

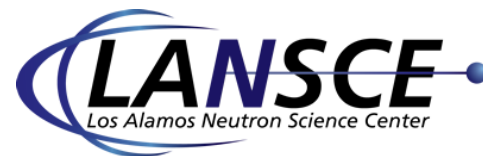
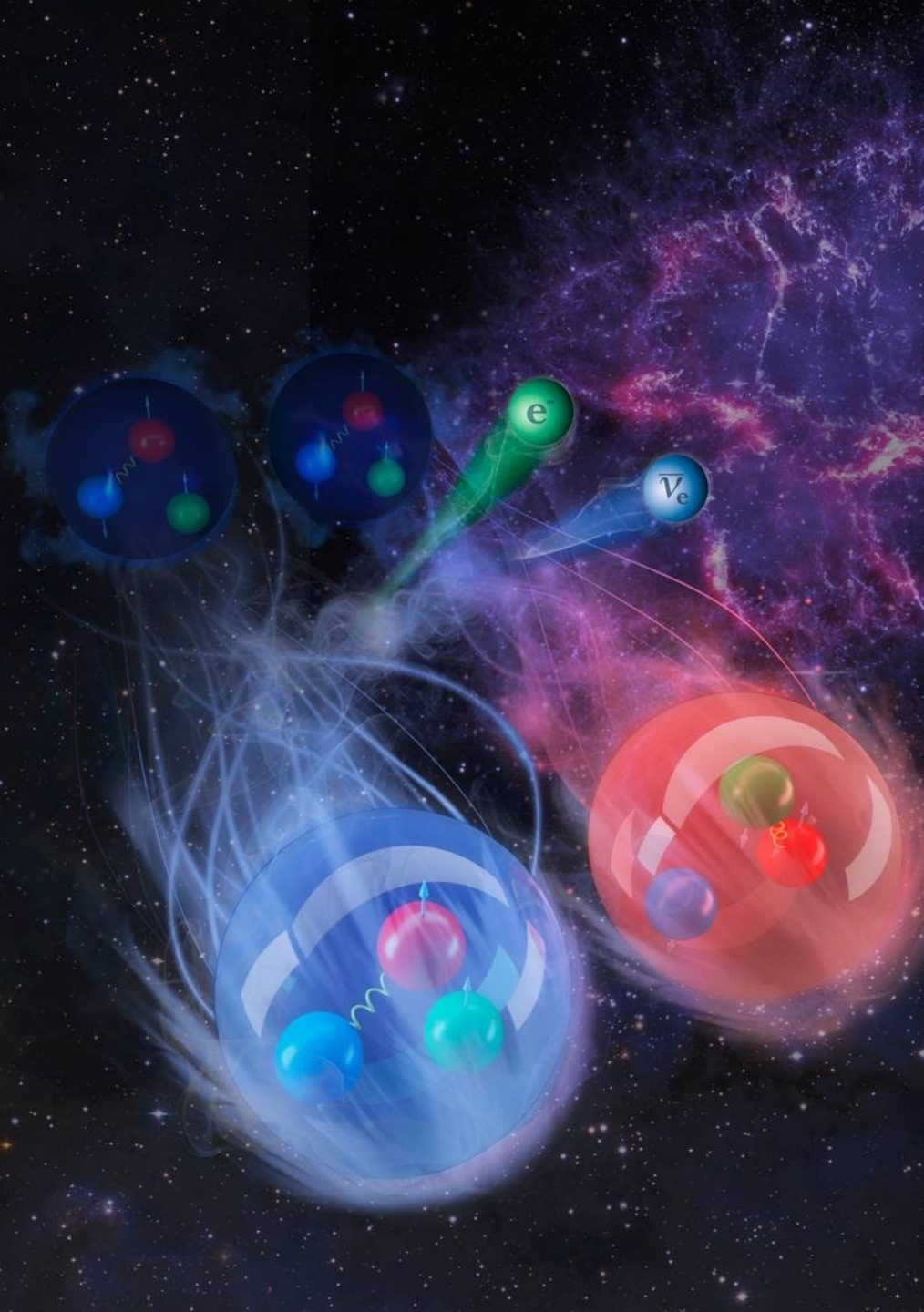
Recent Results from UCNtau and implication for V_{ud}

Chen-Yu Liu

Indiana University

11/23/2021

11th International Workshop on the CKM
Unitarity Triangle (CKM2021)

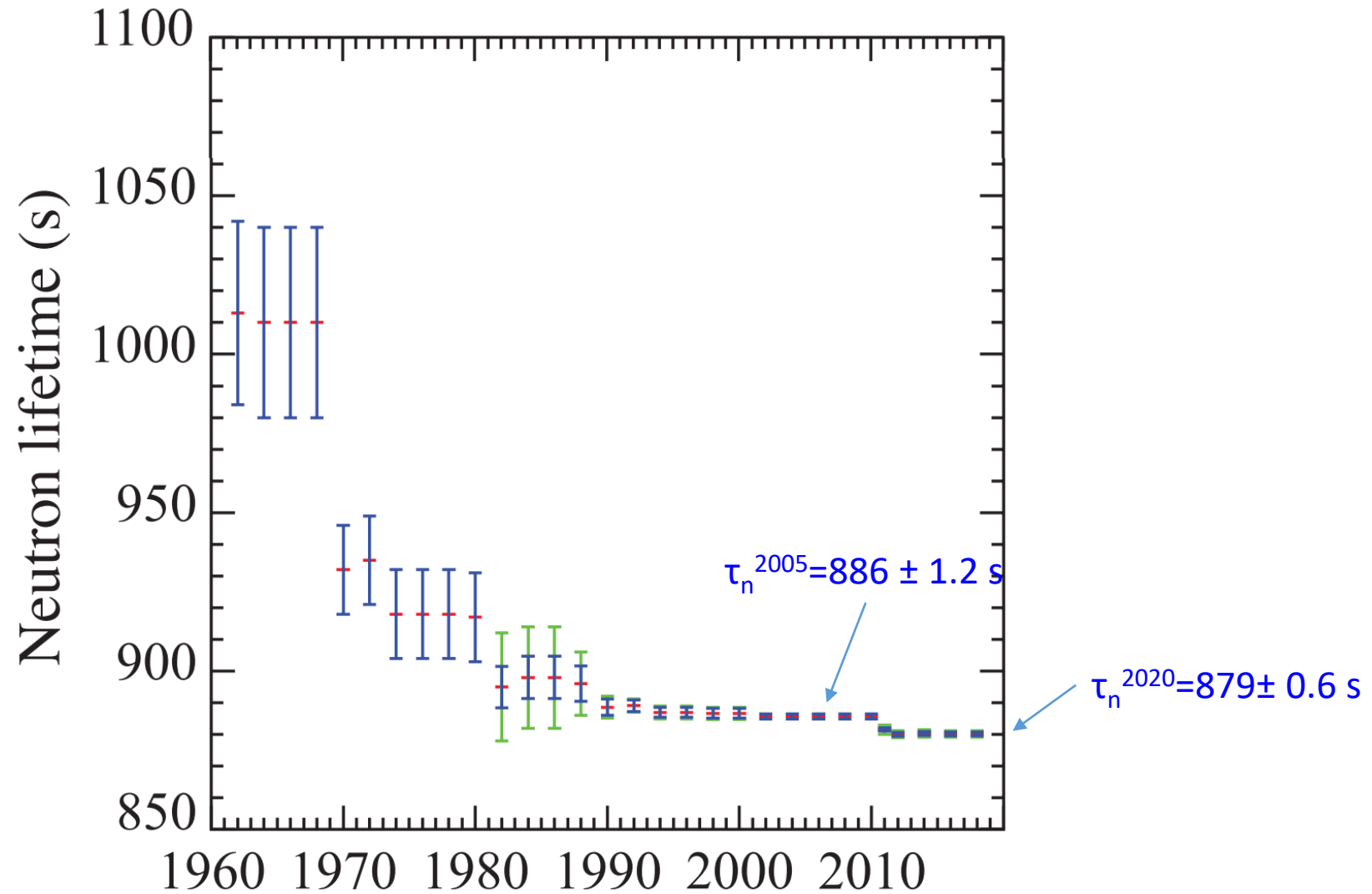


U.S. DEPARTMENT OF
ENERGY

Office of Science

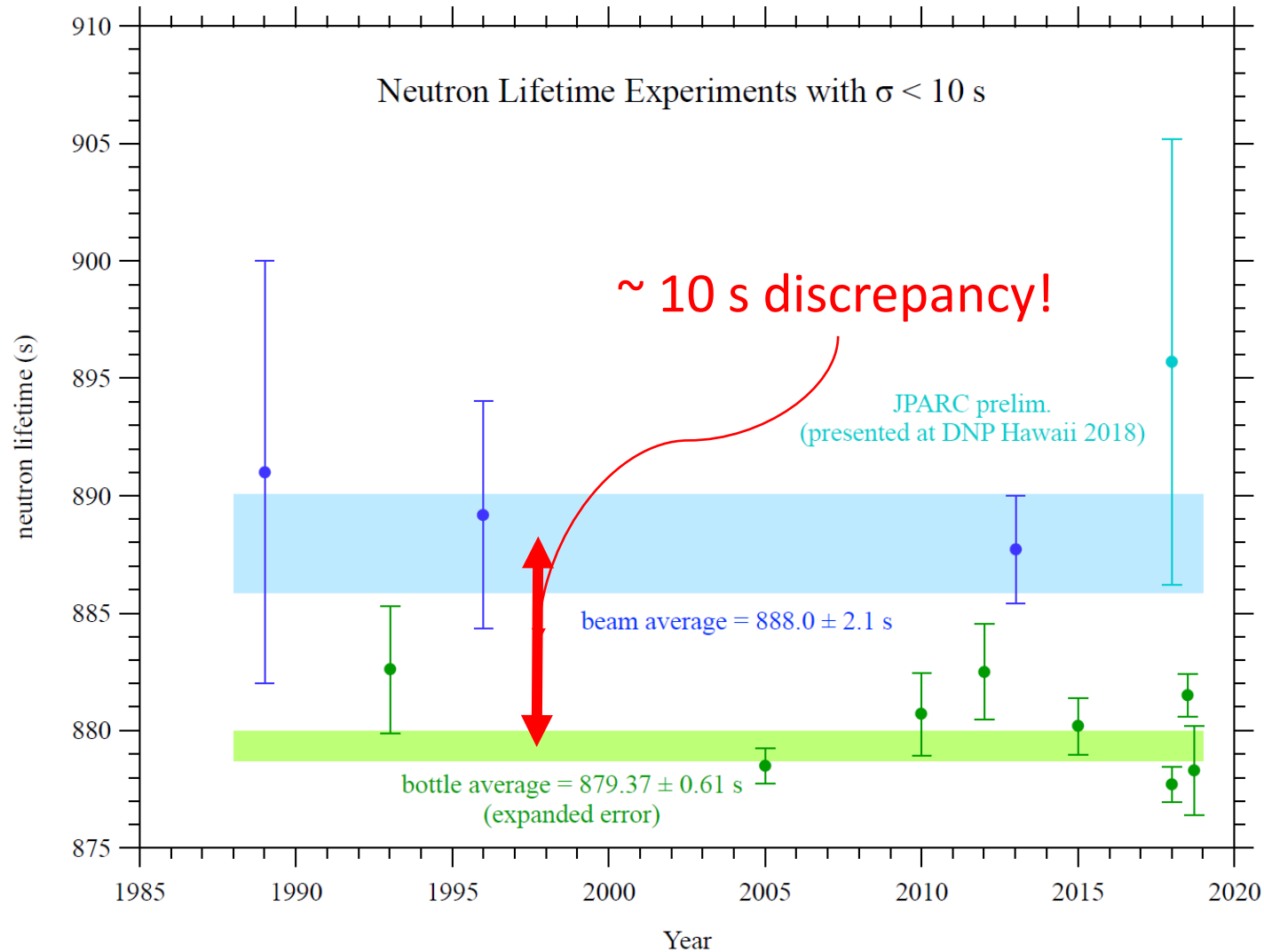


The Particle Data Group (PDG) average of the neutron lifetime:



There is an unresolved discrepancy between two leading methods to measure the neutron lifetime: neutrons in a bottle seem to disappear faster.

The Situation Today - 2019

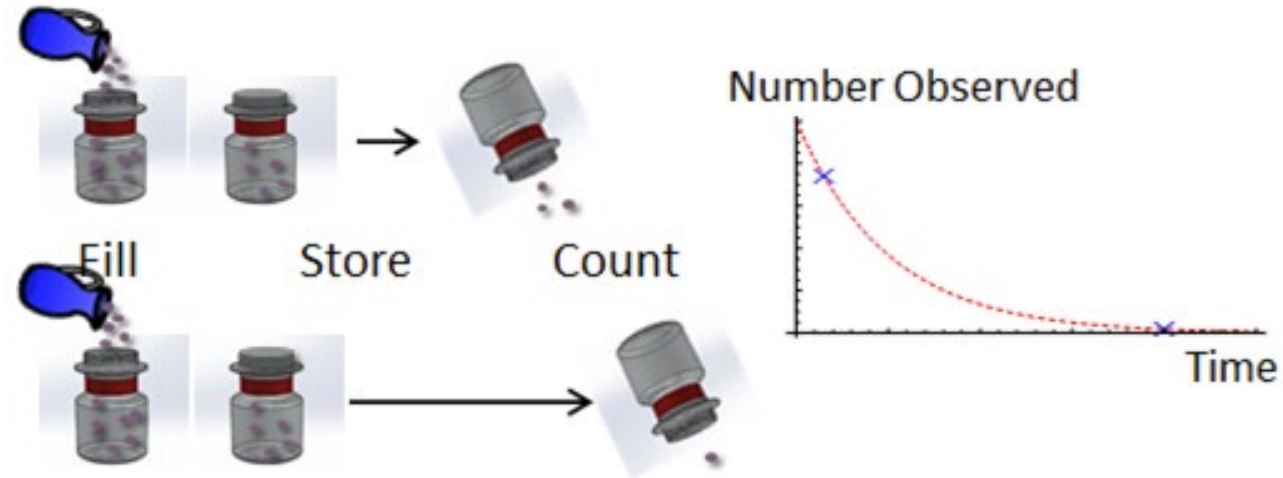
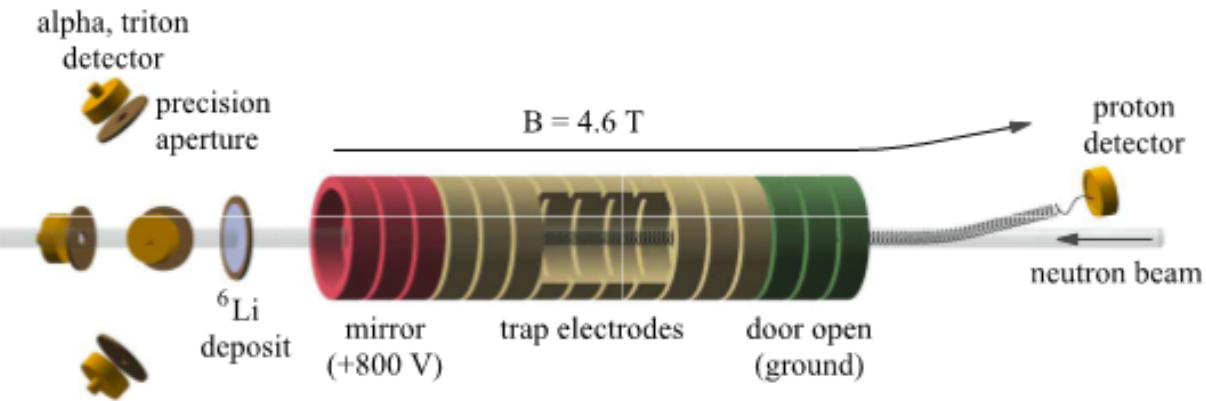


The “beam” and “bottle” techniques

$$\tau_n = \frac{L \dot{N}_n / \epsilon_n}{v_n \dot{N}_p / \epsilon_p}$$

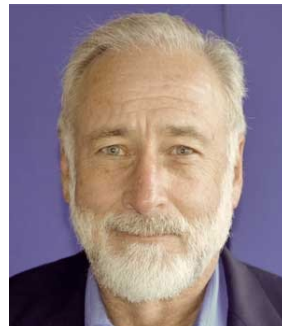
$$Y(t) = Y_0 e^{-t / \tau_{meas}}$$

$$\tau_{meas}^{-1} = \tau_n^{-1} + \tau_{loss}^{-1}$$



“It sounds hard, and it is hard”
said Geoff Greene.

(2021 Bonner Prize Recipient)



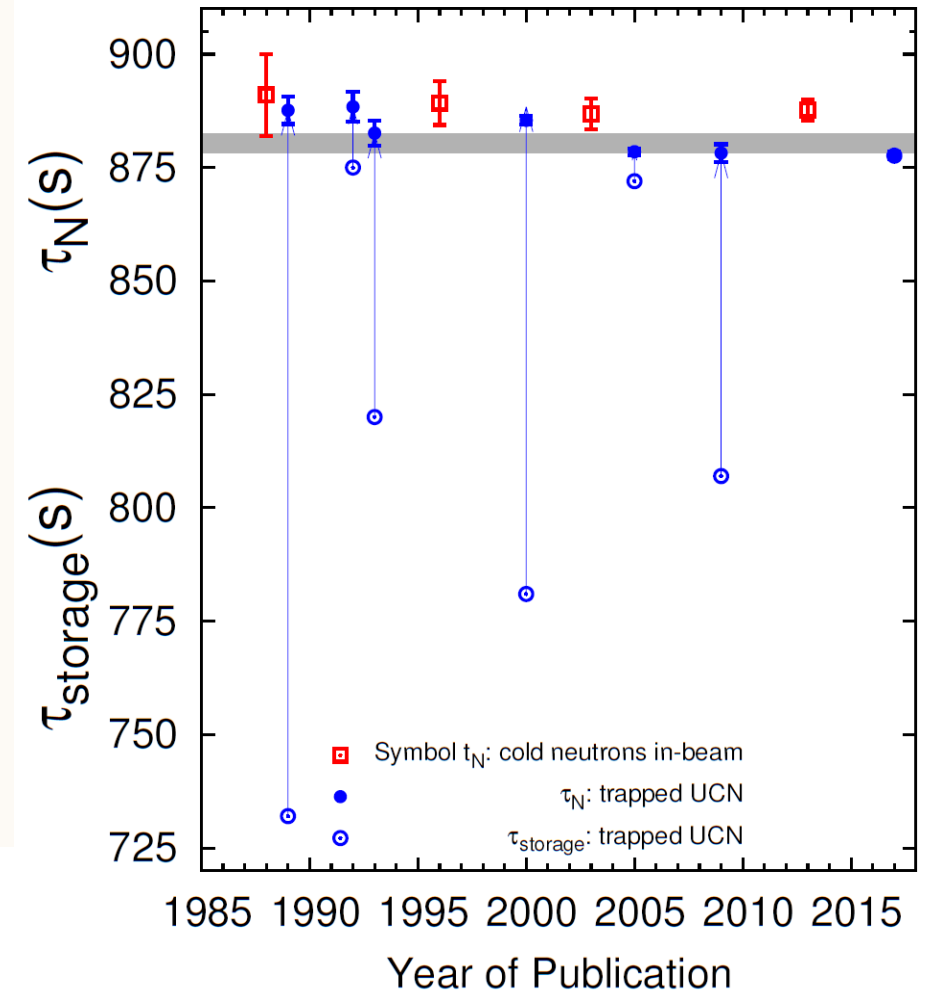
“It sounds easy, and it is hard”
said Geoff Greene.



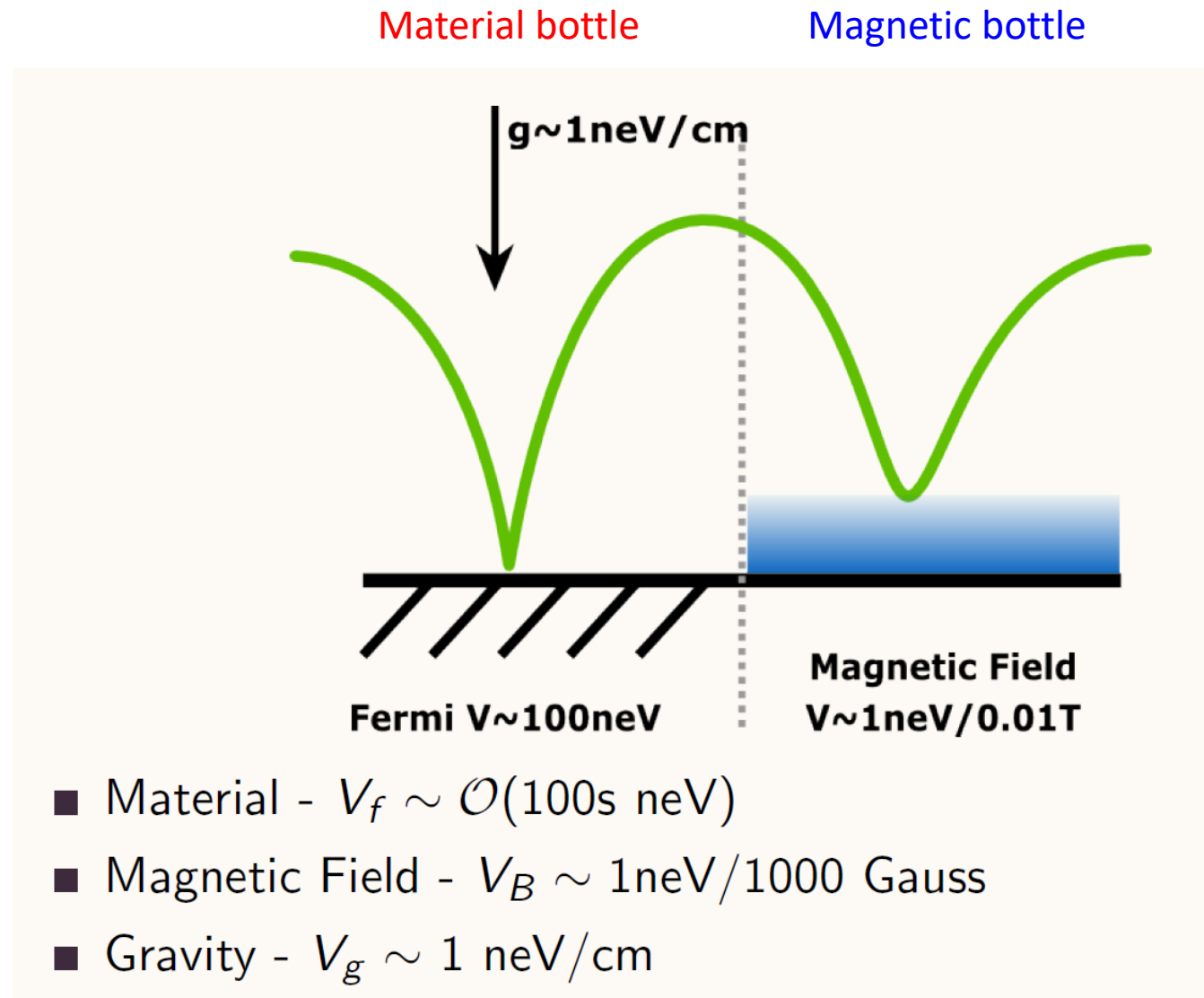
Many experiments need to correct for the systematic effects and extrapolate from the measured lifetime to report the Neutron Lifetime

$$1/\tau_{\text{bottle}} = 1/\tau_n + 1/\tau_{\text{wall}} + 1/\tau_{\text{gas}} + \dots$$

| Author | $\sigma_{\text{stat.}}$ [s] | $\Delta\tau_{\text{sys.}}$ [s] | Extrap. [s] | Method |
|-----------------|-----------------------------|--------------------------------|-------------|--------|
| Arzumanov 2015 | 0.64 | 3.6 | 40-280 | Bottle |
| Steyerl 2012 | 1.4 | ~ 7 | > 200 s | Bottle |
| Pichlmaier 2010 | 1.3 | 1 | 110-300 | Bottle |
| Serebrov 2005 | 0.7 | 0.4 | 10-20 | Bottle |
| Yue 2013 | 1.2 | 1 | 2-15 | Beam |
| Byrne 1996 | 3 | 5.9 | - | Beam |

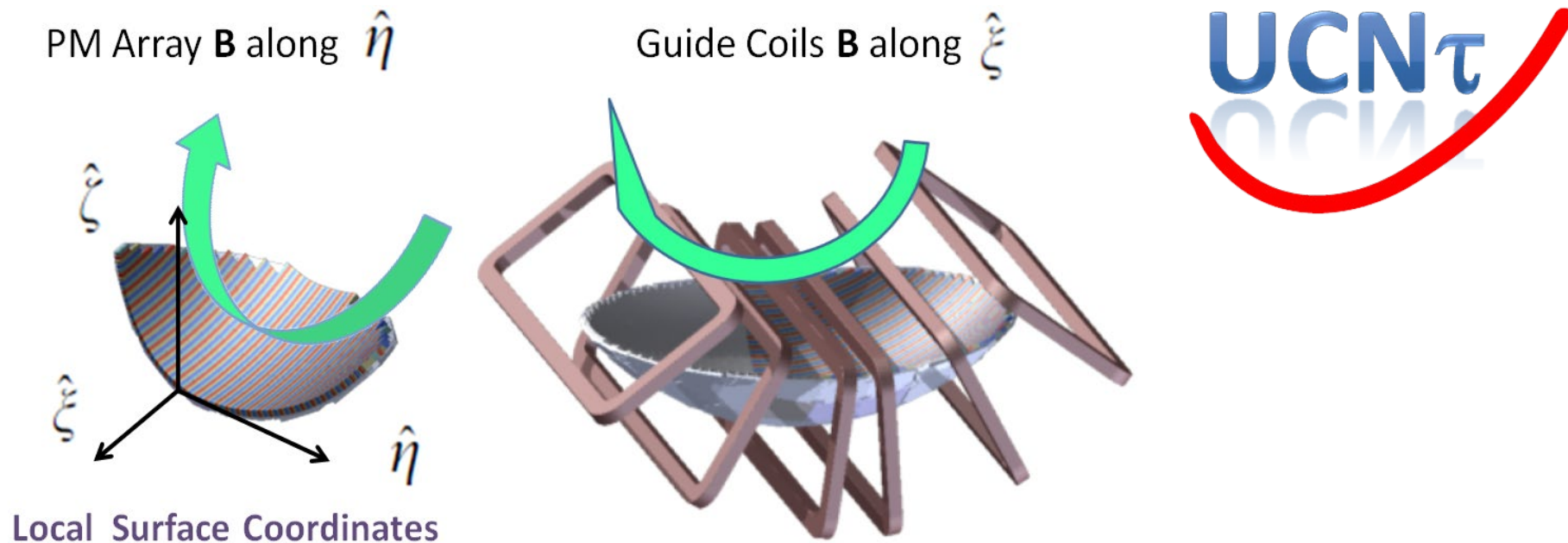


Neutron-wall interactions



A magneto-gravitational trap eliminates neutron losses on the walls

- **Magnetic trapping:** Halbach array of permanent magnets along trap floor repels spin polarized neutrons.
- **Minimize UCN spin-depolarization loss:** EM Coils arranged on the toroidal axis generates holding **B** field throughout the trap (perpendicular to the Halbach array field).



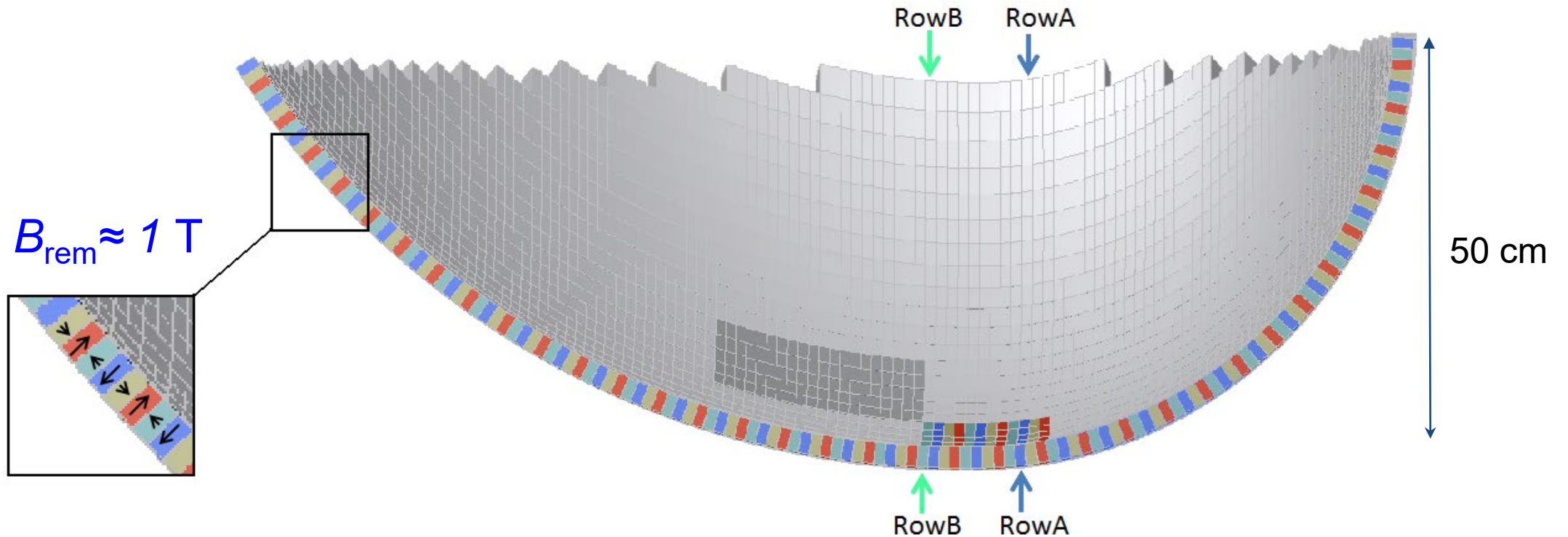
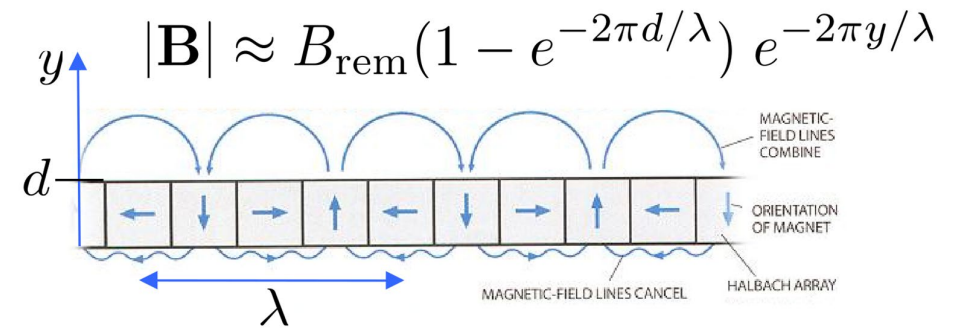
The UCN τ “Halbach” array

DESIGN OF PERMANENT MULTIPOLE MAGNETS
WITH ORIENTED RARE EARTH COBALT MATERIAL*

K. HALBACH

University of California, Lawrence Berkeley Laboratory, Berkeley, CA 94720, U.S.A.

Received 20 August 1979



Bailey inside the Halbach array performing field mapping (before Christmas 2012)

← Tweet



Chris Hadfield 
@Cmdr_Hadfield



That's Bailey in a neutron bottle. She helped discover that neutrons in the wild last 14.629 minutes (in an atom they can last billions of years).

[@LosAlamosNatLab](#)

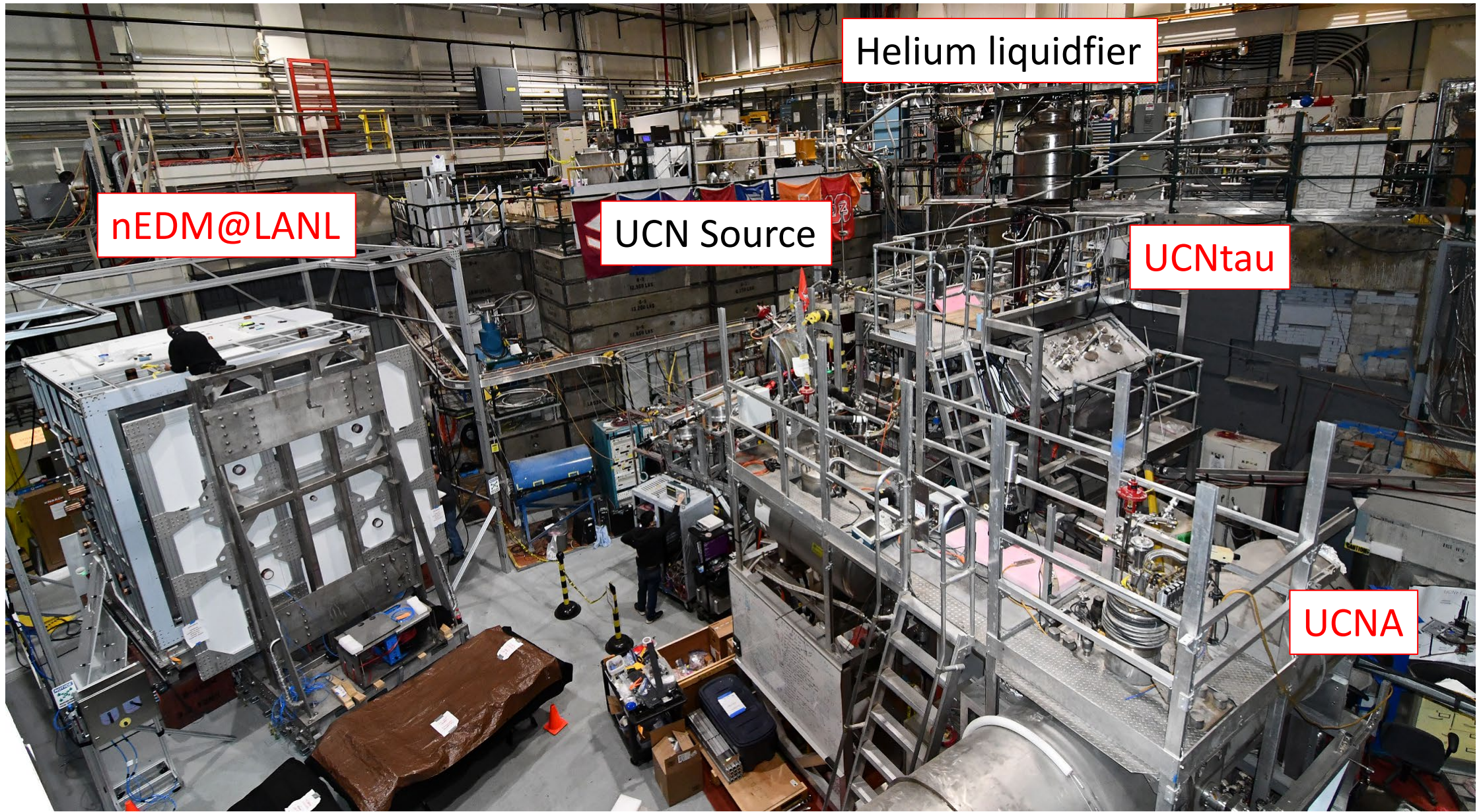
The details: bit.ly/3mBp5Tm



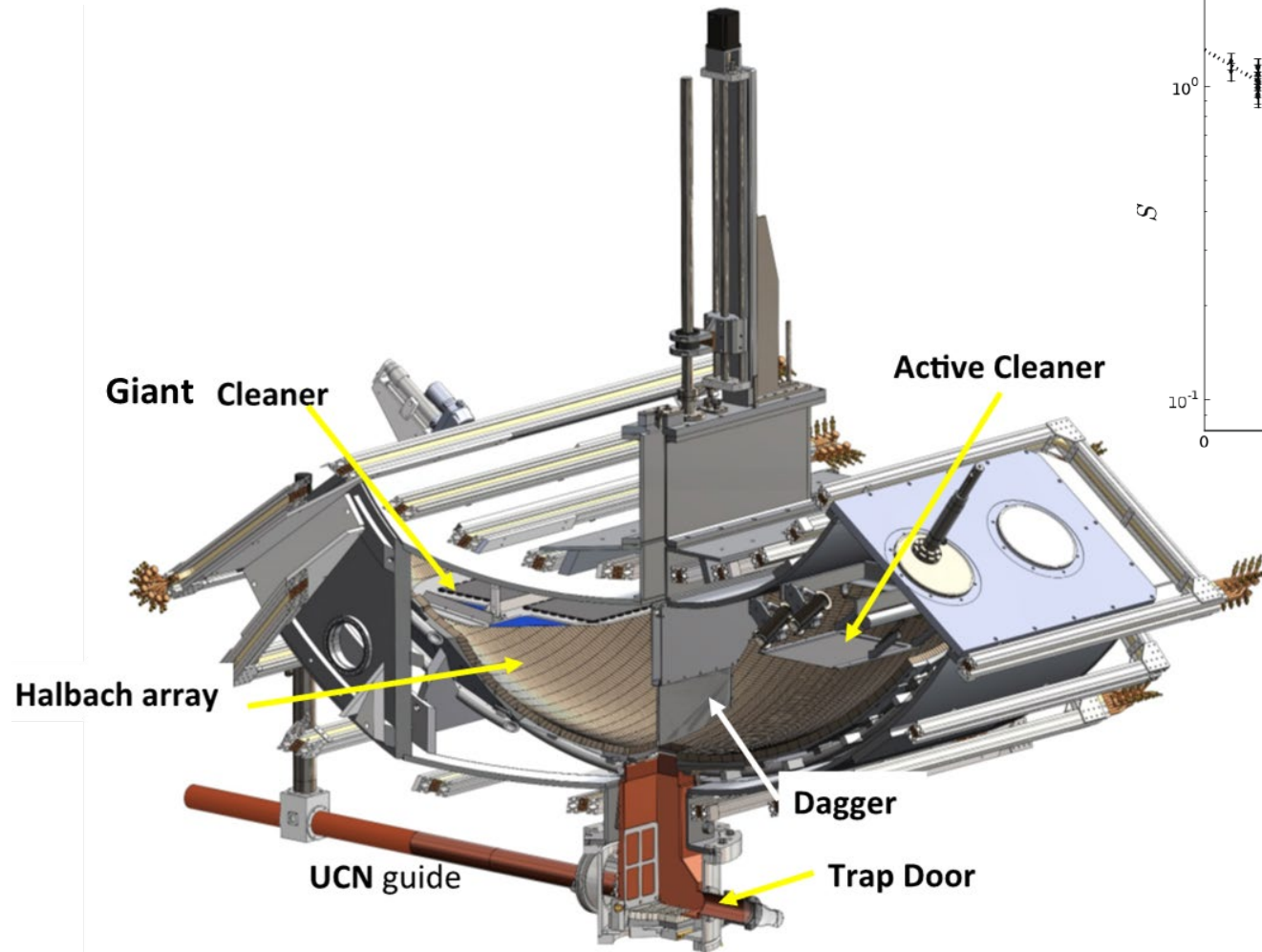
9:20 AM · Nov 2, 2021 · Twitter Web App

Los Alamos Neutron Science Center (LANSCE)

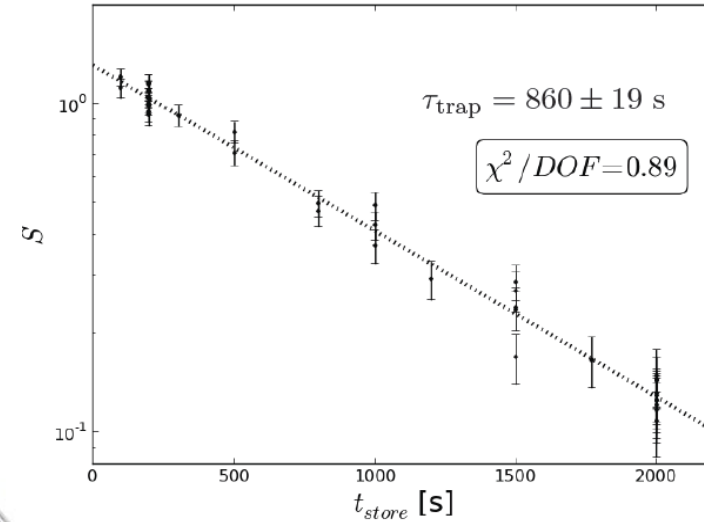
UCN Experimental Area (2021)



The UCN τ Apparatus

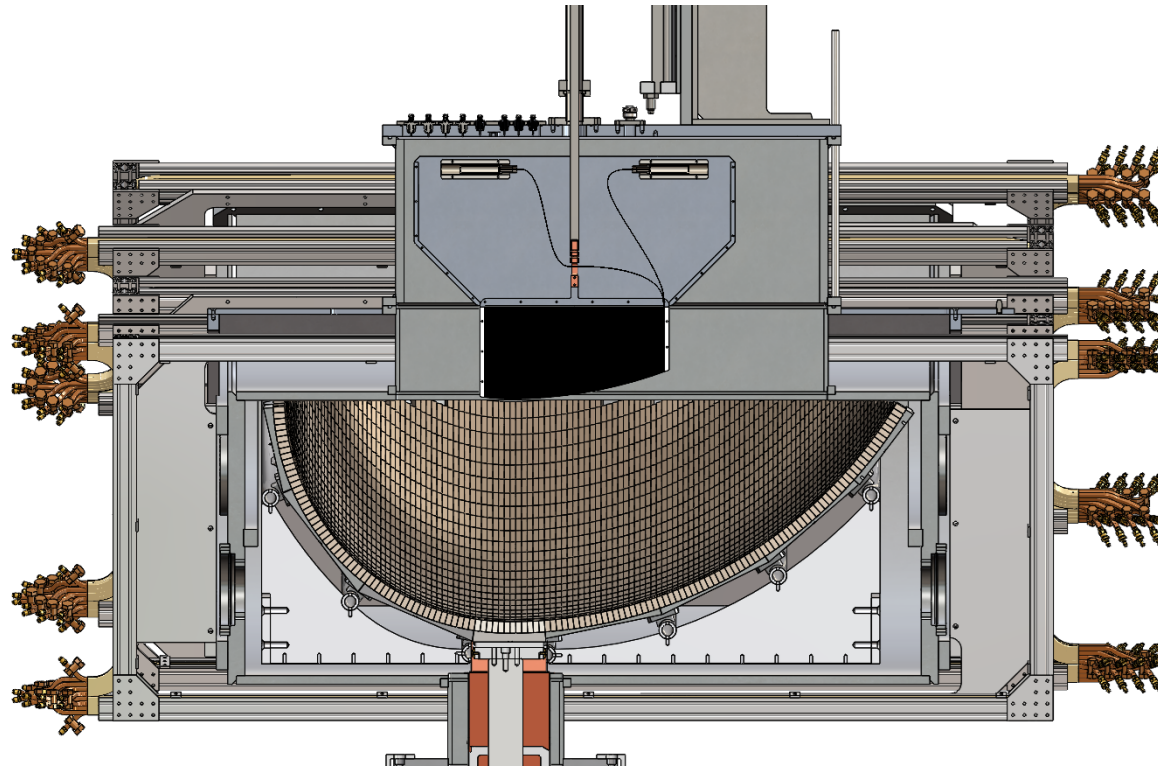


First Physics Data: 2013

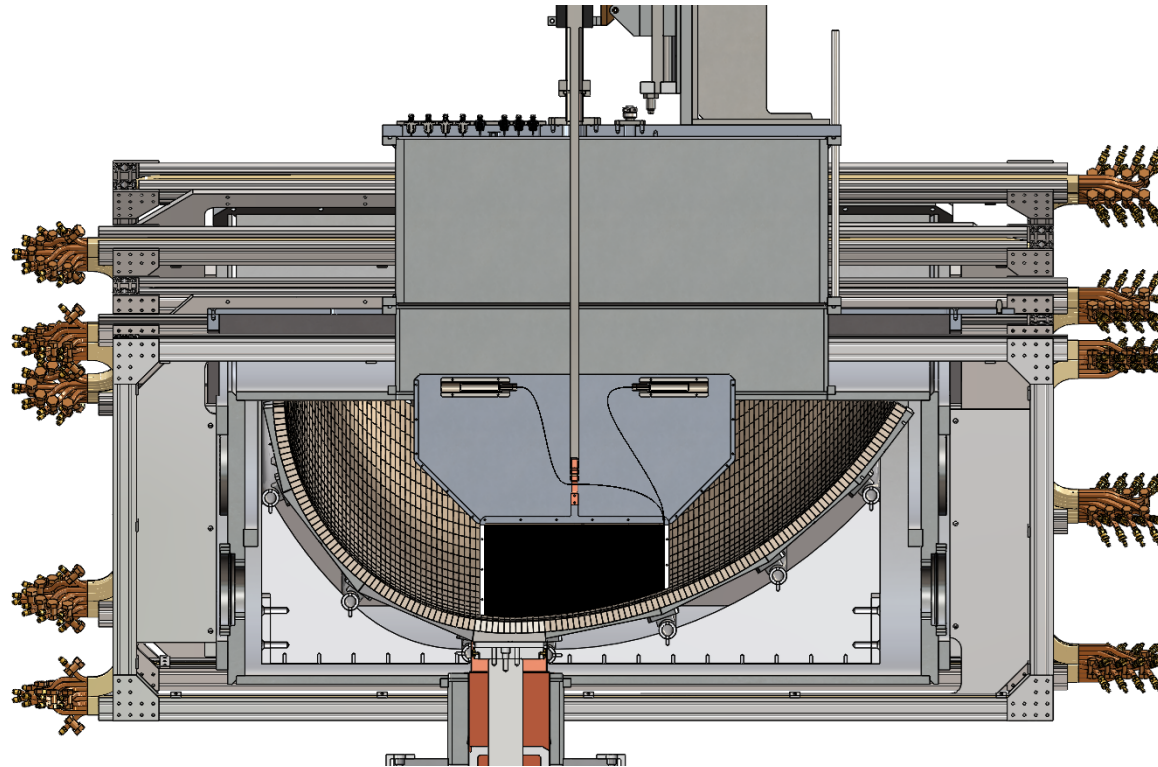


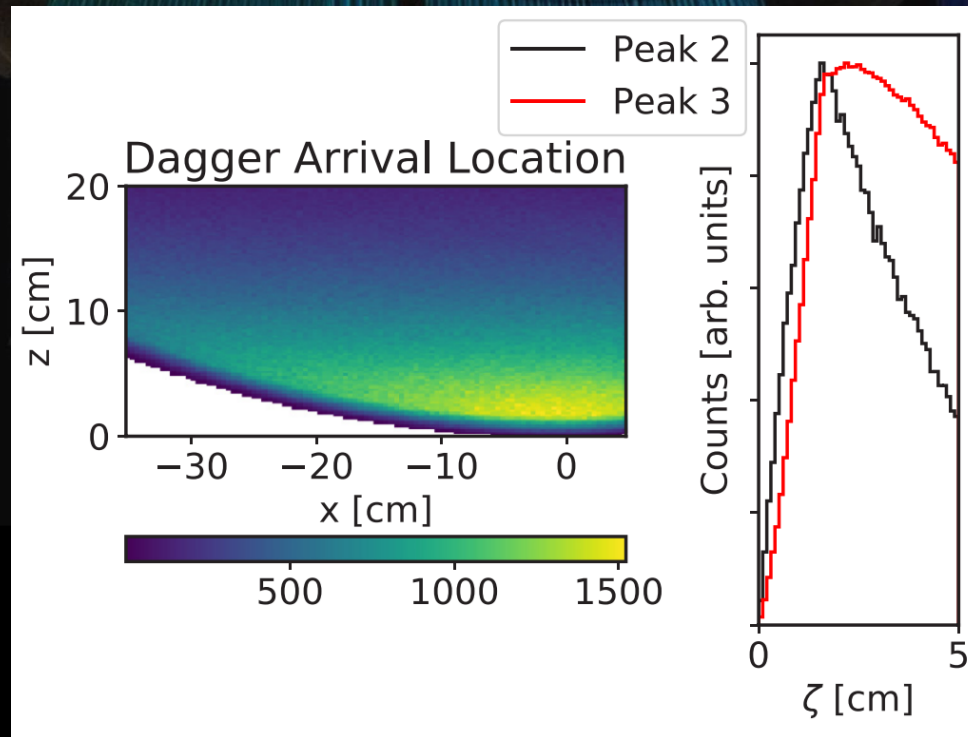
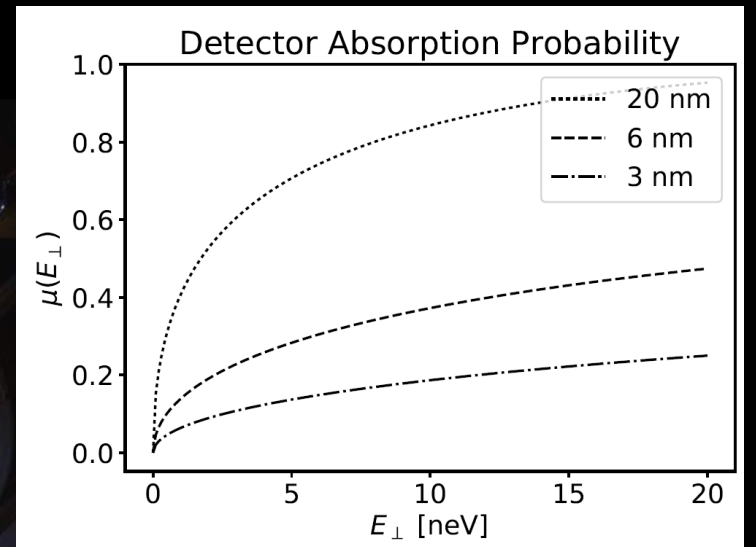
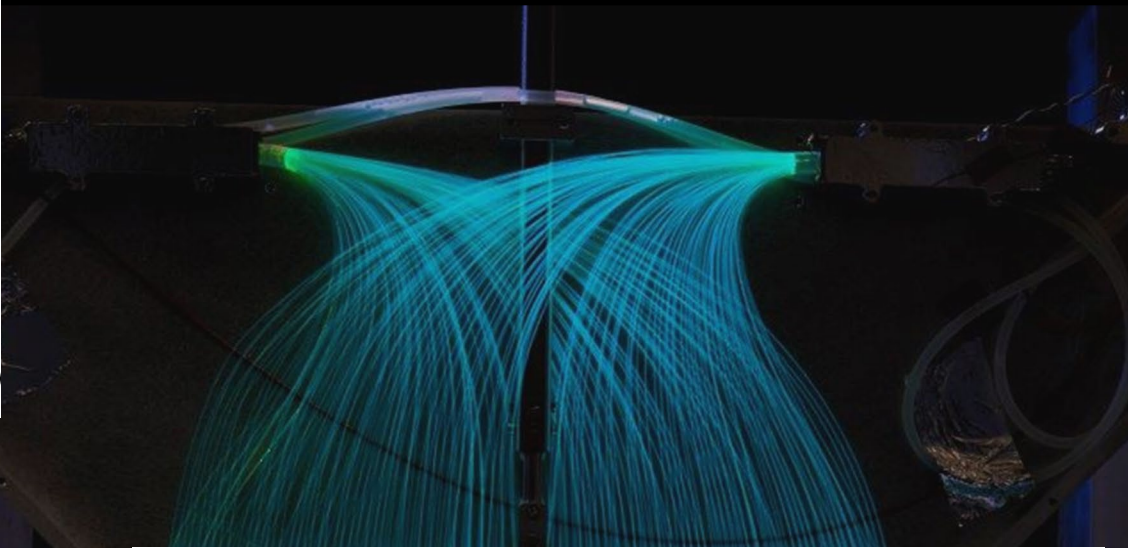
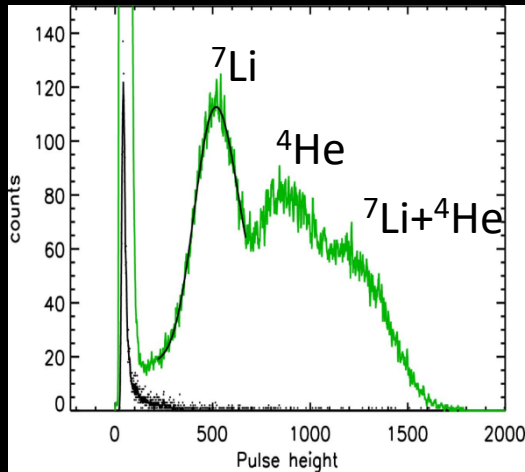
D. Salvat, PRC 89, 052501 (2014)

We also implemented a new way to count the trapped neutrons:



In-situ UCN detection using a “dagger” detector:
detection time ~ 8 s



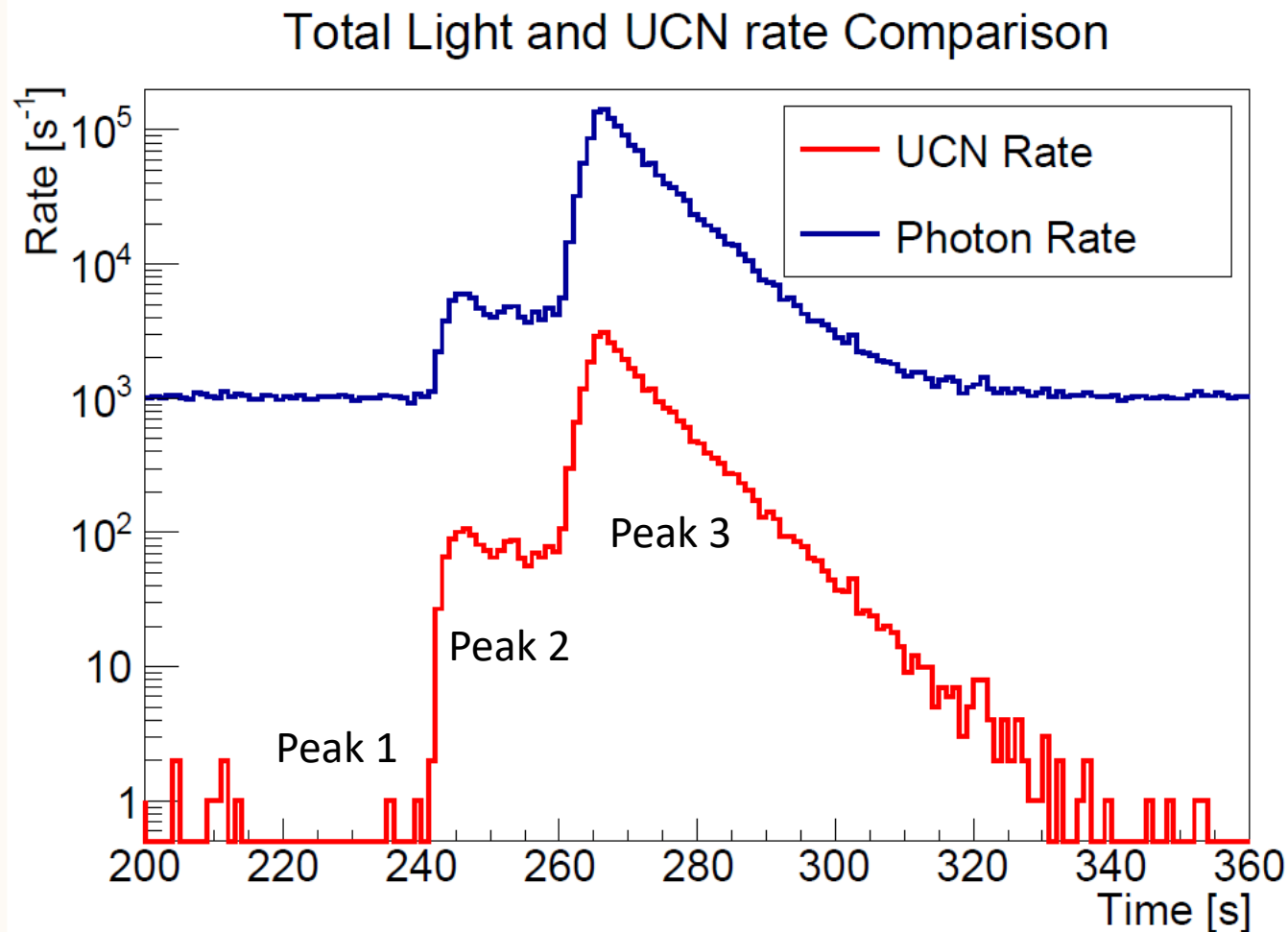


The response function (single collision) of the in-situ detector is highly non-uniform.

It takes ~ 5 collisions before a UCN is detected.

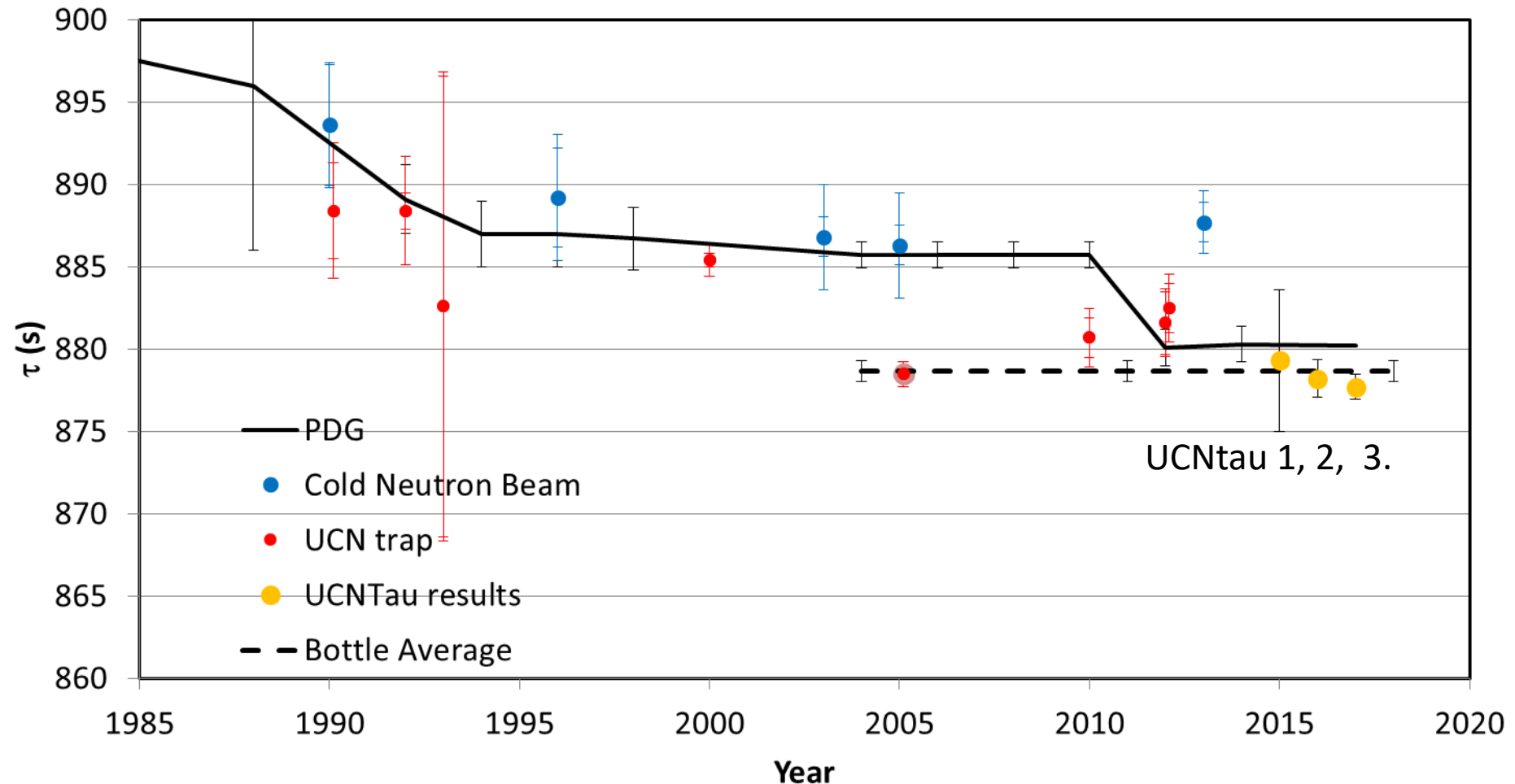
Singles vs Coincidence

- Singles - High Background (100-400 s correction)
- Coincidence - Pileup/Deadtime (as high as 3 s)



UCNtau results (2018)

1. 2015 commission data (RSI)
2. 2015-2016 data
3. 2016-2017 data (Science, 2018)



With UCNtau, we have made a measurement of τ_n for the first time with **no extrapolation**: 877.7 ± 0.7 (stat) $+0.3/-0.1$ (sys) s.

This gives confidence in previous bottle lifetime experiments

2019 Review of Particle Physics.

M. Tanabashi *et al.* (Particle Data Group), Phys. Rev. D **98**, 030001 (2018) and 2019 update.

n MEAN LIFE

INSPIRE search

Limits on lifetimes for *bound* neutrons are given in the section "p PARTIAL MEAN LIVES."

We average seven of the best eight measurements, those made with ultracold neutrons (UCN's). If we include the one in-beam measurement with a comparable error (YUE 2013), we get 879.6 ± 0.8 s, where the scale factor is now 2.0.

For a recent discussion of the long-standing disagreement between in-beam and UCN results, see CZARNECKI 2018 (Physical Review Letters 120 202002 (2018)). For a full review of all matters concerning the neutron lifetime until about 2010, see WIETFELDT 2011, F.E. Wietfeldt and G.L. Greene, "The neutron lifetime," Reviews of Modern Physics 83 1173 (2011).

| VALUE (s) | DOCUMENT ID | TECN | COMMENT | |
|---|--|------|---------|---------------------------|
| 879.4 ± 0.6 | OUR AVERAGE Error includes scale factor of 1.6. | | | |
| 878.3 ± 1.6 ± 1.0 | EZHOV | 2018 | CNTR | UCN magneto-gravit. trap |
| 877.7 ± 0.7 ^{+0.4} _{-0.2} | 1 PATTIE | 2018 | CNTR | UCN asym. magnetic trap |
| 881.5 ± 0.7 ± 0.6 | SEREBROV | 2018 | CNTR | UCN gravitational trap |
| 880.2 ± 1.2 | 2 ARZUMANOV | 2015 | CNTR | UCN double bottle |
| 882.5 ± 1.4 ± 1.5 | 3 STEYERL | 2012 | CNTR | UCN material bottle |
| 880.7 ± 1.3 ± 1.2 | PICHLMAIER | 2010 | CNTR | UCN material bottle |
| 878.5 ± 0.7 ± 0.3 | SEREBROV | 2005 | CNTR | UCN gravitational trap |
| ••• We do not use the following data for averages, fits, limits, etc. ••• | | | | |
| 887.7 ± 1.2 ± 1.9 | 4 YUE | 2013 | CNTR | In-beam n , trapped p |
| 881.6 ± 0.8 ± 1.9 | 5 ARZUMANOV | 2012 | CNTR | See ARZUMANOV 2015 |
| 886.3 ± 1.2 ± 3.2 | NICO | 2005 | CNTR | See YUE 2013 |
| 886.8 ± 1.2 ± 3.2 | DEWEY | 2003 | CNTR | See NICO 2005 |
| 885.4 ± 0.9 ± 0.4 | ARZUMANOV | 2000 | CNTR | See ARZUMANOV 2012 |

New data 2017-2018 data sets, 3 Independent Analyses



Frank Gonzalez
(Indiana)

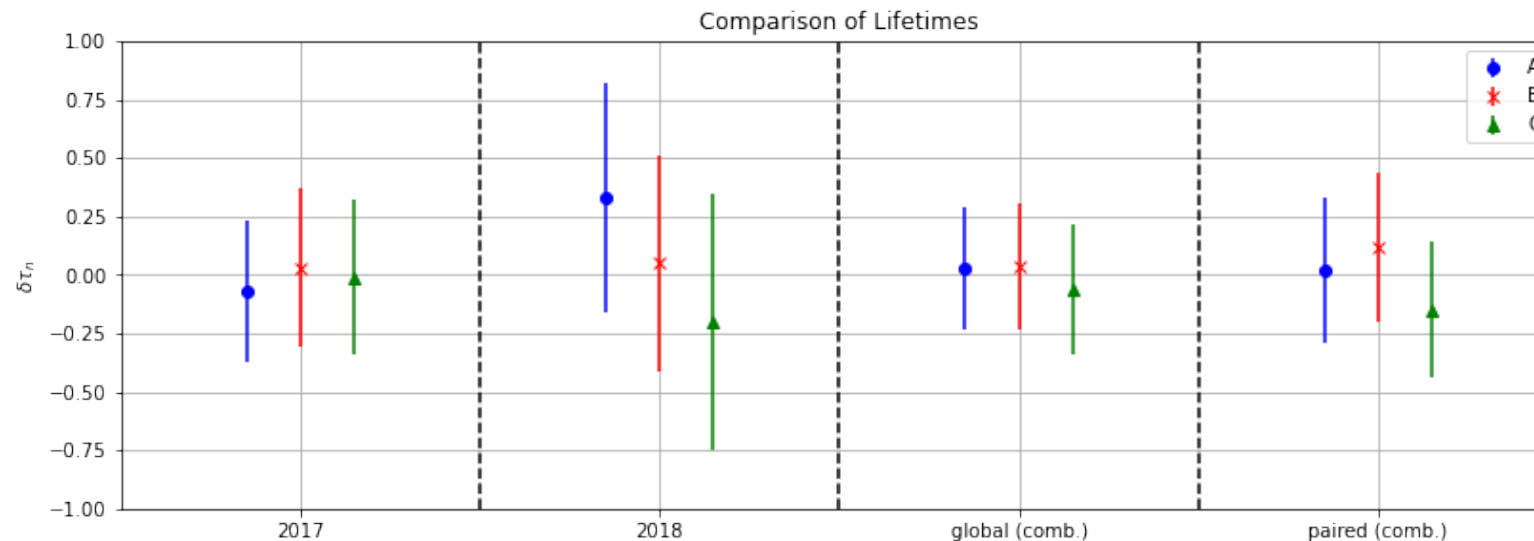


Eric Fries
(Caltech)



Chris Morris
(LANL)

- Blinded data:
 - Holding time is modified
 - Measured lifetime blinded by up to ± 15 s
- Unblinding Criteria:
 - Three complete (statistical and systematic) analyses
 - After cross-checking analyses, lifetimes combined via unweighted average, using largest uncertainties

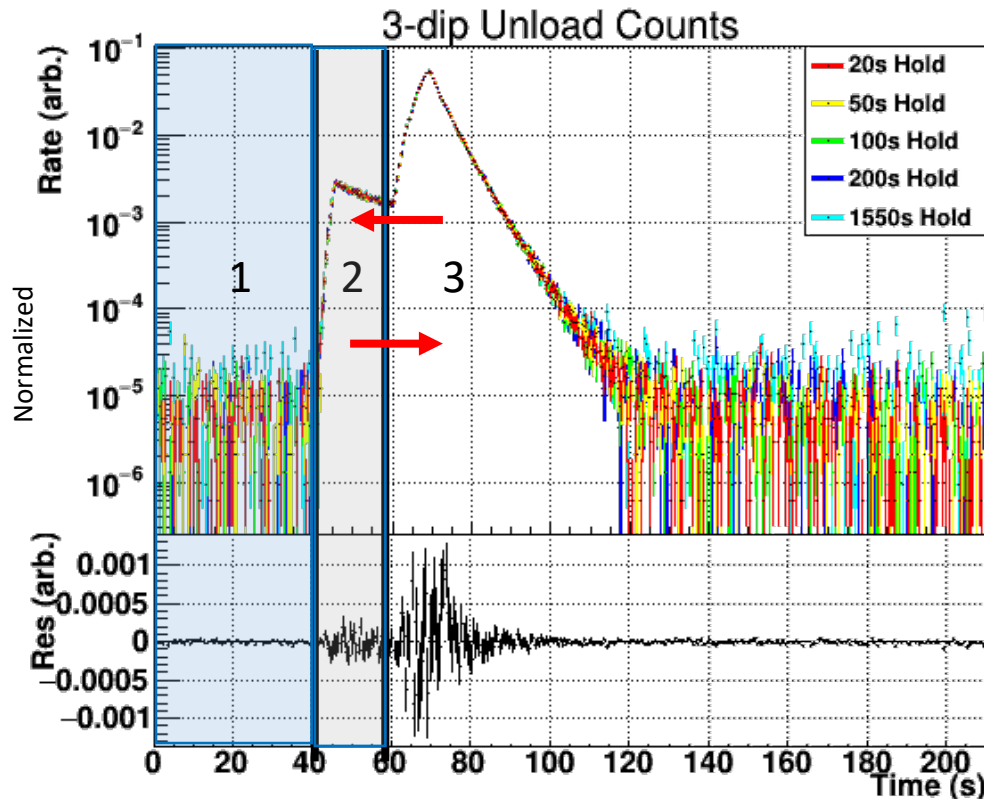


New Result: $\tau_n = 877.75 \pm 0.28^{+0.22}_{-0.16}$ s

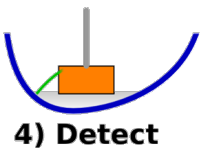
| Effect | Previous Reported Value (s) | New Reported Value (s) | Notes |
|-------------------------|---|---|---|
| τ_{meas} | 877.5 ± 0.7 | 877.58 ± 0.28 | Uncorrected Value! |
| UCN Event Definition | 0 ± 0.04 | 0 ± 0.13 | Single photon analysis vs. Coincidence analysis |
| Normalization Weighting | -- | 0 ± 0.06 | Previously unable to estimate |
| Depolarization | $0 + 0.07$ | $0 + 0.07$ | |
| Uncleaned UCN | $0 + 0.07$ | $0 + 0.11$ | |
| Heated UCN | $0 + 0.24$ | $0 + 0.08$ | |
| Phase Space Evolution | 0 ± 0.10 | -- | Now included in stat. uncertainty |
| Al Block | -- | 0.06 ± 0.05 | Accidentally dropped into trap... |
| Residual Gas Scattering | 0.16 ± 0.03 | 0.11 ± 0.06 | |
| Sys. Total | $0.16^{+0.4}_{-0.2}$ | $0.17^{+0.22}_{-0.16}$ | |
| TOTAL | $877.7 \pm 0.7^{+0.4}_{-0.2}$ | $877.75 \pm 0.28^{+0.22}_{-0.16}$ | |



Quantify the systematic effects using the 3-Step unload scheme

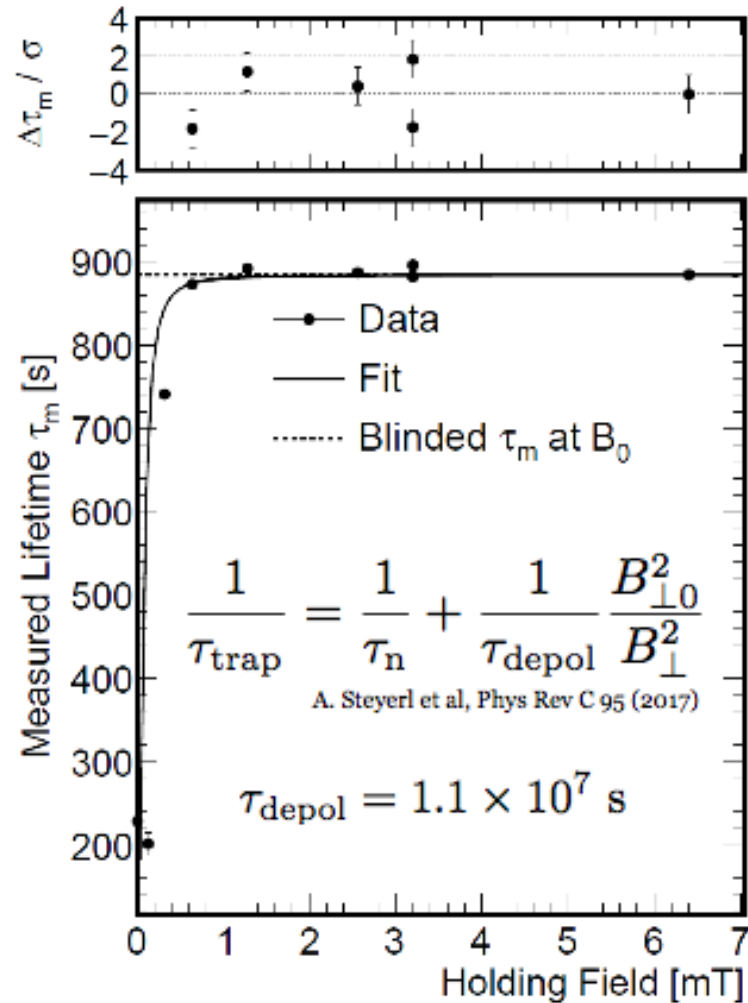


- Overthreshold UCN above cleaning height
 - “Heated” UCN (at long times): $\Delta\tau_{heat} = 0 + 0.08$ s
 - “Uncleaned” UCN (at short times): $\Delta\tau_{unc} = 0 + 0.11$ s
- Overthreshold neutrons $< 2 \times 10^{-5}$!
- Require constant detection efficiency
 - Phase space evolution couples to counting time
 - Use mean arrival time during unload
 - $\Delta\tau_{PSE} = 0.02 \pm 0.01$ s



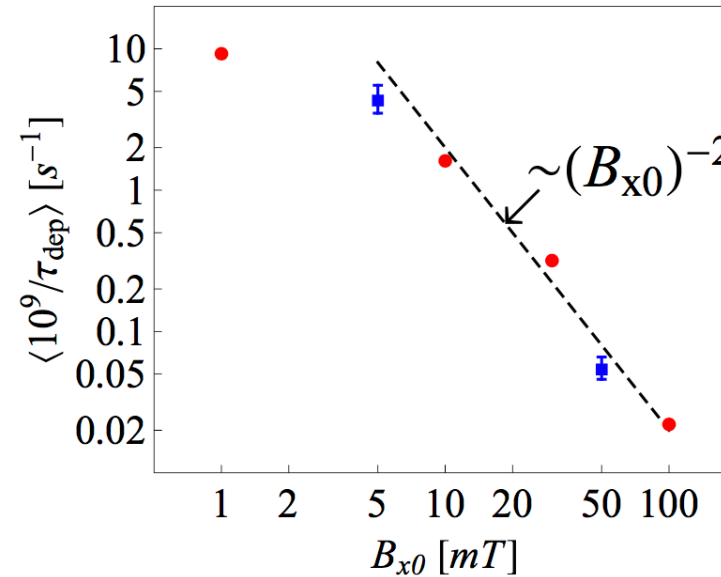
$$\tau_n^{-1} = \tau_{meas}^{-1} - \tau_{loss}^{-1}$$

UCN Depolarization

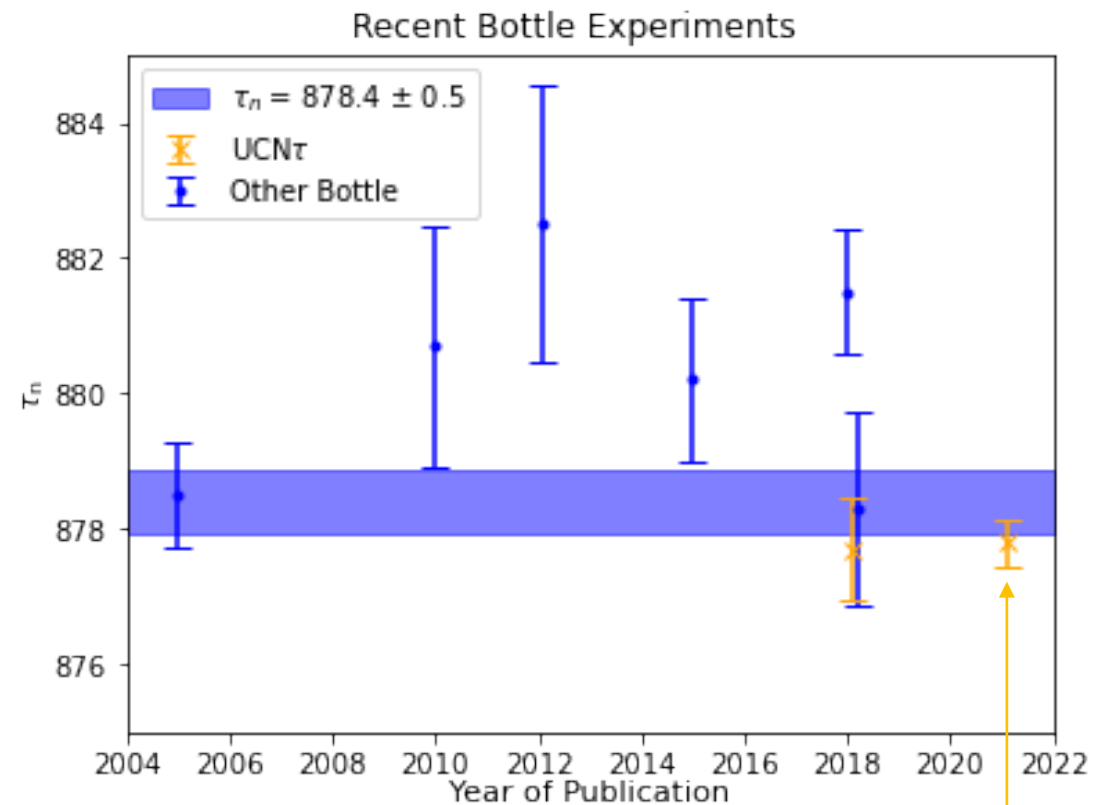
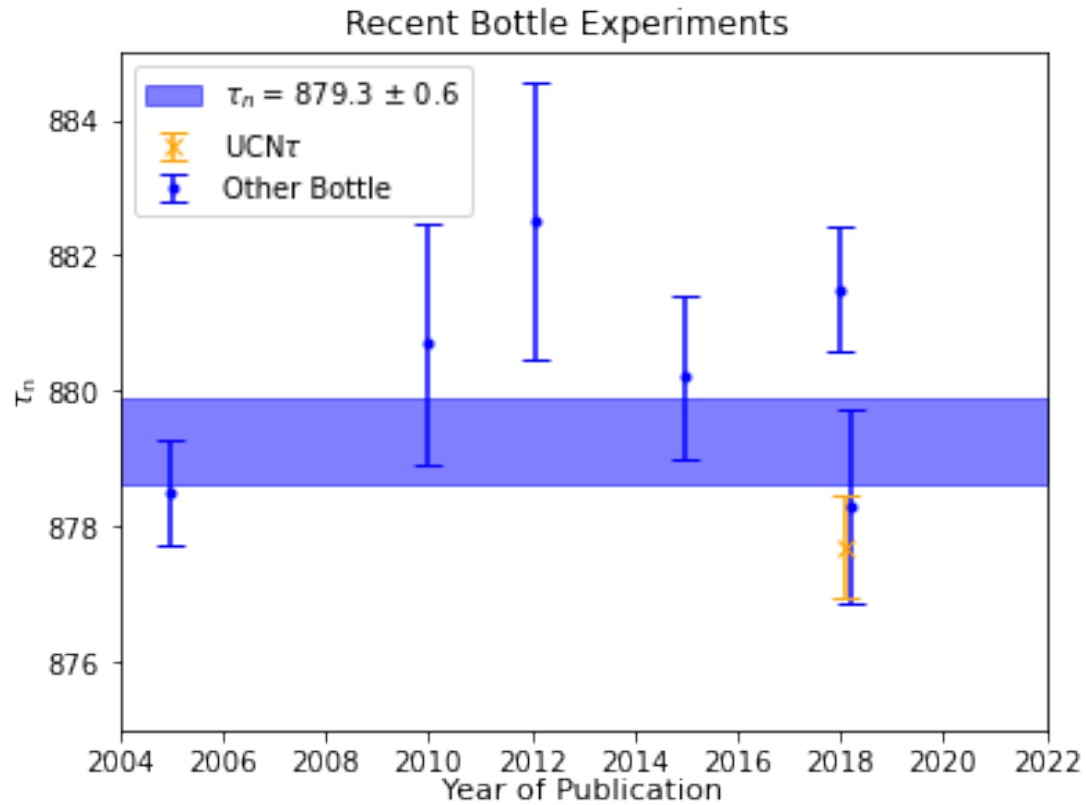


Three things to check:

1. Scaling behavior at low holding field.
2. Ideal vs. actual magnetic field.
3. Where to assess spin flip probability.



A new lifetime result (2021):



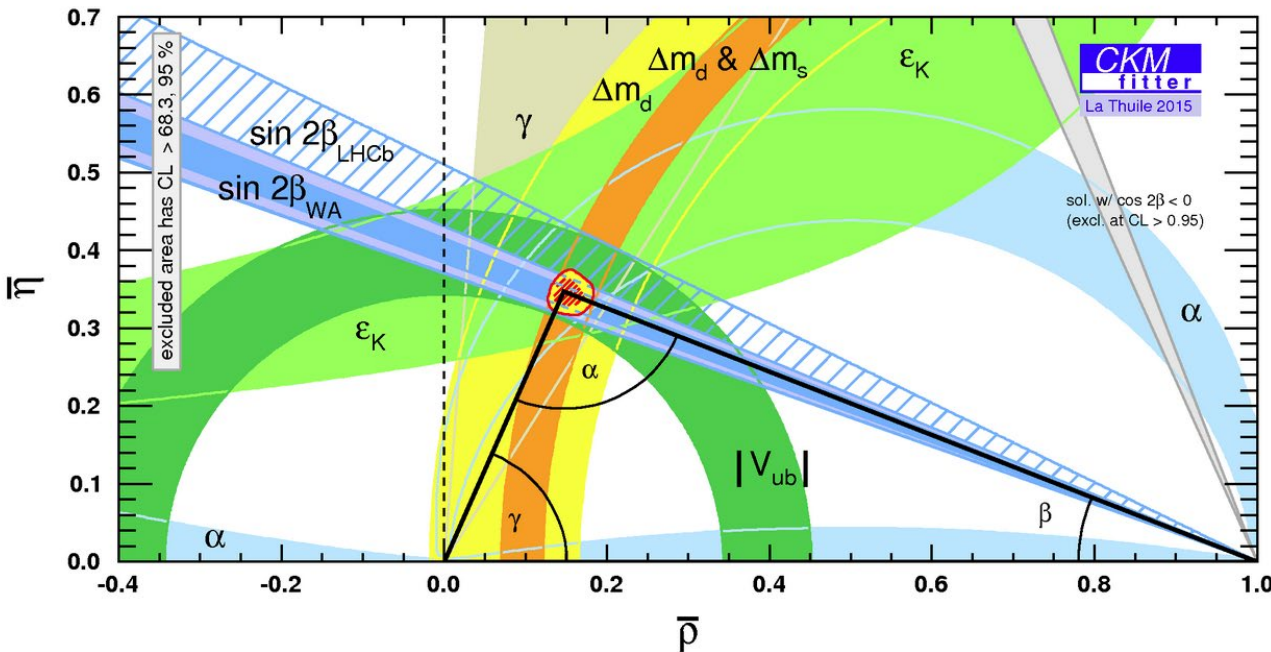
$$\tau_n = 877.75 \pm 0.28^{+0.22}_{-0.16} \text{ s}$$

The CKM matrix quantifies the quark flavor mixing.

$$\begin{pmatrix} d \\ s \\ b \end{pmatrix}_f = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}_m$$

Wolfenstein parameterization – expansion in $\lambda = \sin \theta_c \sim 0.22$

$$V = \begin{pmatrix} 1 - \frac{1}{2}\lambda^2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \frac{1}{2}\lambda^2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + \mathcal{O}(\lambda^4)$$



Area of the triangle \rightarrow CP-violating phase

Precision Test on the CKM Unitarity

First Row: $|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1 + \Delta_{BSM}$

$V_{ub} \ll V_{ud}$ and V_{us} , so negligible contribution

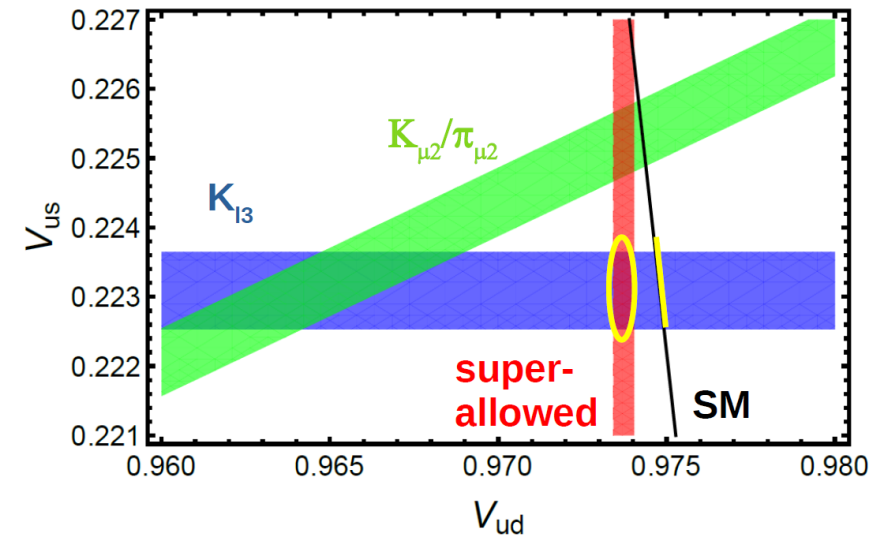
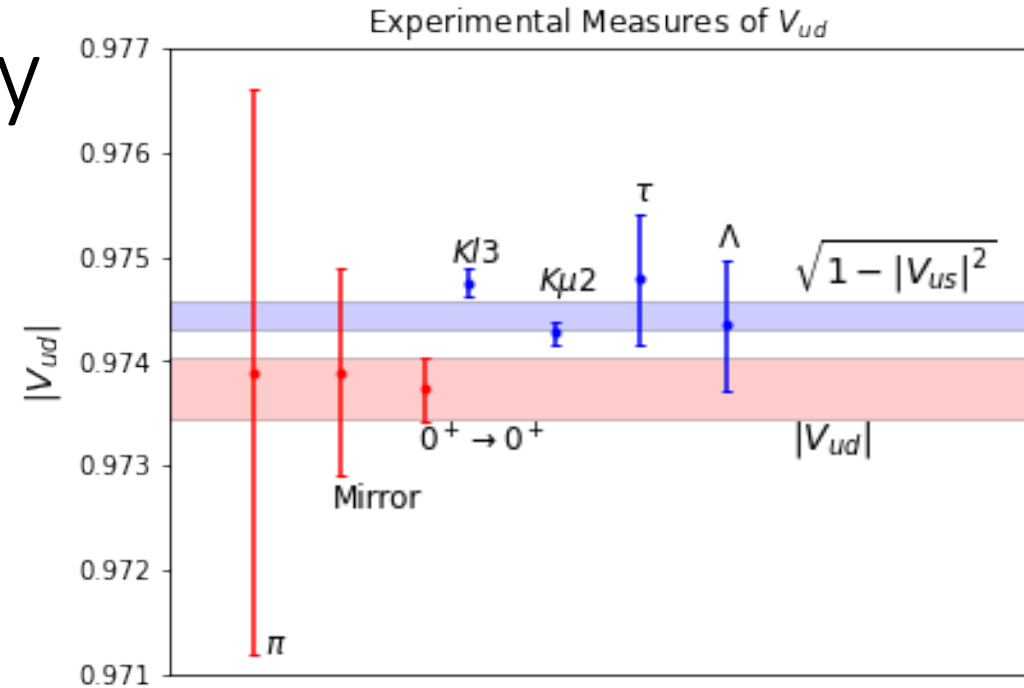
Measurements of V_{ud} :

- Most precise “Superallowed” $0^+ \rightarrow 0^+$ decays
- Mirror nuclei and Pions less precise
- Large theoretical uncertainties from radiative corrections and nuclear structure

Measurements of V_{us} :

- Most precise from Kaon decays
- Cabibbo angle anomaly ($V_{us} = \lambda = \sin \Theta_c$) between different decay channels
- Also limits from τ and Λ hyperons

Most precise measurements disagree (up to 3σ)!



“Cabibbo Angle Anomaly (CAA)” $\sim 3\sigma$

Discovery potential of the beta decay anomalies

A concrete example: First-row CKM unitarity with $|V_{ud}|$ from 0^+ beta decay and $|V_{us}|$ from $K_{\ell 3}$ decay

$$|V_{ud}|_{0^+}^2 + |V_{us}|_{K_{\ell 3}}^2 + \cancel{|V_{ub}|^2} - 1 = -0.0021(7)$$

SOURCES OF UNCERTAINTY:

$\delta|V_{ud}|_{0^+}^2$, **RC:**

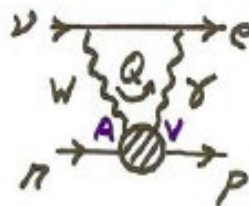
Theory uncertainties in the single-nucleon radiative corrections (RC)



| | |
|--|-----------------------|
| $ V_{ud} _{0^+}^2 + V_{us} _{K_{\ell 3}}^2 - 1$ | -2.1×10^{-3} |
| $\delta V_{ud} _{0^+}^2$, exp | 2.1×10^{-4} |
| $\delta V_{ud} _{0^+}^2$, RC | 1.8×10^{-4} |
| $\delta V_{ud} _{0^+}^2$, NS | 5.3×10^{-4} |
| $\delta V_{us} _{K_{\ell 3}}^2$, exp+th | 1.8×10^{-4} |
| $\delta V_{us} _{K_{\ell 3}}^2$, lat | 1.7×10^{-4} |
| Total uncertainty | 6.5×10^{-4} |
| Significance level | 3.2σ |

Extracting V_{ud} with neutron decays

f: Phase space factor=1.6886
(Fermi function, nuclear mass, size,
recoil)



$$1/\tau_n = f G_F^2 |V_{ud}|^2 m_e^5 (1+3g_A^2)(1+RC)/2\pi^3$$

From μ -decay: 0.6 ppm (MuLan 2011)

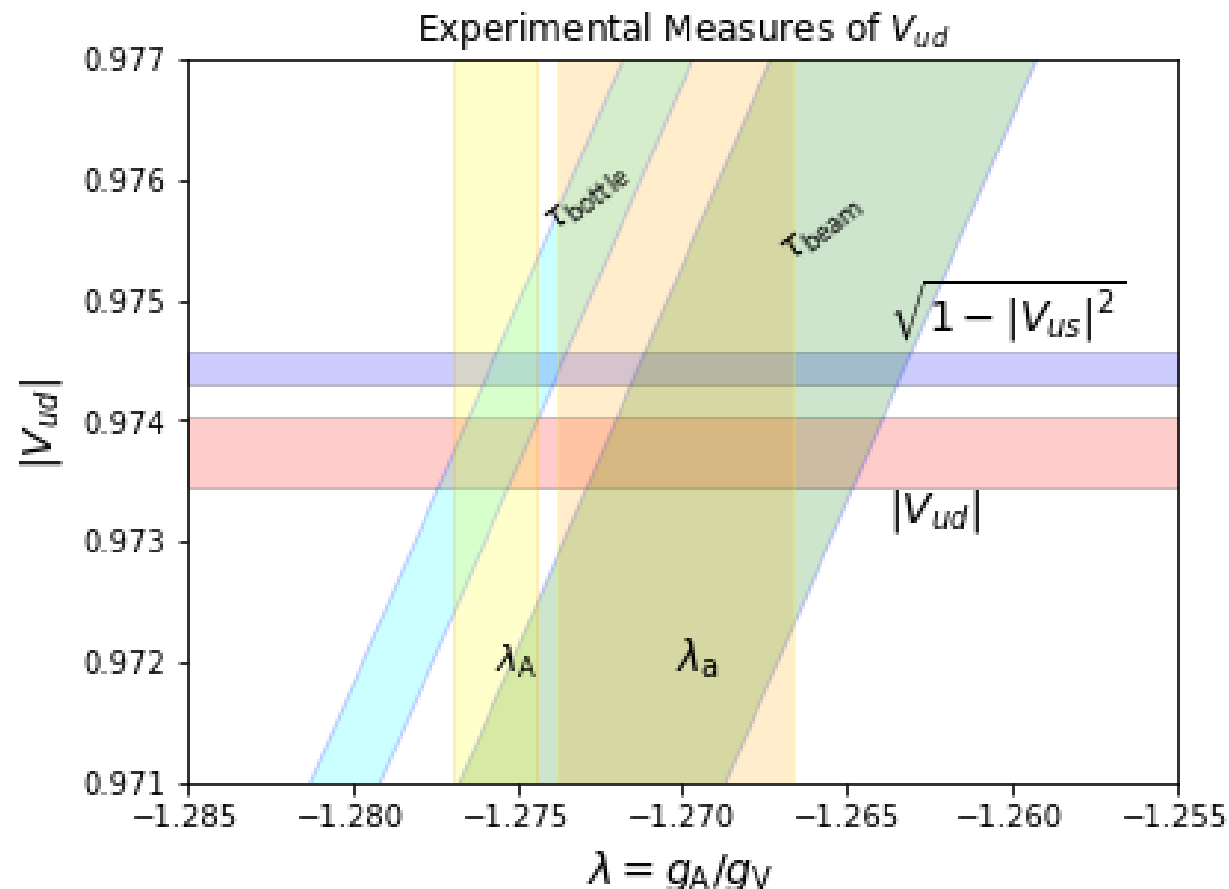
$$|V_{ud}|^2 = \frac{4905.7 \pm 1.7 \text{ s}}{\tau_n (g_V + 3g_A^2)}$$

Marciano & Sirlin, PRL 96, 032002 (2006)

Seng et al, PRL 121 (2018); Seng et al, PRD 100 (2019);

Czarnecki, Marciano & Sirlin, PRD 100 (2019)

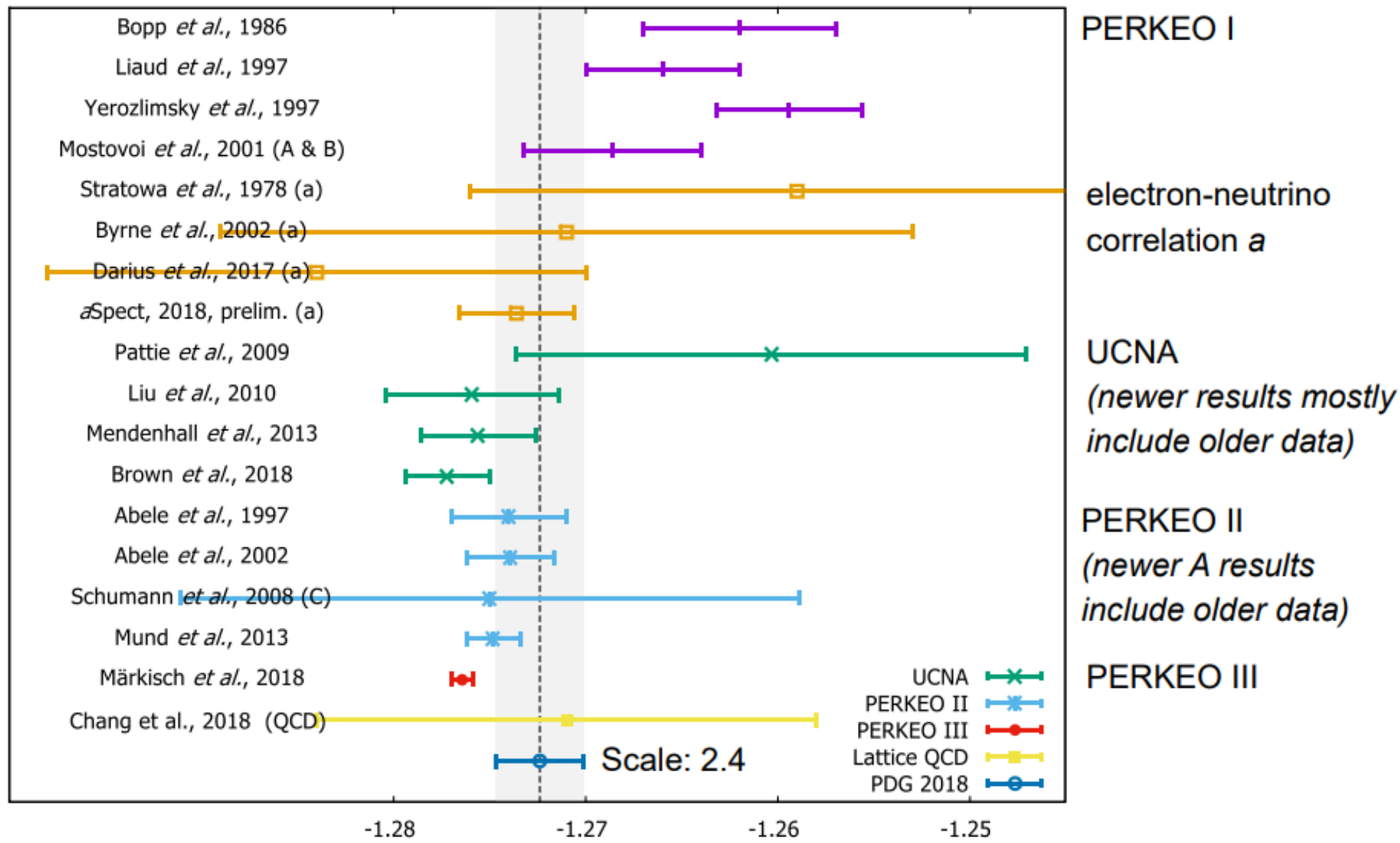
To match the theoretical uncertainty: 3.5×10^{-4} , it requires experimental uncertainties of: $\Delta A/A = 4\Delta\lambda/\lambda < 2 \times 10^{-3}$ and $\Delta\tau/\tau = 3.5 \times 10^{-4}$.



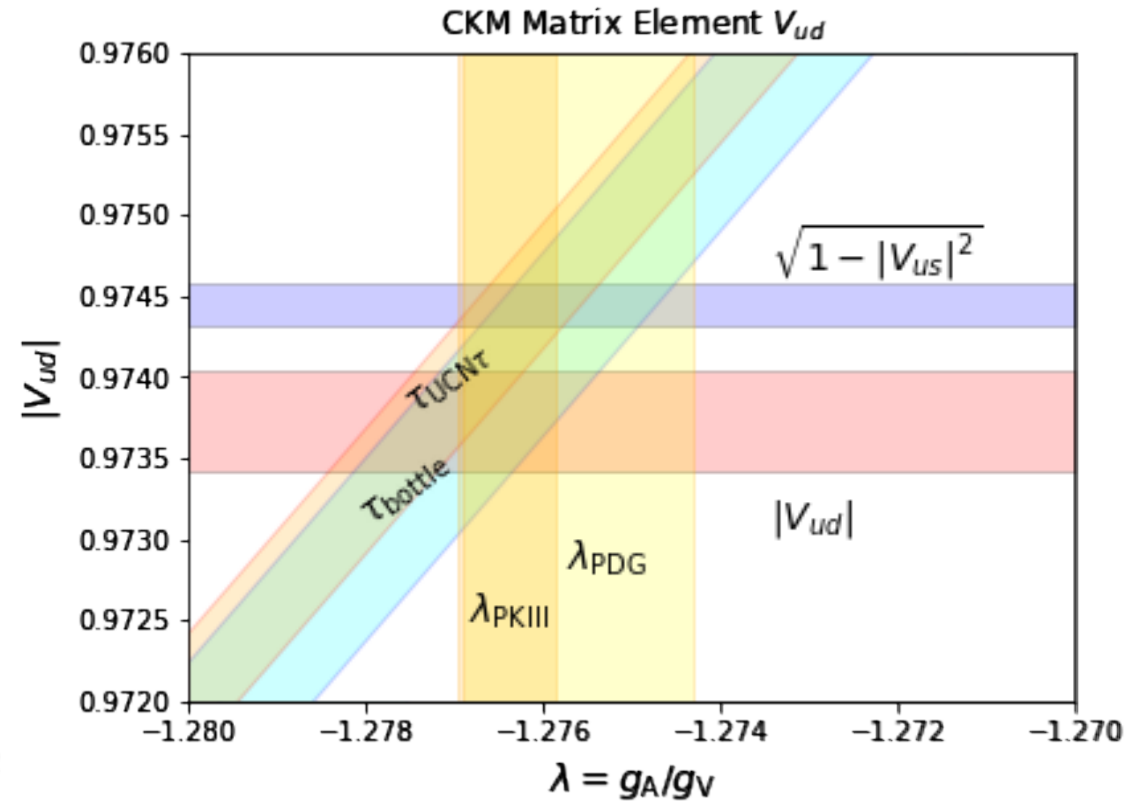
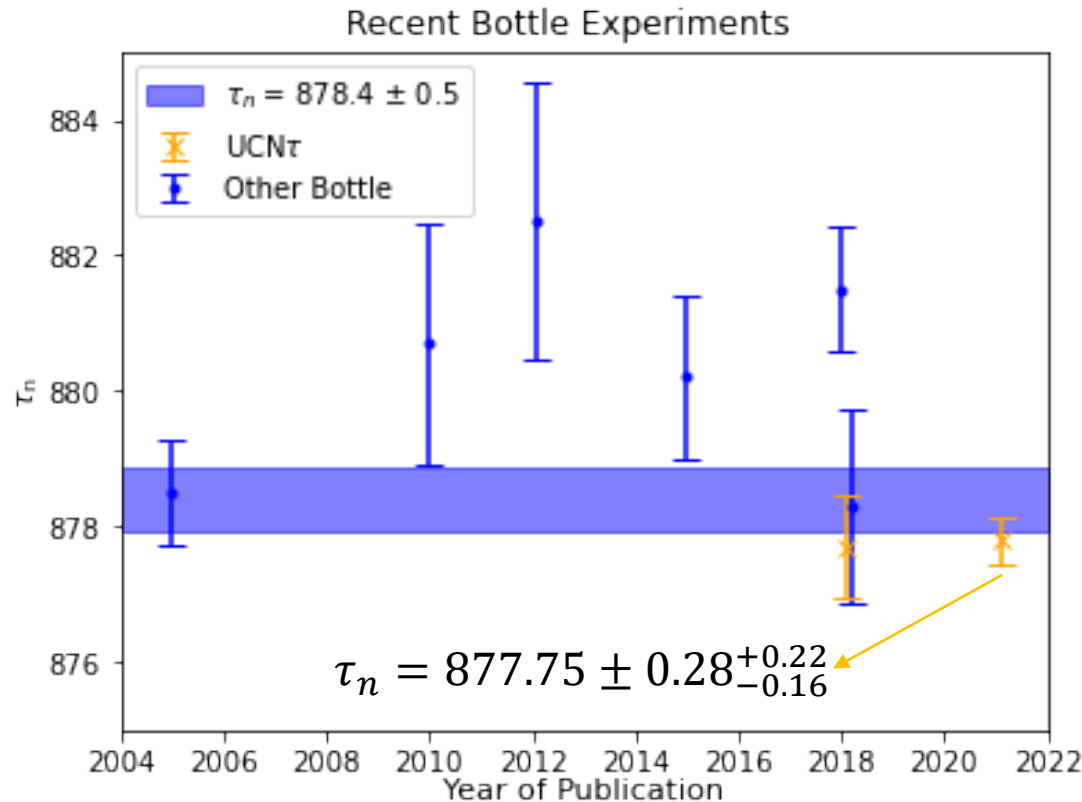
To be consistent with CKM unitarity, it requires a smaller $|g_A|$, or a shorter τ_n .

Axial Coupling: Status

Results from beta asymmetry A , unless where noted otherwise



With new UCNtau lifetime result (+ Perkeo III), the extracted V_{ud} agrees with the CKM unitarity.



We report a measurement of τ_n with 0.34 s (0.039%) uncertainty, improving upon our past results by a factor of 2.25 using two blinded datasets from 2017 and 2018. The new result incorporates improved experimental and analysis techniques over our previous result [Science **360**, 627 (2018)].

This is the first neutron lifetime measurement precise enough to confront SM theoretical uncertainties.

Summary

Low-energy neutrons are useful in testing the Standard Model of particle physics. Storage of UCN allows for the long observation times needed for precision measurement of many neutron observables. High-precision measurements, confronted with theoretical predictions, probe high-energy physics.

Precision measurements on the neutron lifetime ($\delta t < 0.1$ s), combined with the beta-decay asymmetry ($\delta A/A < 0.1\%$), test the unitarity of the CKM matrix (to $1e-4$ level of precision) and probe physics beyond the Standard Model. With UCN τ , all systematic uncertainties have been quantified by measurements.

- $\tau_n = 877.7 \pm 0.7 \begin{smallmatrix} +0.3 \\ -0.1 \end{smallmatrix}$ s (Science 2018)
- $\tau_n = 877.75 \pm 0.28 \begin{smallmatrix} +0.22 \\ -0.16 \end{smallmatrix}$ s (PRL 2021)

To be consistent with CKM unitarity, it requires a smaller $|g_A|$, or a shorter τ_n .

Discrepancy with CKM unitarity is an opportunity for new physics.

Moving forward:

- UCN τ + (immediate future): elevator loading, reaching $\delta t = 0.1$ s
- UCN τ 2 (future): superconducting coils (conceptual design), reaching $\delta t = 0.01$ s

The UCN τ Collaboration



Argonne National Laboratory

N. B. Callahan

California Institute of Technology

M. Blatnik, B. Filippone, E. M. Fries, K. P. Hickerson,
S. Slutsky, V. Su, X. Sun, C. Swank, W. Wei

DePauw University

A. Komives

East Tennessee State University

R. W. Pattie, Jr.

Indiana University/CEEM

M. Dawid, W. Fox, C.-Y. Liu, D. J. Salvat,
J. Vanderwerp, G. Visser

Institute Laue-Langevin

P. Geltenbort

Joint Institute for Nuclear Research

E. I. Sharapov

Los Alamos National Laboratory

S. M. Clayton (co-spokesperson), S. A. Currie,
M. A. Hoffbauer, T. M. Ito, M. Makela, C. L. Morris,
C. O'Shaughnessy, Z. Tang, W. Uhrich,
P. L. Walstrom, Z. Wang

North Carolina State University

T. Bailey, J. H. Choi, C. Cude-Woods, E.B. Dees,
L. Hayen, R. Musedinovic, A. R. Young, B. A. Zeck

Oak Ridge National Laboratory

L. J. Broussard, F. Gonzalez, J. Ramsey, A. Saunders

Tennessee Technological University

R. Colon, D. Dinger, J. Ginder, A. T. Holley (co-
spokesperson),
M. Kemp, C. Swindell

