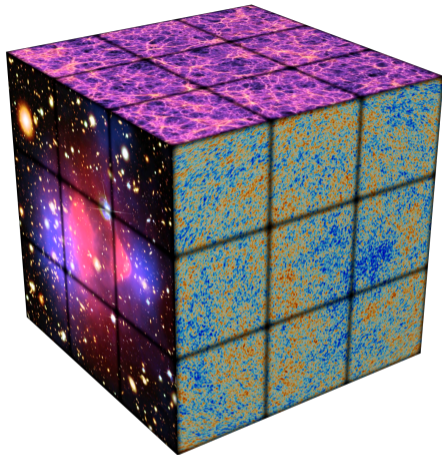


26/08/22

COSMOS 22



Dark Matter from Preheating

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with

Mathias Pierre (DESY)

Sarunas Verner (Florida)

2206.08940



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1. Preheating



2. Weak coupling



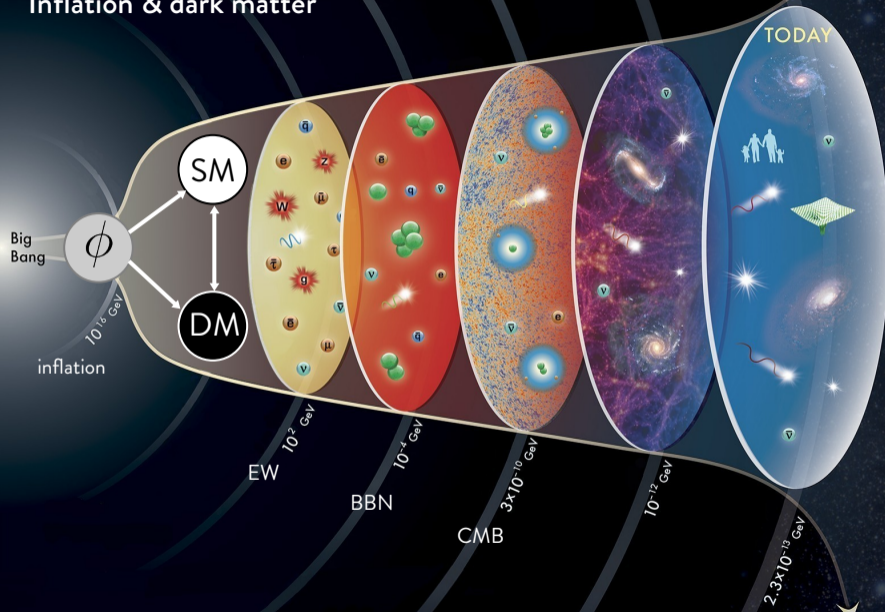
3. Strong coupling



4. Constraints



Inflation & dark matter



1. Preheating



2. Weak coupling



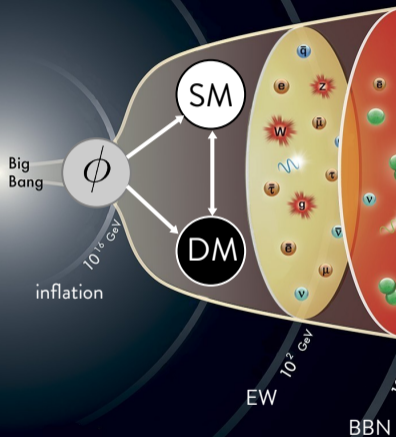
3. Strong coupling



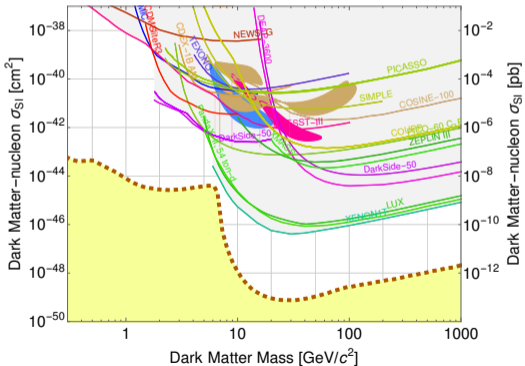
4. Constraints



Inflation & dark matter



No WIMP?



1. Preheating



2. Weak coupling



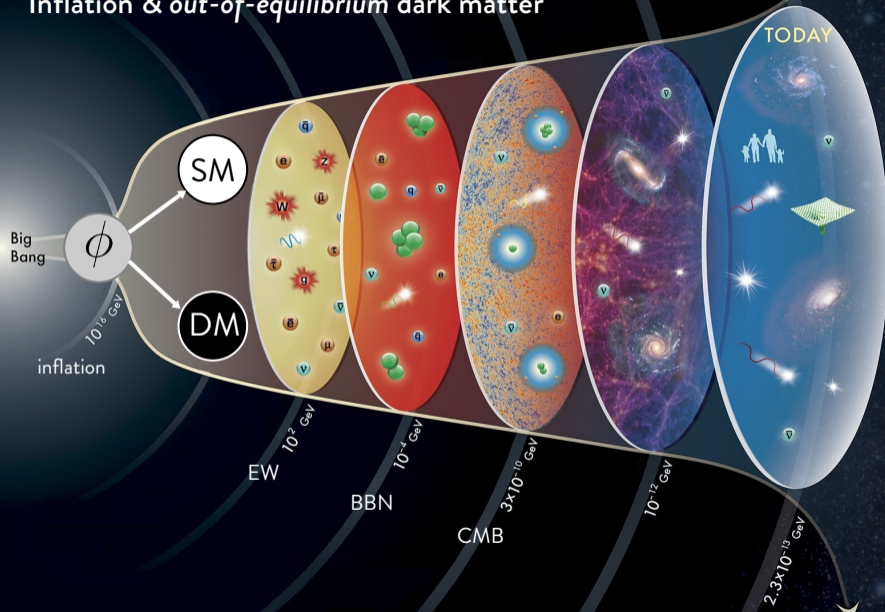
3. Strong coupling



4. Constraints



Inflation & out-of-equilibrium dark matter



1. Preheating



2. Weak coupling



3. Strong coupling



4. Constraints



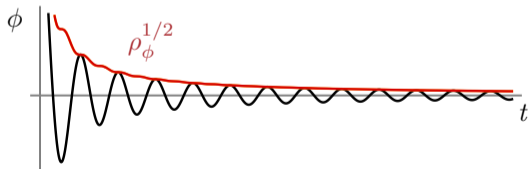
The inflaton and its decay products

$$\mathcal{S} = \int d^4x \sqrt{-g} \left[\frac{1}{2} (\partial_\mu \phi)^2 - 6\lambda M_P^4 \tanh^2 \left(\frac{\phi}{\sqrt{6} M_P} \right) \right. \\ \left. + \frac{1}{2} (\partial_\mu \chi)^2 - \frac{1}{2} (m_\chi^2 + \sigma \phi^2) \chi^2 \right. \\ \left. + \bar{\psi} i \bar{\gamma}^\mu \nabla_\mu \psi - y \phi \bar{\psi} \psi + \dots \right]$$

inflaton

dark matter

radiation



$$\langle p_\phi \rangle = \frac{1}{2} \langle \dot{\phi}^2 + m_\phi^2 \phi^2 \rangle \simeq 0$$

(matter)

1. Preheating



2. Weak coupling



3. Strong coupling



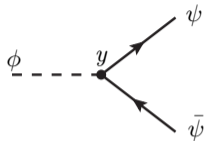
4. Constraints



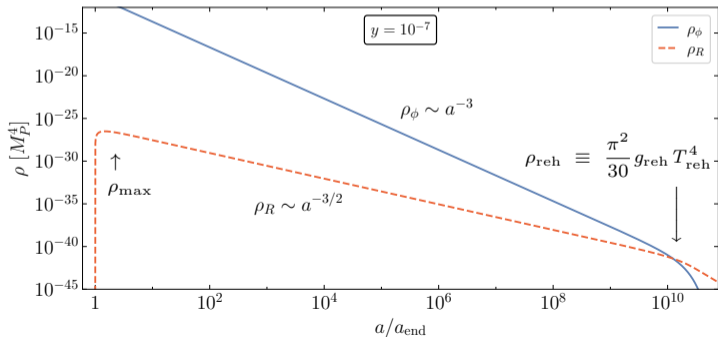
Perturbative decay of the inflaton

Decay into visible sector is assumed to be perturbative

$$\mathcal{L} \supset -y\phi\bar{\psi}\psi$$



$$\Gamma_\phi = \frac{y^2}{8\pi} m_\phi$$



$$\dot{\rho}_\phi + 3H\rho_\phi = -\Gamma_\phi\rho_\phi$$

$$\dot{\rho}_R + 4H\rho_R = \Gamma_\phi\rho_\phi$$

$$\rho_\phi + \rho_R = 3H^2 M_P^2$$

1. Preheating



2. Weak coupling



3. Strong coupling



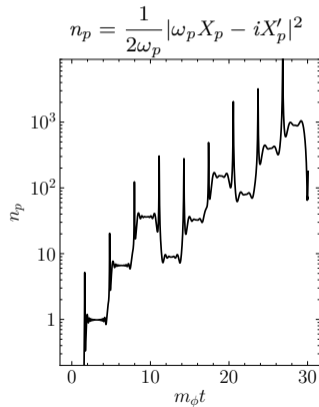
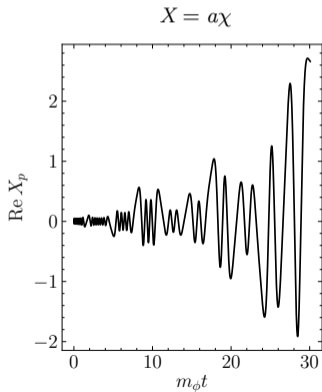
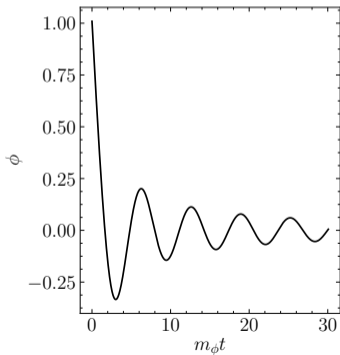
4. Constraints



Scalar preheating

$$\left(\frac{d^2}{dt^2} + 3H \frac{d}{dt} + \frac{p^2}{a^2} + m_\chi^2 + \sigma \phi^2 \right) \chi_p = 0$$

Resonant growth of fluctuations



1. Preheating



2. Weak coupling



3. Strong coupling



4. Constraints



The Boltzmann approximation

Dissipation of fluctuations of the inflaton condensate into χ quanta

$$\phi(t) \simeq \phi_0(t)\mathcal{P}(t) = \phi_0(t) \sum_n \mathcal{P}_n e^{-in\omega t}$$

$$\begin{aligned} \frac{\partial f_\chi}{\partial t} - H|\mathbf{P}| \frac{\partial f_\chi}{\partial |\mathbf{P}|} &= \frac{1}{P^0} \sum_{n=1}^{\infty} \int \frac{d^3\mathbf{K}}{(2\pi)^3 n_\phi} \frac{d^3\mathbf{P}'}{(2\pi)^3 2P'^0} (2\pi)^4 \delta^{(4)}(K_n - P - P') |\overline{\mathcal{M}}_n|^2 \\ &\quad \times \left[f_\phi(K)(1 + f_\chi(P))(1 + f_\chi(P')) - f_\chi(P)f_\chi(P')(1 + f_\phi(K)) \right] \\ &= \frac{\pi |\overline{\mathcal{M}}_2|^2}{2m_\phi^2 \beta(t)} \delta(|\mathbf{P}| - m_\phi \beta(t)) (1 + 2f_\chi(|\mathbf{P}|)) \end{aligned}$$

Condensate PSD

$$f_\phi(P, t) = (2\pi)^3 n_\phi(t) \delta^{(3)}(\mathbf{P})$$

Effective mass

$$m_{\text{eff}}^2 = m_\chi^2 + \sigma\phi^2$$

Mode momenta

$$K_n = (E_n, \mathbf{0}) = (nm_\phi, \mathbf{0})$$

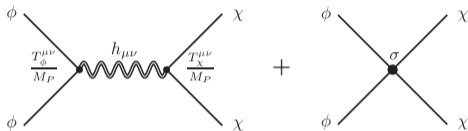
Kinematic factor

$$\beta(t) \equiv \sqrt{1 - \frac{m_{\text{eff}}^2(t)}{m_\phi^2}}$$

1. Preheating



2. Weak coupling



$$\mathcal{L}_I = -\frac{1}{M_P} h_{\mu\nu} (T_\phi^{\mu\nu} + T_\chi^{\mu\nu}) - \frac{\sigma}{2} \phi^2 \chi^2$$

$$\mathcal{M} = -\frac{1}{M_P^2} \left[1 + \frac{2m_{\text{eff}}^2}{s} \right] V(\phi) + \sigma \phi^2$$

3. Strong coupling



4. Constraints



Here $m_{\text{eff}}^2 = m_\chi^2 + \sigma\phi^2$. For quadratic $V(\phi)$ only the second mode contributes

$$|\overline{\mathcal{M}}_2|^2 = \frac{\phi_0^4}{32} \left[\sigma - \lambda \left(1 + \frac{m_{\text{eff}}^2}{2m_\phi^2} \right) \right]^2 \equiv \frac{1}{8} \frac{\rho_\phi^2}{m_\phi^4} \hat{\sigma}^2$$

- For $\sigma/\lambda > 1$, direct decay suppressed by graviton exchange
- For $\sigma/\lambda < 1$, gravitational production suppressed by direct coupling
- For $\sigma/\lambda = 1$, complete interference at $m_{\text{eff}} = 0$

1. Preheating



2. Weak coupling



3. Strong coupling



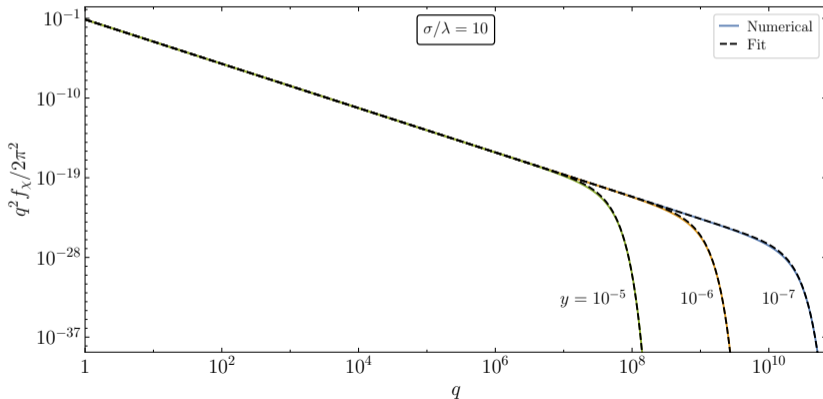
4. Constraints



The Boltzmann approximation

With $f_{\chi}(|\mathbf{P}|, t) \equiv \frac{1}{2} \left[\exp(2f_{\chi}^c(|\mathbf{P}|, t)) - 1 \right]$ and $q \equiv \frac{P}{T_{\star}} \left(\frac{a}{a_0} \right)$, $T_{\star} \equiv m_{\phi} \left(\frac{a_{\text{end}}}{a_0} \right)$

$$f_{\chi}^c(q, t) \simeq \frac{\sqrt{3}\pi\hat{\sigma}^2\rho_{\text{end}}^{3/2}M_P}{16m_{\phi}^7} q^{-9/2} e^{-1.56\left(\frac{a_{\text{end}}}{a_{\text{reh}}}\right)^2 q^2} \theta(q-1), \quad (t \gg t_{\text{reh}}, \beta \simeq 1)$$



1. Preheating



2. Weak coupling



3. Strong coupling



4. Constraints



Non-perturbative particle production

Equation of motion for χ

$$\left(\frac{d^2}{dt^2} - \frac{\nabla^2}{a^2} + 3H \frac{d}{dt} + m_\chi^2 + \sigma \phi^2 \right) \chi = 0$$

In terms of conformal time, $dt/d\tau = a$, and the re-scaled field $X = a\chi$

$$X(\tau, \mathbf{x}) = \int \frac{d^3 \mathbf{p}}{(2\pi)^{3/2}} e^{-i\mathbf{p} \cdot \mathbf{x}} \left[X_p(\tau) \hat{a}_{\mathbf{p}} + X_p^*(\tau) \hat{a}_{-\mathbf{p}}^\dagger \right]$$

the equation of motion for the mode functions is

$$X_p'' + \omega_p^2 X_p = 0$$

with

$$\omega_p^2 = p^2 + a^2 m_{\text{eff}}^2 = p^2 + a^2 \left(m_\chi^2 + \sigma \phi^2 + \frac{1}{6} R \right)$$

and Bunch-Davies initial condition $X_p(\tau_0) = \frac{1}{\sqrt{2\omega_p}}$, $X_p'(\tau_0) = -\frac{i\omega_p}{\sqrt{2\omega_p}}$

1. Preheating



2. Weak coupling



3. Strong coupling

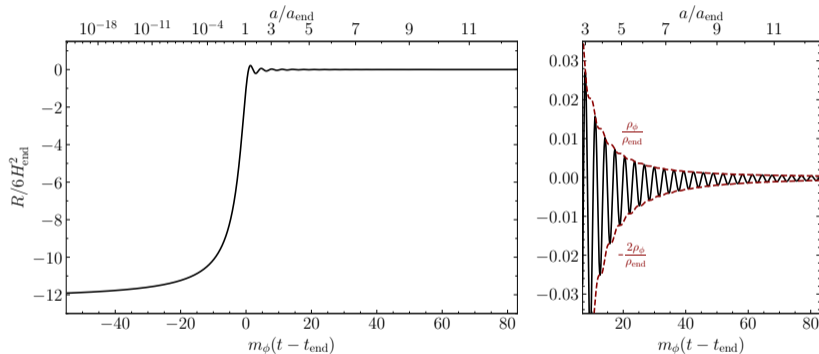


4. Constraints



Non-perturbative particle production

$$\frac{1}{6}R = -\frac{a''}{a^3} = -\frac{1}{6M_P^2} (4V - \dot{\phi}^2)$$



$$\omega_p^2(t_{\text{end}}) = p^2 + a_{\text{end}}^2 (m_\chi^2 + \sigma\phi^2 - H_{\text{end}}^2)$$

For $\sigma/\lambda \lesssim 10^{-1/2}$ superhorizon modes grow during inflation due to tachyonic instability

1. Preheating



2. Weak coupling



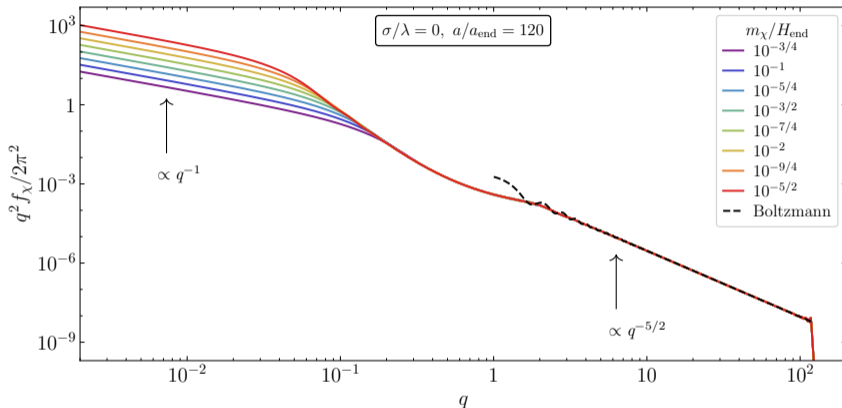
3. Strong coupling



4. Constraints



Pure gravitational production



IR regulated by the present comoving scale $p_0 = a_0 H_0$, or $q_0 = \frac{H_0}{m_\phi} \left(\frac{a_0}{a_{\text{end}}} \right)$

N. Herring, D. Boyanovsky and A. Zentner, PRD 101 (2020), 083516

S. Ling and A. Long, PRD 103 (2021), 103532

1. Preheating



2. Weak coupling



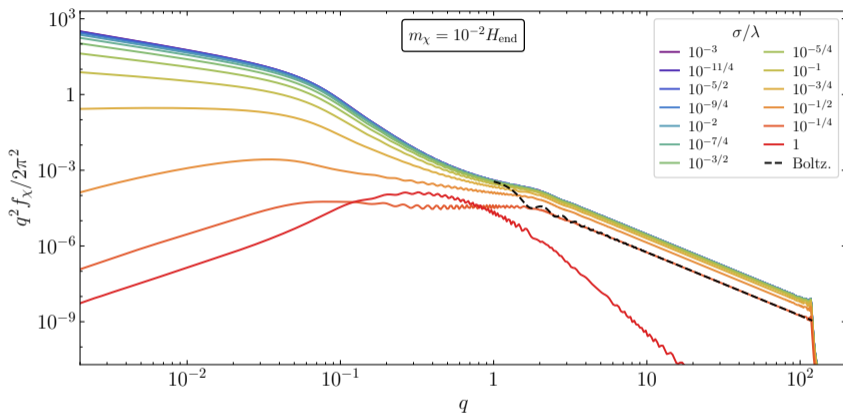
3. Strong coupling



4. Constraints



Weak coupling ($\sigma/\lambda \leq 1$)



Non-perturbatively the interference is not exact for $\sigma/\lambda = 1$, and $f_\chi \sim q^{-15/2}$ in the UV

1. Preheating



2. Weak coupling



3. Strong coupling



4. Constraints



Strong coupling: Hartree

At strong coupling the resonant production of χ can influence the background dynamics

Strong, but not too strong couplings ($\sigma/\lambda \lesssim 10^{7/2}$): **Hartree approximation**

$$\rho_\phi + \rho_\chi = 3H^2 M_p^2,$$

$$\ddot{\phi} + 3H\dot{\phi} + V_\phi + \sigma\langle\chi^2\rangle\phi = 0,$$

where

$$\langle\chi^2\rangle = \frac{1}{(2\pi)^3 a^2} \int d^3\mathbf{p} \left(|X_p|^2 - \frac{1}{2\omega_p} \right).$$

L. Kofman, A. Linde, A. Starobinsky, PRD 56 (1997) 3258

MG, K. Kaneta, Y. Mambrini, K. Olive, S. Verner, JCAP 03 (2022) 016

No backreaction if $\rho_\chi \lesssim 0.1\rho_\phi$

1. Preheating



2. Weak coupling



3. Strong coupling



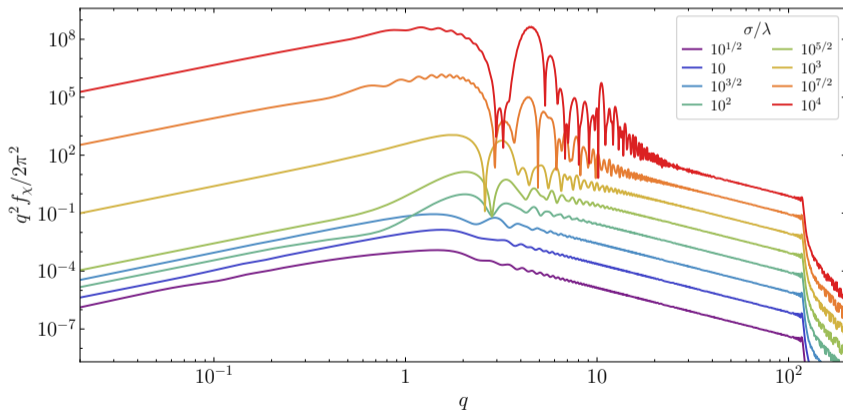
4. Constraints



Strong coupling: Hartree

At strong coupling the resonant production of χ can influence the background dynamics

Strong, but not too strong couplings ($\sigma/\lambda \lesssim 10^{7/2}$): Hartree approximation



1. Preheating



2. Weak coupling



3. Strong coupling



4. Constraints



Strong coupling: Lattice

For stronger couplings the re-scattering of χ into ϕ disrupts the inflaton condensate

Mode-mode couplings of perturbations make spectral codes unsuitable for the task

Solution: **Classical fields on a configuration-space lattice**

$$\ddot{\phi} + 3H\dot{\phi} - \frac{\nabla^2 \phi}{a^2} + V_{,\phi} = 0$$

$$\ddot{\chi} + 3H\dot{\chi} - \frac{\nabla^2 \chi}{a^2} + V_{,\chi} = 0$$

Software of choice: CosmoLattice (v1.0)

D. Figueroa, et al., arXiv:2102.01031 [astro-ph.CO]

Caveat: no metric perturbations

1. Preheating



2. Weak coupling



3. Strong coupling



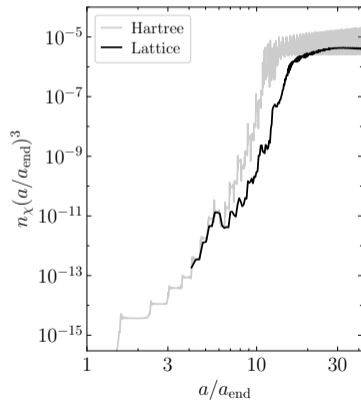
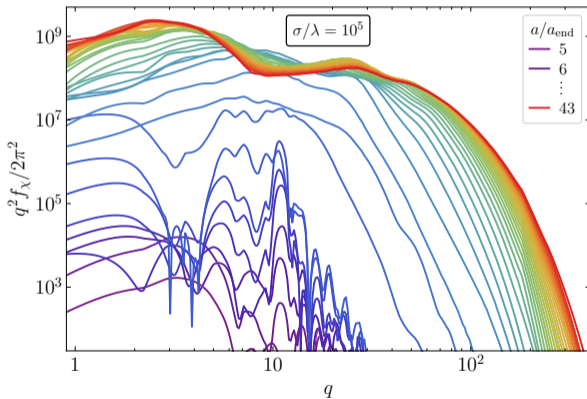
4. Constraints



Strong coupling: Lattice

Re-scattering leads to a broadening distribution with pseudo-thermal tail for ϕ and χ

$$f_\chi \sim e^{-\alpha(\sigma/\lambda;t)q} \quad \text{in the UV}$$



1. Preheating



2. Weak coupling



3. Strong coupling

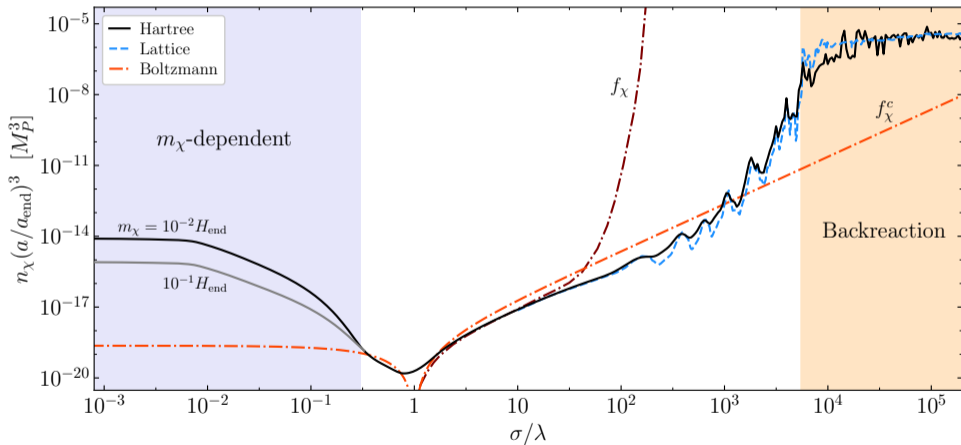


4. Constraints



Comoving number densities

Boltzmann has a very limited range of applicability



1. Preheating



2. Weak coupling



3. Strong coupling

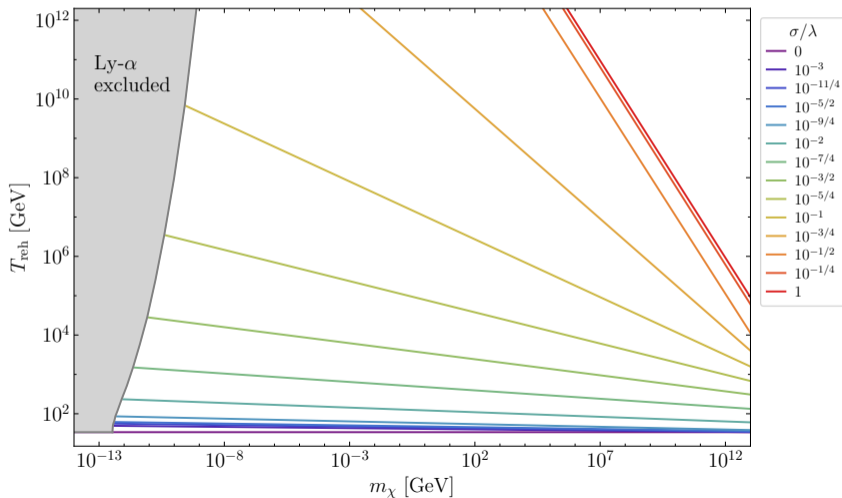


4. Constraints



Relic abundance at weak coupling

Saturating the DM relic abundance (curves of $\Omega_\chi h^2 = 0.12$):



1. Preheating



2. Weak coupling



3. Strong coupling

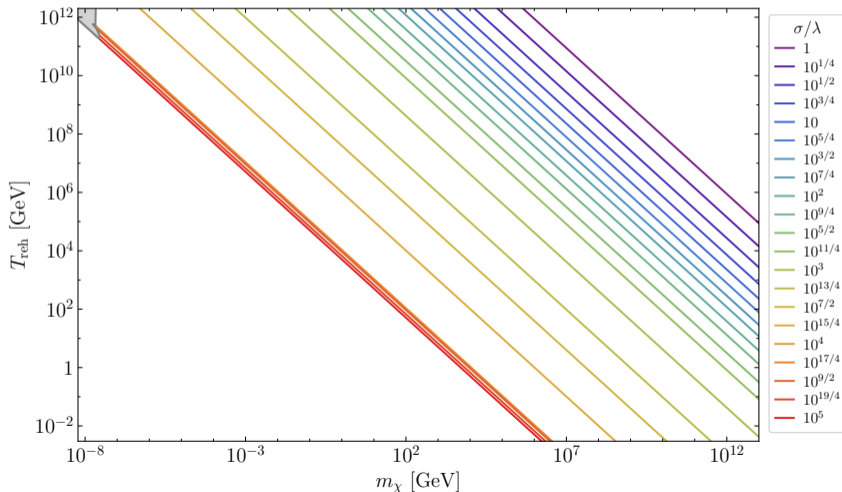


4. Constraints



Relic abundance at strong coupling

Saturating the DM relic abundance (curves of $\Omega_\chi h^2 = 0.12$):



1. Preheating



2. Weak coupling



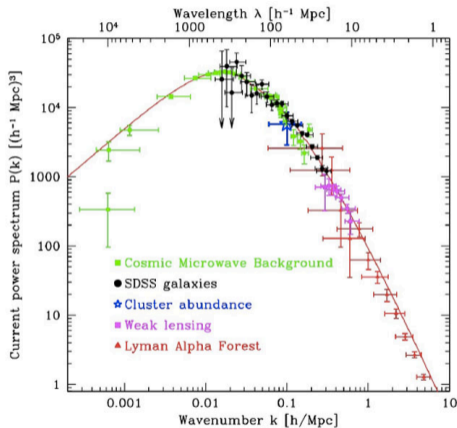
3. Strong coupling



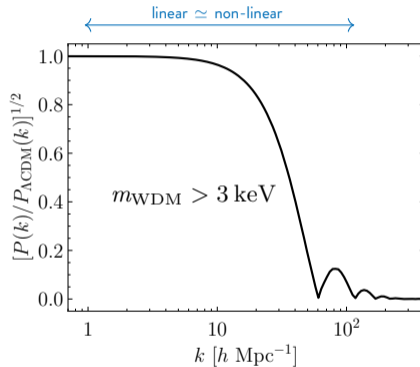
4. Constraints



How warm is out-of-equilibrium dark matter?



R. Murgia, V. Iršič and M. Viel, PRD 98 (2018), 083540



G. Ballesteros, MG and M. Pierre, JCAP 03 (2021), 101

1. Preheating



2. Weak coupling



3. Strong coupling



4. Constraints



How warm is out-of-equilibrium dark matter?

R. Murgia, V. Iršič and M. Viel, PRD 98 (2018), 083540

$$f(\mathbf{x}, \mathbf{p}, \tau) = f_0(|\mathbf{p}|, \tau) [1 + \Psi(\mathbf{x}, \mathbf{p}, \tau)]$$



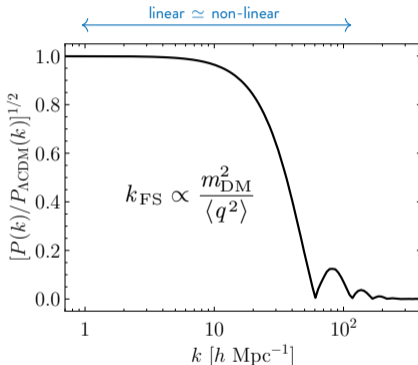
$$\ddot{\delta} + \mathcal{H}\dot{\delta} - \frac{3}{2}\mathcal{H}^2 \left(1 - w \frac{10}{9} \frac{k^2}{\mathcal{H}^2}\right) \delta = 0$$

$$k_{\text{FS}}^2(a) = \frac{9}{10} \frac{\mathcal{H}^2}{w}$$

$$k_{\text{H}}(a) = \left[\int_0^a \frac{d\tilde{a}}{\tilde{a} k_{\text{FS}}(\tilde{a})} \right]^{-1}$$

$$w \simeq \frac{\delta P}{\delta \rho} = \frac{T_\star^2}{3m_{\text{DM}}^2} \frac{\langle q^2 \rangle}{a^2}$$

$$w(m_{\text{DM}}) = w_{\text{WDM}}(m_{\text{WDM}}) \longrightarrow$$



G. Ballesteros, MG and M. Pierre, JCAP 03 (2021), 101

$$m_{\text{DM}} = m_{\text{WDM}} \left(\frac{T_\star}{T_{\text{WDM}}} \right) \sqrt{\frac{\langle q^2 \rangle}{\langle q^2 \rangle_{\text{WDM}}}}$$

1. Preheating



2. Weak coupling



3. Strong coupling



4. Constraints



Lyman- α constraint

From Boltzmann f_{χ}^c ,

$$\langle q^2 \rangle \simeq 0.641 \left(\frac{a_{\text{reh}}}{a_{\text{end}}} \right)^2 \frac{\Gamma(1/4, 1.56(a_{\text{end}}/a_{\text{reh}})^2)}{\Gamma(-3/4, 1.56(a_{\text{end}}/a_{\text{reh}})^2)} \simeq 2.433 \sqrt{\frac{a_{\text{reh}}}{a_{\text{end}}}}$$

for $a_{\text{reh}} \gg a_{\text{end}}$, and

$$\begin{aligned} m_{\text{DM}} &> 15.78 \text{ keV} \left(\frac{m_{\text{WDM}}}{3 \text{ keV}} \right)^{4/3} m_{\phi} \rho_{\text{end}}^{-1/4} g_{\text{reh}}^{-1/12} \\ &\simeq 32.4 \text{ eV} \left(\frac{m_{\text{WDM}}}{3 \text{ keV}} \right)^{4/3} \left(\frac{\lambda}{2.05 \times 10^{-11}} \right)^{1/4} \left(\frac{427/4}{g_{\text{reh}}} \right)^{1/12} \end{aligned}$$

Weaker than WDM one, and without dependence on duration of reheating!

1. Preheating



2. Weak coupling



3. Strong coupling



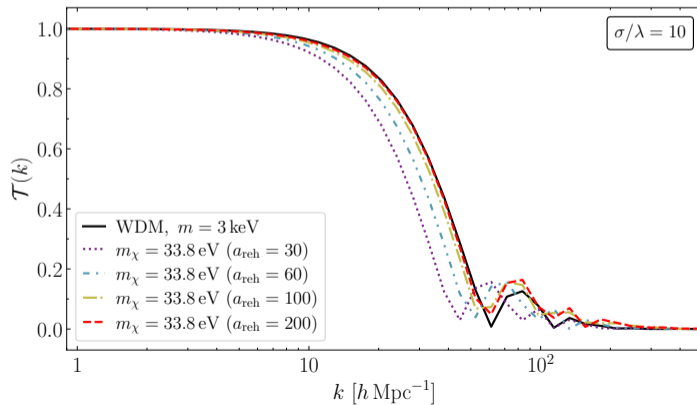
4. Constraints



Lyman- α constraint

For all non-lattice, $\sigma \neq \lambda$ cases, the perturbative tail is present

$$(m_{\text{DM}})_{\text{non-pert}} = (m_{\text{DM}})_{\text{pert}} \sqrt{\frac{\langle q^2 \rangle_{\text{non-pert}}}{\langle q^2 \rangle_{\text{pert}}}}$$



1. Preheating



2. Weak coupling



3. Strong coupling



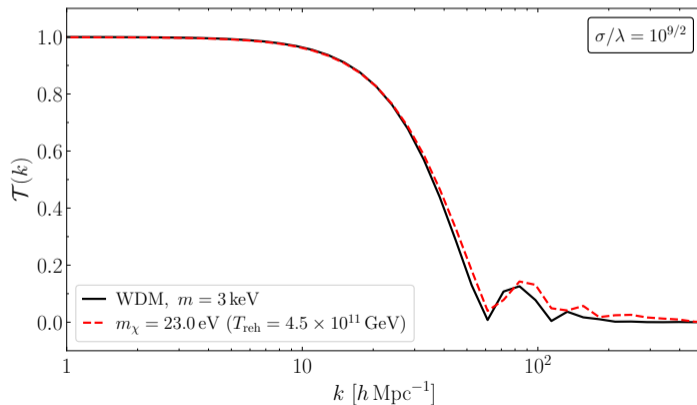
4. Constraints



Lyman- α constraint

For strong backreaction, or $\sigma = \lambda, \langle q^2 \rangle \rightarrow \text{const.}$ during reheating

$$m_{\text{DM}} > 9.58 \text{ keV} \left(\frac{m_{\text{WDM}}}{3 \text{ keV}} \right)^{4/3} \sqrt{\langle q^2 \rangle} \frac{m_\phi T_{\text{reh}}^{1/3}}{\rho_{\text{end}}^{1/3}}.$$



1. Preheating



2. Weak coupling



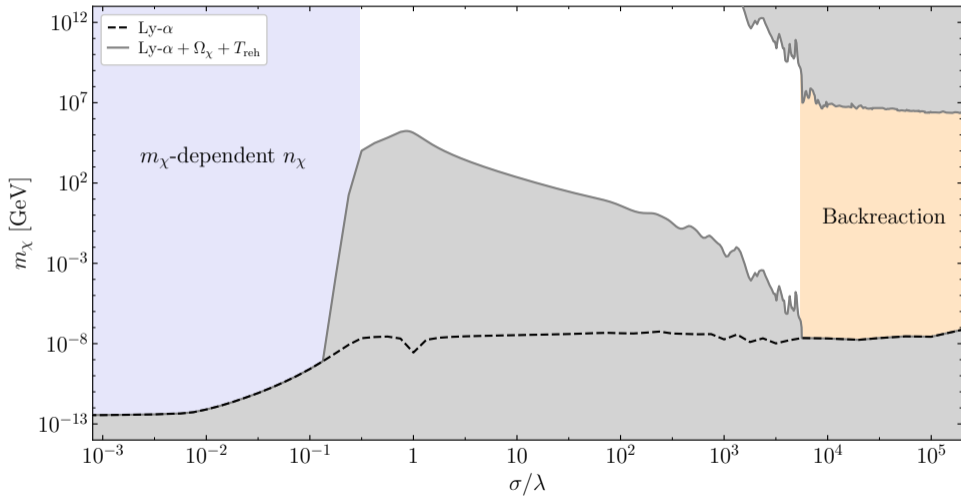
3. Strong coupling



4. Constraints



The allowed parameter space



1. Preheating



2. Weak coupling



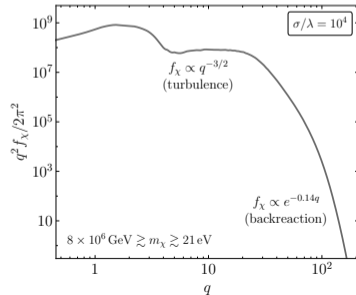
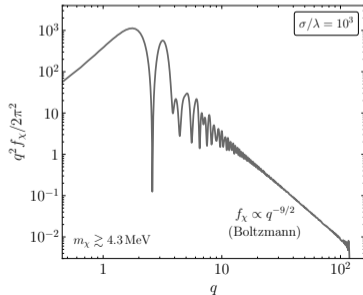
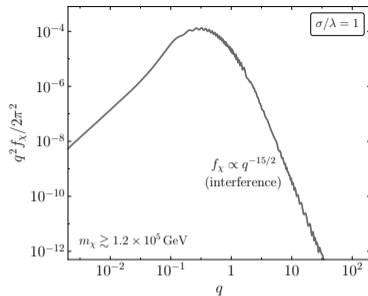
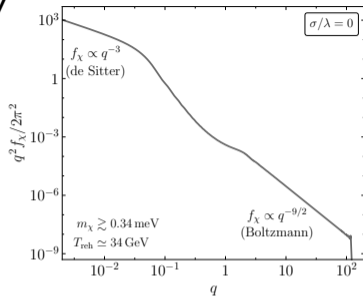
3. Strong coupling



4. Constraints



Summary



1. Preheating



2. Weak coupling



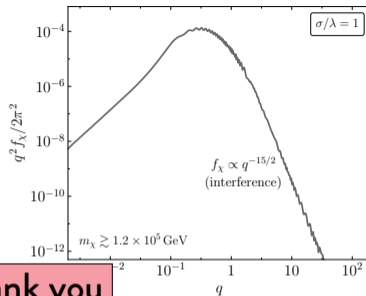
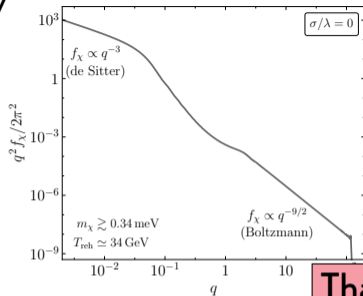
3. Strong coupling



4. Constraints



Summary



Thank you

