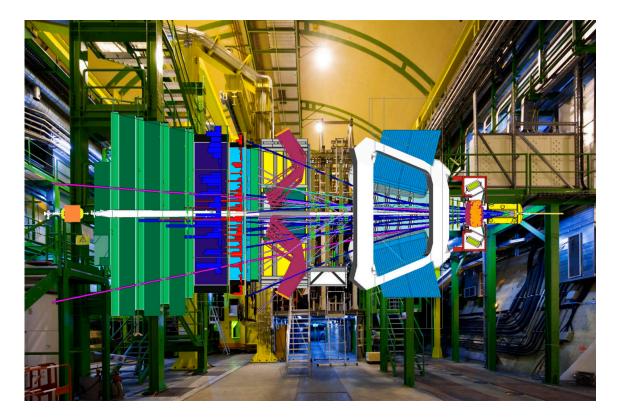


This Run 2 search for low-mass dimuon pairs provides a nice example of trigger scouting opening up phase space

Similarly, trigger-level analysis (TLA @ ATLAS) has extended the reach, shown here for Z'



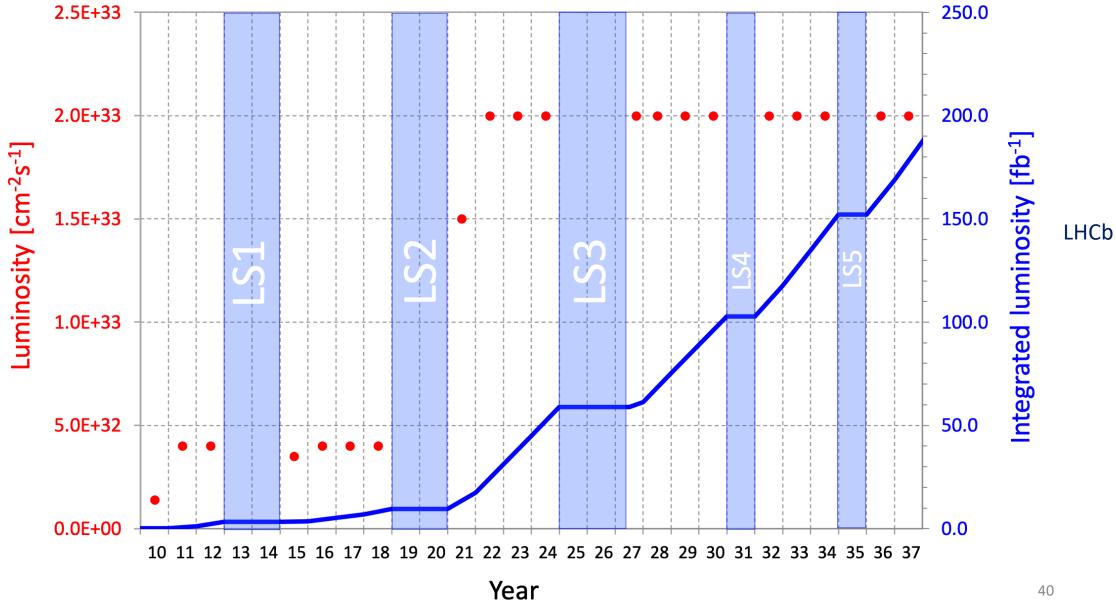
LHCb: Run 3 and beyond

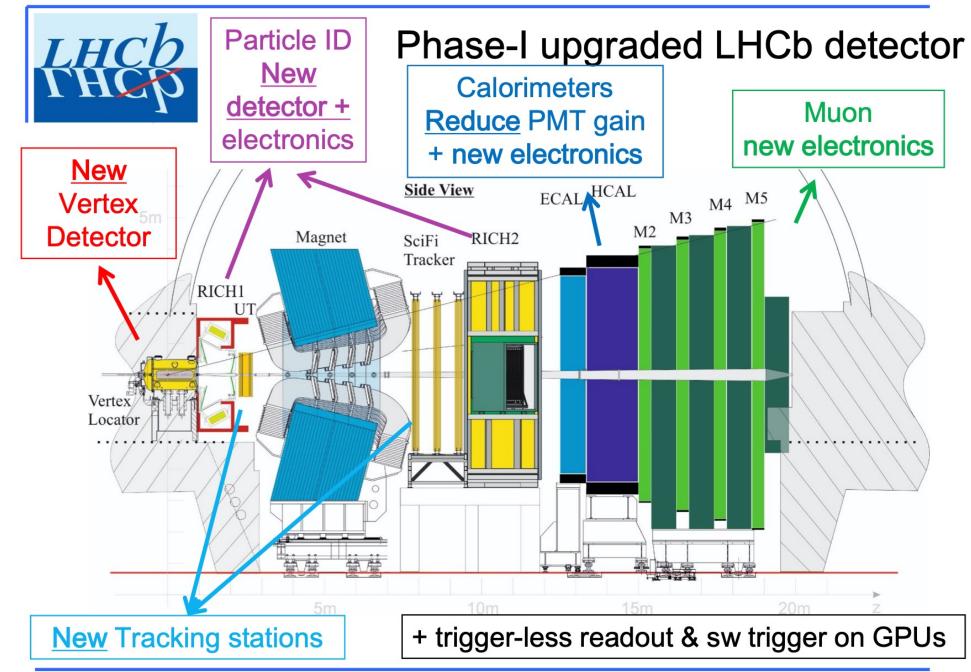


- LHCb is a forward arm spectrometer, working with particles from collisions in the forward region (2 $<\eta<$ 5)
- Targets precision measurements in the b and charm quark sectors
- Has extended to access to exploring QCD and EW physics
- Participates in pp running, with lower luminosities, and heavy ion runs

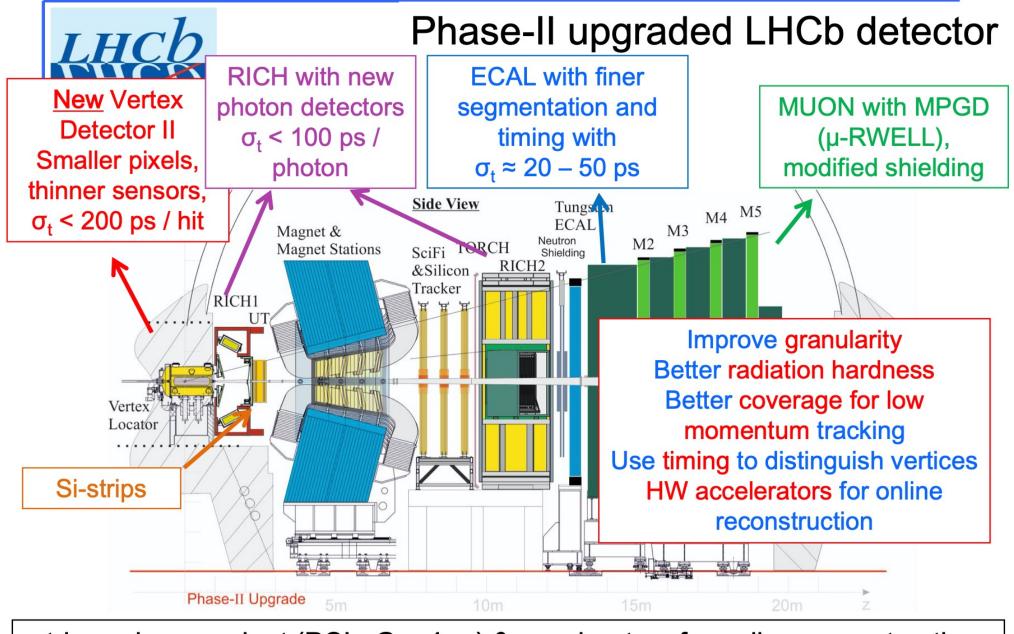
NOTE: LS3 has been shifted one year forward (see next slide)

• Peak luminosity —Integrated luminosity



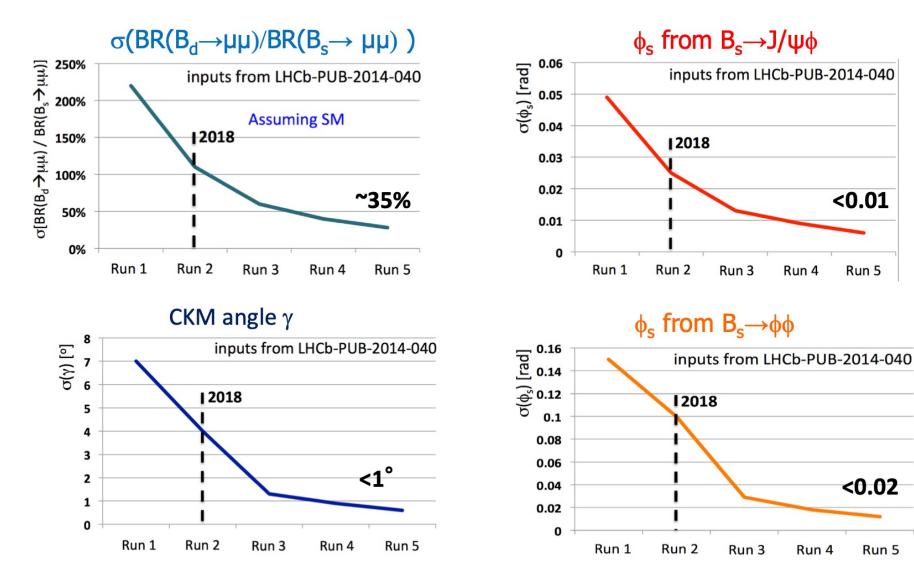


See talk by F. Alessio for more details



+ trigger-less readout (PCIe Gen4++) & accelerators for online reconstruction

Sensitivity of several high-priority LHCb measurements





<0.01

Run 5

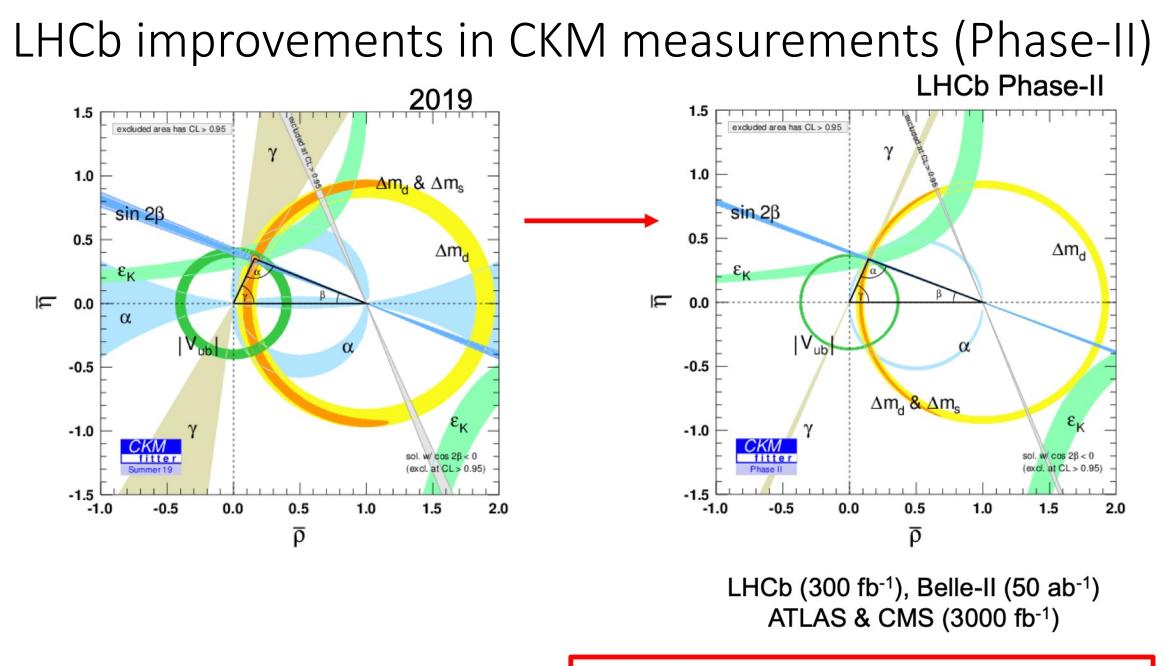
<0.02

Run 5

Run 4

Run 3

Run 4



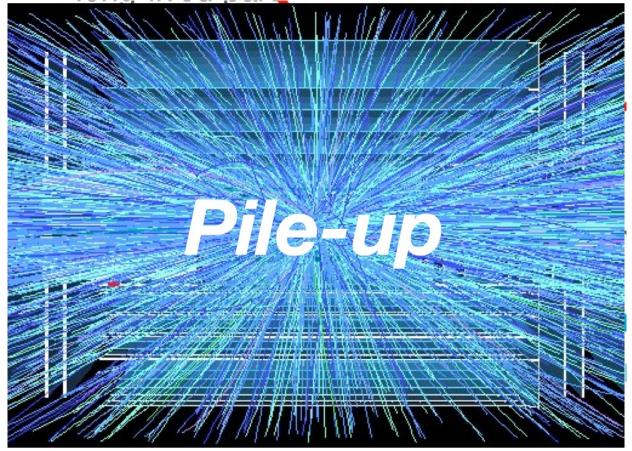
 $\sigma_v \approx 5^\circ (2019) \rightarrow 1^\circ (\text{Phase 1}) \rightarrow 0.35^\circ (\text{Phase 2})$ 44

Observable	Current LHCb	LHCb 2025	Belle II	Upgrade II
EW Penguins				
$\overline{R_K \ (1 < q^2 < 6 \mathrm{GeV}^2 c^4)}$	0.1	0.025	0.036	0.007
$R_{K^*} \ (1 < q^2 < 6 \mathrm{GeV}^2 c^4)$	0.1	0.031	0.032	0.008
$R_{\phi}, R_{pK}, R_{\pi}$	-	0.08, 0.06, 0.18	-	0.02, 0.02, 0.05
CKM tests				
$\overline{\gamma}$, with $B_s^0 \to D_s^+ K^-$	$\binom{+17}{-22}^{\circ}$	4°	-	1°
γ , all modes	$(^{+17}_{-22})^{\circ}_{(^{+5.0}_{-5.8})^{\circ}}$	1.5°	1.5°	0.35°
$\sin 2\beta$, with $B^0 \to J/\psi K_s^0$	0.04	0.011	0.005	0.003
ϕ_s , with $B_s^0 \to J/\psi\phi$	49 mrad	14 mrad	-	$4 \mathrm{mrad}$
ϕ_s , with $B_s^0 \to D_s^+ D_s^-$	170 mrad	35 mrad	-	$9 \mathrm{mrad}$
$\phi_s^{s\bar{s}s}$, with $B_s^0 \to \phi\phi$	154 mrad	39 mrad	-	11 mrad
$a_{\rm sl}^s$	$33 imes 10^{-4}$	$10 imes 10^{-4}$	-	$3 imes 10^{-4}$
$ ec{V}_{ub} / V_{cb} $	6%	3%	1%	1%
$B^0_s, B^0{ ightarrow}\mu^+\mu^-$				
$\overline{\mathcal{B}(B^0 \to \mu^+ \mu^-)} / \mathcal{B}(B^0_s \to \mu^+ \mu^-)$	90%	34%	_	10%
$\tau_{B^0_s \to \mu^+ \mu^-}$	22%	8%	-	2%
$S_{\mu\mu}$	_	-	-	0.2
$b \to c \ell^- \bar{\nu_l} \text{ LUV studies}$				
$\overline{R(D^*)}$	0.026	0.0072	0.005	0.002
$R(J/\psi)$	0.24	0.071	-	0.02
Charm				
$\overline{\Delta A_{CP}(KK - \pi\pi)}$	$8.5 imes10^{-4}$	$1.7 imes10^{-4}$	$5.4 imes10^{-4}$	$3.0 imes 10^{-5}$
$A_{\Gamma} \ (\approx x \sin \phi)$	$2.8 imes 10^{-4}$	$4.3 imes 10^{-5}$	$3.5 imes 10^{-4}$	$1.0 imes 10^{-5}$
$x\sin\phi$ from $D^0 \to K^+\pi^-$	$13 imes 10^{-4}$	$3.2 imes 10^{-4}$	$4.6 imes10^{-4}$	$8.0 imes 10^{-5}$
$x\sin\phi$ from multibody decays	_	$(K3\pi)~4.0\times 10^{-5}$	$(K_{ m S}^0\pi\pi)~1.2 imes10^{-4}$	$(K3\pi) 8.0 \times 10^{-6}$

LHCb Run 3 goal Run 4 goal

http://arxiv.org/abs/arXiv:1808.08865

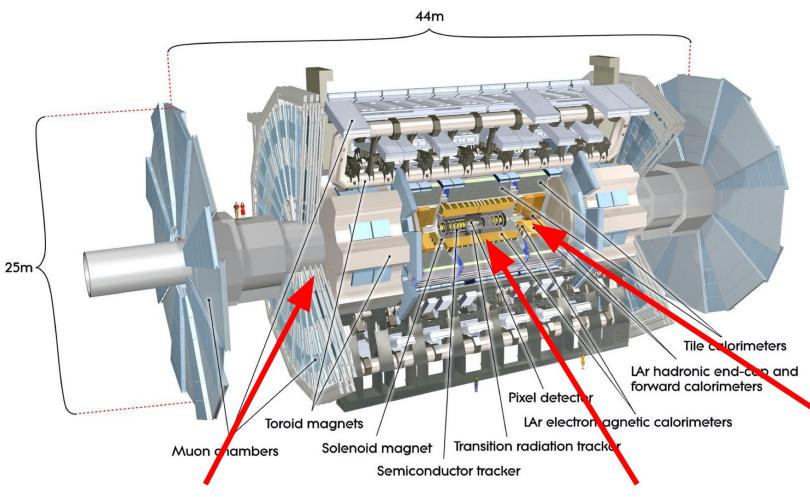
HL-LHC



What is new in the HL-LHC era?

- Higher energy, with $\sqrt{s} = 14 \text{ TeV}$
- Significant increase in instantaneous luminosity (up to 140 200 overlapping interactions)
- Increasing dataset by factor of ten
- Many upgraded or new detector components, including timing detectors

Overview of ATLAS Phase-II Upgrades



New Muon Chambers

 Inner barrel region with new Resistive Plate Chambers and new Monitored Drift Tubes (sMDT) detectors

New Inner Tracking Detector (ITk)

• All silicon (9 layers), up to $|\eta| = 4$

Upgraded Trigger and Data Acquisition system

- Level-0 Trigger at 1 MHz
- Improved High-Level Trigger
- (150 kHz full-scan tracking)

Electronics Upgrades

- On-detector and off-detector electronics upgrades of:
- LAr Calorimeter
- Tile Calorimeter
- Muon Detectors

High Granularity Timing Detector (HGTD)

- Forward region
- Precision time recon. (30 ps) with Low-Gain Avalanche Detectors (LGAD)

Additional small upgrades

- Luminosity detectors (1% precision)
- HL-ZDC (Heavy Ion physics)

Overview of CMS Phase-II Upgrades

L1-Trigger/HLT/DAQ

https://cds.cern.ch/record/2283192 https://cds.cern.ch/record/2283193

- Tracks in L1-Trigger at 40 MHz for 750 kHz PFlow-like selection rate
- HLT output 7.5 kHz

Calorimeter Endcap

https://cds.cern.ch/record/2293646

- Si, Scint+SiPM in Pb-W-SS
- 3D shower topology with precise timing

Tracker https://cds.cern.ch/record/2272264

- Si-Strip and Pixels increased granularity
- Design for tracking in L1-Trigger
- Extended coverage to $\eta\simeq 3.8$

Barrel Calorimeters

https://cds.cern.ch/record/2283187

- ECAL crystal granularity readout at 40 MHz with precise timing for e/γ at 30 GeV
- ECAL and HCAL new Back-End boards

Muon systems

https://cds.cern.ch/record/2283189

- DT & CSC new FE/BE readout
- New GEM/RPC 1.6 < η < 2.4
- Extended coverage to $\eta \simeq 3$

Beam Radiation Instr. and Luminosity, and Common Systems and Infrastructure https://cds.cern.ch/record/2020886

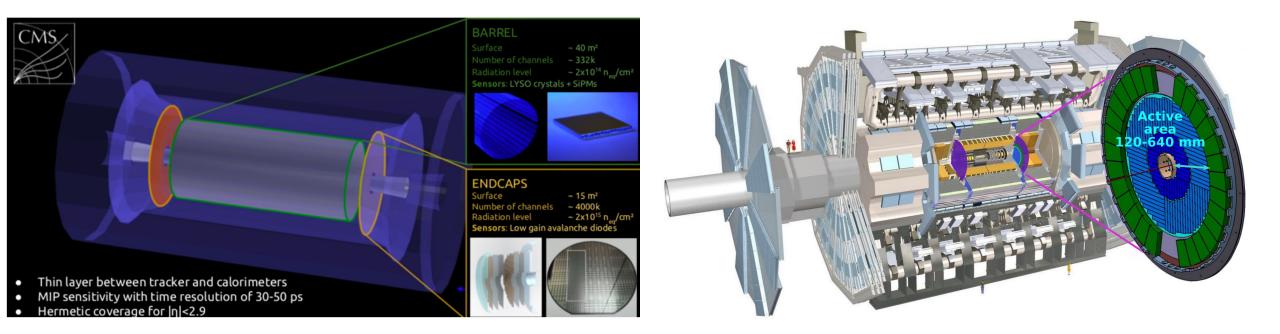
MIP Timing Detector https://cds.cern.ch/record/2296612

- ~ 30 ps resolution
- Barrel layer: Crystals + SiPMs
- Endcap layer: Low Gain Avalanche Diodes

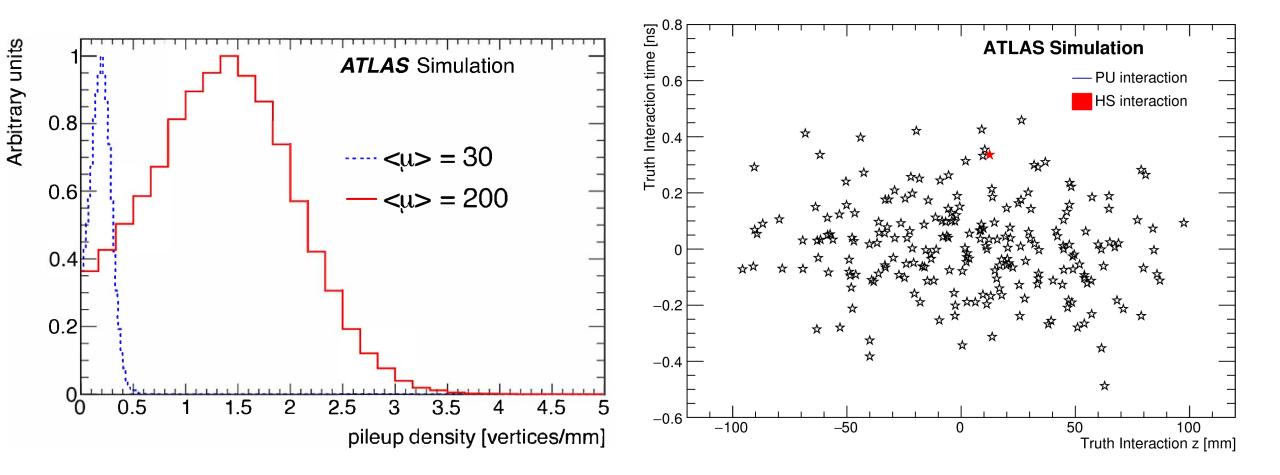
A quick aside regarding timing...

CMS MIP Timing Detector (MTD)

ATLAS High-Granularity Timing Detector (HGTD)

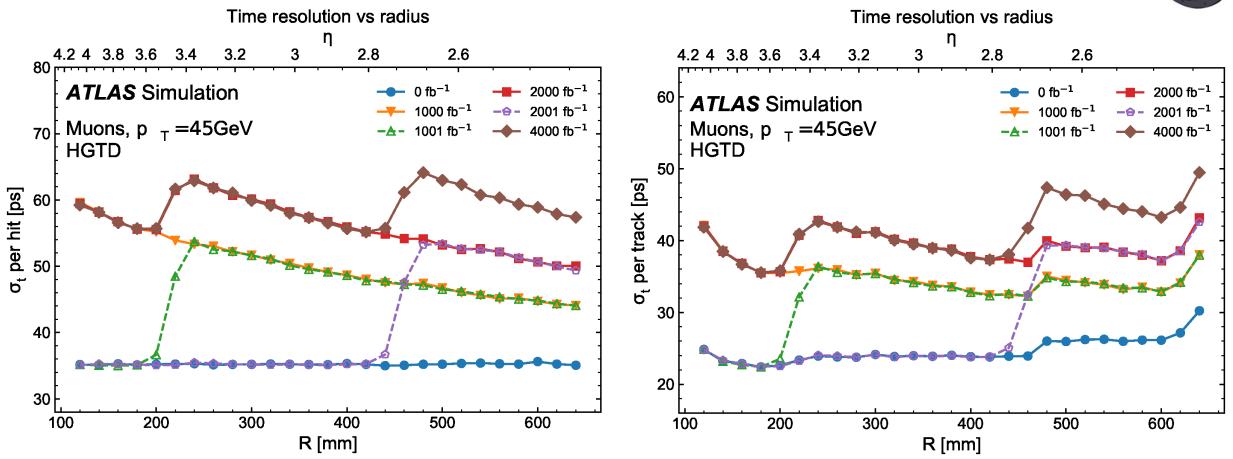


The potential of timing information...





How well can we measure time?

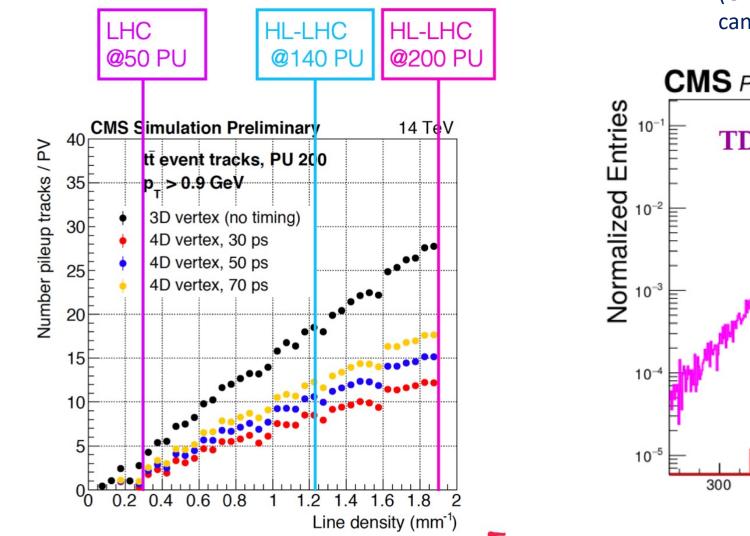


How well can we use it?

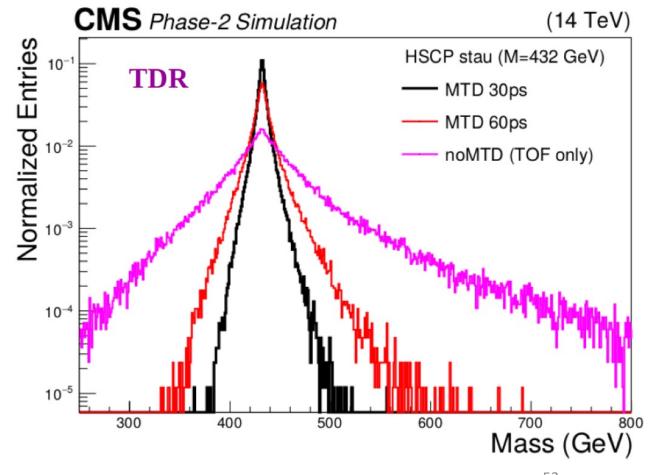


- If each hit in your timing detector has a time associated with it, you could
 - use the times to improve your tracking, reducing fakes from combinatorics
 - assign a time to each track.
- The track times could be used to assign times with reconstructed vertices so that pile-up vertices can be separated from the hard-scatter vertex of interest.
- It isn't trivial to use the timing information an maintain high efficiency, particularly for ATLAS where we only have access to timing from forward objects which primarily helps with VBF-like signatures.
- The HL-LHC will provide a tremendous opportunity to learn how to exploit this additional dimension, on the route to 5D!

CMS MTD performance



long-lived staus moving slowly through detector (GMSB model) – for the same background they can achieve a factor of four in signal gain



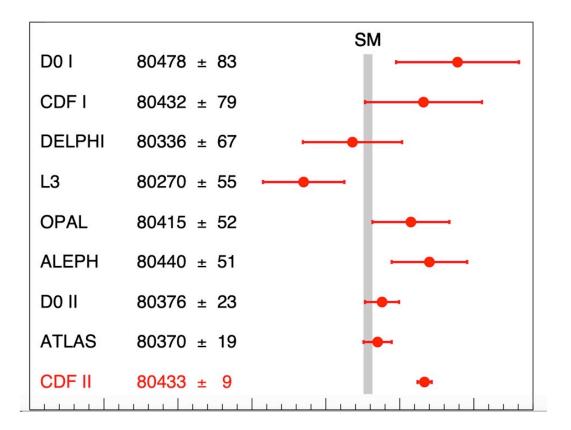
More results in a beautiful talk by Livia Soffi here

HL-LHC Physics Program

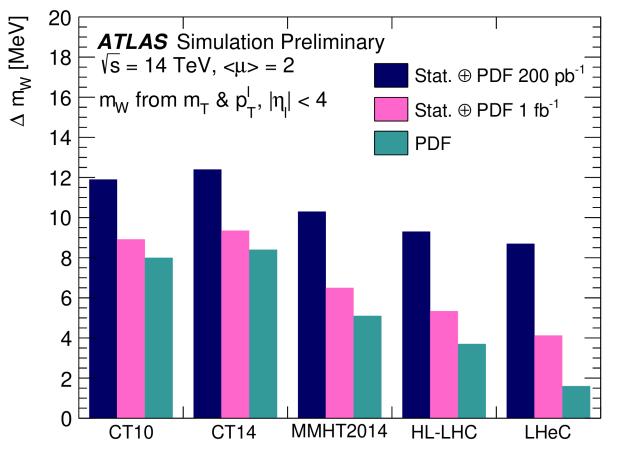
- Rediscover the SM, and measure properties with higher precision
- Improve measurements of Higgs production and branching ratios, with a target of percent-level uncertainties for range of decay modes
- Some Holy Grails:
 - Improve measurements of WW scattering
 - Access the Higgs self coupling
- Use the Higgs boson to search for BSM physics
- Extend searches for
 - SUSY, long-lived particles, Heavy new bosons (W'/Z'), exotic Higgs bosons, and much, much more!

Precision Standard Model: W boson mass

There has been renewed interest in the W mass measurement due to the recent CDF result.



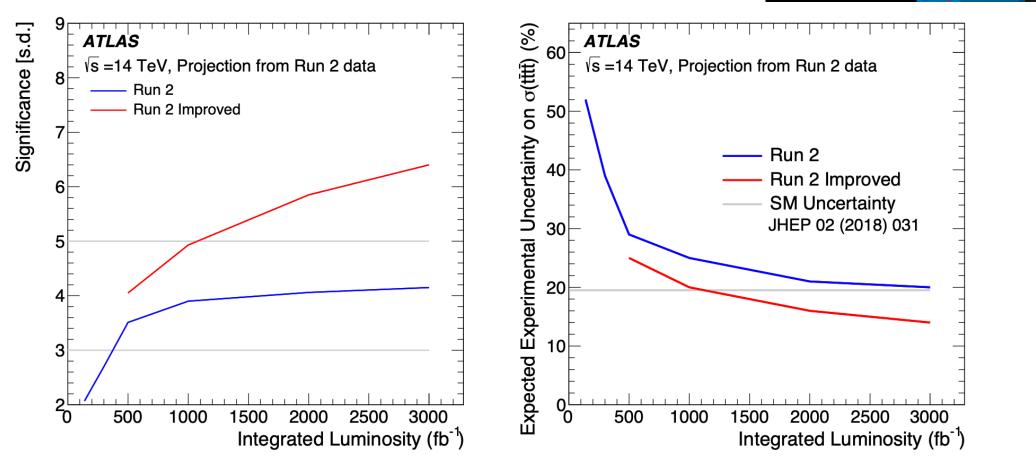
At HL-LHC, requires a dedicated low pile-up dataset. The analysis benefits from improved forward lepton ID, and is very dependent on PDF knowledge.

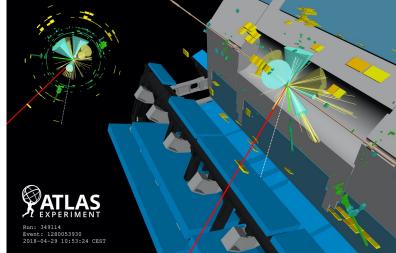


Four top cross section

Run 2 analysis improvements enabled ATLAS to update our HL-LHC projection

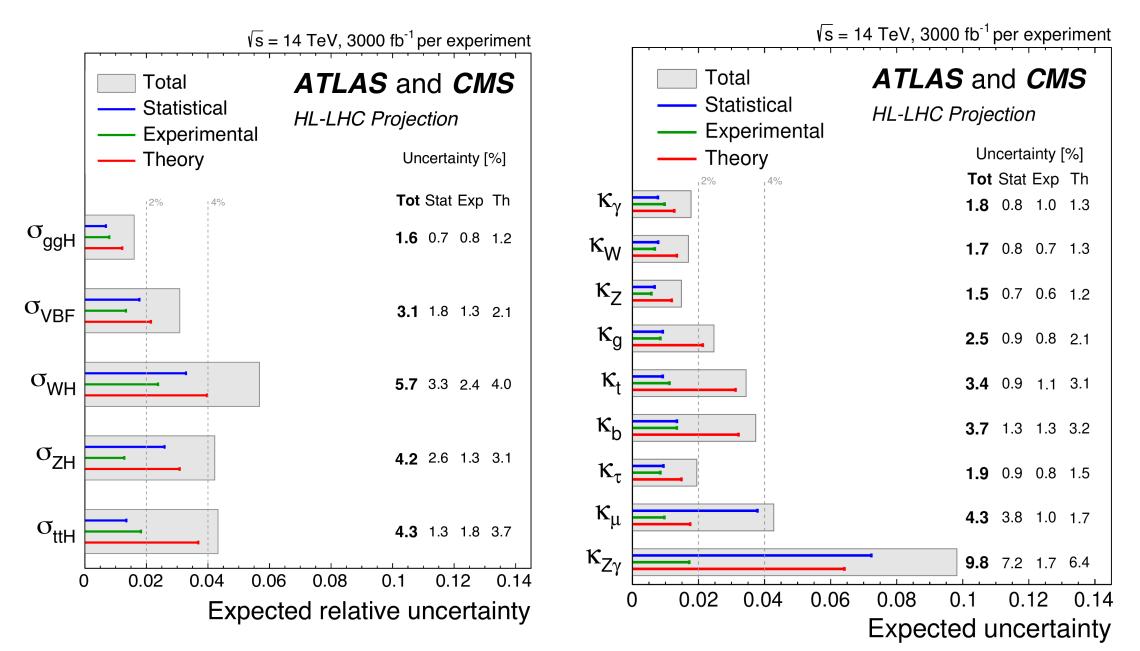
Our experimental uncertainty approaches, or dips below, the current uncertainty from theory.





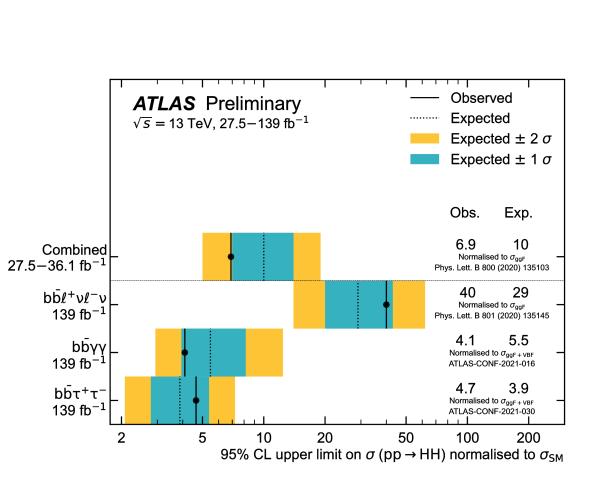
56

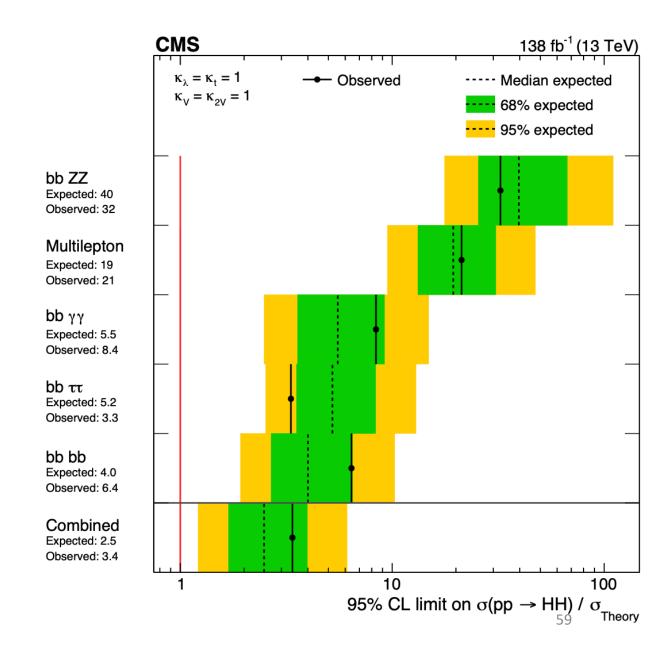
Higgs boson Production and Decay Modes



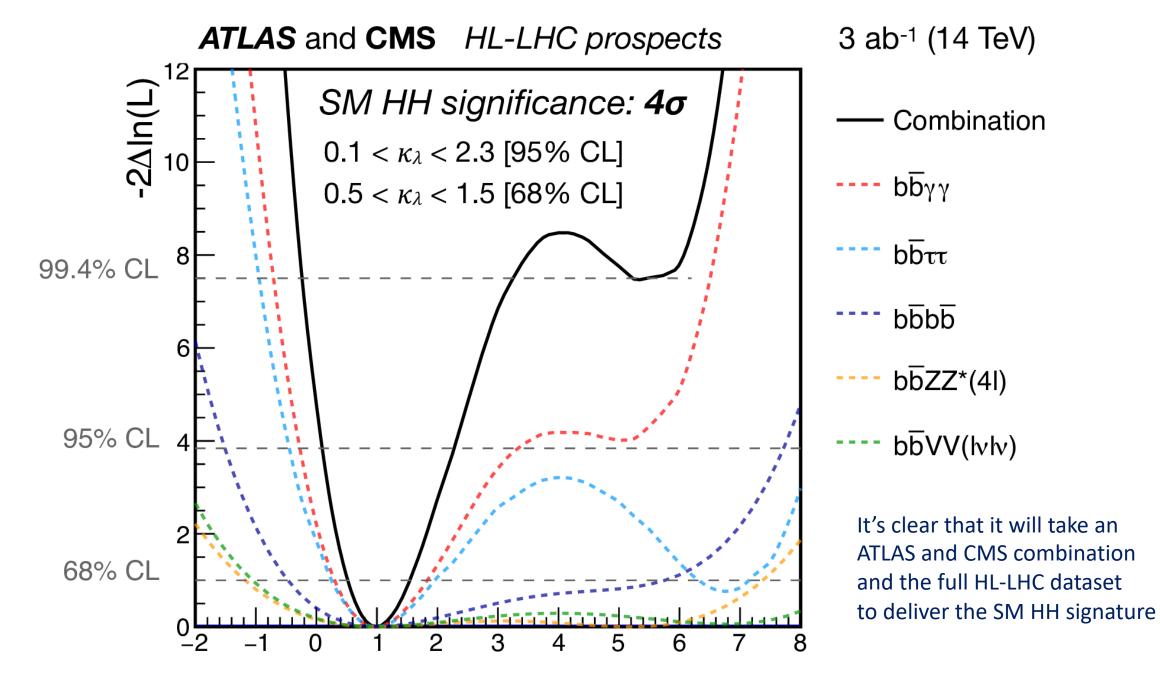
Higgs boson self-interaction

- Accessing the Higgs boson self-interaction is a primary goal of the HL-LHC
 - Central to our understanding of the vacuum instability and inflation
- Measuring λ_{HHH} would constrain the shape of the Higgs potential and contribute to our understanding of EW symmetry breaking
 - new dynamics at higher energy scales could modify the Higgs self coupling
- HH production provides our most sensitive window at the HL-LHC
- Extremely challenging measurement, since the rate is a factor of 1000 smaller than single Higgs production. We need to exploit all modes and combine ATLAS and CMS measurements to reach sensitivity.

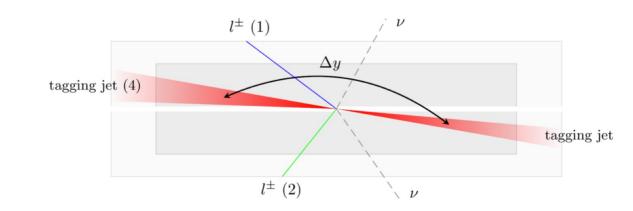


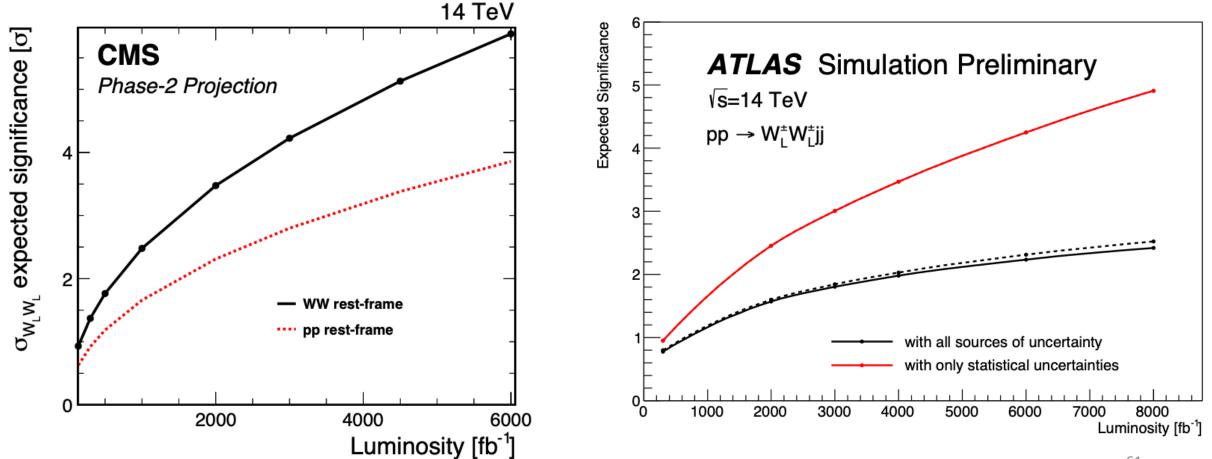


HH Current Picture



Vector boson Scattering: longitudinal mode

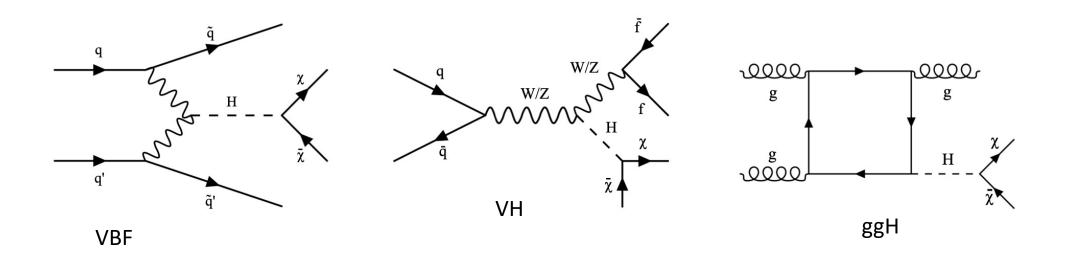




Invisible Width of the Higgs

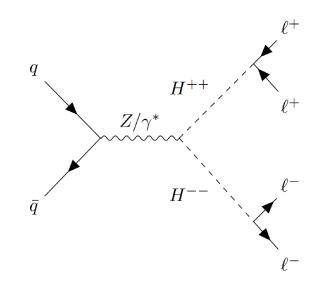
A challenge at a hadron collider is the inability to measure the full width of objects produced. Without a known collision-by-collisions center-of-mass energy, we can't know the # of Higgs bosons produced.

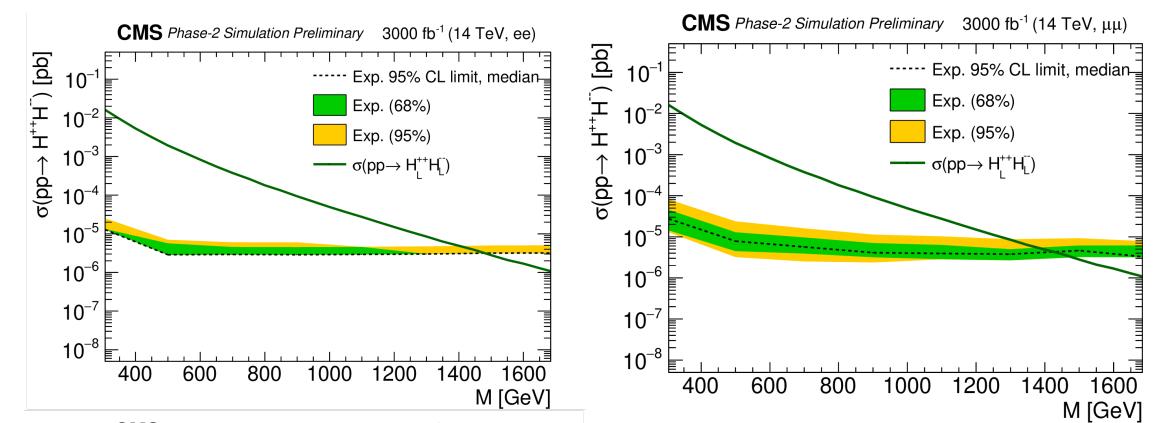
Current limits are at the level of 10%. We expect, with the HL- LHC, to measure $B(H \rightarrow inv) < 2.5$ %.



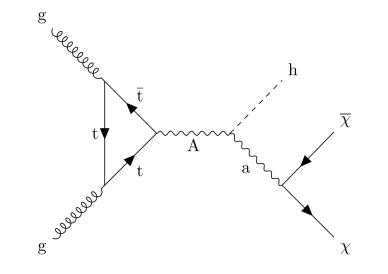
CMS Search for H⁺⁺/H⁻⁻

We can continue the hunt for a doubly-charged Higgs boson. This search benchmarks scenarios where decays are exclusively to electrons (left) or to muons (right)





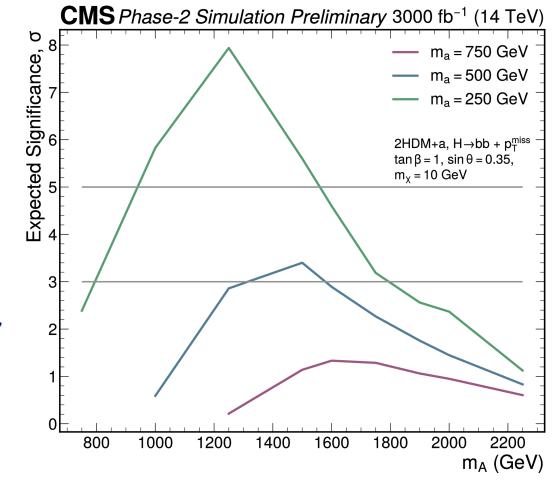
Using the Higgs boson: searches for dark matter

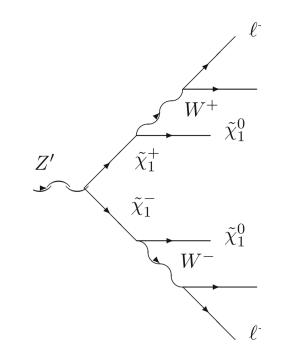


Using boosted events where the higgs decays to a pair of b quarks, with that final state reconstructed in a large radius jet

Interpretation in the 2HDM+a model

Important to remember that we have clear discovery potential for a wide range of new physics with the HL-LHC



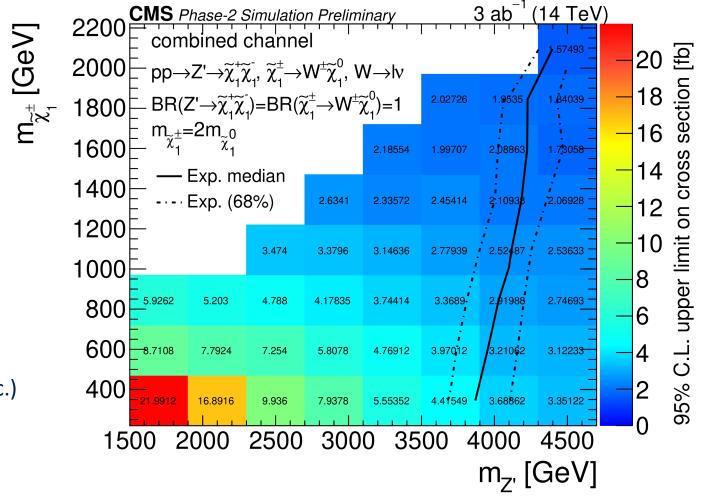


Search motivated by extensions to SUSY

Communicating the results is a challenge, as is optimizing the analysis, due to the huge potential phase space covered (BRs, masses, etc.)

Deep NN used for signal extraction

Search for a "leptophobic" Z'



Conclusions (1)



- Run 3 is well underway. We expect our Run 2 dataset to be significantly expanded.
- Improvements for all LHC experiments will extend the physics reach beyond what a simple extension of Run-2-like data would enable
 - Be on the lookout for combinations of Run 2 and Run 3 data, in both searches and measurements!
- The challenging HL-LHC high-pileup conditions will be supported by many new detector components from the LHC experiments.
 - More sophisticated triggers, with tracking information available early and more analysis moved online, will help select the optimal dataset for physics exploration within the constraints.
 - Hallmarks include higher granularity information and improved resolutions, along with a new category of information via dedicated timing detectors.

Conclusions (2)







- There is uncertainty in our near-term, and long-term, futures.
- We are not isolated from our environments:
 - CERN, like the rest of Europe, is impacted by the invasion on the continent and the energy crisis. This impacts our integrated luminosity.
- Long term: we carry the future of the energy frontier with us, collectively. We must extract as much as we can from LHC and HL-LHC data as early as we can to best inform the path forward.
- The experience of the pandemic has influenced all of us.
 - We learned some lessons about taking care of ourselves and each other that we should not forget as things return more to "normal".
- It's an incredibly exciting time at CERN, as we launch into the third run of the LHC and have the HL-LHC to look forward to. It is useful to keep in mind that there are few better generators of creativity than challenges and constraints.

Thank you for your attention, and nice to be with you in person!



And thank you to colleagues across many experiments for the beautiful papers/notes/talks that informed or populated these slides!

Additional Material

- Machine Learning
 - New <u>PDG</u> Chapter
 - Living Review: https://iml-wg.github.io/HEPML-LivingReview/
- Run 3 Prospects
 - CMS Briefing
- HL-LHC Prospects
 - "<u>Yellow Report</u>" input for European Strategy
 - ATLAS & CMS <u>Snowmass input</u>