

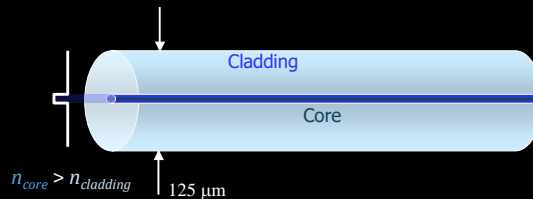
Microstructured optical fibres: a novel tool to make light interact with gas efficiently

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Group for Fibre Optics



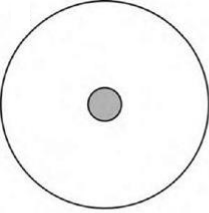
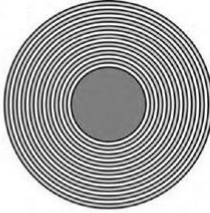
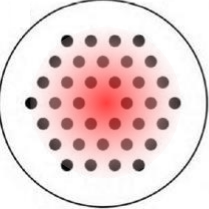
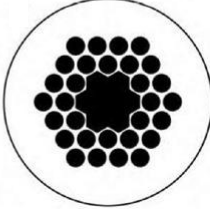
At the origin, there is the optical fibre...



- * Fibres are **small** and can be **seamlessly integrated** in a structure and the environment
- * Fibres are **chemically inert** and **electrically insulating**
- * Fibres are **extremely transparent** (50% light remaining after 15km)
- * Fibres **do not distort signals** (1 THz over 1000km)
- * Fibres are **green**: they are made from a very abundant and widespread raw material: **silica**



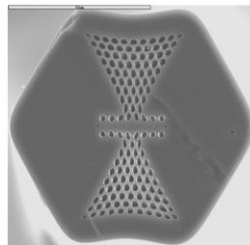
Different types of guiding have emerged

Index guiding		Antiresonant guiding	<ul style="list-style-type: none"> • Zero dispersion in the visible → Supercontinuum generation • Endlessly single mode fibres • Small mode area → Enhanced nonlinearities • Large mode area → Decreased nonlinearities • Guiding in air, fluids and vacuum → Absence of dispersion → Absence of nonlinearities → Interaction in gases & liquids • Prospect to overcome the loss floor of silica
			
Effective index guiding		Photonic bandgap guiding	
			

Fibres dedicated to pressure sensing

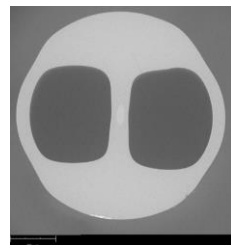
The isotropic hydrostatic pressure causes an asymmetric deformation of a fibre designed with an anisotropic profile → **Birefringence is modified by pressure**

Photonic crystal fibre designed for optimised response



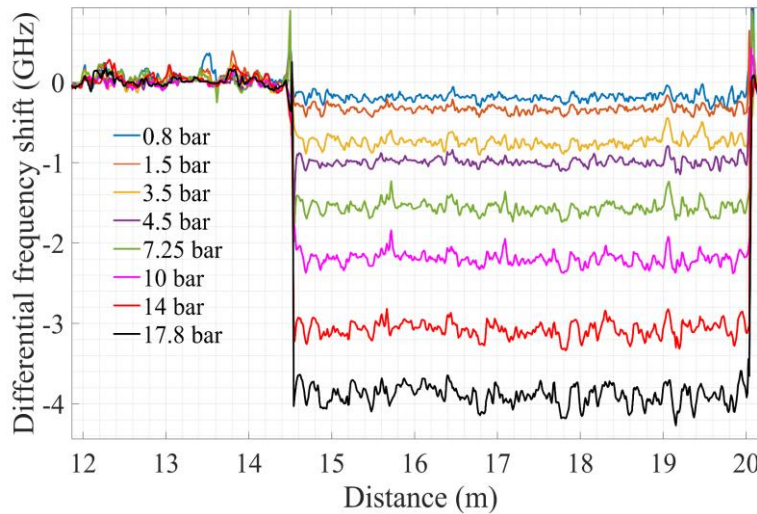
- Strong response
- High loss
- Poor uniformity
- High cost

Side-hole fibre designed for simplified fabrication

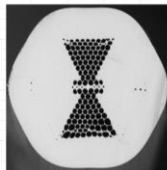
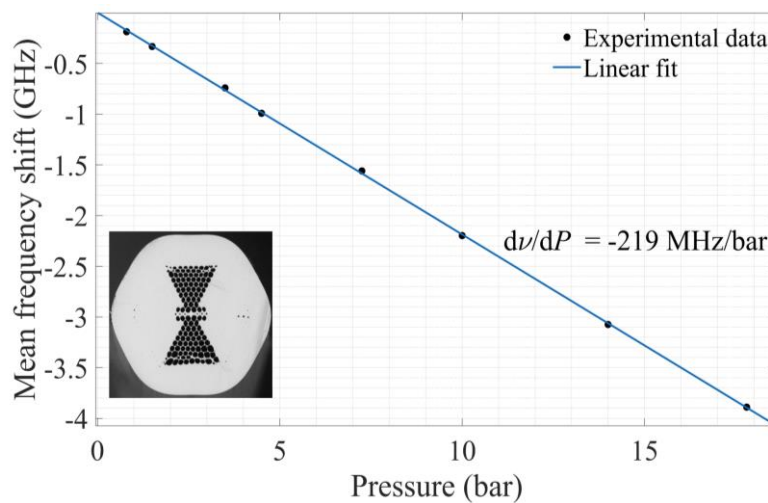


- Good response
- Low loss
- Good uniformity
- Low cost

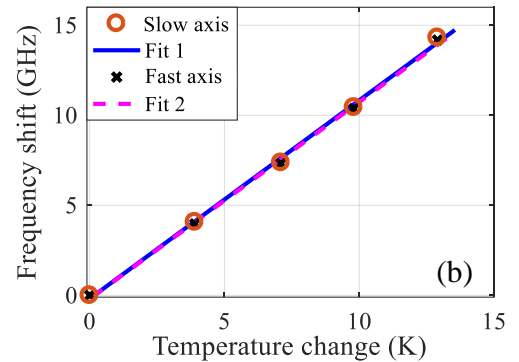
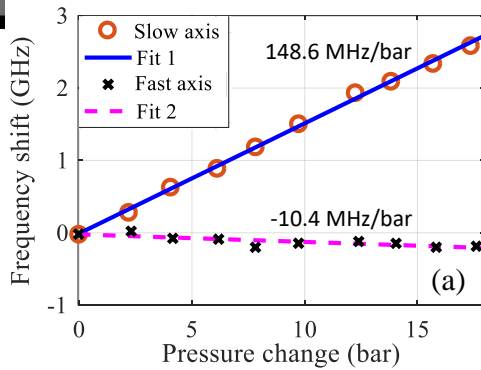
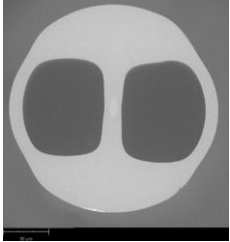
Distributed birefringence change vs Pressure



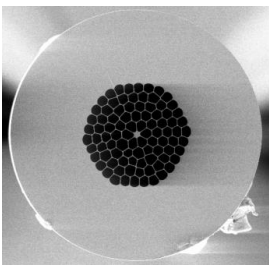
Distributed birefringence change vs Pressure



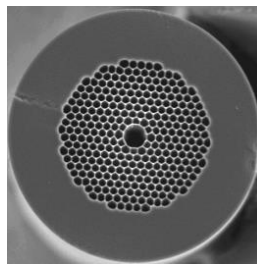
Pressure and temperature response along the 2 polarisations



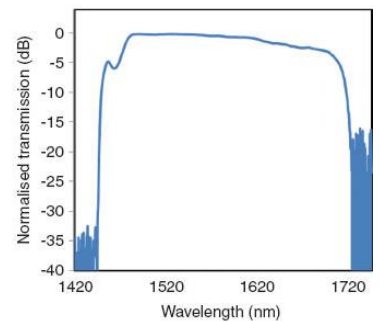
Index Guiding vs Photonic Bandgap Guiding



Light is essentially in SiO_2

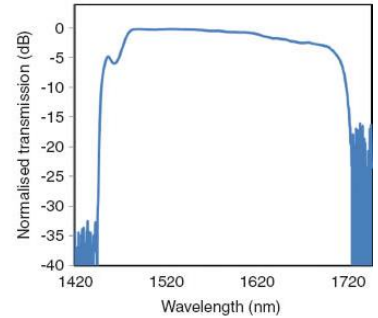
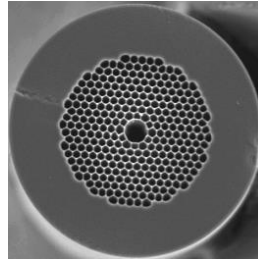
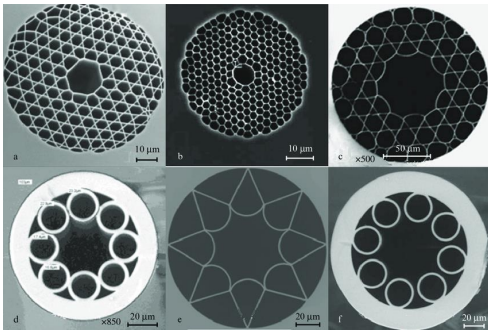


Light is essentially in air
 A tiny fraction propagates in SiO_2 ($\sim 1\%$)
 Loss limited by surface roughness



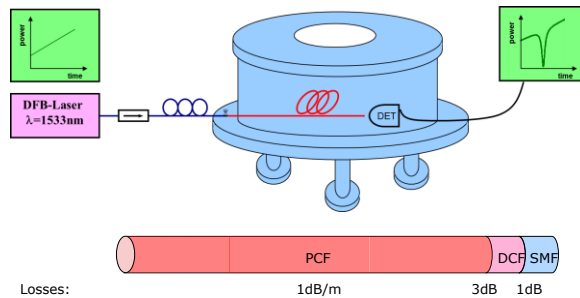
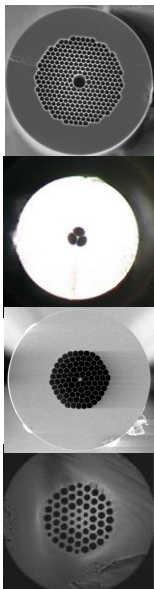
Index Guiding vs Photonic Bandgap Guiding

Structures have evolved to decrease the fraction of light in SiO_2



Light is essentially in air
A tiny fraction propagates in SiO_2 ($\sim 1\%$)
Loss limited by surface roughness

Preparation of all-fibre absorption cells



The arc-fusion splice requires **splicing at atmospheric pressure**, but this means that the **low-pressure gas** inside the PCF will be **contaminated by air**.

The **solution** is to use the **high permeability of helium through silica**: we insert high-pressure (~ 2 bar) helium before making the splice. Helium will diffuse out of the silica walls in 1-2 hours.

P. S. Light, F. Couny, and F. Benabid, "Low optical insertion-loss and vacuum-pressure all-fiber acetylene cell based on hollow-core photonic crystal fiber," Opt. Lett. **31**, 2538-2540 (2006)

Preparation of all-fibre absorption cells

The diagram illustrates the preparation of an all-fibre absorption cell. At the top, a DFB-Laser with a wavelength of $\lambda = 1533\text{nm}$ is connected to a fiber optic system. The system includes a circulator (C), a gas cell (G), and a detector (DET). Two graphs show power vs. time: one shows a linear increase, and the other shows a step-like increase. Below this, a schematic of the fiber assembly shows segments of SMF, DCF, PCF, DCF, and SMF. Losses are specified as 1dB for SMF, 4dB for DCF, 1dB/m for PCF, 3dB for DCF, and 1dB for SMF.

The preparation process is divided into three steps:

- STEP 1:** Fiber purging (HCPCF) and Active gas loading (C_2H_2).
- STEP 2:** He loading and Splicing.
- STEP 3:** He evacuation through permeation process and Gas cell.

P. S. Light, F. Couny, and F. Benabid, "Low optical insertion-loss and vacuum-pressure all-fibre acetylene cell based on hollow-core photonic crystal fiber," Opt. Lett. 31, 2538-2540 (2006)

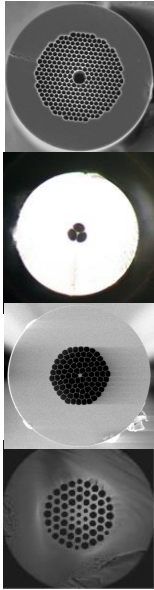
Preparation of all-fibre absorption cells

The diagram shows the fiber assembly schematic from the previous slide, with losses: 1dB (SMF), 4dB (DCF), 1dB/m (PCF), 3dB (DCF), and 1dB (SMF).

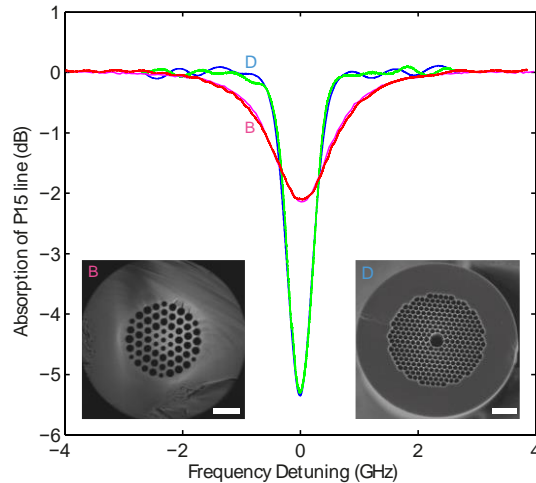
The transmission spectrum of the PCF is shown as a graph of Relative transmitted intensity (dB) vs. Wavelength (nm). The spectrum shows a series of periodic dips. Two specific dips are labeled P14 and P15. The x-axis ranges from 1525 to 1540 nm, and the y-axis ranges from 0 to -20 dB.

A photograph on the right shows the physical fiber assembly, which is a long, thin, clear fiber with connectors at both ends, resting on a green surface next to a coin for scale.

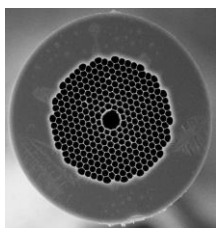
Preparation of all-fibre absorption cells



Long term response: Nov. 2008 – Sep. 2011

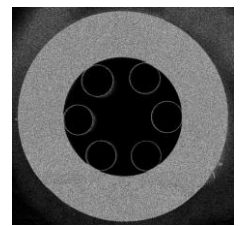


Clear interest for opto-acoustic interactions in gases



Hollow core fibres have a transparent and bright future:

- Will probably outperform silica fibres in term of loss (currently < 1 dB/km)
- Higher power handling capability (no nonlinear effect)
- No dispersion, broader spectral range
- Low latency



➔ **The possibilities for optical signal processing are much limited, accordingly!**

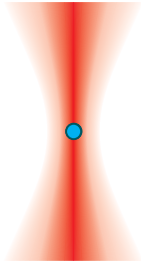
If **stimulated Brillouin scattering** could be activated in gas

➔ a wide choice of all-optical interactions could be implemented:

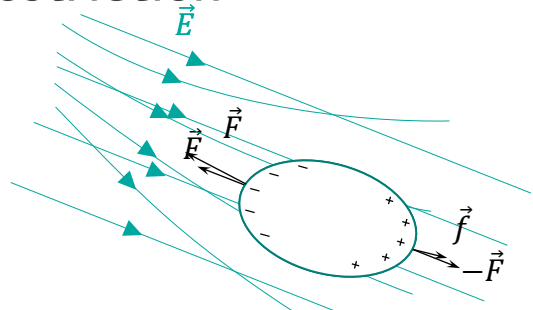
- Amplification & lasing
- Slow & fast light
- Selective spectral filtering
- Optical storage
- All-optical calculus
- Distributed sensing, etc...

If observable, will the interaction be large enough to offer any interest?

Polarisation force & Electrostriction



Optical trapping & tweezers
Nobel Prize 2018



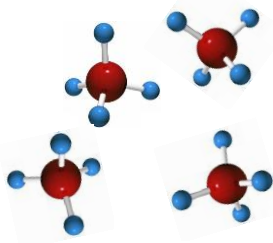
A dielectric object in a non-uniform field feels a force toward regions of higher field strength.



$$\vec{F} \sim \vec{E} \nabla \vec{E}$$

Induced charges Non-uniformity strength Force strength

The Clausius-Mossotti relation - Elasto-optic effect



$$\frac{n^2 - 1}{n^2 + 2} = N_A \frac{\alpha}{3M} \rho$$

with ρ : Material density
 N_A : Avogadro's constant
 α : Molecular polarisability
 M : Molar mass

For a tenuous medium like a gas ($n \sim 1$) :

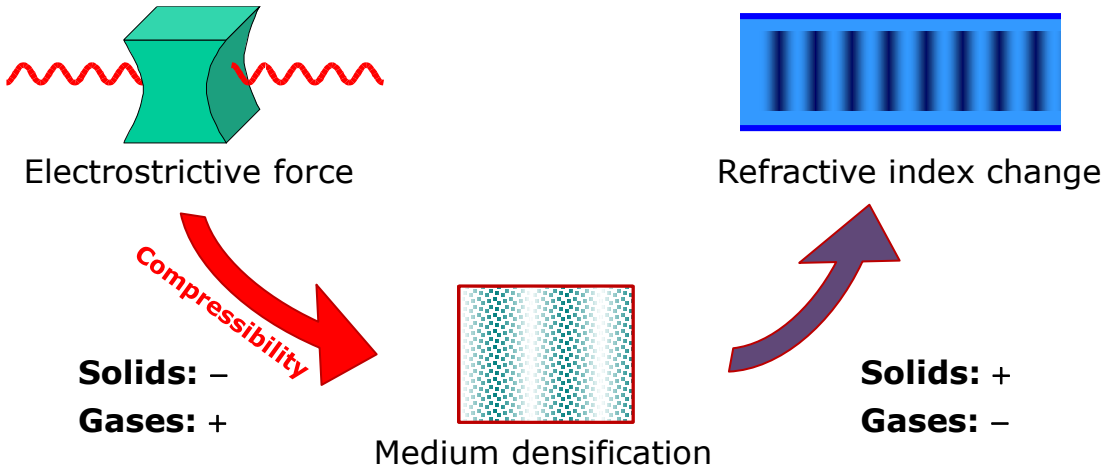
$$n \approx 1 + N_A \frac{\alpha}{2M} \rho$$



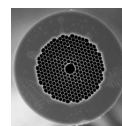
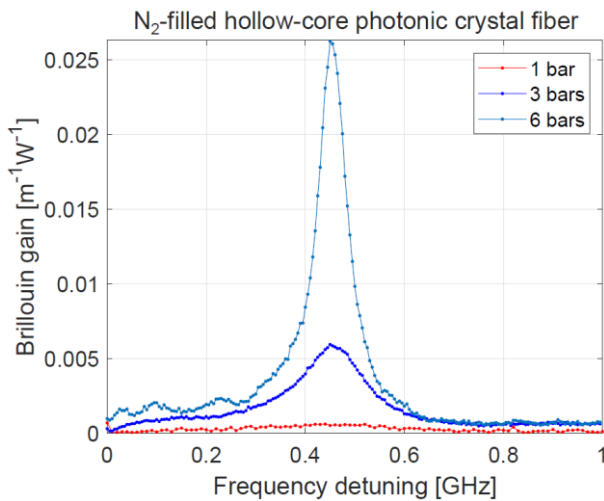
$$\Delta n \approx N_A \frac{\alpha}{2M} \Delta \rho$$

A large **density change** leads to a proportionally large **change in refractive index**

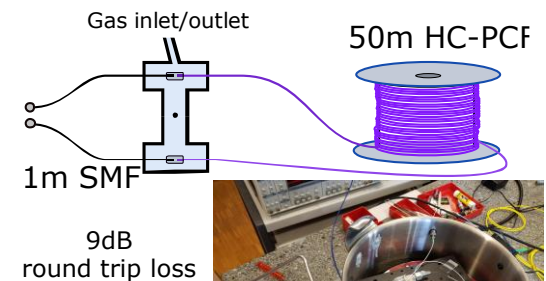
From electrostriction to dynamic grating



Observation of stimulated Brillouin scattering in gas



All results were obtained in a 50 metre gas-filled hollow-core photonic bandgap fibre (core size 10µm)



Brillouin linear gain in silica fibres is 100X larger!

Pressure dependence

Opto-acoustic strength depends linearly on pressure:

$$\Delta n \approx N_A \frac{\alpha}{2M} \Delta \rho \sim P$$

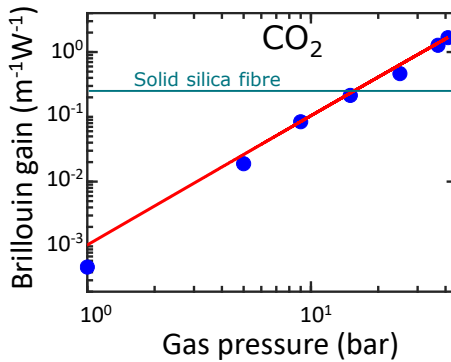
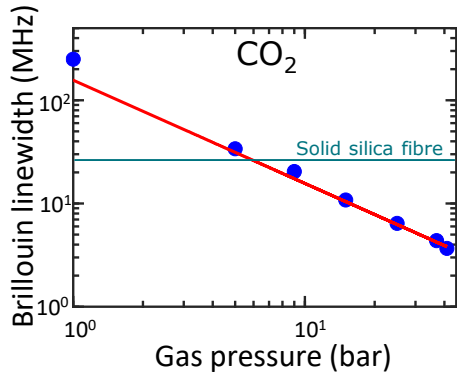
Acoustic lifetime also depends linearly on pressure:

$$\tau_A \sim P$$

→ Gain

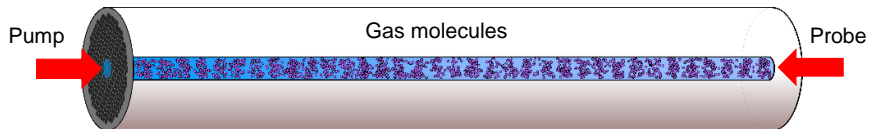
$$g_B \sim P^2$$

Pressure squared!

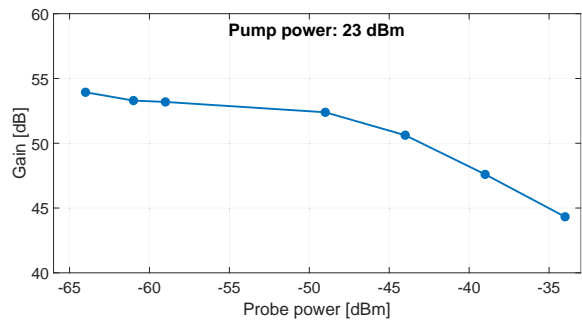
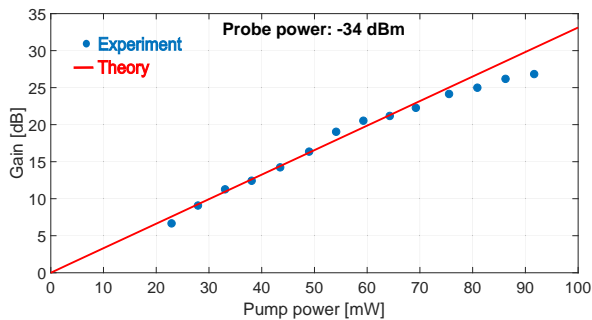


A hollow-core fibre can sustain a pressure up to 1000 bar and more.

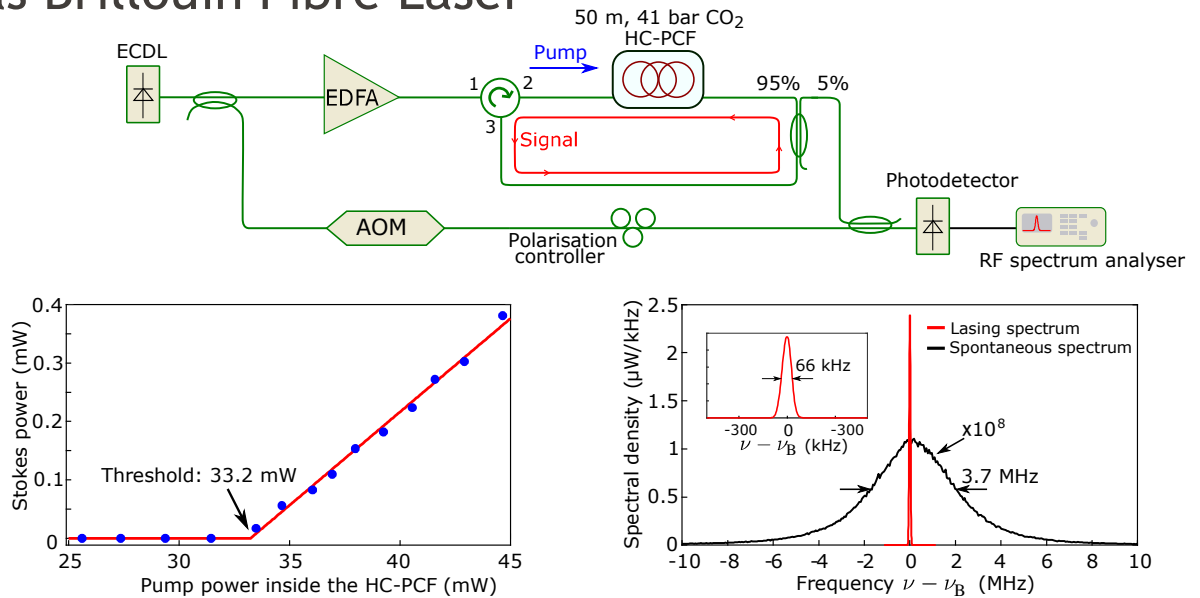
Amplification



Gain obtained along a 50m hollow-core fibre filled by CO₂ at 41 bar



Gas Brillouin Fibre Laser



Influence of the gas species

The Brillouin gain depends on several material parameters:

- Molecular mass \nearrow
- Molecular size \nearrow
- Viscosity \searrow
- Polarisability \nearrow

Gas	Name	V_a [m/s]	m [g/mol]	d [pm]	η_s [10^{-5} Pa · s]
CH ₄	Methane	466	16	400	1.88
N ₂	Nitrogen	354	28	370	2.86
CO ₂	Carbon dioxide	280	44	232	2.71

Measured Brillouin gains at 24°C and 10 bar:

Gas	Name	ν_B [MHz]	$\Delta\nu_B$ [MHz]	G [$m^{-1}W^{-1}$]
N ₂	Nitrogen	451	41	0.025
CH ₄	Methane	580	40	0.034
CO ₂	Carbon dioxide	351	21	0.105

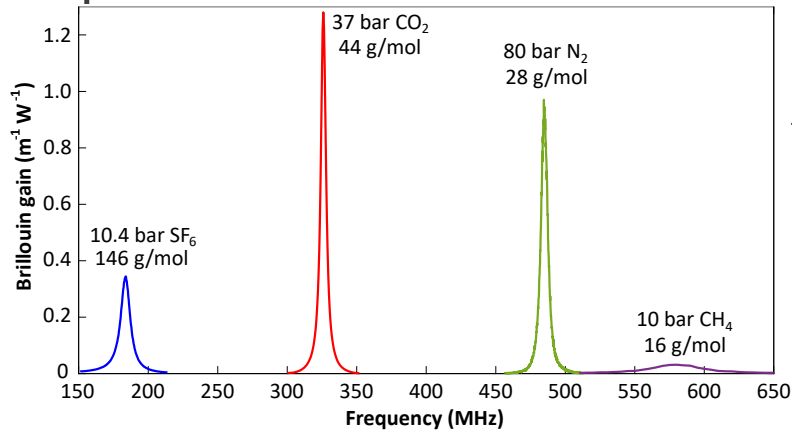
At first glance, gases with **complex and heavy molecules** should deliver more gain.

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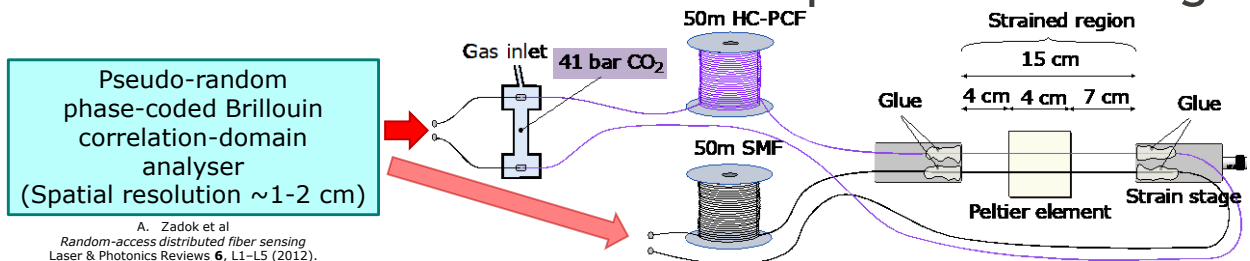
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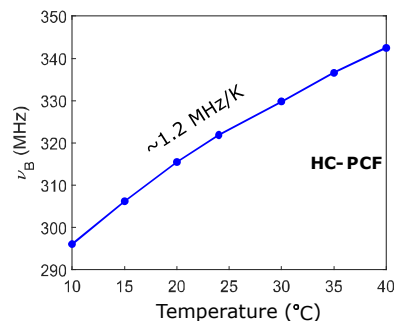
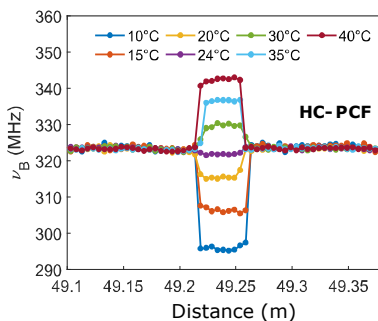
- But:
- May quickly turn into **liquid phase** at higher pressure.
 - May show **spectral absorption lines**, moreover significantly broadened at high pressure

Strain-Insensitive Distributed Temperature Sensing

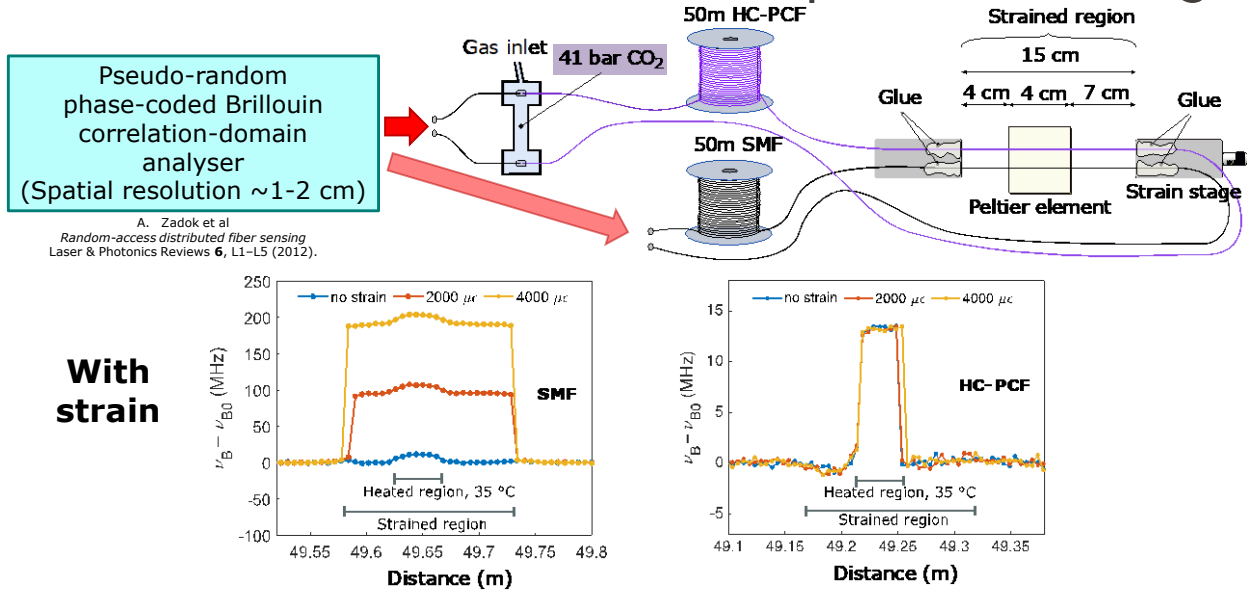


A. Zadok et al
Random-access distributed fiber sensing
Laser & Photonics Reviews **6**, L1-L5 (2012).

No strain

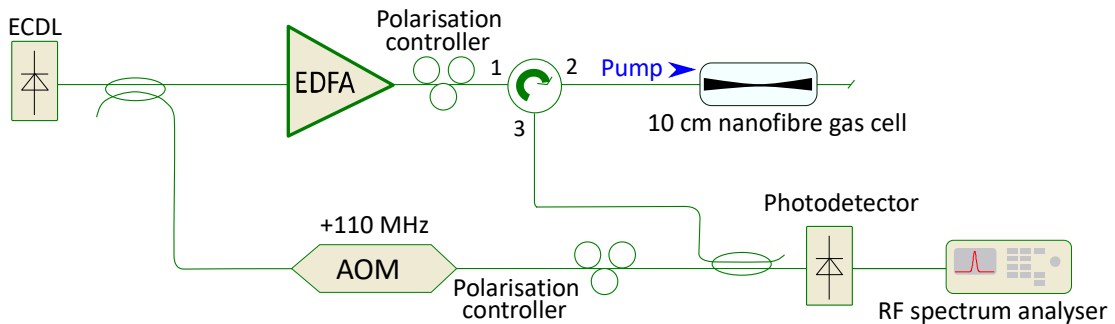
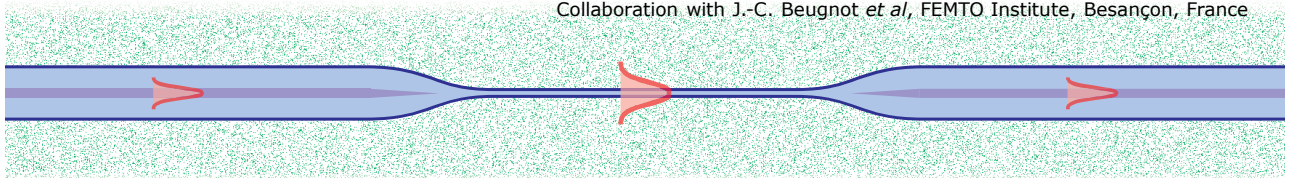


Strain-Insensitive Distributed Temperature Sensing



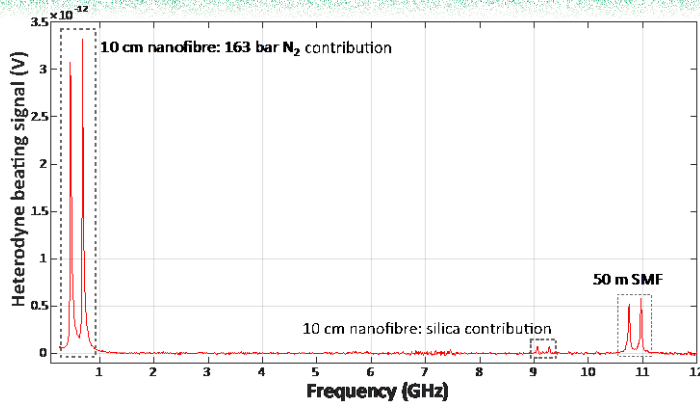
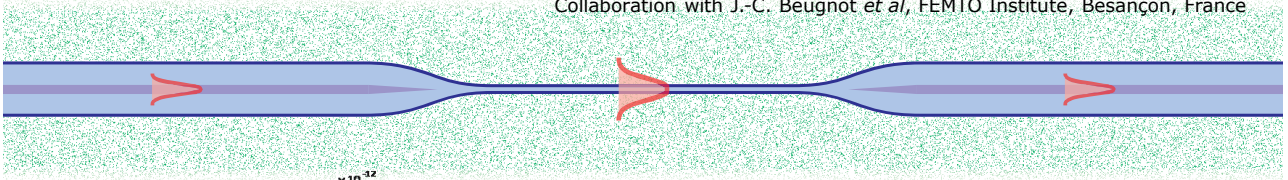
Brillouin amplification using evanescent field in gas

Collaboration with J.-C. Beugnot *et al*, FEMTO Institute, Besançon, France



Brillouin amplification using evanescent field in gas

Collaboration with J.-C. Beugnot *et al*, FEMTO Institute, Besançon, France



Final words

	SBS in silica	SBS in chalcogenide	Raman gain in H ₂ at 10 bar	Forward Brillouin gain in air	SBS in N ₂ at 1 bar	SBS in N ₂ at 1000 bars
Linear gain (m/W)	3×10^{-11}	6×10^{-9}	4.2×10^{-12}	4×10^{-14}	3×10^{-14}	3×10^{-8}

Stimulated Brillouin scattering in gases can outperform any existing nonlinear amplification!

