Null test BSM searches with rare charm baryon decays

Marcel Golz

based on
JHEP 09 (2021) 208; [arxiv:2107.13010] MG, G. Hiller, T. Magorsch

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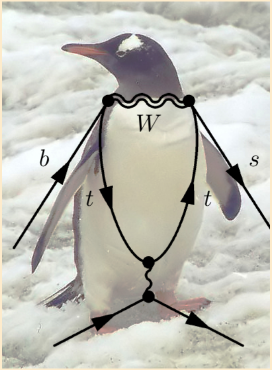
Outline

- Part I - Introduction to rare charm decays
  - Effective field theory framework
  - Null test overview
  - Baryon modes
- Part II - Null tests with baryons
  - Forward-backward asymmetry
  - F_L
  - LU ratios
  - CP–violation
- Outlook

Part I - Introduction rare charm decays

EFT at the charm scale de Boer, (2017), PhD thesis, TU Dortmund

What is the difference between $b \to s$ and $c \to u$ penguins?



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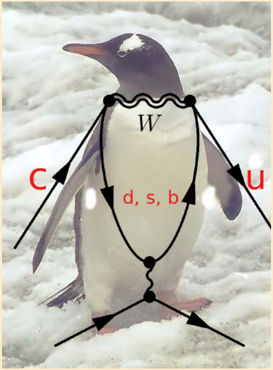

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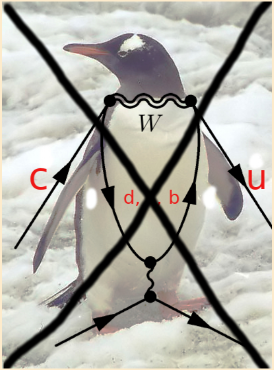
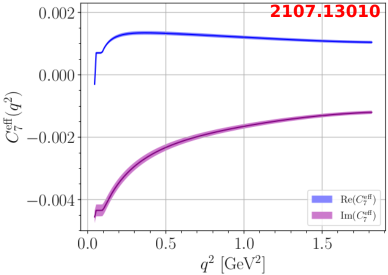
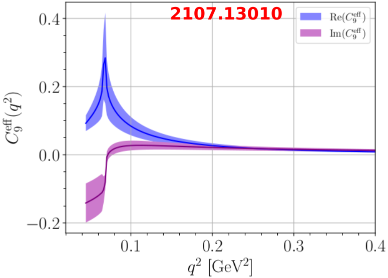


- Light quark masses need to be set to zero at μ_W
- Effective GIM-mechanism kills $C_{7,9,10}$ at μ_W
- $C_{7,9}^{\mathrm{eff}}$ are induced by RG running to μ_c
- $C_{10}(\mu_c) = 0$

Introduction





$$\mathcal{H}_{\text{eff}} \supset -\frac{4G_F}{\sqrt{2}}\frac{\alpha_e}{4\pi}\left[\sum_{i=7,9,10,S,P}(C_iO_i + C_i'O_i') + \sum_{i=T,T5}C_iO_i\right],$$

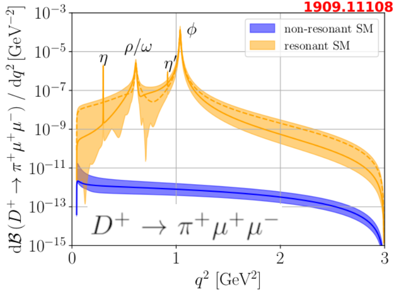
$$O_7 = \frac{m_c}{e}(\bar{u}_L\sigma_{\mu\nu}c_R)F^{\mu\nu},$$

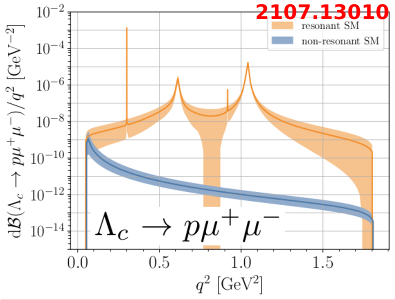
$$O_9 = (\bar{u}_L\gamma_\mu c_L)(\bar{\ell}\gamma^\mu\ell), \quad O_{10} = (\bar{u}_L\gamma_\mu c_L)(\bar{\ell}\gamma^\mu\gamma_5\ell),$$

$$O_S = (\bar{u}_Lc_R)(\bar{\ell}\ell), \quad O_P = (\bar{u}_Lc_R)(\bar{\ell}\gamma_5\ell),$$

$$O_T = \frac{1}{2}(\bar{u}\sigma_{\mu\nu}c)(\bar{\ell}\sigma^{\mu\nu}\ell), \; O_{T5} = \frac{1}{2}(\bar{u}\sigma_{\mu\nu}c)(\bar{\ell}\sigma^{\mu\nu}\gamma_5\ell).$$

- Primed operators obtained with $L \leftrightarrow R$
- $C_7^{\text{eff}} = \mathcal{O}(10^{-3})$, $C_9^{\text{eff}} = \mathcal{O}(10^{-2})$.
- $C_{10}^{(\prime)} = C_S^{(\prime)} = C_P^{(\prime)} = 0$ and $C_T = C_{T5} = 0$



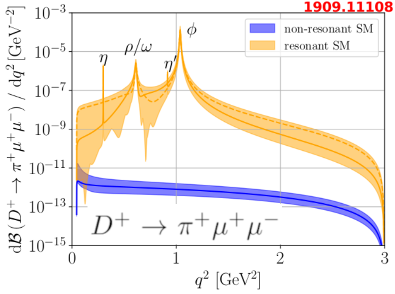


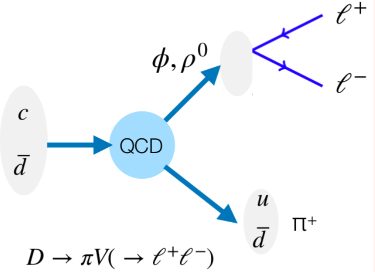
- SM contributions dominated by long range dynamics
  $\mathcal{B}^{\mathrm{SM}}(D \to \pi\ell^+\ell^-) \approx \mathcal{B}(D \to \pi V(\to \ell^+\ell^-))$
- Parametrized by a sum of Breit-Wigner contributions (fit from data)

$$C_9^R(q^2) = a_\omega e^{i\delta_\omega}\left(\frac{1}{q^2 - m_\omega^2 + i m_\omega \Gamma_\omega} - \frac{3}{q^2 - m_\rho^2 + i m_\rho \Gamma_\rho}\right) + \frac{a_\phi e^{i\delta_\phi}}{q^2 - m_\phi^2 + i m_\phi \Gamma_\phi}.$$

$$C_P^R(q^2) = \frac{a_\eta e^{i\delta_\eta}}{q^2 - m_\eta^2 + i m_\eta \Gamma_\eta} + \frac{a_{\eta'}}{q^2 - m_{\eta'}^2 + i m_{\eta'} \Gamma_{\eta'}}$$

Introduction





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Introduction

- Parameters in C_9^R, C_P^R are main source of uncertainties
- Process dependent 1909.11108, 1805.08516

	a_ρ/ GeV^2	a_ϕ/ GeV^2	a_η/ GeV^2	$a_{\eta'}$/ GeV^2
$D^+ \to \pi^+$	~ 0.2	~ 0.2	$\sim 6 \times 10^{-4}$	$\sim 8 \times 10^{-4}$
$D^0 \to \pi^0$	~ 0.9	~ 0.3	$\sim 5 \times 10^{-4}$	$\sim 8 \times 10^{-4}$
$D_s^+ \to K^+$	~ 0.5	~ 0.1	$\sim 6 \times 10^{-4}$	$\sim 7 \times 10^{-4}$
$D^0 \to \pi^+\pi^-$	~ 0.7	~ 0.3	$\sim 1 \times 10^{-3}$	$\sim 1 \times 10^{-3}$
$D^0 \to K^+K^-$	~ 0.7	–	$\sim 3 \times 10^{-4}$	–

- Parameters in C_9^R, C_P^R are main source of uncertainti[illegible]
- Process dependent 1909.11108, 1805.[illegible]

	[illegible]	[illegible]	[illegible]	[illegible] GeV^2
$D^+ \to \pi^+$	[illegible]	[illegible]	[illegible]	$\sim 8 \times 10^{-4}$
$D^0 \to \pi^0$	[illegible]	[illegible]	$\sim 5 \times 10^{-4}$	$\sim 8 \times 10^{-4}$
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	$\Lambda_c \to p$	$\Xi_c^+ \to \Sigma^+$	$\Xi_c^0 \to \Sigma^0$	$\Xi_c^0 \to \Lambda^0$	$\Omega_c^0 \to \Xi^0$
a_ω	0.062 ± 0.009	~ 0.06	~ 0.06	~ 0.06	~ 0.05
a_ϕ	0.110 ± 0.008	~ 0.1	~ 0.1	~ 0.1	~ 0.09

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- Parameters in C_9^R, C_P^R are main source of uncertainties
- Process dependent 1909.11108, 1805.0...

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	$\Lambda_c \to p$	$\Xi_c^+ \to \Sigma^+$	$\Xi_c^0 \to \Sigma^0$	$\Xi_c^0 \to \Lambda^0$	$\Omega_c^0 \to \Xi^0$
a_ω	0.062 ± 0.009	~ 0.06	~ 0.06	~ 0.06	~ 0.05
a_ϕ	0.110 ± 0.008	~ 0.1	~ 0.1	~ 0.1	~ 0.09

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- With more data model parameters (a_i, δ_i) can be constrained and the model can be improved.

Overview null test strategies

- Rare charm decays observed only at the resonances, for the rest of the signal region U.L. are available at 90 % C.L. (see 2011.09478)
  - $\mathcal{B}(D^0 \to \mu^+\mu^-) < 6.2 \times 10^{-9}$
  - $\mathcal{B}(D^+ \to \pi^+\mu^+\mu^-) < 6.7 \times 10^{-8}$
  - $\mathcal{B}(D^0 \to \pi^+\pi^-\mu^+\mu^-) < 5.5 \times 10^{-7}$
  - $\mathcal{B}(\Lambda_c \to p\mu^+\mu^-) < 7.7 \times 10^{-8}$
  - **New results from LHCb 2111.03327**

Direct bounds on WC's from $\mathcal{B}(D^+ \to \pi^+\mu^+\mu^-)$, $\mathcal{B}(D^0 \to \mu^+\mu^-)$ imply

- $C_{S,P} \lesssim \mathcal{O}(0.1)$ $C_{7,9,10,T,T5} \lesssim \mathcal{O}(1)$

Overview null test strategies

NP searches in branching ratios are challenging.

Instead define null test observables where

Signal $\leftrightarrow$ NP

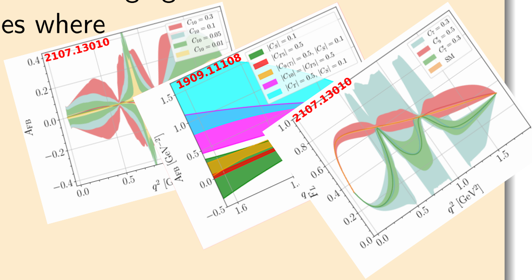
Overview null test strategies

NP searches in branching ratios are challenging.

Instead define null test observables where

Signal ↔ NP

- Angular observables

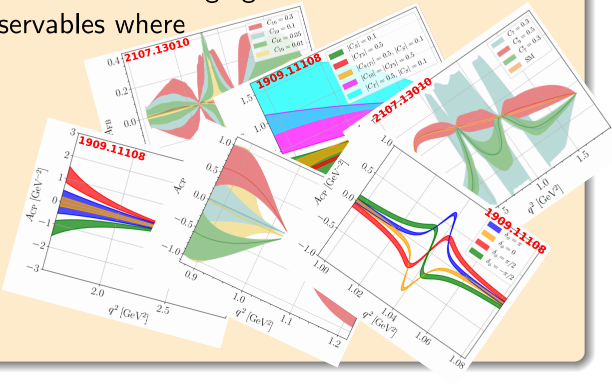


NP searches in branching ratios are challenging.

Instead define null test observables where

Signal $\leftrightarrow$ NP

- Angular observables
- CP–asymmetries



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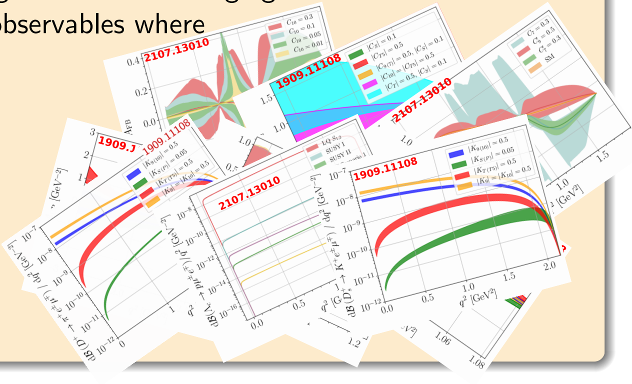
Instead define null test observables where

Signal $\leftrightarrow$ NP

- Angular observables
- CP–asymmetries
- LFV



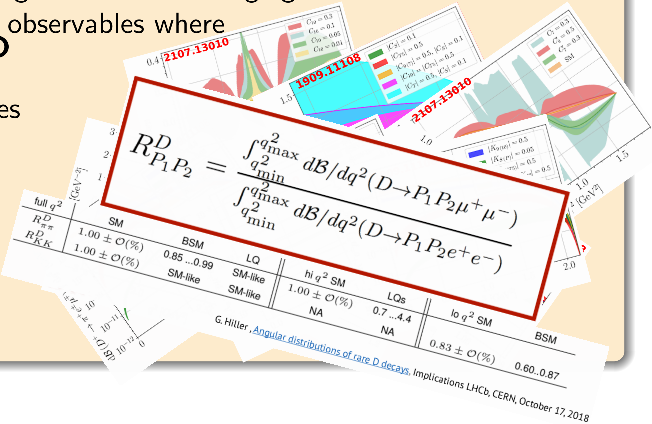
Overview null test strategies

NP searches in branching ratios are challenging.

Instead define null test observables where

Signal $\leftrightarrow$ NP

- Angular observables
- CP–asymmetries
- LFV
- LU ratios



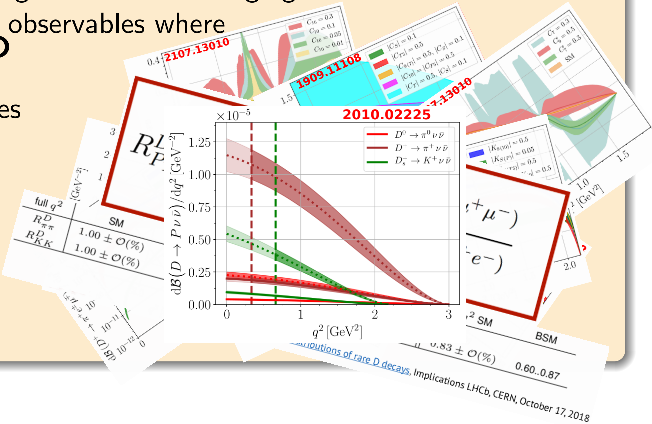
Overview null test strategies

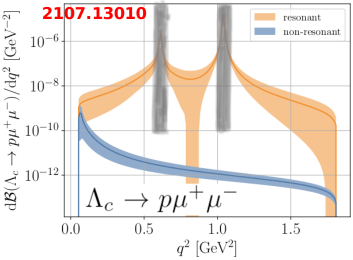
NP searches in branching ratios are challenging.

Instead define null test observables where

Signal $\leftrightarrow$ NP

- Angular observables
- CP–asymmetries
- LFV
- LU ratios
- Dineutrino modes





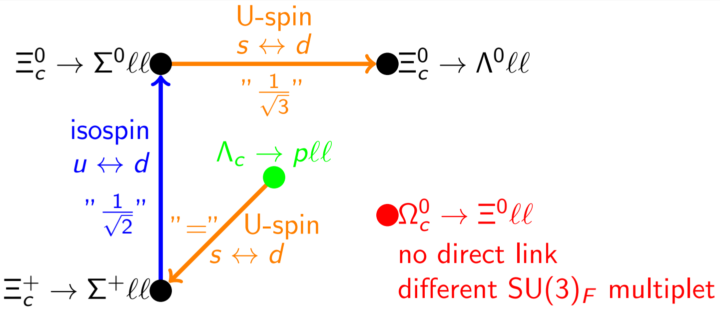
- LHCb upper limit at 90 % C.L. with ±40 MeV cuts around known resonance masses, then extrapolated to full q^2 region

$$\mathcal{B}_{\text{LHCb}}(\Lambda_c \to p\mu^+\mu^-) < 7.7 \times 10^{-8}$$

- Including form factor and resonance uncertainties and integrating the LHCb search region

$$\mathcal{B}^{\text{SM}}(\Lambda_c \to p\mu^+\mu^-) = (1.9^{+1.8}_{-1.5}) \times 10^{-8}$$
$$\mathcal{B}^{\text{SM}}(\Xi_c^+ \to \Sigma^+\mu^+\mu^-) \sim 1.8 \times \mathcal{B}^{\text{SM}}(\Lambda_c \to p\mu^+\mu^-)$$
$$\mathcal{B}^{\text{SM}}(\Xi_c^0 \to \Sigma^0\mu^+\mu^-) \sim 0.4 \times \mathcal{B}^{\text{SM}}(\Lambda_c \to p\mu^+\mu^-)$$
$$\mathcal{B}^{\text{SM}}(\Omega_c^0 \to \Xi^0\mu^+\mu^-) \sim 1.3 \times \mathcal{B}^{\text{SM}}(\Lambda_c \to p\mu^+\mu^-)$$

- Form factors are available for $\Lambda_c \to p\mu^+\mu^-$ from Lattice QCD (see 1712.05783)
- No results for other decay modes (from the lattice) available yet.
- Use $SU(3)_F$ flavor symmetries to relate modes



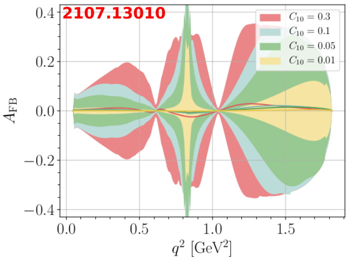
Part II - Null tests with baryons

$\Lambda_c \to p\mu^+\mu^-$ decay distribution:

$$\frac{d\Gamma}{dq^2 d\cos\theta_\ell} = \frac{3}{2}\left(K_{1ss}\sin^2\theta_\ell + K_{1cc}\cos^2\theta_\ell + K_{1c}\cos\theta_\ell\right),$$

$$A_{FB} = \frac{1}{d\Gamma/dq^2}\left[\int_0^1 - \int_{-1}^0\right]\frac{d\Gamma}{dq^2 d\cos\theta_\ell}\, d\cos\theta_\ell = \frac{3}{2}\frac{K_{1c}}{2K_{1ss}+K_{1cc}}.$$

$K_{1c} \sim C_9 C_{10}$, $C_9' C_{10}'$ and $C_7^{(\prime)} C_{10}^{(\prime)}$ interference terms $\Rightarrow$ **no SM**.



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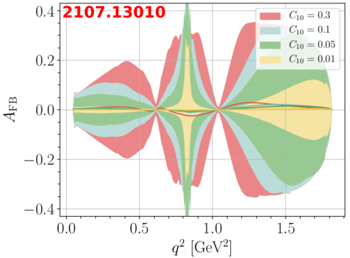
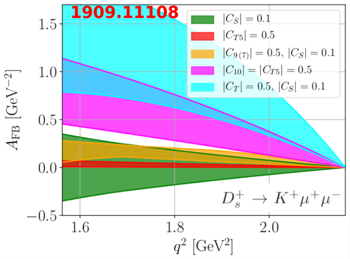
$K_{1c} \sim C_9 C_{10},\ C_9' C_{10}'$ and $C_7^{(\prime)} C_{10}^{(\prime)}$ interference terms $\Rightarrow$ **no SM**.

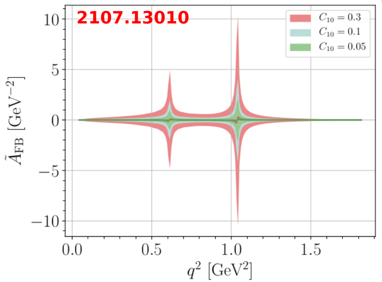
Complementary results in $D \to P\mu^+\mu^-$ (Normalized to integrated rate)

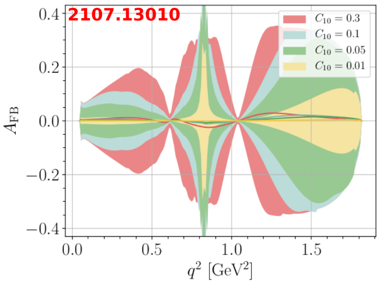




Normalization matters

- $\tilde{A}_{FB} = \frac{3}{2} \frac{K_{1c}}{\Gamma}$ with $\Gamma = \int_{q^2_{min}}^{q^2_{max}} (2\,K_{1ss} + K_{1cc})$
- $A_{FB} = \frac{3}{2} \frac{K_{1c}}{d\Gamma/dq^2}$ with $d\Gamma/dq^2 = 2\,K_{1ss} + K_{1cc}$
- Resonance enhancement

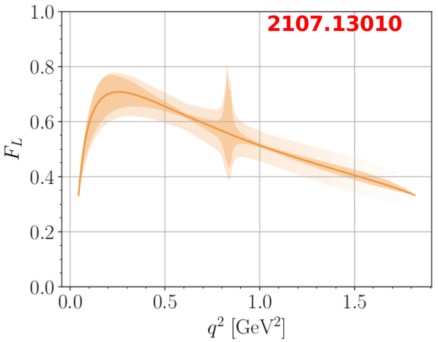


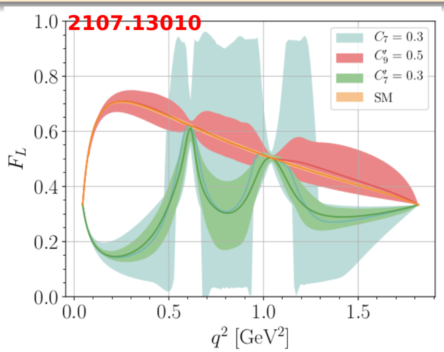


Fraction of longitudinally polarized dimuons

$$F_L = \frac{2\,K_{1ss} - K_{1cc}}{2\,K_{1ss} + K_{1cc}}$$

No null test, but cancellation of hadronic uncertainties in the SM and highly sensitive to $C_7^{(\prime)}$.





LU ratios

$$R_\pi^D = \int_{q^2_{\min}}^{q^2_{\max}} \frac{d\mathcal{B}(D \to \pi \mu^+ \mu^-)}{dq^2} dq^2 \Big/ \int_{q^2_{\min}}^{q^2_{\max}} \frac{d\mathcal{B}(D \to \pi e^+ e^-)}{dq^2} dq^2$$

R_π^D	SM	$\lvert C_9\rvert = 0.5$	$\lvert C_{10}\rvert = 0.5$	$\lvert C_9\rvert = \pm\lvert C_{10}\rvert = 0.5$
full q^2	$1.00 \pm \mathcal{O}(\%)$	SM-like	SM-like	SM-like
low q^2	$0.95 \pm \mathcal{O}(\%)$	$\mathcal{O}(100)$	$\mathcal{O}(100)$	$\mathcal{O}(100)$
high q^2	$1.00 \pm \mathcal{O}(\%)$	$0.2 \ldots 11$	$3 \ldots 7$	$2 \ldots 17$

1909.11108

LU ratios

$$R_\pi^D = \int_{q^2_{\min}}^{q^2_{\max}} \frac{d\mathcal{B}(D \to \pi\mu^+\mu^-)}{dq^2} dq^2 \Big/ \int \dots dq^2$$

$$R_p^{\Lambda_c} = \int_{q^2_{\min}}^{q^2_{\max}} \frac{d\mathcal{B}(\Lambda_c \to p\mu^+\mu^-)}{dq^2} dq^2 \Big/ \int_{q^2_{\min}}^{q^2_{\max}} \frac{d\mathcal{B}(\Lambda_c \to pe^+e^-)}{dq^2} dq^2$$

	SM	$\lvert C_9^\mu \rvert = 0.5$	$\lvert C_{10}^\mu \rvert = 0.5$	$\lvert C_9^\mu \rvert = \lvert C_{10}^\mu \rvert = 0.5$	$\lvert C_9'^\mu \rvert = 0.5$	$\lvert C_{10}'^\mu \rvert = 0.5$	$\lvert C_9'^\mu \rvert = \lvert C_{10}'^\mu \rvert = 0.5$
full q^2	$1.00 \pm \mathcal{O}(\%)$	SM-like	SM-like	SM-like	SM-like	SM-like	SM-like
low q^2	$0.94 \pm \mathcal{O}(\%)$	$7.5 \dots 20$	$4.4 \dots 13$	$11 \dots 32$	$4.6 \dots 14$	$4.4 \dots 13$	$8.2 \dots 26$
high q^2	$1.00 \pm \mathcal{O}(\%)$	$\mathcal{O}(100)$	$\mathcal{O}(100)$	$\mathcal{O}(100)$	$\mathcal{O}(100)$	$\mathcal{O}(100)$	$\mathcal{O}(100)$

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$$R^D_\pi = \int_{q^2_{\text{min}}}^{q^2_{\text{max}}} \ldots$$

$$R^D_{P_1P_2} = \frac{\int_{q^2_{\text{min}}}^{q^2_{\text{max}}} d\mathcal{B}/dq^2(D\to P_1P_2\mu^+\mu^-)}{\int_{q^2_{\text{min}}}^{q^2_{\text{max}}} d\mathcal{B}/dq^2(D\to P_1P_2e^+e^-)}$$

$$\int_{q^2_{\text{min}}}^{q^2_{\text{max}}} \frac{d\mathcal{B}(\Lambda_c\to pe^+e^-)}{dq^2}dq^2$$

full q^2	SM	BSM	LQ	hi q^2 SM	LQs	lo q^2 SM	BSM
$R^D_{\pi\pi}$	$1.00\pm\mathcal{O}(\%)$	0.85 ...0.99	SM-like	$1.00\pm\mathcal{O}(\%)$	0.7 ...4.4	$0.83\pm\mathcal{O}(\%)$	0.60..0.87
R^D_{KK}	$1.00\pm\mathcal{O}(\%)$	SM-like	SM-like	NA	NA		

	SM			$\lvert C_9'^{\mu}\rvert = \lvert C_{10}'^{\mu}\rvert = 0.5$
full q^2	$1.00\pm\mathcal{O}(\%)$	$7.5\ldots20$	$\mathcal{O}(\ldots$	SM-like
low q^2	$0.94\pm\mathcal{O}(\%)$	$\mathcal{O}(100)$		$8.2\ldots26$
high q^2	$1.00\pm\mathcal{O}(\%)$			$\mathcal{O}(100)$

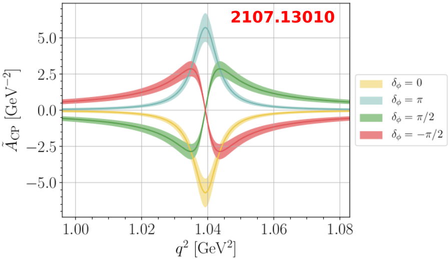
G. Hiller, Angular distributions of rare D decays, Implications LHCb, CERN, October 17, 2018

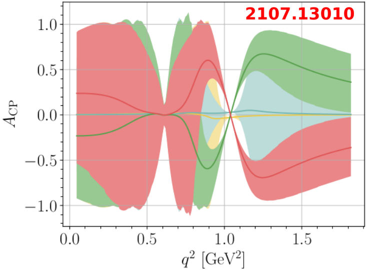
1805.08516

CP–asymmetries

$$\tilde{A}_{\mathrm{CP}}(q^2) = \frac{1}{\Gamma + \bar{\Gamma}}\left(\frac{\mathrm{d}\Gamma}{\mathrm{d}q^2} - \frac{\mathrm{d}\bar{\Gamma}}{\mathrm{d}q^2}\right) \quad \text{with} \quad \Gamma = \int_{q^2_{\mathrm{min}}}^{q^2_{\mathrm{max}}} \frac{\mathrm{d}\Gamma}{\mathrm{d}q^2}\mathrm{d}q^2 \quad \text{or} \quad A_{\mathrm{CP}}(q^2) = \frac{\frac{\mathrm{d}\Gamma}{\mathrm{d}q^2} - \frac{\mathrm{d}\bar{\Gamma}}{\mathrm{d}q^2}}{\frac{\mathrm{d}\Gamma}{\mathrm{d}q^2} + \frac{\mathrm{d}\bar{\Gamma}}{\mathrm{d}q^2}}$$

- Strong phase from resonances needed $\longrightarrow$ measurement around the ϕ resonance. S. Fajfer and N. Kosnik 1208.0759
- NP weak phase needed. $\longrightarrow$ benchmark $C_9 = 0.5\ \exp(i\pi/4)$.
- Binning necessary.



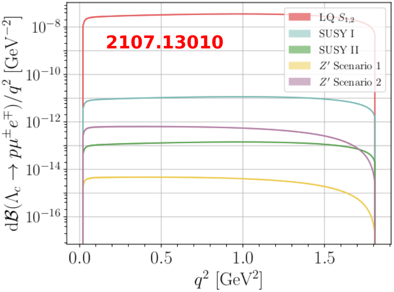


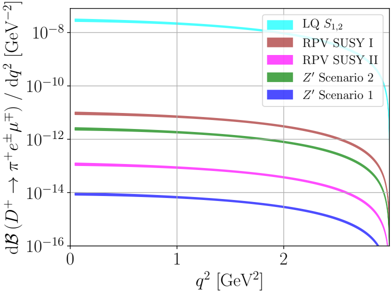
Outlook

- Rare charm decays are dominated by long-distance resonance contributions
- NP can be probed in **clean** null test observables
- Plenty of opportunities with baryon modes
- Resonance enhancement turns the bug into a feature
- Different observables test different Wilson coefficients $\rightarrow$ complementarity

BACKUP

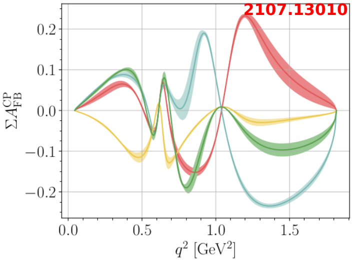
Lepton flavor violation

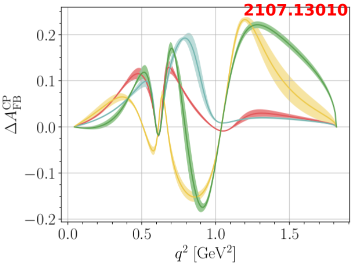


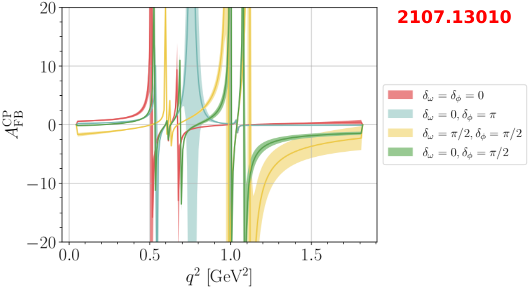


LQ's ($K_9' = K_{10}' = 0.5$),
SUSY + R-parity violation ($K_9 = -K_{10} = 0.009$),
SUSY no R-parity violation ($K_9 = -K_{10} = 0.001$),
Z' 1 ($K_9 = K_9' = -K_{10} = -K_{10}' = 1.4 \cdot 10^{-4}$),
Z' 2 ($K_9 = K_9' = -K_{10} = -K_{10}' = 2.3 \cdot 10^{-4}$).

A_{FB}^{CP} for baryons







Angular observables in $D \to P\ell\ell$