

Measurement of the Branching Ratio of $B^0_{(s)}$ mesons to $\mu^+\mu^-(\gamma)$ at LHCb

Flavour Physics

[Phys. Rev. D105 \(2022\) 012010](#)

[Phys. Rev. Lett. 128 \(2022\) 041801](#)

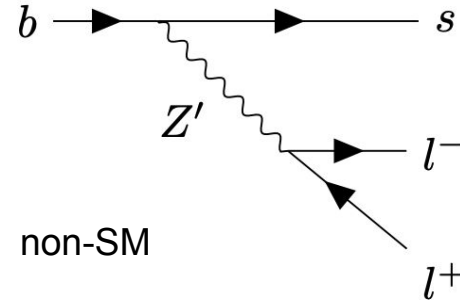
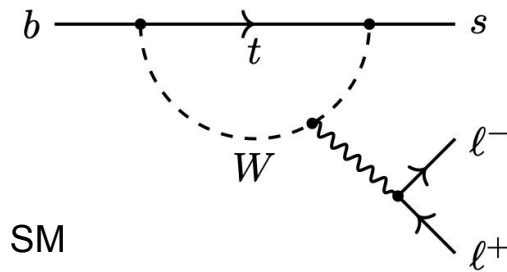
Jayashri on behalf of Group B

Group B: Quentin, Jiwoong, Laszlo, Lakshmi, Anežka, Aravind, Mike, Eszter, Ashish, Jinyoung, Francesco, Ritik, Madeeha, Jonas, Malinda



Introduction

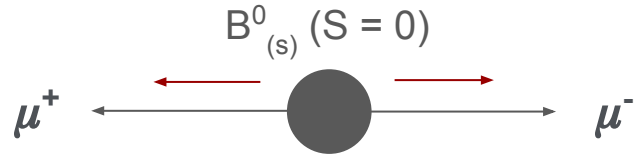
- Flavour changing neutral currents (FCNC) forbidden at tree level in the Standard Model (SM)
- $b \rightarrow sll$ transition can only occur via loop \rightarrow rare decays



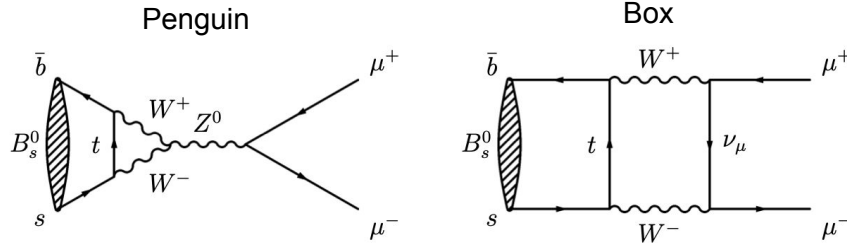
- Process could be mediated by **new virtual particles** \rightarrow probe for New Physics (NP)
 - Z' gauge boson, leptoquarks, non-SM Higgs

Indirect searches for NP using precision measurements of FCNC processes.

$B_{(s)}^0 \rightarrow \mu^+ \mu^-$ Decay



- $B_{(s)}^0$ decays to 2 muons: **loop** and **helicity suppressed**



- Effective field theory (EFT) gives SM prediction for branching ratios (BR)

$$\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-)_{\text{SM}} = (3.66 \pm 0.14) \times 10^{-9}$$

[Phys. Rev. Lett. 112, 101801 \(2014\)](#)

$$\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-)_{\text{SM}} = (1.03 \pm 0.05) \times 10^{-10}$$

[JHEP10\(2019\)232](#)

- Previous measurement:

- LHCb + ATLAS + CMS combined measurement at 95% CL, consistent with the SM within 2σ

$$\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-) = (2.69_{-0.35}^{+0.37}) \times 10^{-9}$$

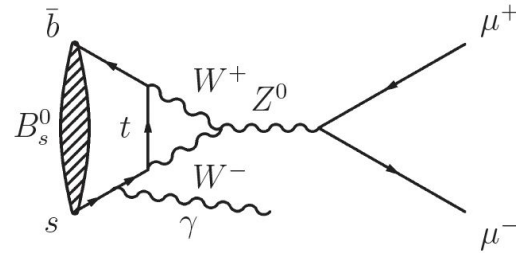
[ATLAS-CONF-2020-049](#)

$$\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-) < 1.9 \times 10^{-10}$$

Analysis Goals and $B_s^0 \rightarrow \mu^+ \mu^- \gamma$

- Measure BRs for

- $B^0 \rightarrow \mu^+ \mu^-$
- $B_s^0 \rightarrow \mu^+ \mu^-$ *
- $B_s^0 \rightarrow \mu^+ \mu^- \gamma$



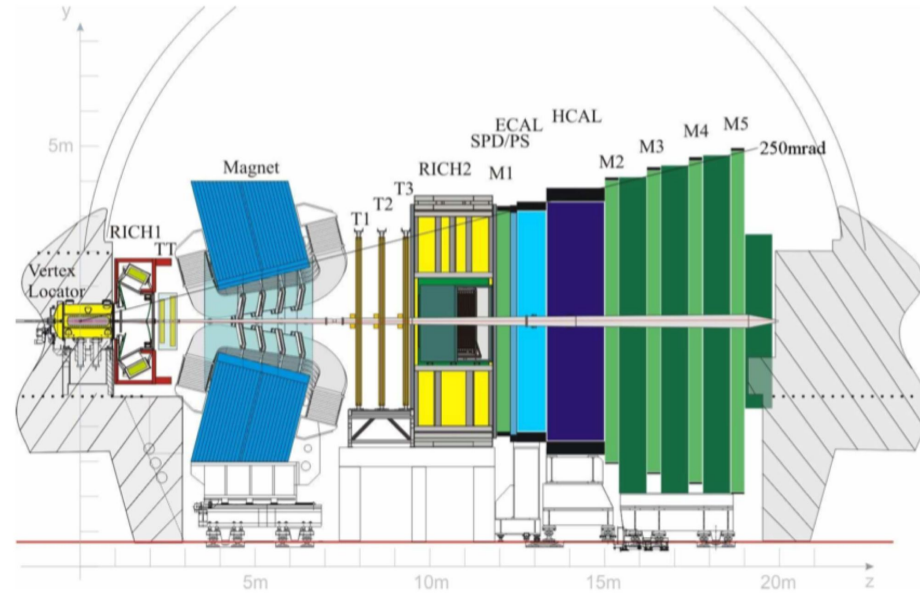
- B_s^0 decays to 2 muons can be accompanied by the emission of photon **
 - Not helicity suppressed
- Sensitivity to initial-state radiation (ISR) from the valence quarks at high dimuon mass
 - SM BR prediction $O(10^{-10})$
 - Final-state radiation (FSR) included in B_s^0 BR

*The paper also reports the measurement of the effective lifetime for $B_s^0 \rightarrow \mu^+ \mu^-$ which is not included in this talk.

**Note that B^0 ISR is negligible due to additional CKM suppression.

LHCb Detector

- **Forward-spectrometer** specialized on b and c decays, located at the LHC interaction point 8
- **Vertex Locator at interaction region**
 - Only 7 mm distance to beam
 - Excellent impact parameter resolution to identify secondary vertices
 - Decay-time resolution < 50 fs (decay-time measurement)
- **Ring-Imaging-Cherenkov detectors** for particle identification (relevant for normalization channels)
- Tracking system with momentum resolution $< 1.0\%$
→ **good mass resolution**
- Muon system to identify and trigger muons



[JINST 3 \(2008\) S08005](#)

Signal Selection

Signal Region

$$B^0_{(s)} \rightarrow \mu^+ \mu^- (\gamma)$$

Trigger: Number of muon > 0

Muon Selection

- $0.25 < p_T < 40$ GeV
- $p < 500$ GeV
- Inconsistent with PV
- Muon identification

B^0_S Selection

- Decay time < 13.25ps
- $p_T > 0.5$ GeV
- Originated at least one PV
- J/ψ veto^[1]
- $4900 \leq M(\mu\mu) \leq 6000$ MeV

Control Region (For normalization)

$$B^0_s \rightarrow hh^{[2]}$$

Triggered independently

Replacing muon selection to hadron selection

Same but no J/ψ veto

$$B^+ \rightarrow J/\psi^{[3]} K^+(\phi)$$

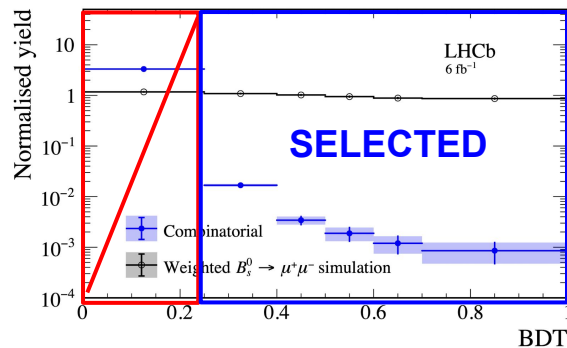
of muon > 0

Same muon selection as signal

- Same but $M(\mu\mu)$ close to J/ψ mass
- Kaon inconsistent with PV

*Additional loose BDT based cut is applied to all regions

Selection using boosted decision tree (BDT)



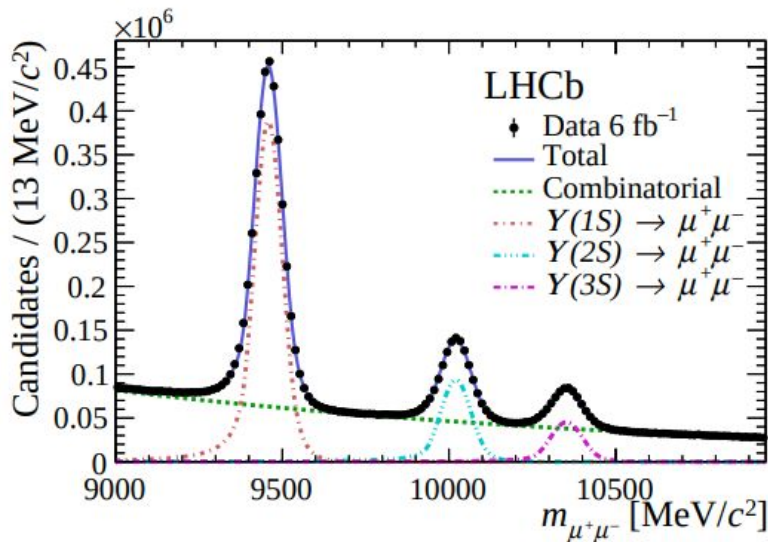
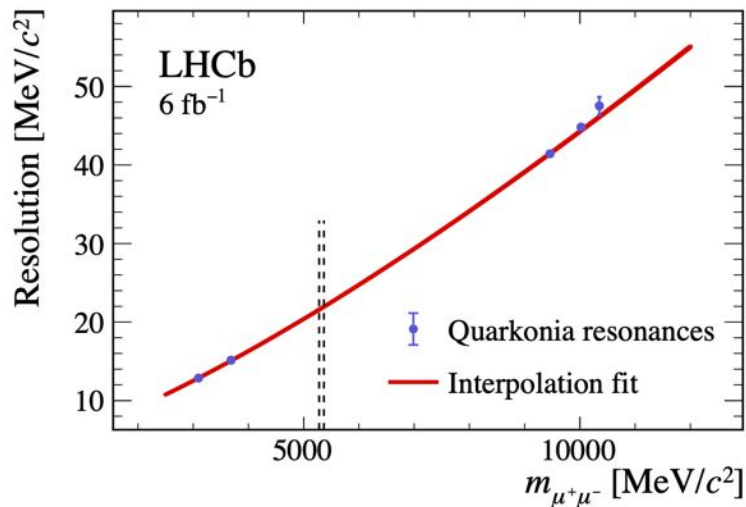
- Discriminating combinatorial background: $bb\bar{b} \rightarrow \mu\mu$
- Optimized using MC samples ($B^0_S \rightarrow \mu\mu$ vs $bb\bar{b} \rightarrow \mu\mu X$)
- Input features: topological variables*

*[Phys. Rev. Lett. 128 \(2022\) 041801](https://arxiv.org/abs/2108.08111)

[1] J/ψ veto : $|M(\mu\mu) - M(J/\psi)| > 30$ MeV [2] h = K or π [3] J/ψ decays to muon pair

Signal Calibration: Mass Shape

- Signal described with a double-sided Crystal-Ball function
- **Mean** of $B_s^0 \rightarrow \mu^+\mu^-$ and $B^0 \rightarrow \mu^+\mu^-$ signal peaks calibrated with $B_s^0 \rightarrow K^+K^-$ and $B^0 \rightarrow K^+\pi^-$ data
- **Resolution** interpolated from mass fits to cc and bb resonances
- **Tail parameters** extracted from simulation



Relative BR Measurement

$$N_{sig} = N_{Tot} f_{sig} B_{sig} \epsilon_{sig}$$

- BR based on
 - Total number of events N_{Tot}
 - Signal yield N_{sig}
 - Selection efficiency ϵ_{sig}
 - Hadronization probability f_{sig}
- **Large uncertainties cancel out in the ratio, e.g., $\sigma(pp \rightarrow bb)$ and luminosity**

$$\frac{B_{sig}}{B_{norm}} = \frac{f_{norm}}{f_{sig}} \frac{\epsilon_{norm}}{\epsilon_{sig}} \frac{N_{sig}}{N_{norm}}$$

Normalization

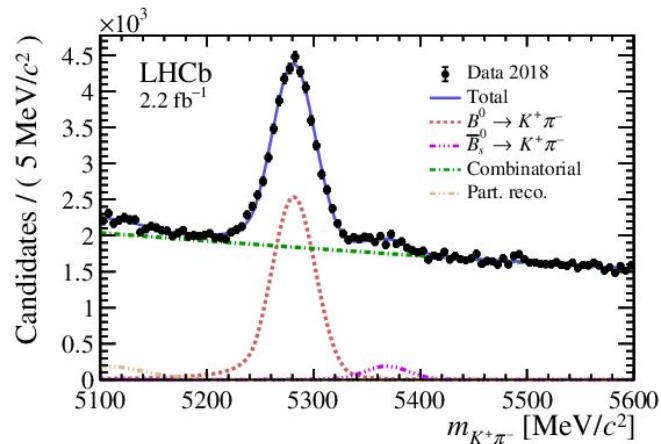
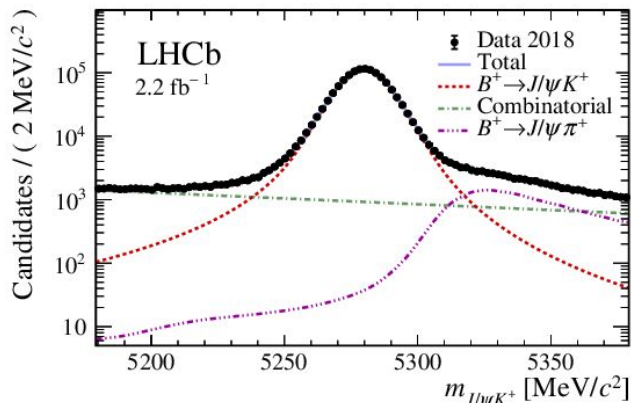
$$\mathcal{B}(B_{(s)}^0 \rightarrow \mu^+ \mu^- (\gamma)) = \frac{f_{\text{norm}}}{f_{\text{sig}}} \frac{\varepsilon_{\text{norm}}}{\varepsilon_{\text{sig}}} \frac{N_{\text{sig}}}{N_{\text{norm}}} \mathcal{B}_{\text{norm}} = \alpha_{\text{sig}} N_{\text{sig}}$$

$\alpha_{\text{sig}} \rightarrow$ single event sensitivity

- BRs measured relative to well known normalization channels with similar selections
 - $B^+ \rightarrow J/\psi (\rightarrow \mu^+ \mu^-) K^+$
 - $B_0 \rightarrow K^+ \pi^-$
- Extraction of normalization channel yields via maximum likelihood fit
- Efficiencies calculated for signal and normalization channels**
- External input for hadronization ratios f_s/f_d (f_s/f_u) \rightarrow large systematic

$$f_s/f_d (13 \text{ TeV}) = 0.254 \pm 0.008$$

YIELD	Period	$B^+ \rightarrow J/\psi K^+$	$B^0 \rightarrow K^+ \pi^-$
	Run 1	$(11.259 \pm 0.015) \times 10^5$	$(14.05 \pm 0.26) \times 10^3$
Run 2	$(36.072 \pm 0.026) \times 10^5$	$(7.99 \pm 0.10) \times 10^4$	



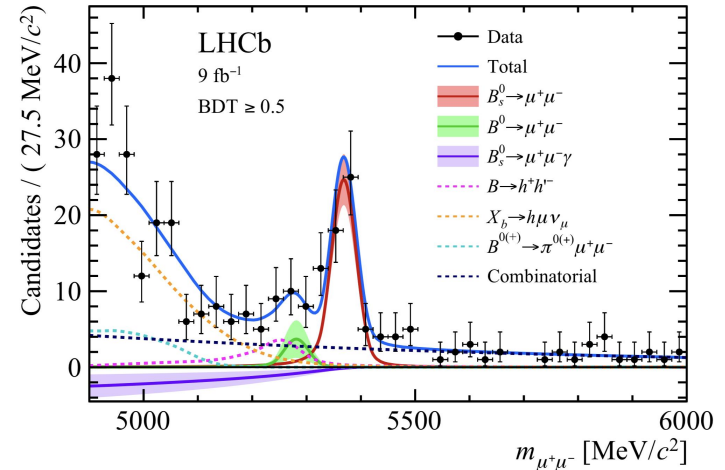
Backgrounds

Physical

- One or more misidentified final state particles (protons, pions, and kaons)
- Use stringent PID requirement to suppress it
- One particle is not reconstructed
 - $B^{(0)+} \rightarrow \pi^{(0)+} \mu^+ \mu^-$
- One hadron is misidentified as muon
 - $X_b \rightarrow h \mu \nu_\mu$
- Two hadrons are misidentified as muons
 - $B^0_{(S)} \rightarrow h^+ h^-$
- **Dominant in low-mass region**

Combinatorial

- Random combination of two muons not originating from a common decay
- Suppress by forming B candidate using topology differences and relative isolation of the muon tracks
- **Dominant in high-mass region**



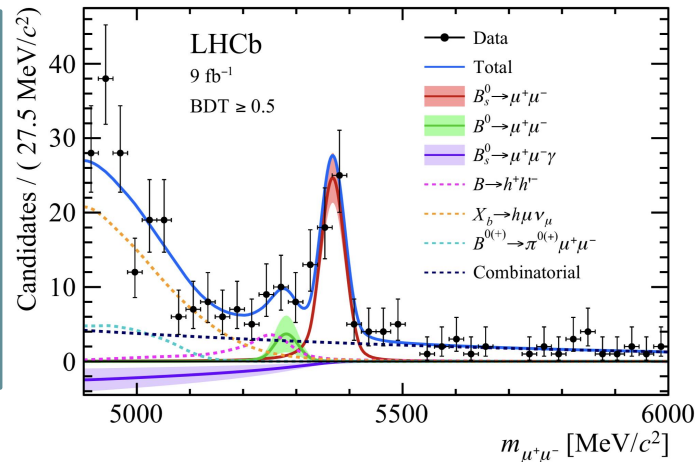
Measurement of Signal BR

Signal yields extracted by **unbinned extended maximum-likelihood fit** with constraint background components

Resulting BRs

$$\begin{aligned}\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-) &= (3.09_{-0.43}^{+0.46} {}_{-0.11}^{+0.15}) \times 10^{-9}, \\ \mathcal{B}(B^0 \rightarrow \mu^+ \mu^-) &= (1.20_{-0.74}^{+0.83} \pm 0.14) \times 10^{-10}, \\ \mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^- \gamma) &= (-2.5 \pm 1.4 \pm 0.8) \times 10^{-9} \text{ with } m_{\mu\mu} > 4.9 \text{ GeV}/c^2.\end{aligned}$$

- Stat. by repeating the fit with all nuisance parameters fixed
- Syst. by subtracting in quadrature the stat. uncertainties from the total



Main contribution to the systematic uncertainties

$$\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-) : f_s/f_d \text{ (3\%)}$$

$$\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-) : B^0 \rightarrow h^+ h^- \text{ and semileptonic b-hadron background (9\%)}$$

Limit Setting of Signal BR

Significance

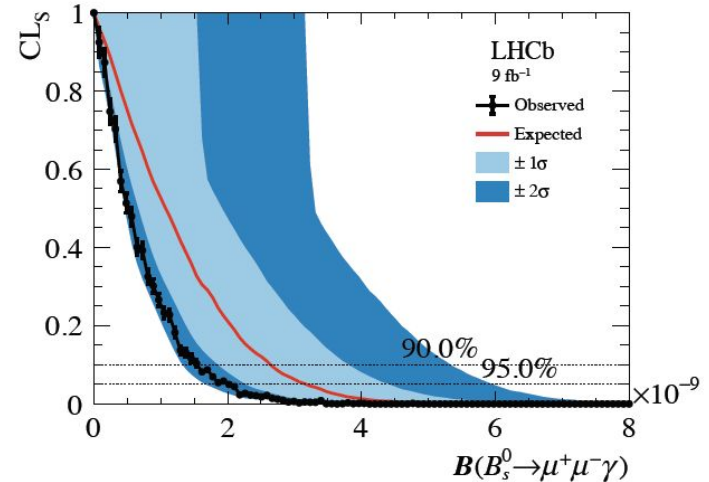
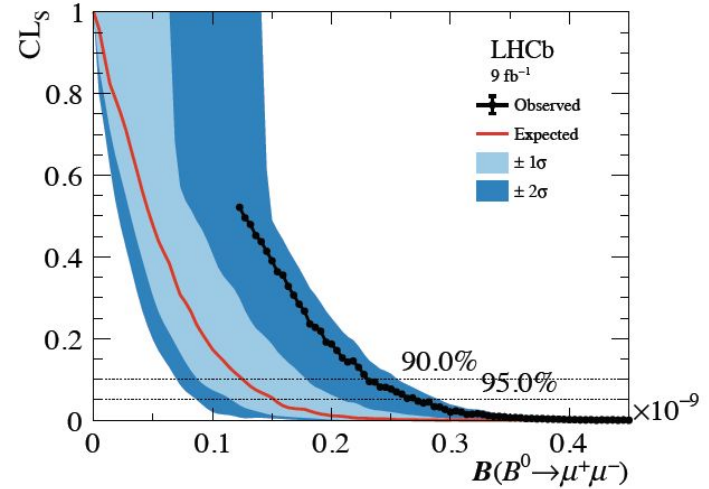
- More than 10σ for $\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-)$
- 1.7σ for $\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-)$
- 1.5σ for $\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^- \gamma)$

No evidence \rightarrow “Limit”

CL_s method with a one-sided test statistic is used for determining upper limits.

Limit on the BRs

$$\begin{aligned}\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-) &< 2.3 (2.6) \times 10^{-10}, \\ \mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^- \gamma) &< 1.5 (2.0) \times 10^{-9} \text{ with } m_{\mu\mu} > 4.9 \text{ GeV}/c^2, \\ &\text{at } 90\% (95\%) \text{ CL.}\end{aligned}$$



Limit Setting of Signal BR

Significance

- More than 10σ for $\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-)$
- 1.7σ for $\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-)$
- 1.5σ for $\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^- \gamma)$

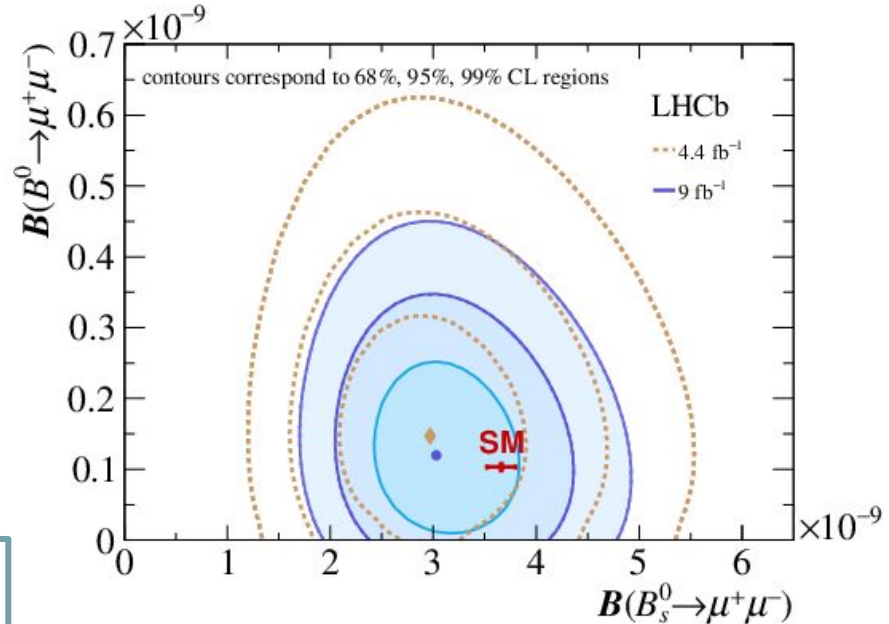
No evidence → “**Limit**”

CL_s method with a one-sided test statistic is used for determining upper limits.

Limit on the BRs

$$\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-) < 2.3 (2.6) \times 10^{-10},$$
$$\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^- \gamma) < 1.5 (2.0) \times 10^{-9} \text{ with } m_{\mu\mu} > 4.9 \text{ GeV}/c^2,$$

at 90% (95%) CL.



In agreement with the SM predictions

Conclusion

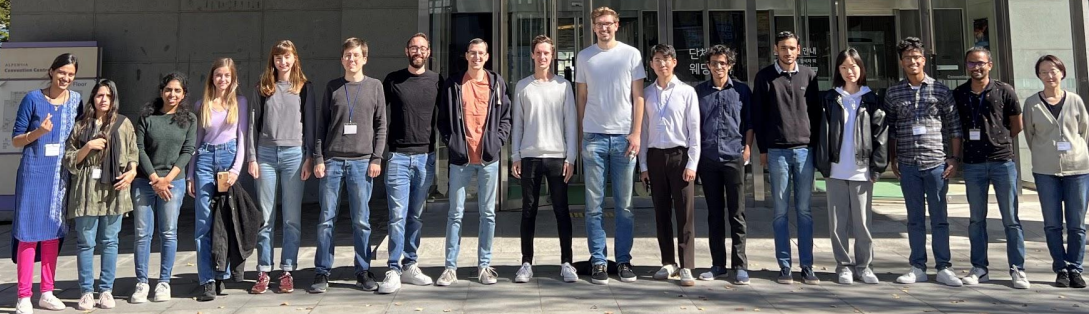
- Rare decays of $B^0_{(s)}$ to two muons \rightarrow powerful probe for NP
- Most precise single experiment measurement
 - BR of $B^0_s \rightarrow \mu^+\mu^-$ measured ($> 10\sigma$ significance)
 - $\mathcal{B}(B^0_s \rightarrow \mu^+\mu^-) = (3.09^{+0.46+0.15}_{-0.43-0.11}) \times 10^{-9}$,
 - Improved limits set for $B^0 \rightarrow \mu^+\mu^-$ and $B^0_s \rightarrow \mu^+\mu^-\gamma$
 - $\mathcal{B}(B^0 \rightarrow \mu^+\mu^-) = (1.20^{+0.83}_{-0.74} \pm 0.14) \times 10^{-10}$,
 $\mathcal{B}(B^0_s \rightarrow \mu^+\mu^-\gamma) = (-2.5 \pm 1.4 \pm 0.8) \times 10^{-9}$ with $m_{\mu\mu} > 4.9 \text{ GeV}/c^2$.

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2022. 10. 05 ~ 2022. 10. 18
Pyeongchang, SOUTH KOREA



Back-up

BR Calculation

Time integrated BR:

$$\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-) = \left[\frac{1 + A_{\Delta\Gamma_s}^{\mu\mu} y_s}{1 - y_s^2} \right] \mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-)_{t=0}$$

where

$$y_s \equiv \Delta\Gamma_s / (2\Gamma_s) = 0.065 \pm 0.003$$

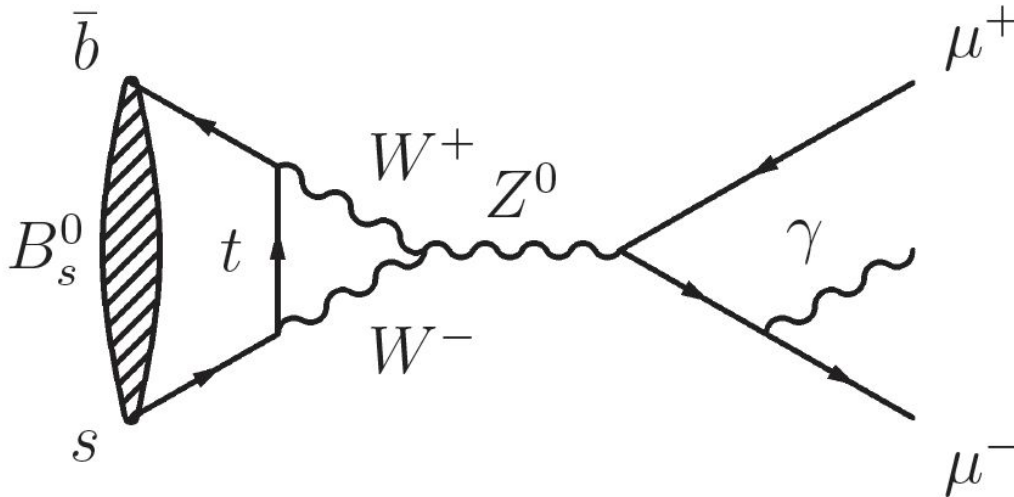
$$A_{\Delta\Gamma_s}^{\mu\mu} \equiv -2\Re(\lambda) / (1 + |\lambda|^2), \text{ with } \lambda = (q/p)(A(\overline{B}_s^0 \rightarrow \mu^+ \mu^-) / A(B_s^0 \rightarrow \mu^+ \mu^-))$$

For the SM, $A_{\Delta\Gamma_s}^{\mu\mu} = 1$ as only CP odd eigenstate of $B_{(s)}^0$ decays two muons $\rightarrow \mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-)_{\text{SM}} = (3.66 \pm 0.14) \times 10^{-9}$

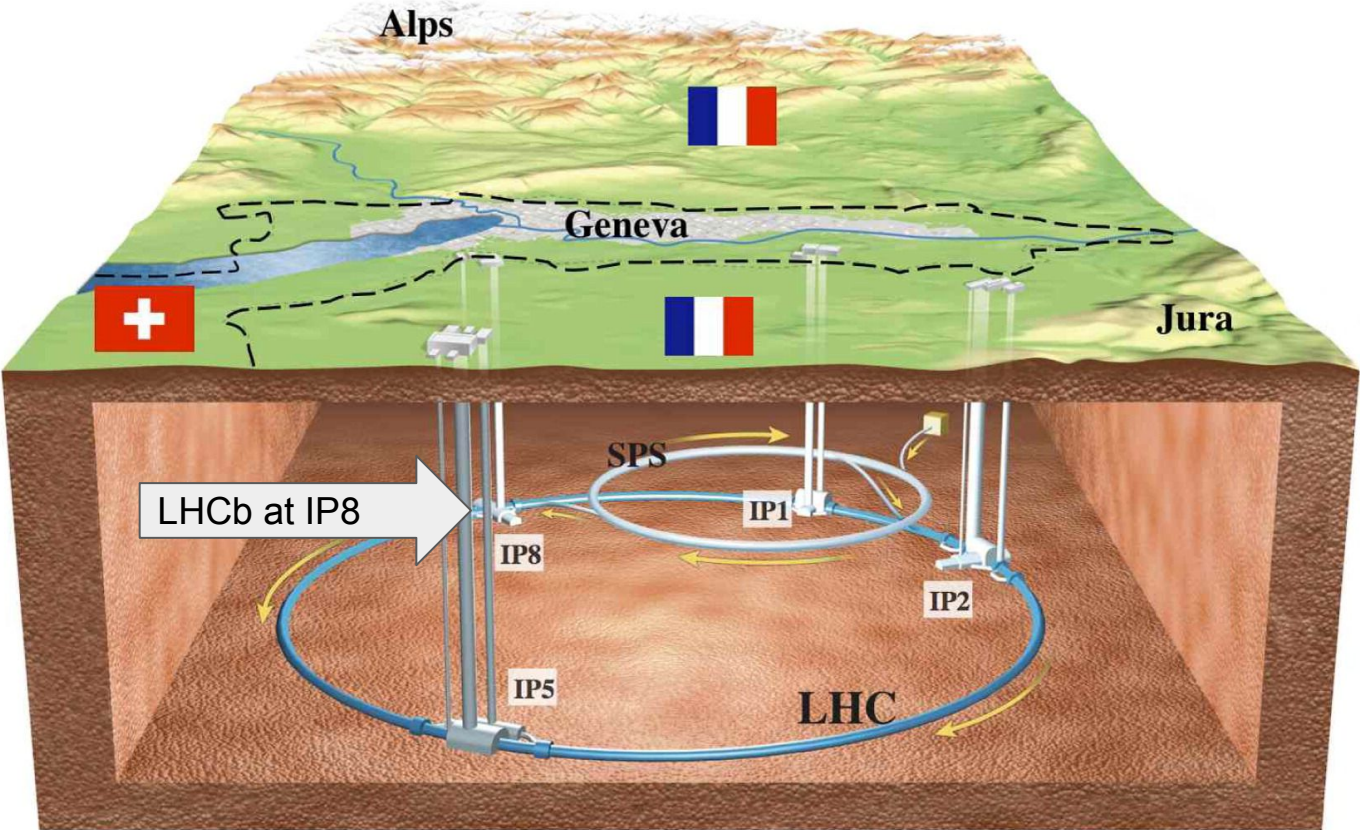
In presence of NP, it can assume any value from -1 to +1.

Final-State Radiation (FSR)

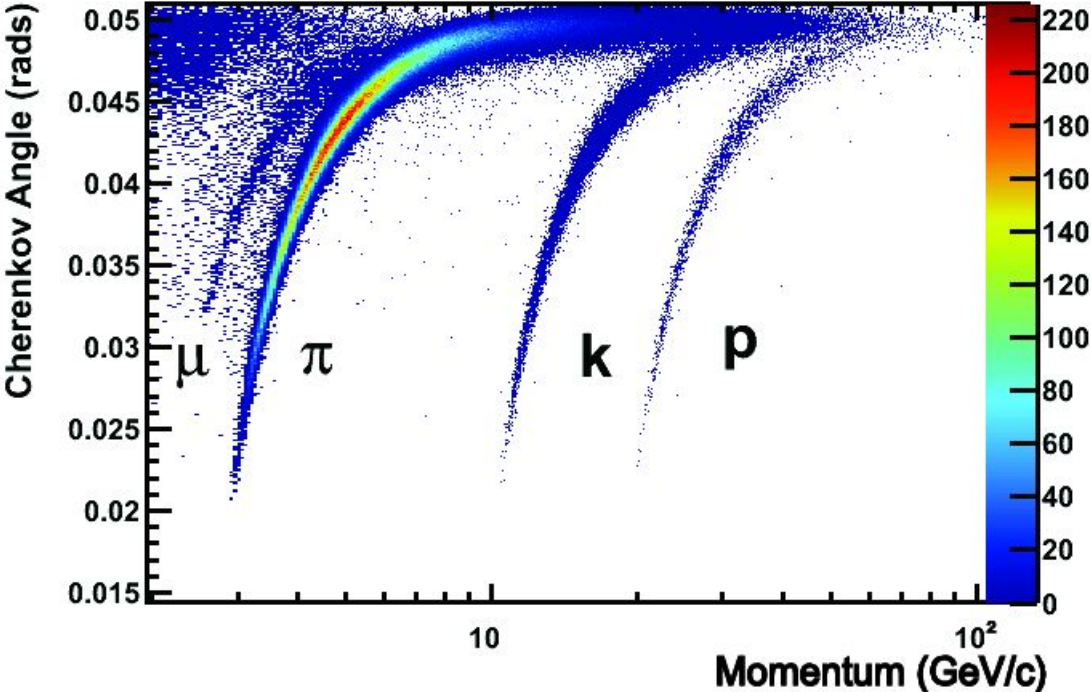
- Included experimentally in $B_s^0 \rightarrow \mu^+ \mu^-$ in the description of the radiative tail in its mass distribution



Detector



Detector

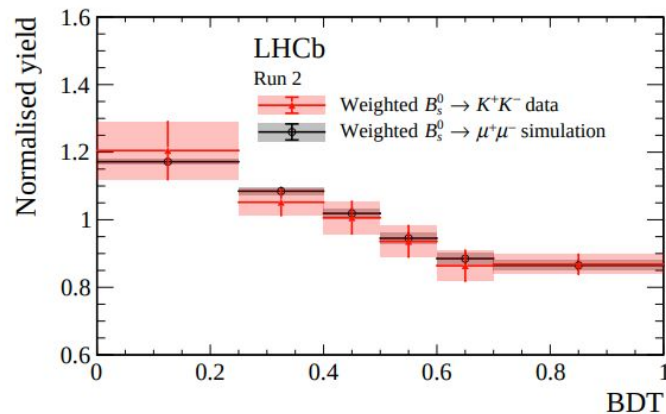
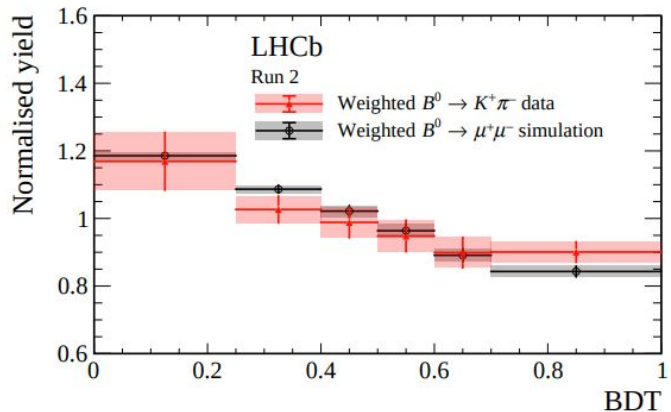


Signal Calibration: BDT

- Events categorised into 6 BDT bins
- Calibration of signal BDT output done using data-corrected simulation*
- Corrections applied to:
 - B meson kinematics and detector occupancy
 - PID and trigger efficiency
 - BDT-lifetime correlations in $B_s^0 \rightarrow \mu^+ \mu^-$
- Fraction of signal decays in each BDT region:

$$f_{\text{BDT},i} = f_{\text{sim},i}^{\mu\mu} \cdot f_{\text{PID},i}^{\mu\mu} \cdot f_{\text{trig},i}^{\mu\mu} (\cdot k_i)$$

- Cross-checked on $B^0 \rightarrow K^+ \pi^-$ and $B_s^0 \rightarrow K^+ K^-$ data
 - Same selection as for the signal, except for trigger and PID requirements



* Reweighting method based on BDT: Extracting $w = \text{data/MC}$ using BDT (arXiv:1608.05806)

BDT Calibration

- **Correction to B meson kinematics:**
 - Boosted decision tree classifier trained to align the data and the simulation
 - p_T , η , and the χ^2_{IP} of the B candidate used as input variables
 - Weights obtained applied to all simulation samples used for calibration and normalisation
- **Correction from event occupancy (measured as the number of tracks in the event):**
 - Weights determined by comparing the relative number of $B^+ \rightarrow J/\psi K^+$ decays in background-subtracted data and simulated samples
- **Corrections to BDT-lifetime correlation for $B^0_s \rightarrow \mu^+ \mu^-$ decays:**
 - Simulation generated using the mean B^0_s lifetime; but effective lifetime may vary between that of light and the heavy mass eigenstates.
 - Correction evaluated for $A^{\mu\mu}_{\Delta\Gamma_s} = -1, 0$ and 1
 - Simulated candidates weighted according to:

$$\omega_j = \frac{\tau_{\text{gen}}}{\tau_{\mu^+\mu^-}} e^{-t_j(1/\tau_{\mu^+\mu^-} - 1/\tau_{\text{gen}})}$$

- Correction factor applied to each BDT region:

$$k_i = \sum_{j=1}^{N_i} \frac{\omega_j}{N_i}$$

Normalization

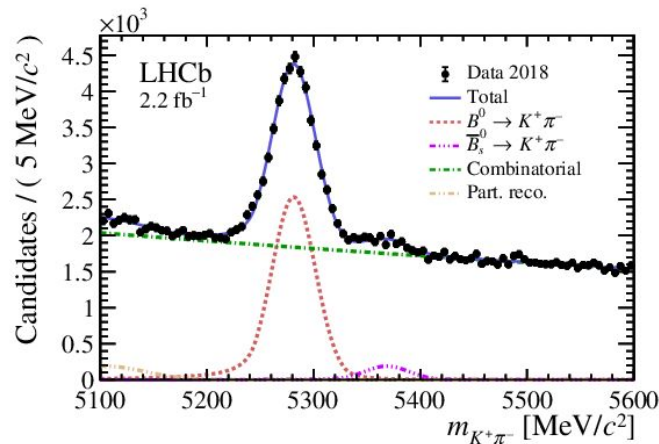
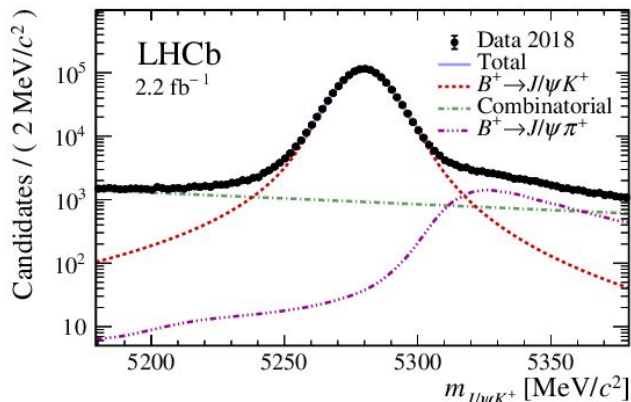
- Branching ratios measured relative to well known normalization channels
 - $B^+ \rightarrow J\psi(\rightarrow \mu^+\mu^-) K^+$ [Br. fr. = $(6.02 \pm 0.17) \times 10^{-5}$]
 - $B_0 \rightarrow K^+ \pi^-$ [Br. fr. = $(1.96 \pm 0.05) \times 10^{-5}$]
- Similar selections, slightly modified (Do we want to give details?)
- Extraction of normalization channel yields via maximum likelihood fit
- External input for hadronization ratios f_s/f_d (f_s/f_u) -> large systematic
- Carefully estimated reconstruction and selection efficiencies

$$\mathcal{B}(B_{(s)}^0 \rightarrow \mu^+\mu^-(\gamma)) = \frac{f_{\text{norm}}}{f_{\text{sig}}} \frac{\varepsilon_{\text{norm}}}{\varepsilon_{\text{sig}}} \frac{N_{\text{sig}}}{N_{\text{norm}}} \mathcal{B}_{\text{norm}} = \alpha_{\text{sig}} N_{\text{sig}}$$

- $\alpha_{\text{sig}} \rightarrow$ single event sensitivity

$$f_s/f_d (13 \text{ TeV}) = 0.254 \pm 0.008$$

YIELD	Period	$B^+ \rightarrow J/\psi K^+$	$B^0 \rightarrow K^+ \pi^-$
	Run 1		$(11.259 \pm 0.015) \times 10^5$
Run 2		$(36.072 \pm 0.026) \times 10^5$	$(7.99 \pm 0.10) \times 10^4$



Efficiencies

For signal only



$$\varepsilon_{\text{norm(sig)}} = \varepsilon_{\text{RecSel}} \cdot \varepsilon_{\text{PID}} \cdot \varepsilon_{\text{Trig}} (\cdot \varepsilon_{\text{BDT}})$$

- Acceptance and reconstruction efficiency
 - Obtained from MC simulation with correction factors applied for simulation modelling
- PID: Particle Identification Efficiency
 - Obtained from high purity sample from data
 - μ from $J/\psi \rightarrow \mu\mu$
 - K^\pm from $D^0 \rightarrow K^\mp \pi^\pm$
 - p from $\Lambda^0 \rightarrow p\pi^-$ and $\Lambda_c^+ \rightarrow p K \pi$
- Trigger efficiency

Efficiencies

$\varepsilon_{\text{RecSel}}$		
	Run 1	Run 2
$B_s^0 \rightarrow \mu^+ \mu^-$	0.0602 ± 0.0003	0.0640 ± 0.0004
$B^0 \rightarrow \mu^+ \mu^-$	0.0594 ± 0.0003	0.0635 ± 0.0004
$B_s^0 \rightarrow \mu^+ \mu^- \gamma$	0.0508 ± 0.0003	0.0546 ± 0.0004
$B^0 \rightarrow K^+ \pi^-$	0.0462 ± 0.0007	0.0500 ± 0.0006
$B^+ \rightarrow J/\psi K^+$	0.0290 ± 0.0003	0.0305 ± 0.0003

ε_{PID}		
	Run 1	Run 2
$B_s^0 \rightarrow \mu^+ \mu^-$	$0.8580 \pm 0.0006 \pm 0.0053$	$0.8822 \pm 0.0003 \pm 0.0039$
$B^0 \rightarrow \mu^+ \mu^-$	$0.8518 \pm 0.0007 \pm 0.0063$	$0.8759 \pm 0.0004 \pm 0.0046$
$B_s^0 \rightarrow \mu^+ \mu^- \gamma$	$0.8487 \pm 0.0006 \pm 0.0088$	$0.8785 \pm 0.0003 \pm 0.0064$
$B^0 \rightarrow K^+ \pi^-$	$0.4741 \pm 0.0049 \pm 0.0010$	$0.5004 \pm 0.0027 \pm 0.0012$
$B^+ \rightarrow J/\psi K^+$	1.0096 ± 0.0005	1.00260 ± 0.00018

$\varepsilon_{\text{Trig}}$		
	Run 1	Run 2
$B_s^0 \rightarrow \mu^+ \mu^-$	$0.9579 \pm 0.0033 \pm 0.0164$	$0.9712 \pm 0.0014 \pm 0.0093$
$B^0 \rightarrow \mu^+ \mu^-$	$0.9570 \pm 0.0032 \pm 0.0176$	$0.9708 \pm 0.0014 \pm 0.0097$
$B_s^0 \rightarrow \mu^+ \mu^- \gamma$	$0.9538 \pm 0.0032 \pm 0.0195$	$0.9694 \pm 0.0013 \pm 0.0111$
$B^0 \rightarrow K^+ \pi^-$	$0.0433 \pm 0.0002 \pm 0.0016$	$0.0727 \pm 0.0002 \pm 0.0020$
$B^+ \rightarrow J/\psi K^+$	$0.8810 \pm 0.0040 \pm 0.0080$	$0.9033 \pm 0.0016 \pm 0.0089$

ε_{BDT}		
	Run 1	Run 2
$B_s^0 \rightarrow \mu^+ \mu^-$	0.723 ± 0.006	0.7071 ± 0.0026
$B^0 \rightarrow \mu^+ \mu^-$	0.720 ± 0.006	0.7036 ± 0.0027
$B_s^0 \rightarrow \mu^+ \mu^- \gamma$	0.656 ± 0.007	0.6531 ± 0.0035

Cross-Check on the Fragmentation Fraction

To cross-check the ratio of the B_s^0 and B^+ fragmentation fractions and its stability over the data taking, the ratio of $B^+ \rightarrow J/\psi K^+$ and $B_s^0 \rightarrow J/\psi \phi$ efficiency-corrected yields,

$$\mathcal{R} = \frac{N_{B_s^0 \rightarrow J/\psi \phi}}{N_{B^+ \rightarrow J/\psi K^+}} \frac{\varepsilon_{B^+ \rightarrow J/\psi K^+}}{\varepsilon_{B_s^0 \rightarrow J/\psi \phi}} = \frac{f_s}{f_u} \frac{\mathcal{B}(B_s^0 \rightarrow J/\psi \phi)}{\mathcal{B}(B^+ \rightarrow J/\psi K^+)}, \quad (12)$$

is also measured, following a similar approach to Ref. [82]. The ratios are found to be similar to those measured in Ref. [82], although the two methods explore different kinematic regions. A dependence on the centre-of-mass energy is seen and found to be consistent with Ref. [82] and the combined analysis of Ref. [61], justifying the use of different f_s/f_d values for the Run 1 and Run 2 data samples.

Reweighting with BDT

$$\text{multiplier}_{\text{bin}} = \frac{w_{\text{bin, RD}}}{w_{\text{bin, MC}}}$$

$$\chi^2 = \sum_{\text{leaf}} \frac{(w_{\text{leaf, MC}} - w_{\text{leaf, RD}})^2}{w_{\text{leaf, MC}} + w_{\text{leaf, RD}}}.$$

Basic weighting method

- Counting histogram
- Limitation: reweighting of one variable often bring disagreement in others



Reweighting with ML

- Define loss function χ^2 as a difference of Data-MC
- Find the region that has maximum χ^2 and minimize the difference

Measurement of Signal BR

