



Stockholm
University

CLASSICAL & QUANTUM

PROBES OF

AXION-LIKE PARTICLES

M.C. DAVID MARSH
STOCKHOLM UNIVERSITY

COSMO '22, RIO DE JANEIRO, 23RD AUGUST, 2022

Stockholm



Pierluca Carenza



Ramkishor Sharma



Axel Brandenburg



Eike Müller

Cambridge



James Matthews



Julia Sisk-Reynes

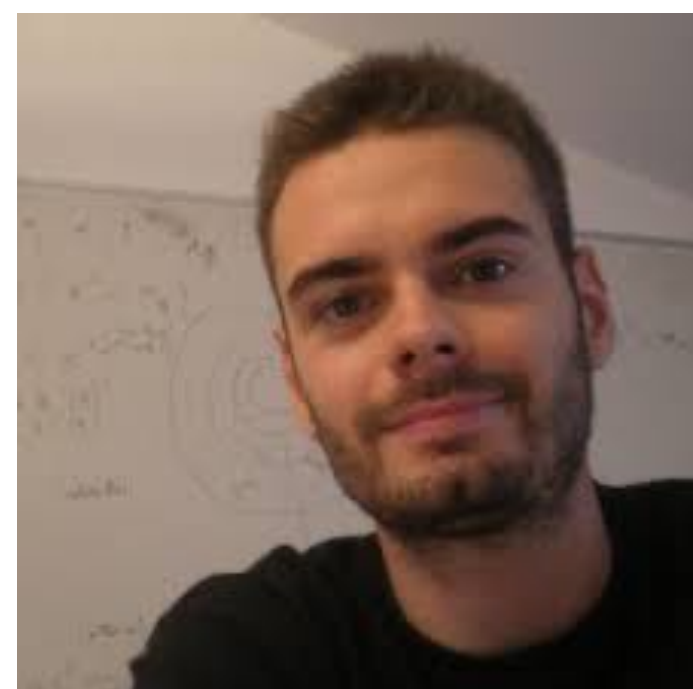


Christopher Reynolds



Helen Russell

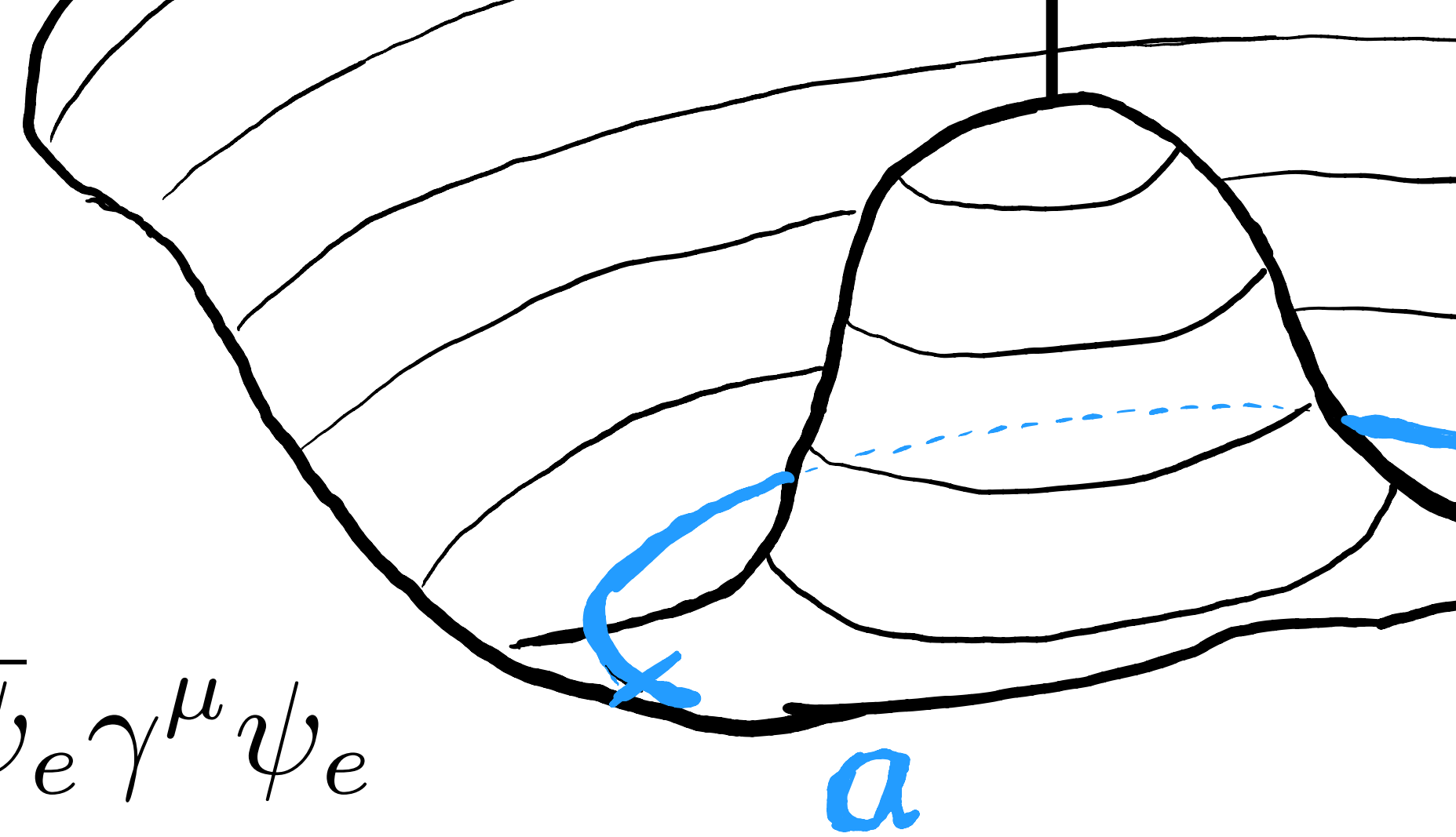
Barcelona



Ricardo Ferreira

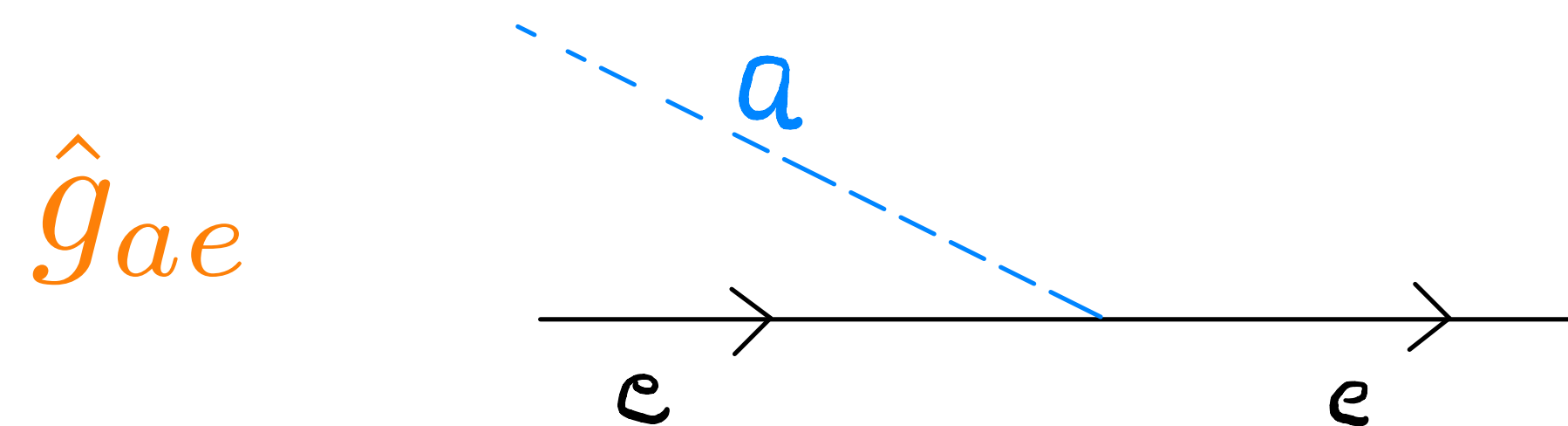
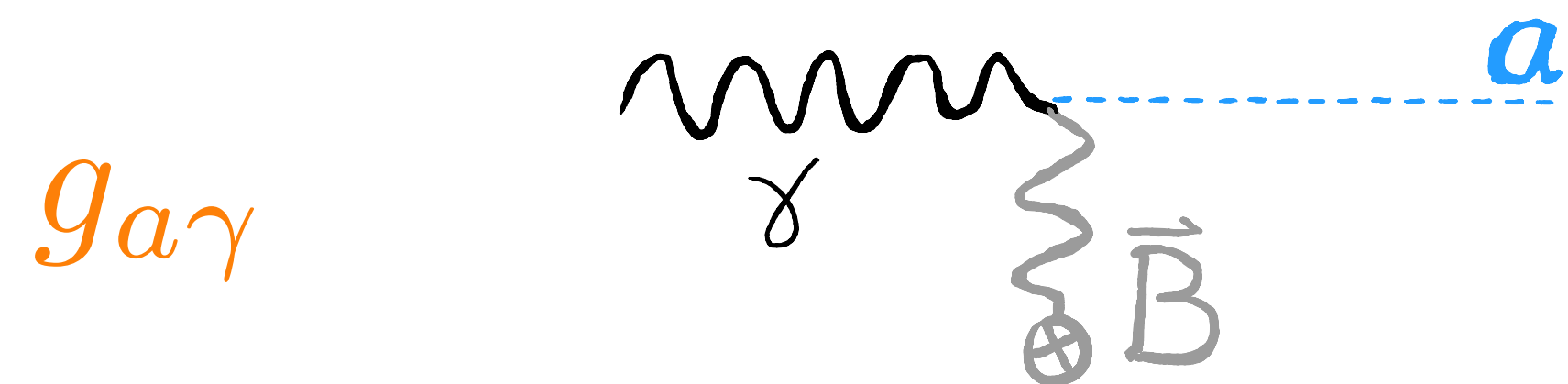
Axion-like particles

[cf. Francesc's talk]



$$\mathcal{L} \supset \frac{1}{2} m_a^2 a^2 + \frac{g_{a\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu} + \hat{g}_{ae} (\partial_\mu a) \bar{\psi}_e \gamma^\mu \psi_e$$

m_a — from negligible to large



How can we use astrophysics to probe this theory?

Classical ALP-photon mixing

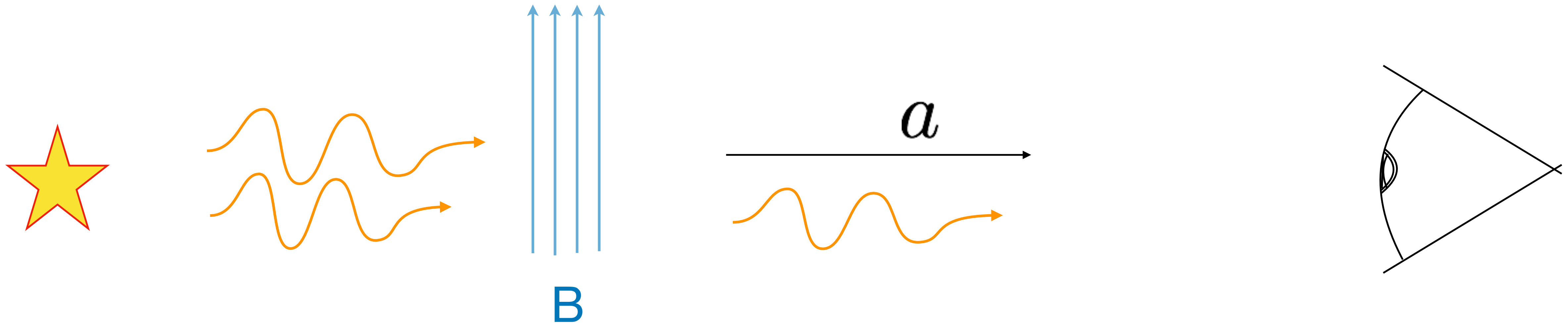
in a magnetised plasma

Schrödinger-like equation
for relativistic ALPs

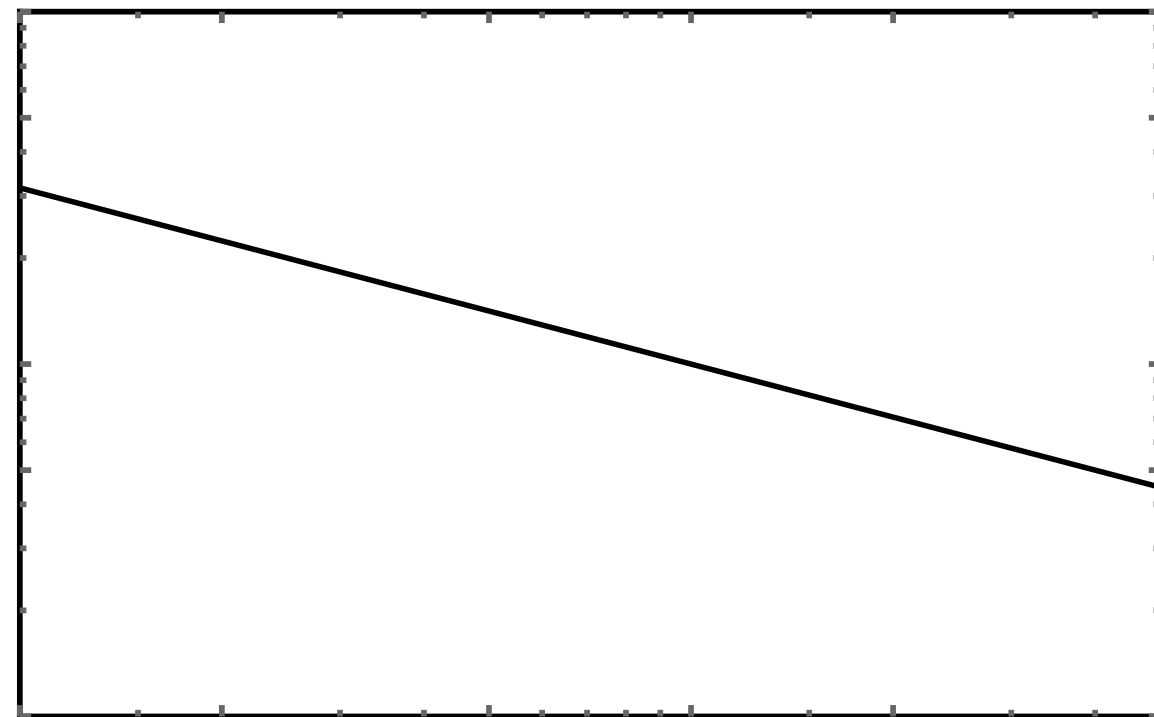
$$i \frac{d}{dz} \Psi(z) = (H_0 + H_I) \Psi(z) ;$$

$$\Psi(z) = \begin{pmatrix} A_x \\ A_y \\ a \end{pmatrix} \quad H_0 = -\frac{1}{2\omega} \begin{pmatrix} \omega_{pl}(z)^2 & 0 & 0 \\ 0 & \omega_{pl}(z)^2 & 0 \\ 0 & 0 & m_a^2 \end{pmatrix} \quad H_I = \frac{ga\gamma}{2} \begin{pmatrix} 0 & 0 & B_x \\ 0 & 0 & B_y \\ B_x & B_y & 0 \end{pmatrix} ;$$

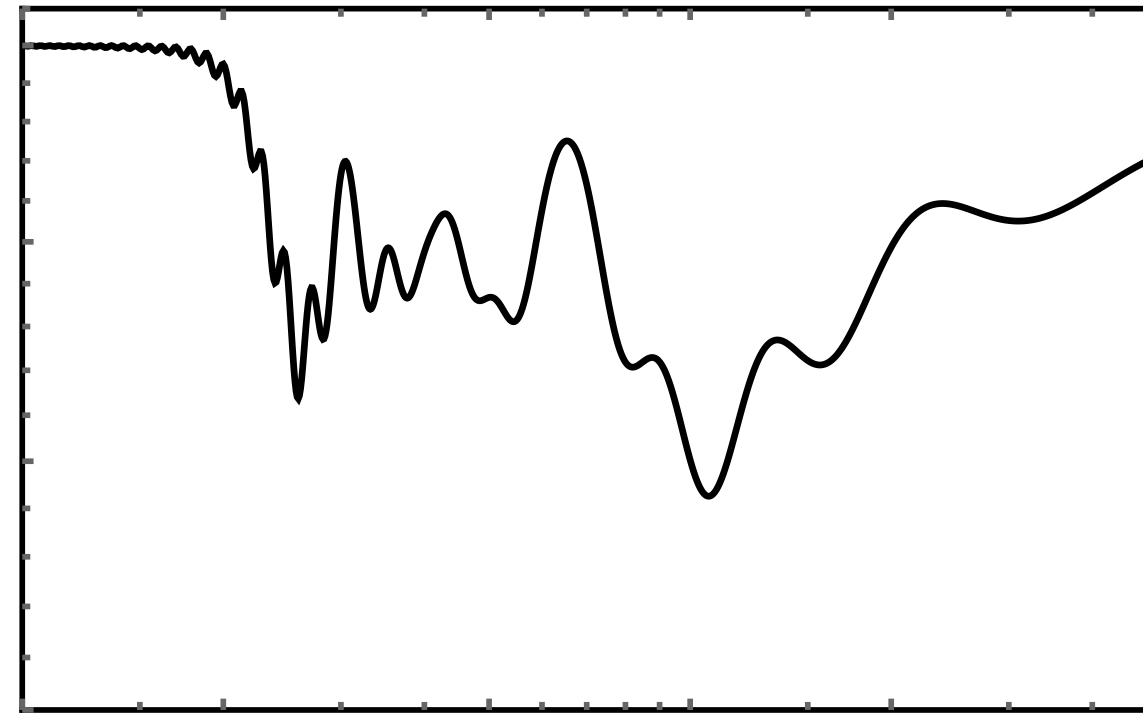
The photon disappearance channel



Initial photon spectrum



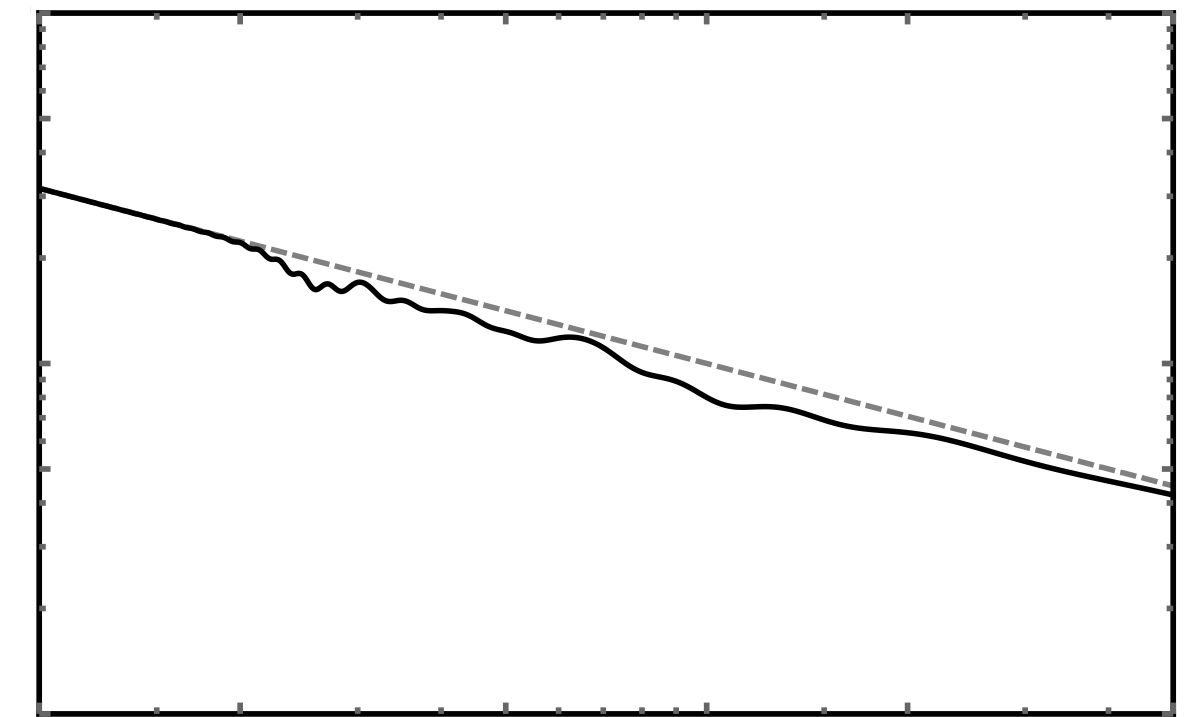
Survival probability



*

=

Final photon spectrum



Energy

Energy

Energy

[Sikivie],
[Raffelt, Stodolsky]

Galaxy clusters are ideal axion-photon converters

Largest gravitational bound objects (~ 100 s kpc).

Magnetised (μG).

Long coherence lengths ($\sim \text{kpc}$).

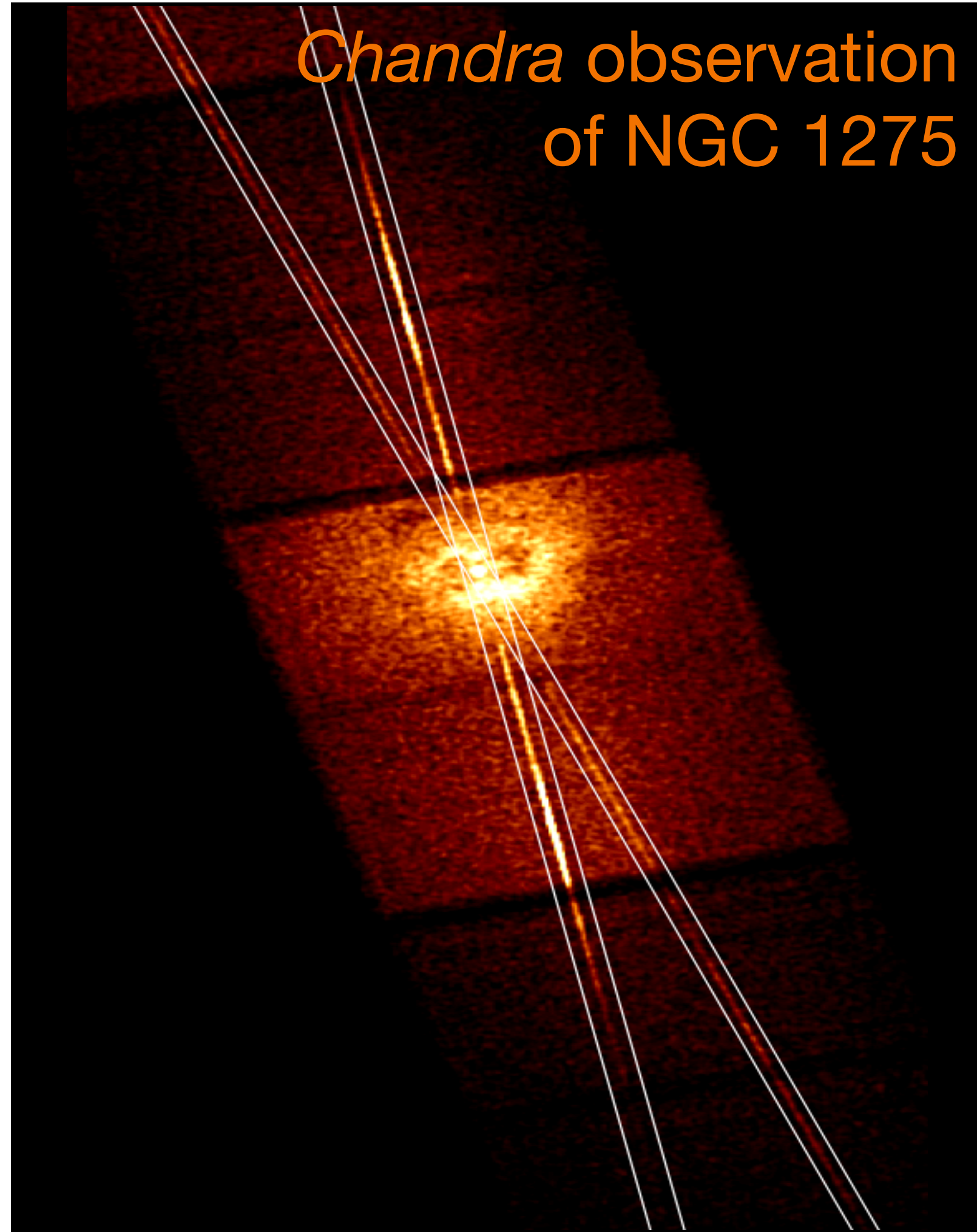
Luminous sources (AGNs, quasars).

Unsuppressed *conversion ratios*:

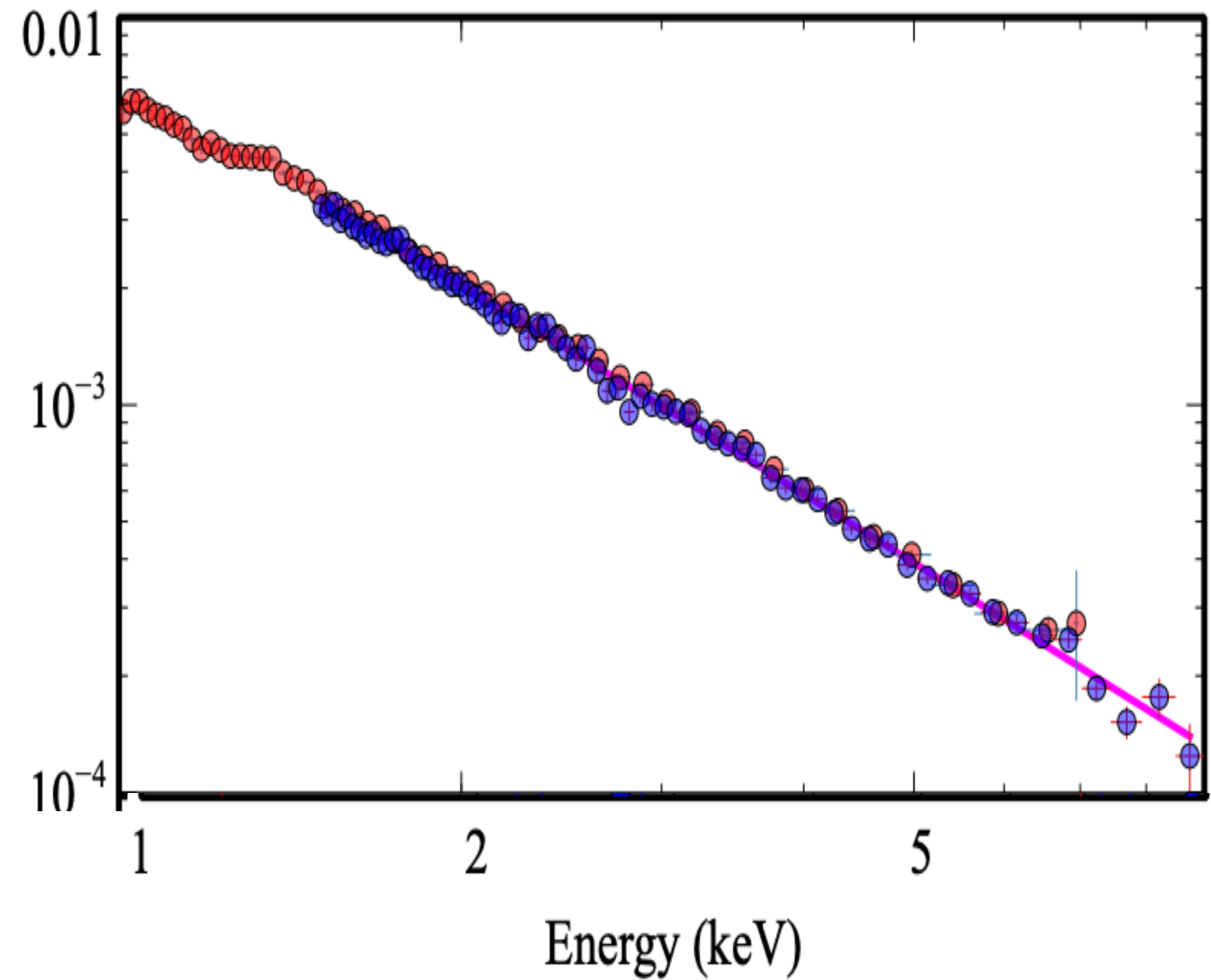
$$P_{\gamma a} \sim \mathcal{O}\left(\frac{1}{2}\right) \times \left(\frac{g_{a\gamma}}{10^{-11} \text{ GeV}}\right)^2$$



Precision spectra

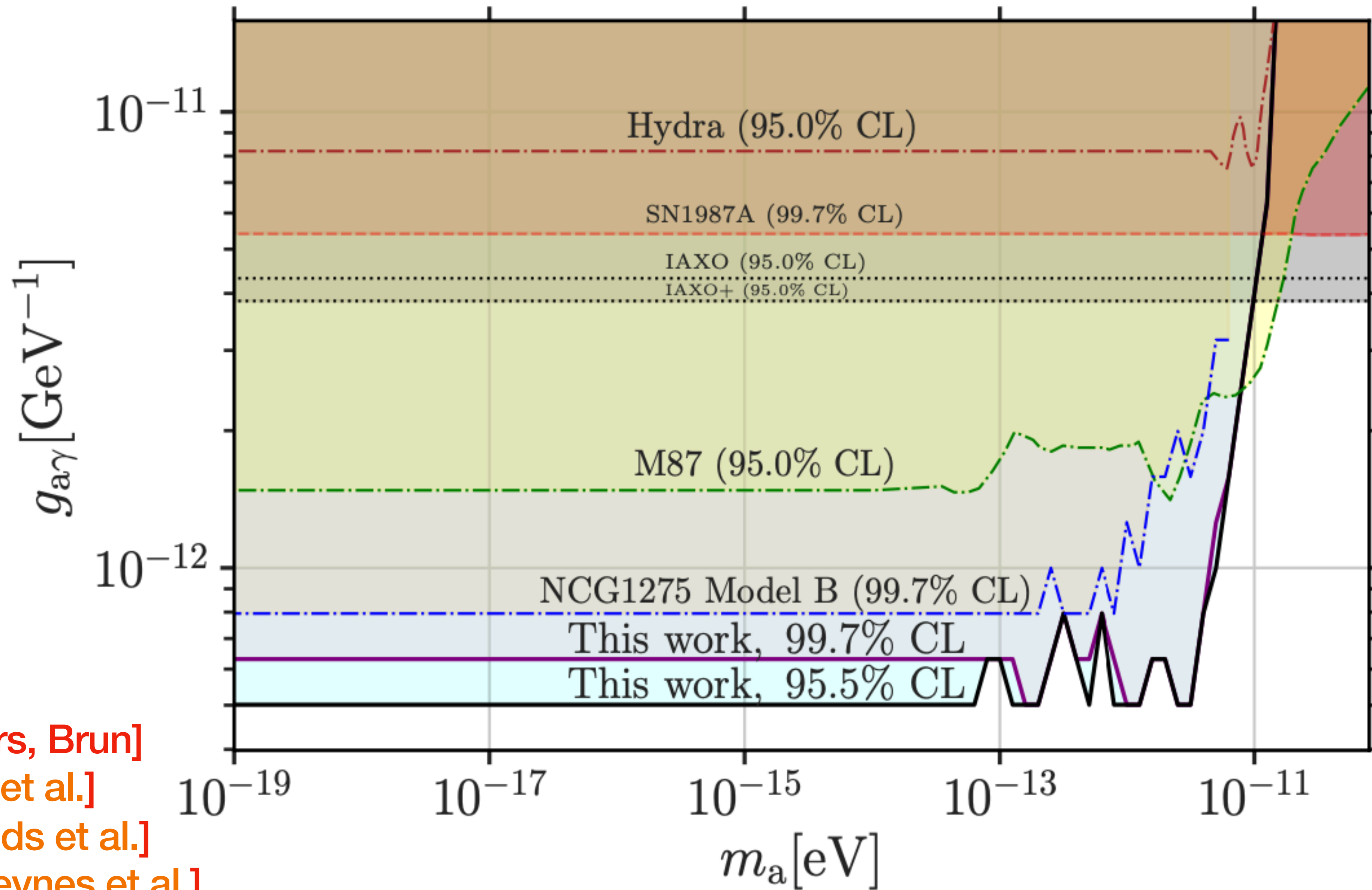


[Reynolds, *DM*, et al.]
[Sisk-Reynes et al.]



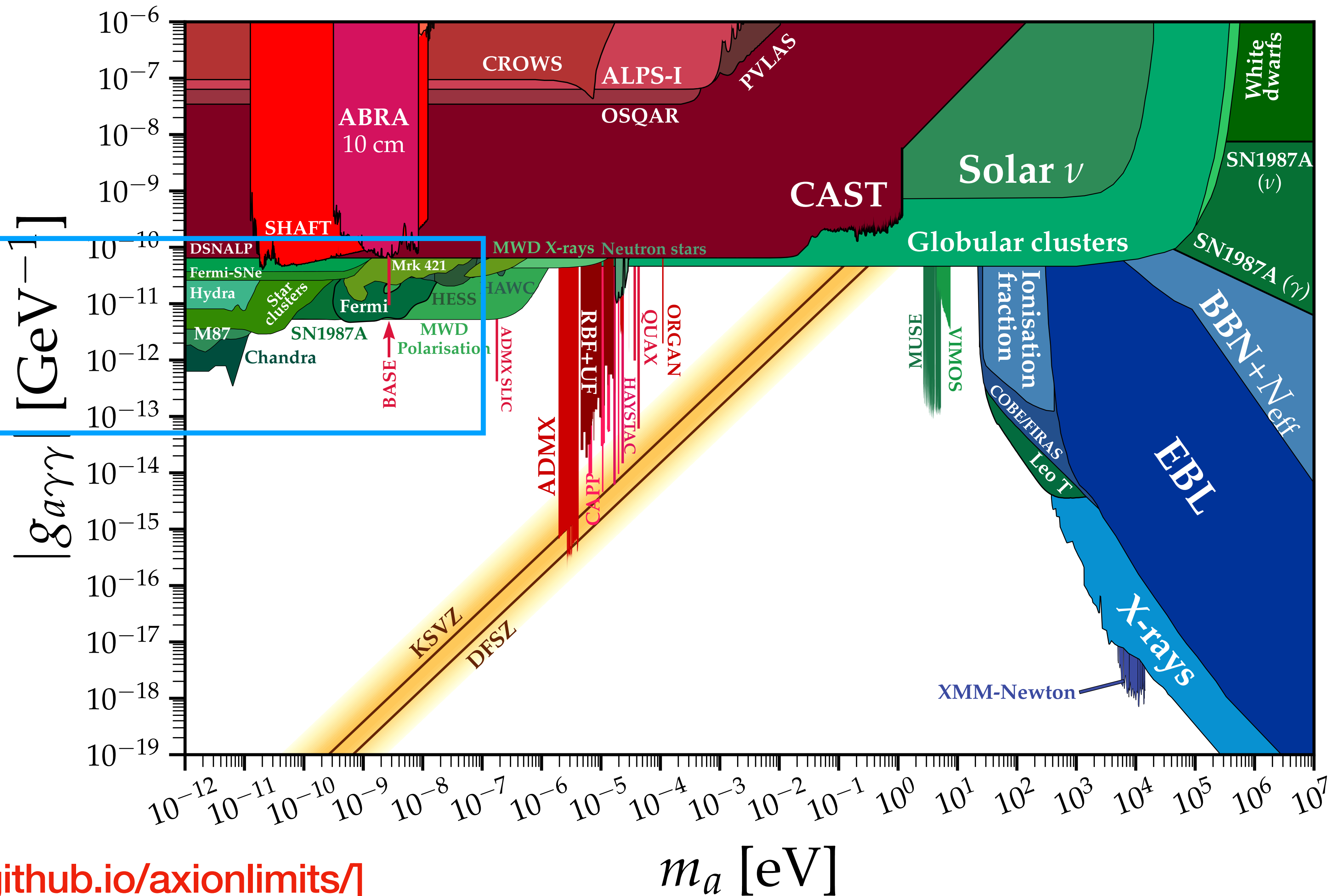
Amplitude of hypothetical
oscillations must be $\lesssim 2.5\%$.

Strongest limits by an order of magnitude



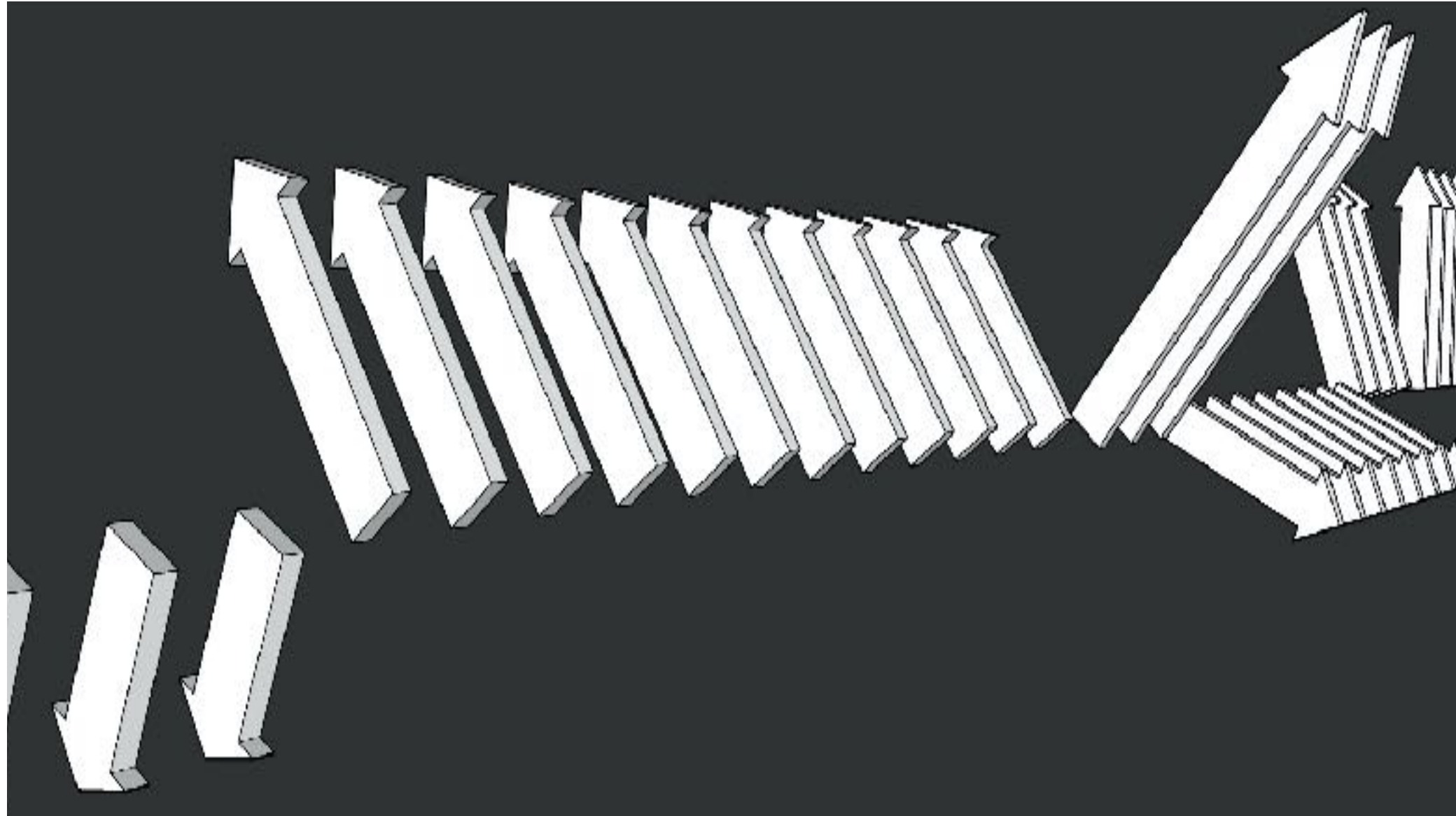
[Wouters, Brun]
[Marsh et al.]
[Reynolds et al.]
[Sisk Reynes et al.]

Context

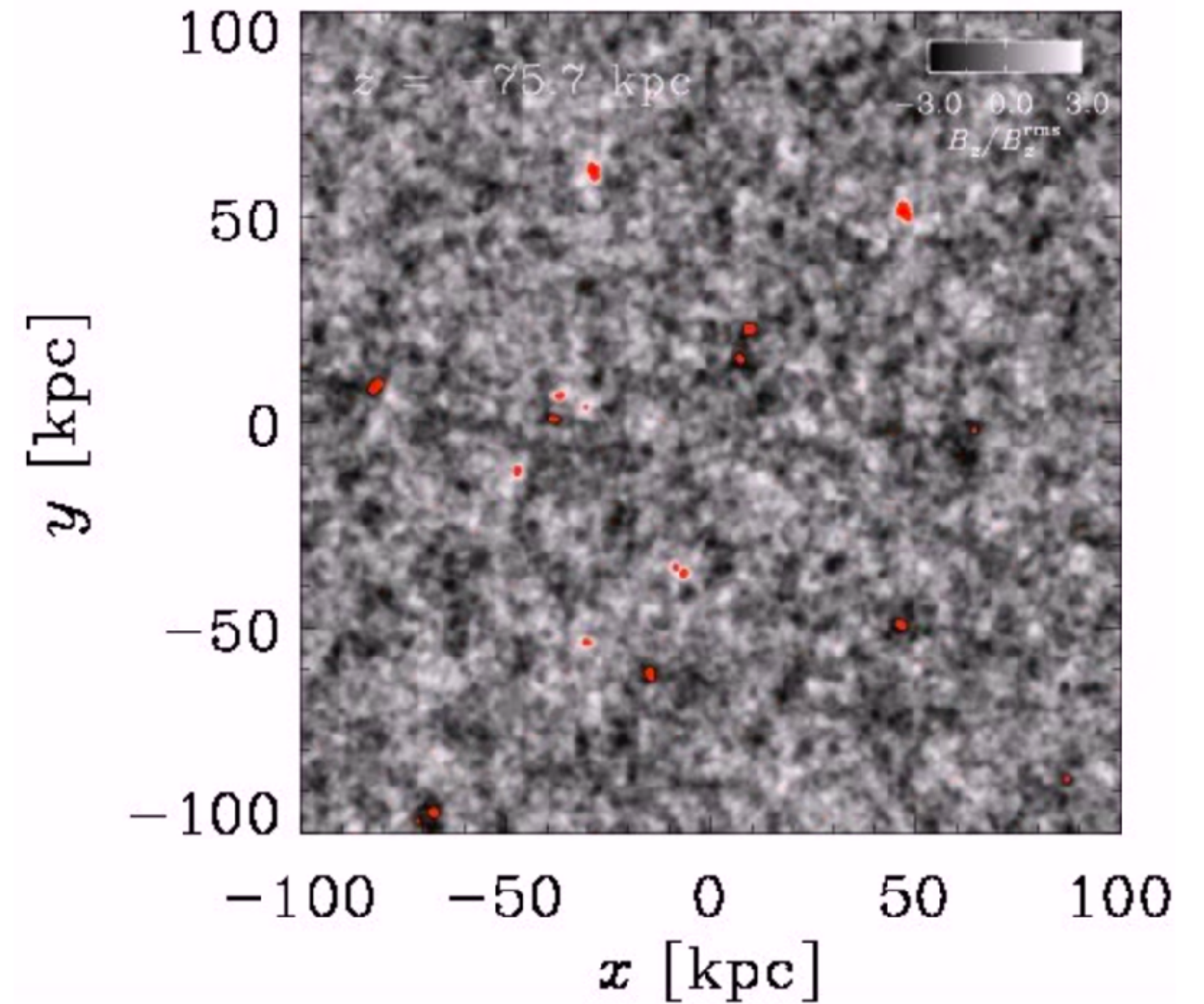


[Wouters, Brun],
 [Conlon et al.],
 [Berg et al.],
 [DM et al.],
 [Reynolds et al.],
 [Chen, Conlon],
 [Day, Krippendorf],
 [Sisk Reynes et al.],
 [Matthews et al.],
 [Schallmoser et al.]

Magnetic field models

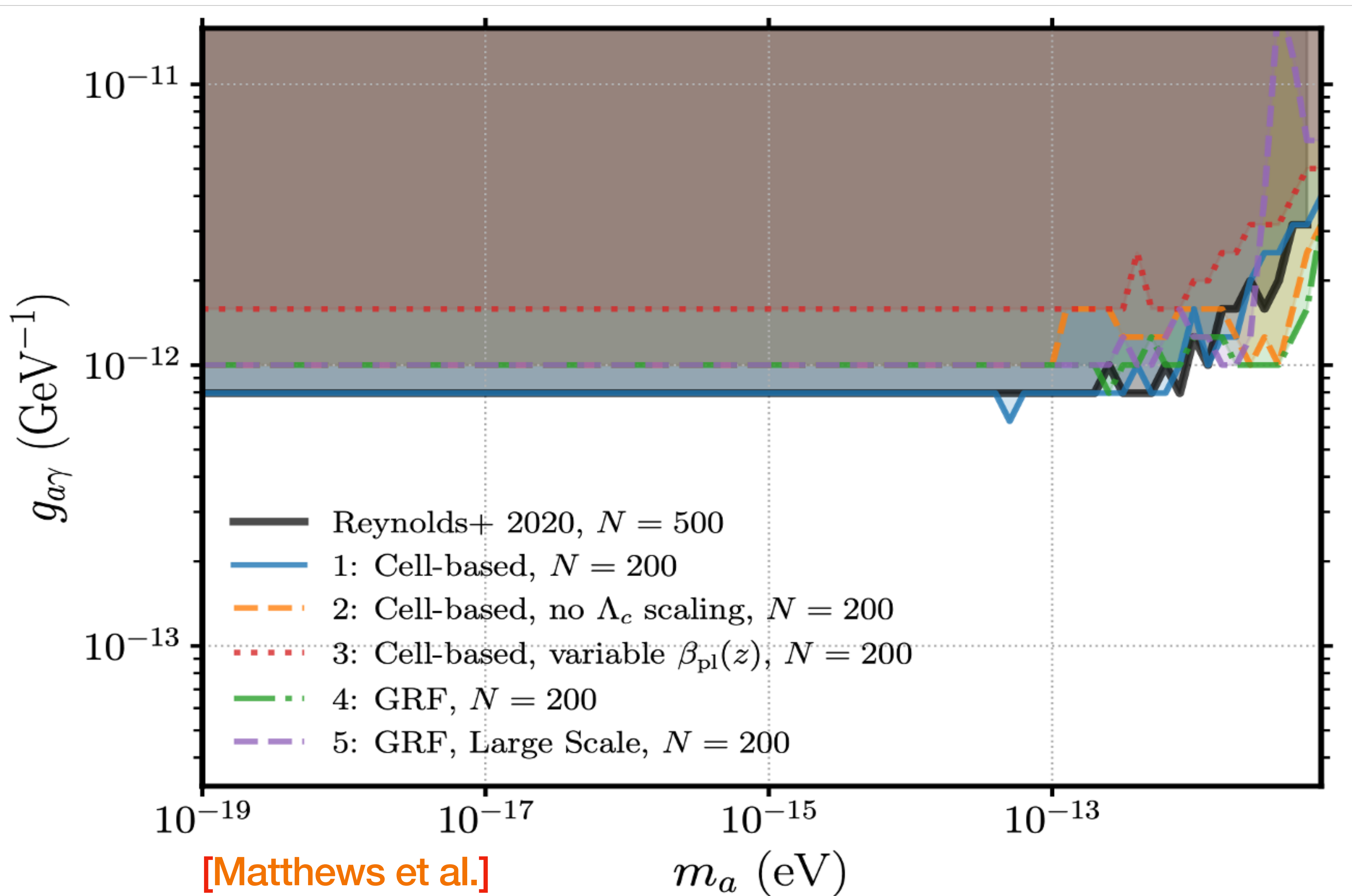


Status: **standard practice**



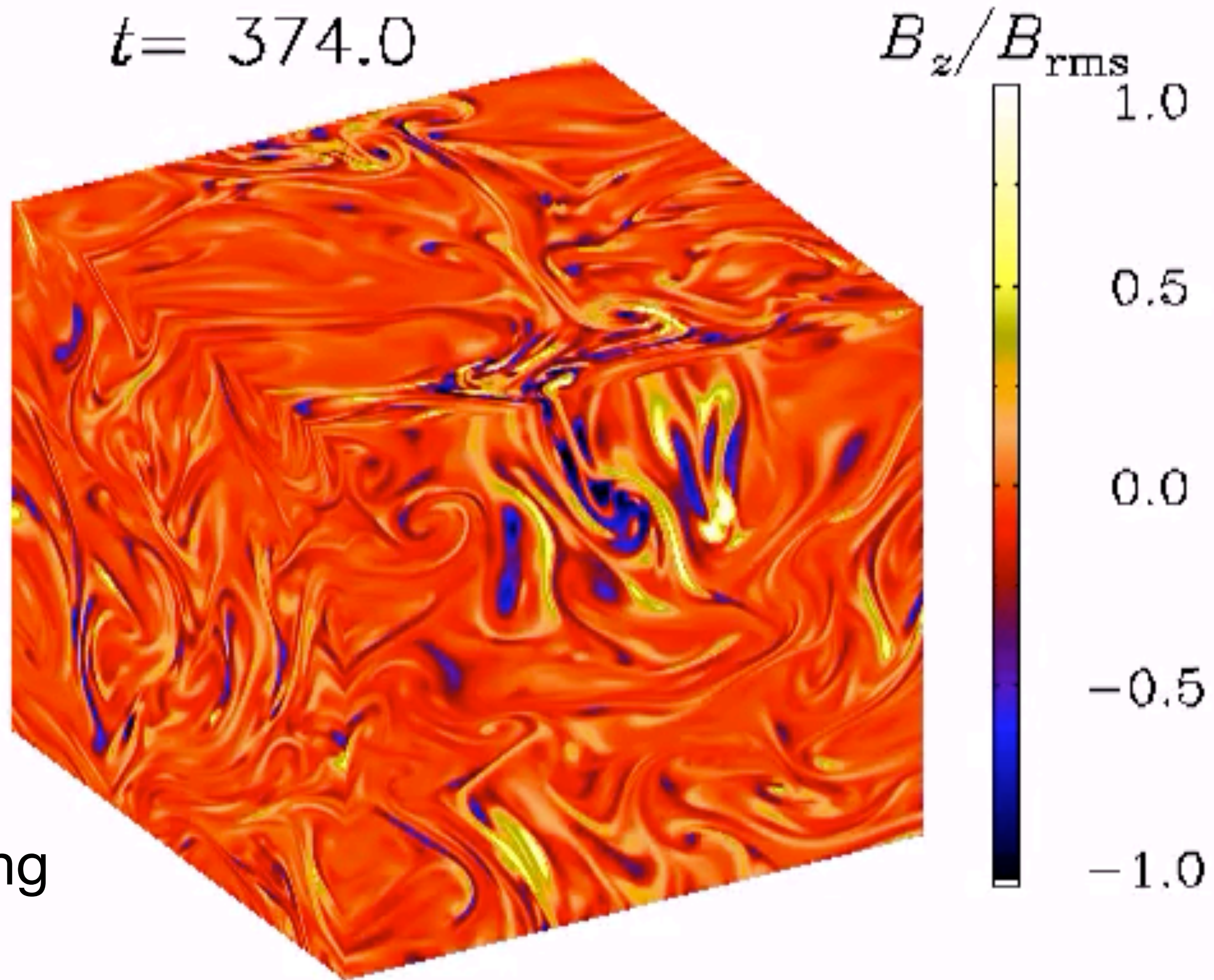
Status: **“state-of-the-art”**

How robust are these limits?



Dedicated MHD simulations: time-evolution

$t = 374.0$



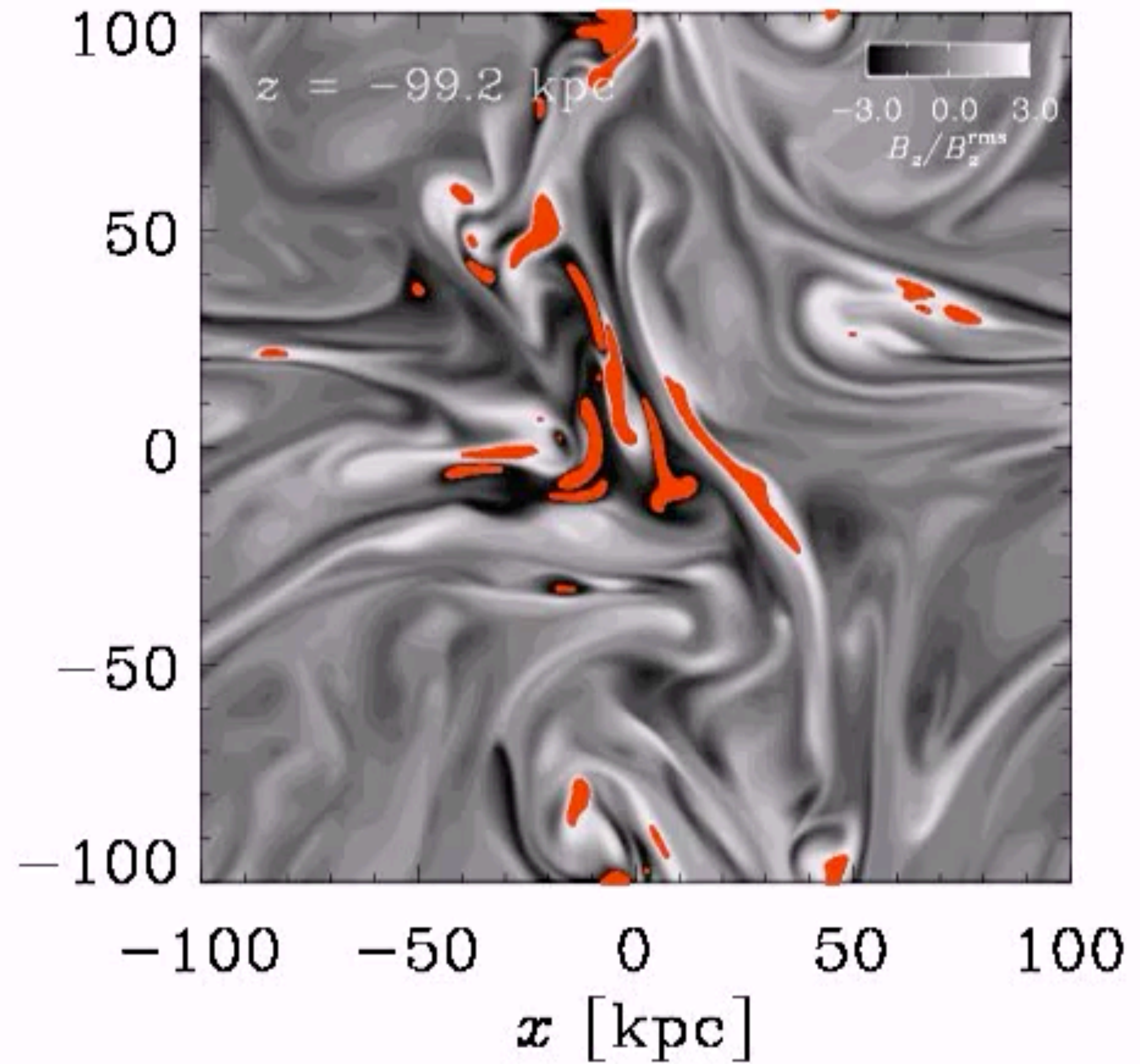
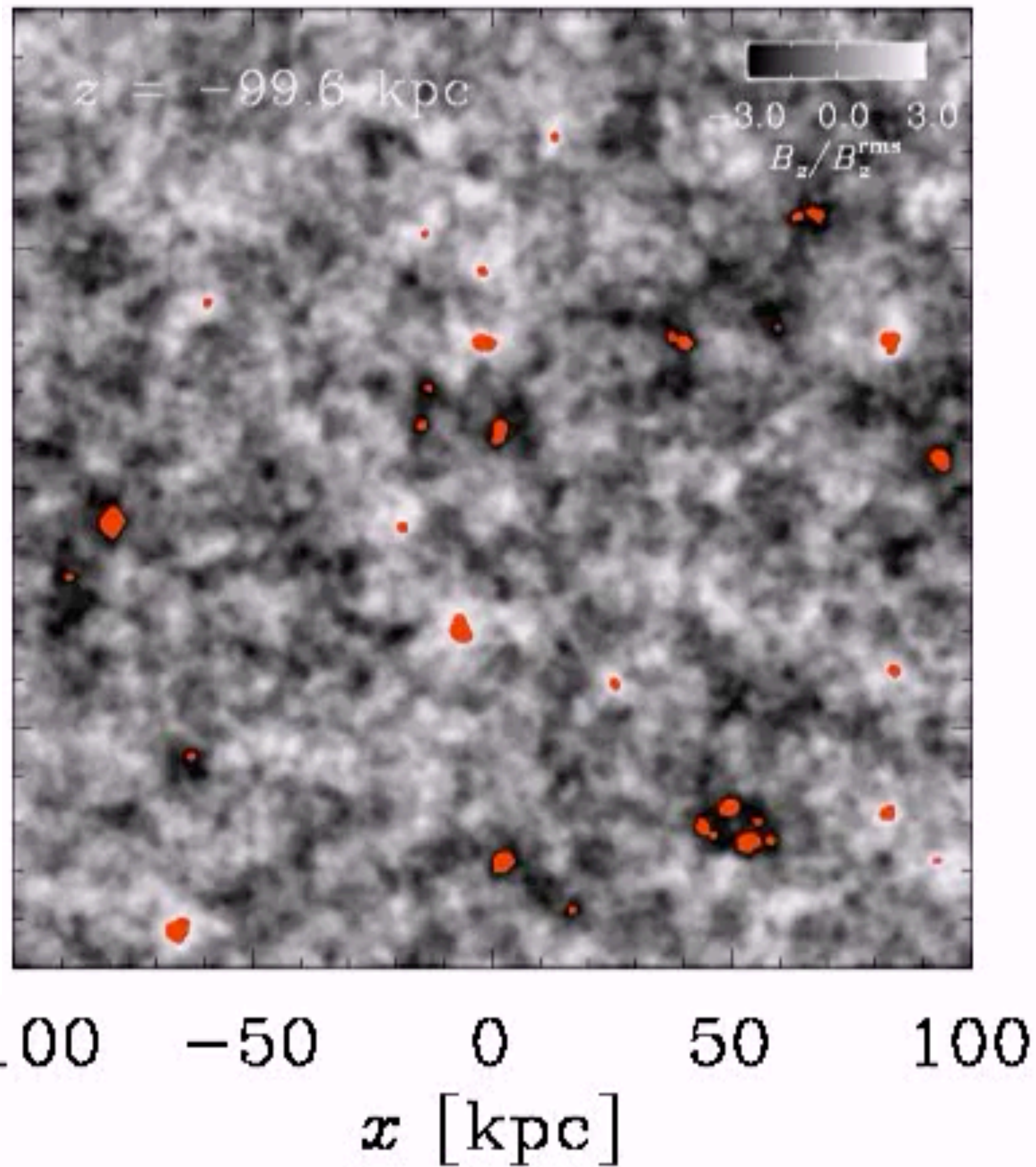
$L^3 = (200 \text{ kpc})^3$
#lattice points = 512^3
periodic bc, external forcing
Dynamo-enhanced,
turbulent magnetic field

[Carenza et al.]

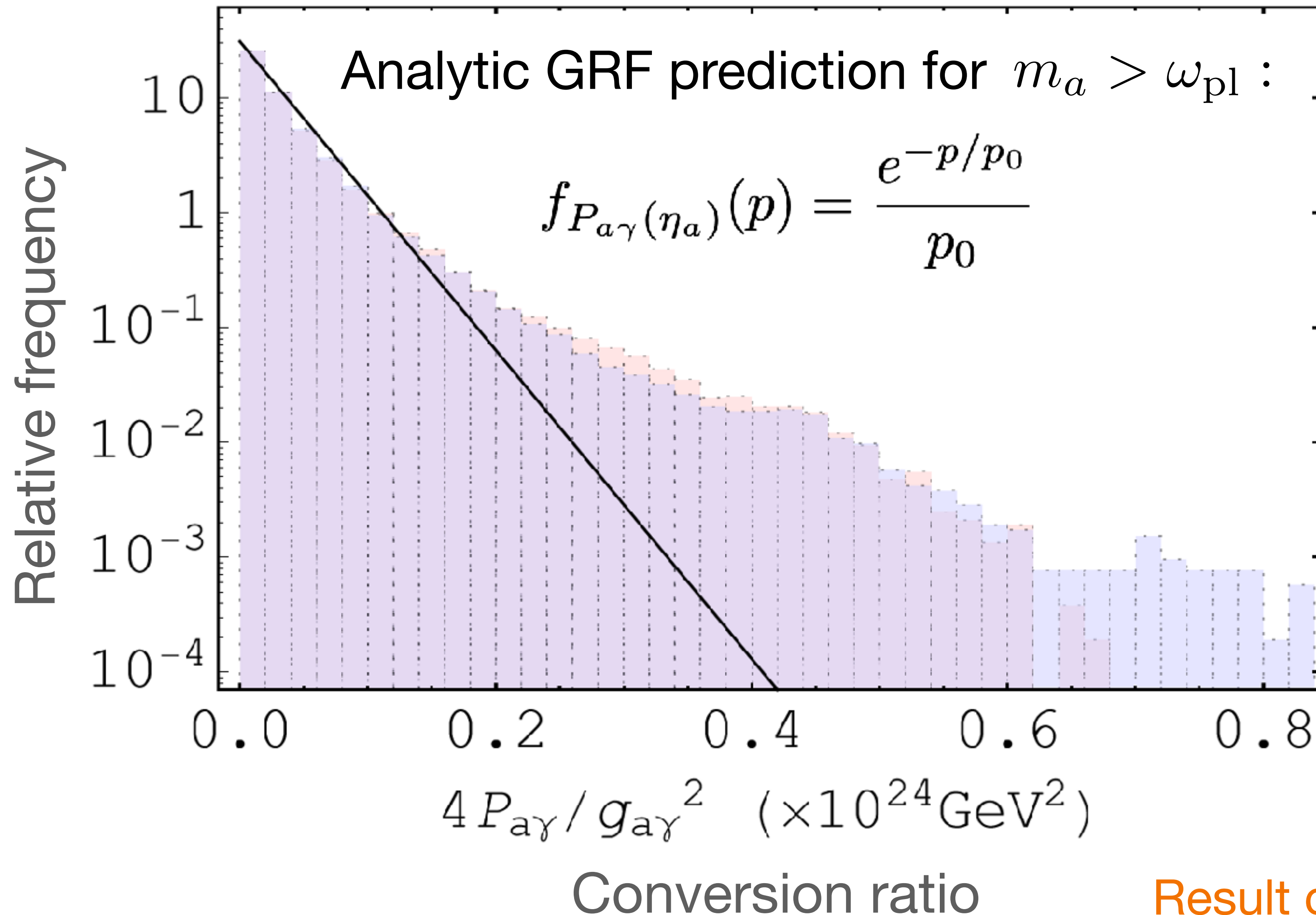
GRF v MHD

(same power spectrum)

Red: $|\mathbf{B}| > 3B_{\text{rms}}$



Heavy-tailed MHD distributions



$$p_0 = \frac{g_{a\gamma}^2}{4} \frac{L}{2\pi} P_{1D}(\eta_a)$$

$$\eta_a = \frac{m_a^2}{2\omega}$$

Skewness & kurtosis:

GRF:

$$S = 2$$

$$K = 9$$

MHD:

$$S = 3.88$$

$$K = 25.56$$

$$S = 4.60$$

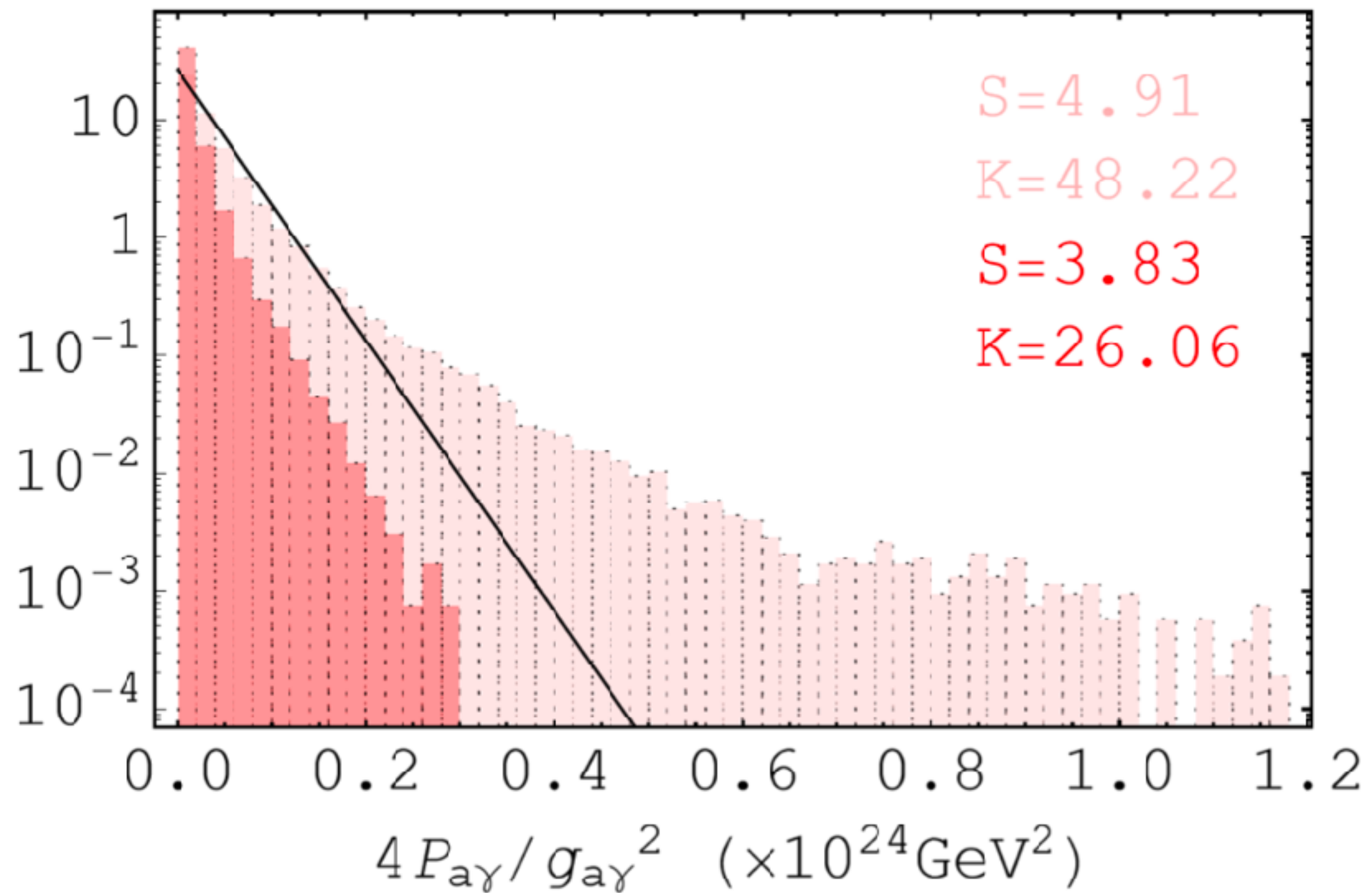
$$K = 41.80$$

Result generalises for arbitrary ALP mass

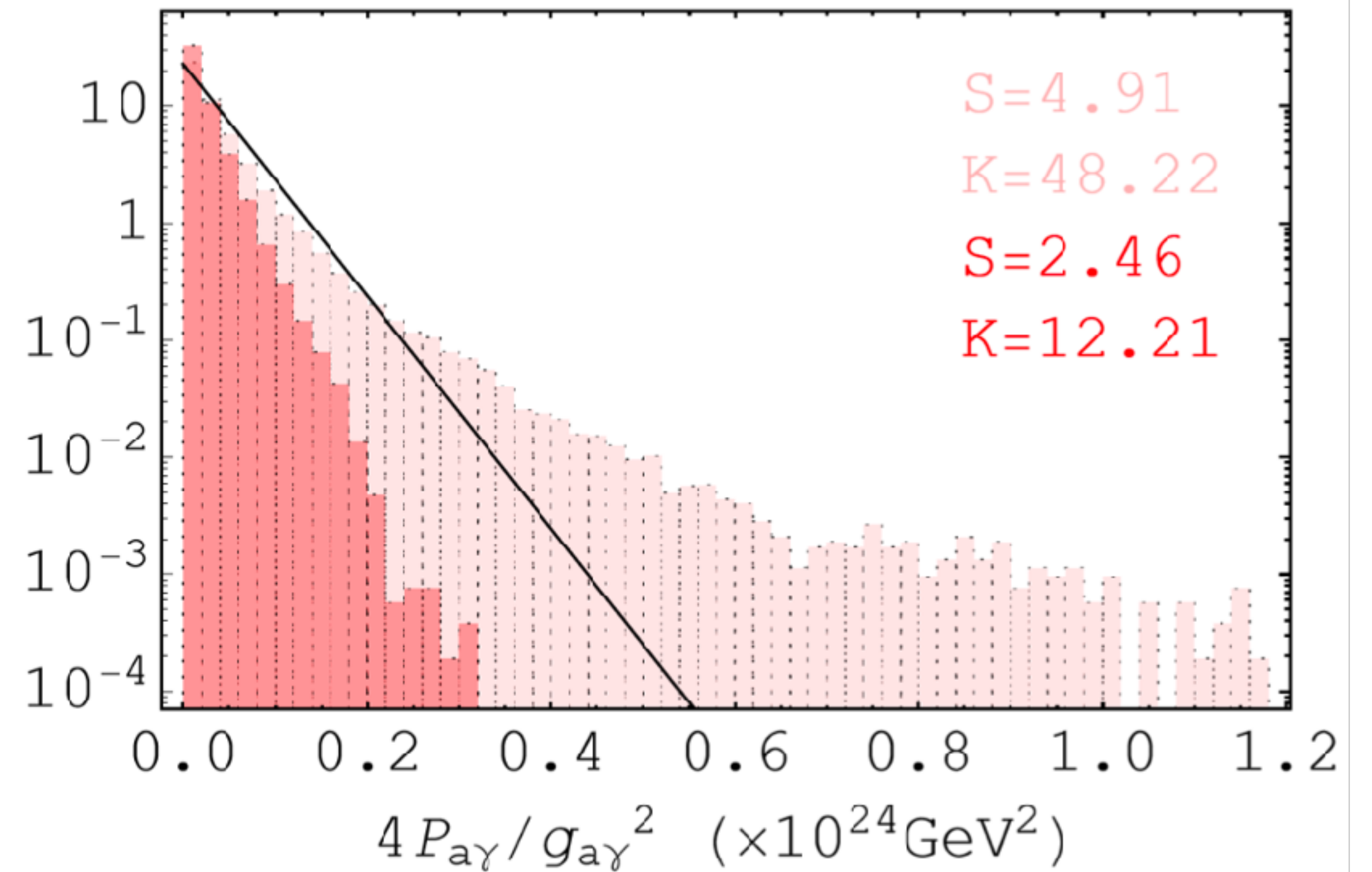
Non-Gaussianity

Two possible sources:

Mask large coherence lengths



Mask high peaks



Non-Gaussianity

Typical predictions essentially set by average:

$$\langle P_{\gamma a}(\eta_a) \rangle = \frac{g_{a\gamma}^2}{4} \langle |\tilde{B}_i(\eta_a)|^2 \rangle = \frac{g_{a\gamma}^2}{4} \frac{L}{2\pi} P_{1D}(\eta_a)$$

Same for MHD and GRF

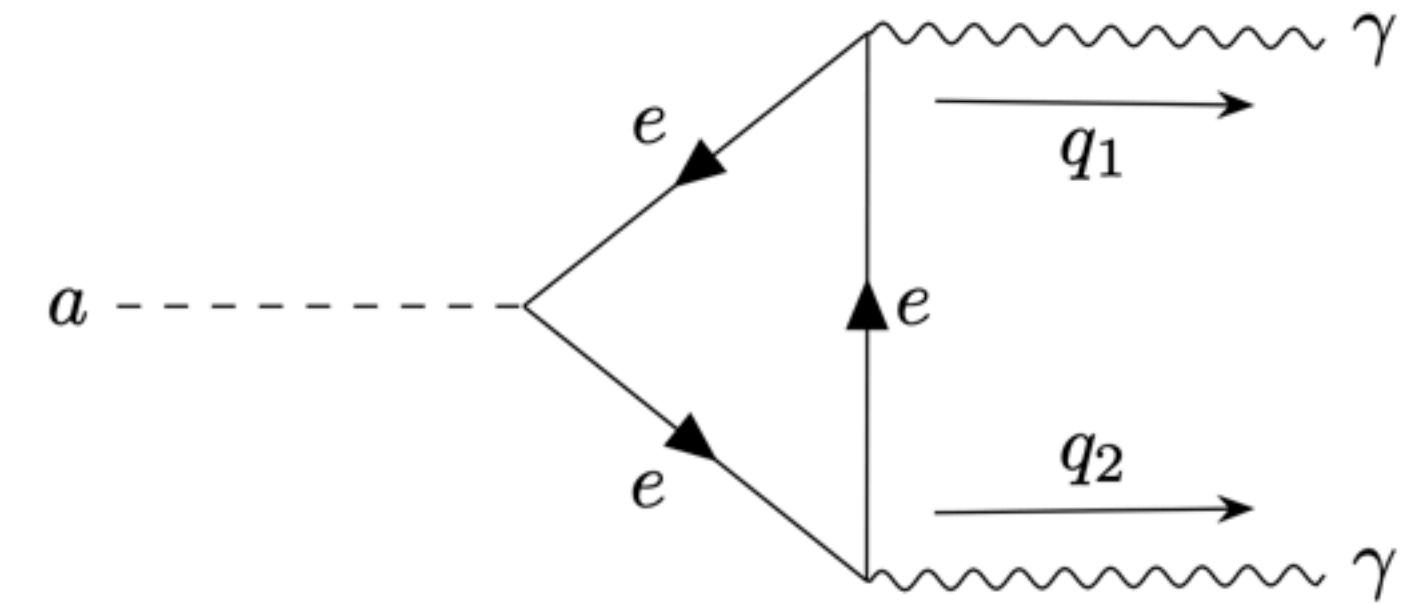
Heavy tails come from larger-than-Gaussian higher-order correlations, i.e.

$$\langle P_{\gamma a}(\eta_a)^2 \rangle, \quad \langle P_{\gamma a}(\eta_a)^3 \rangle, \quad \langle P_{\gamma a}(\eta_a)^4 \rangle \quad \text{etc.}$$

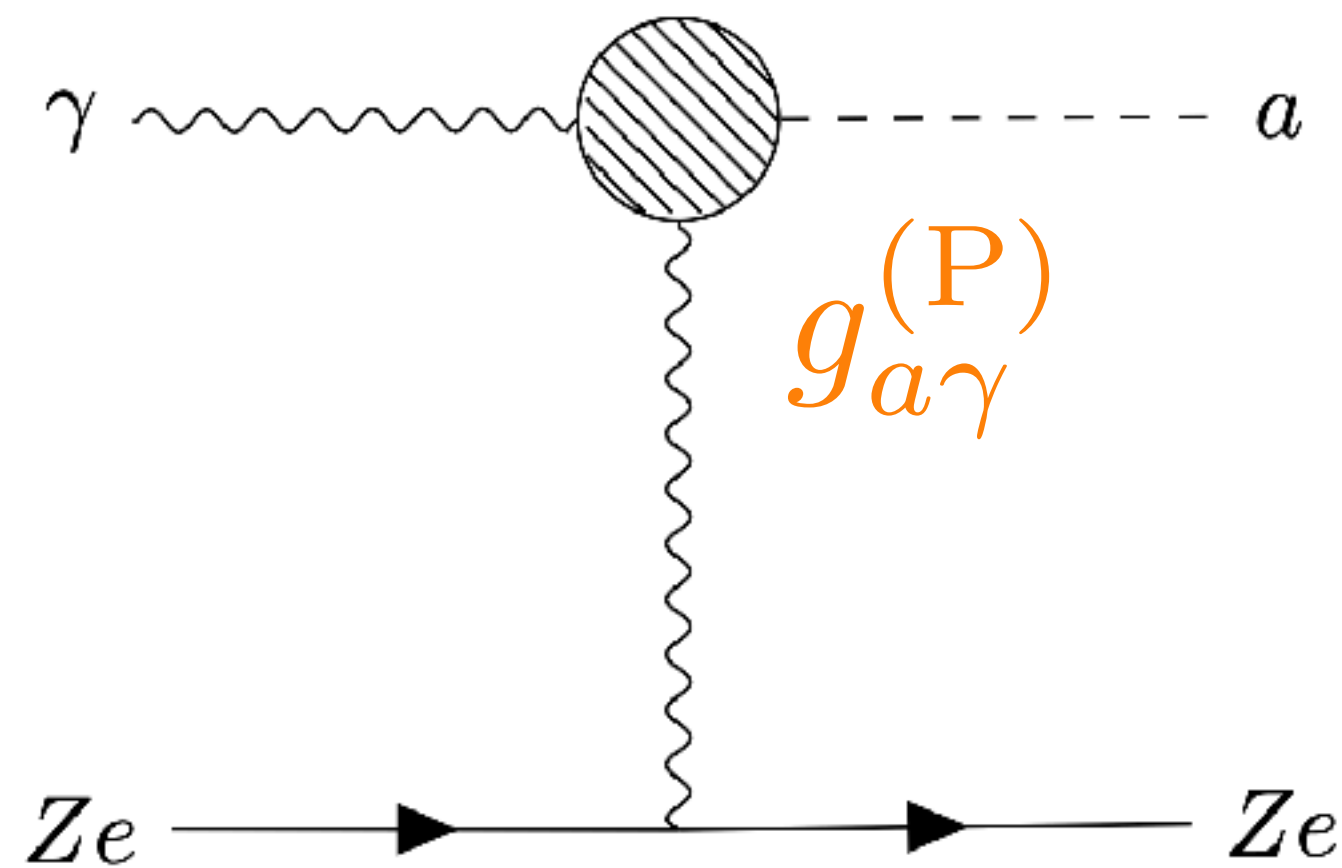
*Larger conversion from MHD
— suggest existing limits conservative*

Quantum ALP-processes

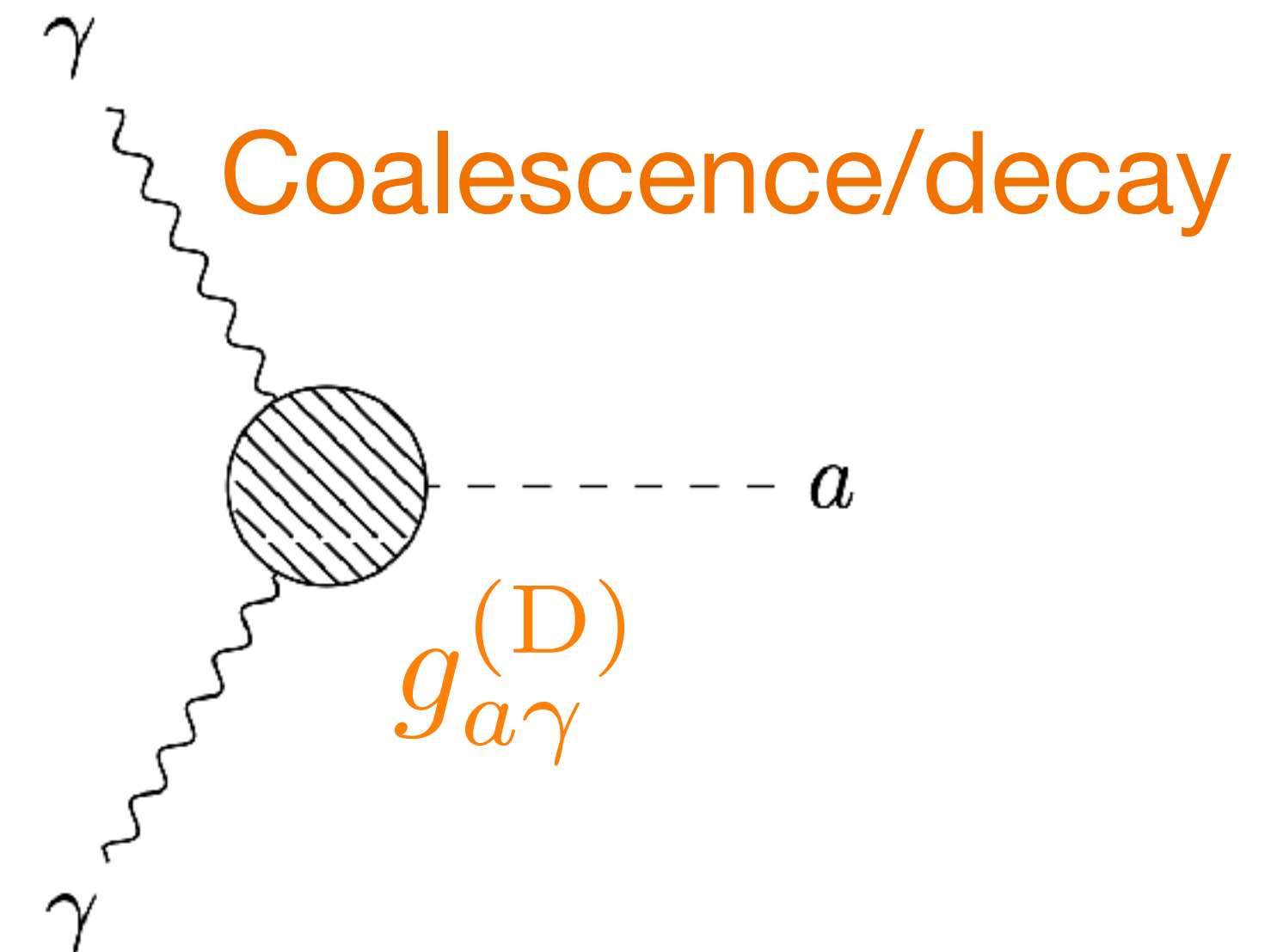
ALPs that don't couple to photons at tree-level still acquire an effective coupling from loops:



Primakoff/inverse Primakoff



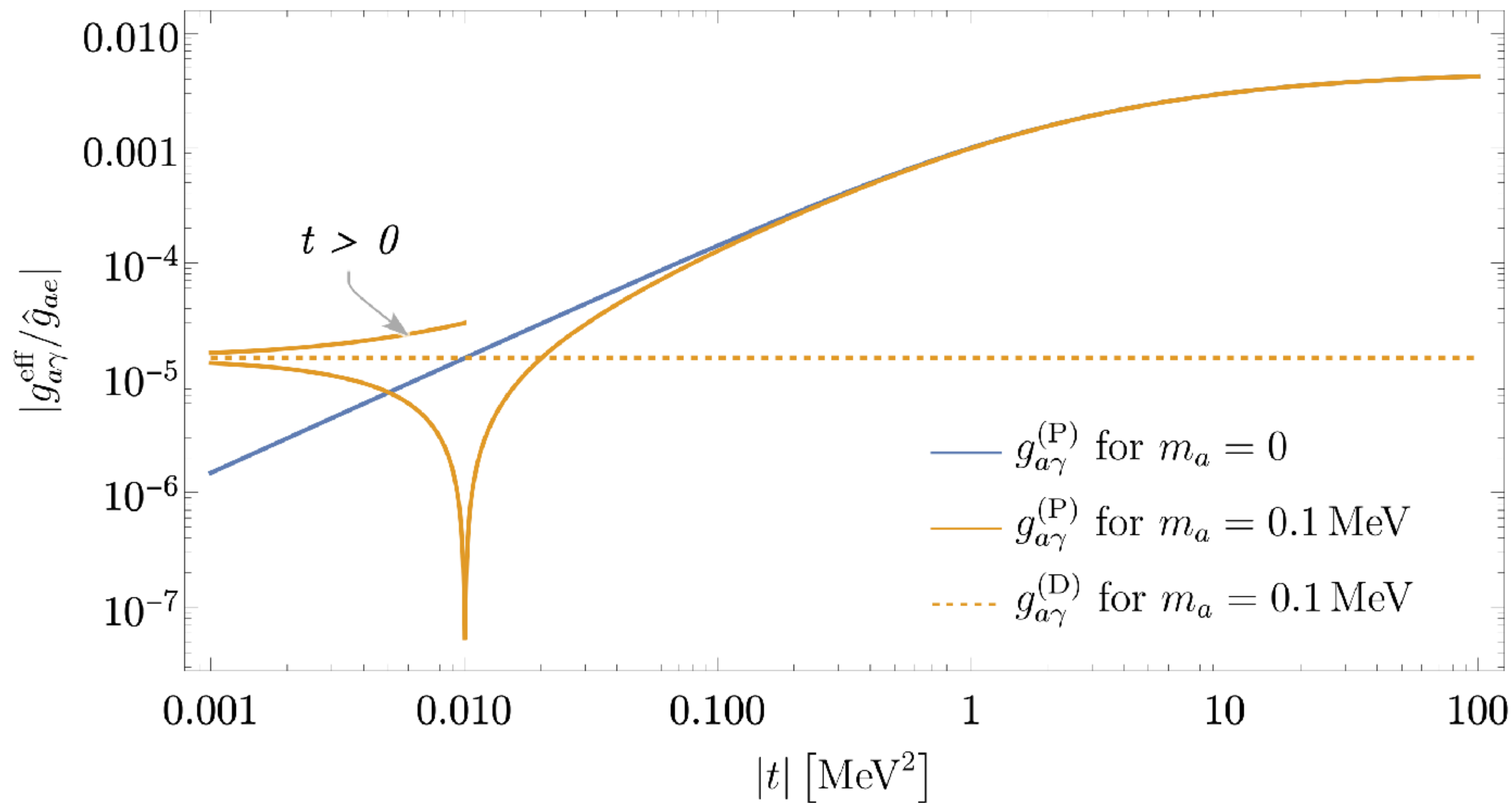
$$g_{a\gamma}^{(P)}$$



Coalescence/decay

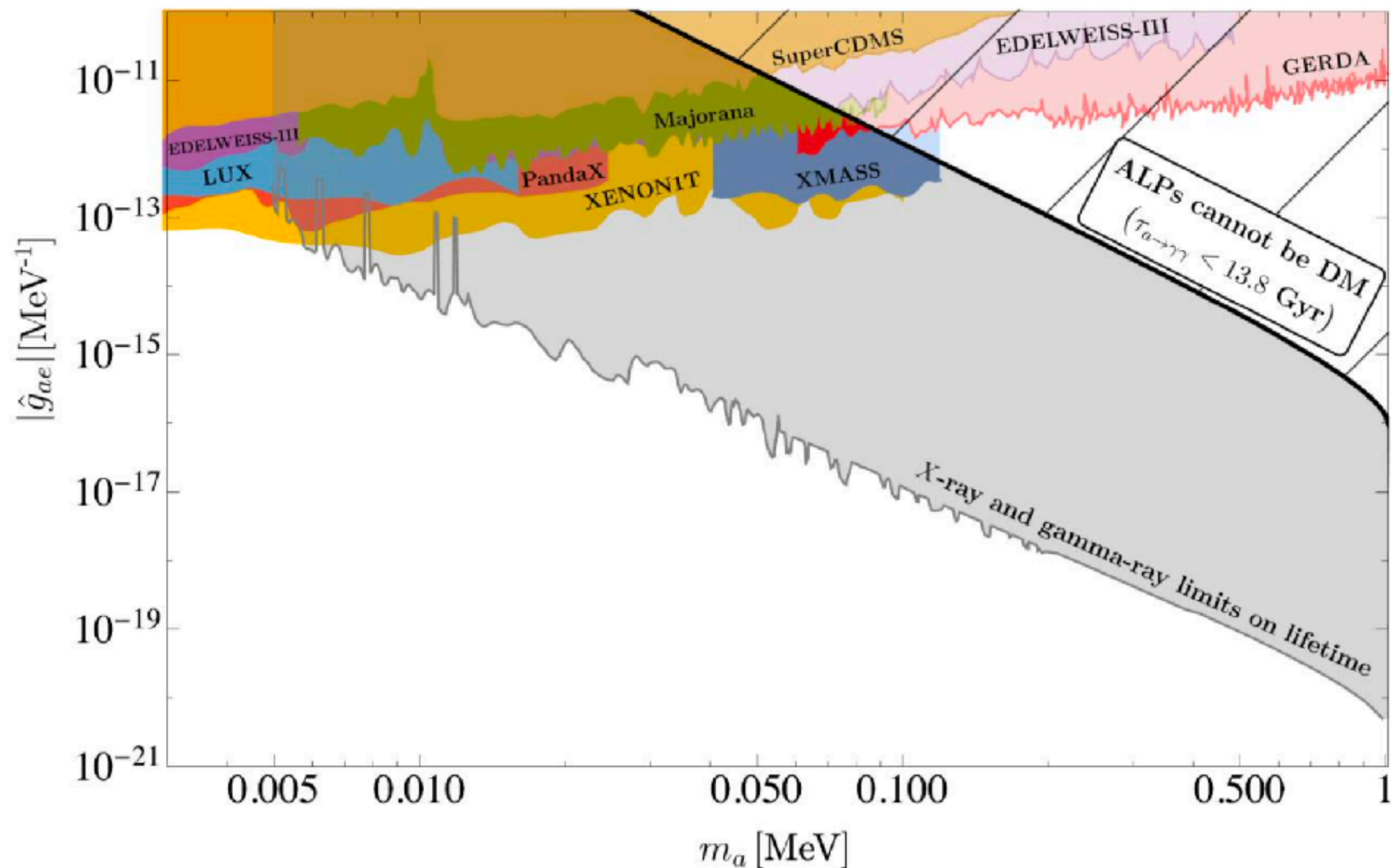
$$g_{a\gamma}^{(D)}$$

Momentum-dependent coupling



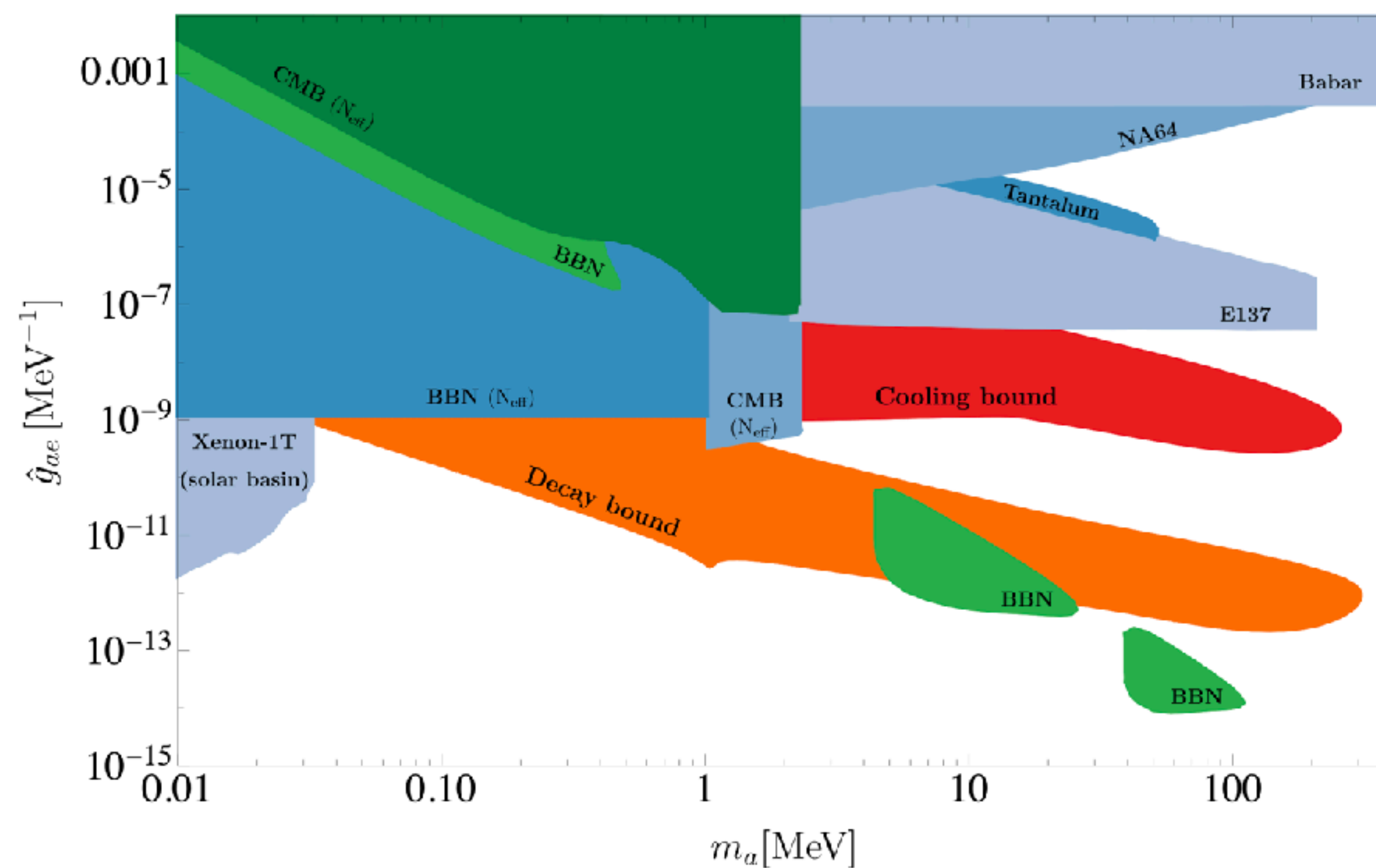
Two applications

ALP dark matter decay



[Ferreira, Marsh, Müller]
Phys.Rev.Lett. 128 (2022) 22

Decay and cooling bound from SN1987A



[Ferreira, Marsh, Müller]
arXiv:2205.07896

Conclusions

Astrophysical probes can be **very sensitive** to ALPs.

MHD models will be the **next state-of-the-art** for ALP-photon conversion.

MHD structure suggests **new observables**.

Quantum processes can dominate — and have drastic consequences.