

# Origin of (anti-)nuclear clusters in hadronic collisions

Starting comments and experimental overview

# Welcome!

- Many thanks from the organisers (Laura, Maximiliano, Kfir, AK) for joining and participating in this special *all-online* TH institute!
- The idea is to have in-depth discussions among experts from the experimental and theory community to understand the properties and the production of anti- and hyper-nuclei in pp, p-Pb, and Pb-Pb collisions at LHC energies.
- Feel free to add discussion points and questions (including plots which were not shown in the talks) in our google-doc:

[https://docs.google.com/presentation/d/1Xr-nMaQCH\\_cSACO112Pqsrls6-3UPK7QrIDC2Q-iu0/edit](https://docs.google.com/presentation/d/1Xr-nMaQCH_cSACO112Pqsrls6-3UPK7QrIDC2Q-iu0/edit) - slide=id.p



# Agenda (1)

TUESDAY, 19 MAY



09:00

→ 10:00

## Starting comments / EXP overview

Speaker: Alexander Philipp Kalweit (CERN)

🕒 1h

📍 Zoom only



10:00

→ 11:00

## The thermal-statistical model applied to nuclear clusters

Speaker: Anton Andronic (Westfaelische Wilhelms-Universitaet Muenster (DE))

🕒 1h

📍 Zoom only



11:00

→ 12:00

## Coalescence as the origin of nuclear clusters

Measurements of the production of light nuclei in relativistic heavy-ion collisions present us with an apparent puzzle: While the momentum spectra of these nuclei indicate a low "kinetic freeze-out" temperature coupled with strong collective flow, their yields reflect a much higher "chemical freeze-out" temperature, consistent with the quark-gluon plasma hadronization temperature. How can these fragile nuclei "survive" the rescattering in the hadronic phase, cooling down in the process and picking up additional collective flow, without getting destroyed and depleted? I will explain how the coalescence model reconciles these observations and how recent kinetic simulations of the process support this model.

Speaker: Ulrich Heinz (The Ohio State University)

🕒 1h

📍 Zoom only



12:00

→ 13:30









Lunch Break

🕒 1h 30m

# Agenda (2)

12:00 → 13:30	<b>Lunch Break</b>	🕒 1h 30m
13:30 → 14:00	<b>Light nuclei yields: EXP vs. TH expectation</b> Speaker: Luca Barioglio (Universita e INFN Torino (IT))	🕒 30m 📍 Zoom only ✎
14:00 → 15:00	<b>Coalescence--correlations relation: B2 vs. HBT/radii</b>  Discussion of the EXP effort to test the predicted relation between pair correlations and coalescence yield. Main questions for TH and for the analysis, results and discussion.  Speaker: Bernhard Hohlweger (TUM)	🕒 1h 📍 Zoom only ✎
15:00 → 16:00	<b>Discussion: coalescence vs. HBT</b>	🕒 1h 📍 Zoom only ✎
16:00 → 16:15	<b>Break</b>	🕒 15m 📍 Zoom only
16:15 → 16:45	<b>The thermal model off-equilibrium</b> Speakers: Volodymyr Vovchenko, Dr Volodymyr Vovchenko (Lawrence Berkeley National Laboratory)	🕒 30m 📍 Zoom only ✎
16:45 → 17:15	<b>Cluster formation in a transport model</b> Speakers: Dr Dmytro Oliinychenko (Lawrence Berkeley National Laboratory), Dmytro Oliinychenko	🕒 30m 📍 Zoom only ✎
17:15 → 18:45	<b>Google Doc Discussion</b>	🕒 1h 30m 📍 Zoom only ✎

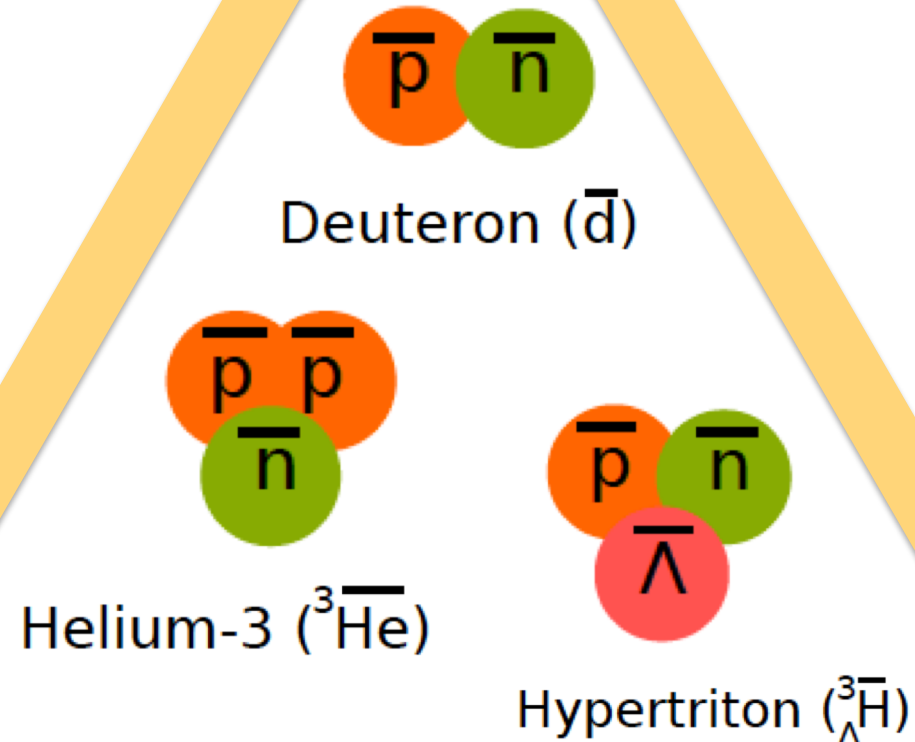
# Agenda (3)

WEDNESDAY, 20 MAY			
09:00	→ 09:30	<b>Hypertriton structure</b> Speaker: Hans-Werner Hammer	🕒 30m 📍 Zoom only 
09:30	→ 10:00	<b>Hypertriton lifetime</b> Speakers: Jean-Marc Rene Richard (Centre National de la Recherche Scientifique (FR)), Jean-Marc Richard, jean-marc Richard (IPNL)	🕒 30m 📍 Zoom only 
10:00	→ 10:45	<b>Hypertriton yield at ALICE: EXP vs. TH</b>	🕒 45m 📍 Zoom only 
10:45	→ 11:30	<b>Discussion: Hypertriton as a probe of cluster formation</b>	🕒 45m 📍 Zoom only 
11:30	→ 12:00	<b>How to distinguish coalescence from thermal production of light nuclei</b> Speaker: Stanislaw Mrowczynski	🕒 30m 📍 Zoom only 
12:00	→ 13:30	<b>Break</b>	🕒 1h 30m 📍 Zoom only
13:30	→ 15:00	<b>Google Doc Discussion &amp; Workshop summary</b>	🕒 1h 30m 📍 Zoom only 
15:00	→ 16:00	<b>TH Colloquium: Jesse Thaler (<a href="https://indico.cern.ch/event/888504/">https://indico.cern.ch/event/888504/</a>)</b>	🕒 1h 📍 Zoom only 

# The LHC as an anti-matter factory

# Production scenarios: thermal vs coalescence

The "golden triangle"  
of anti-nuclei physics at  
the LHC.

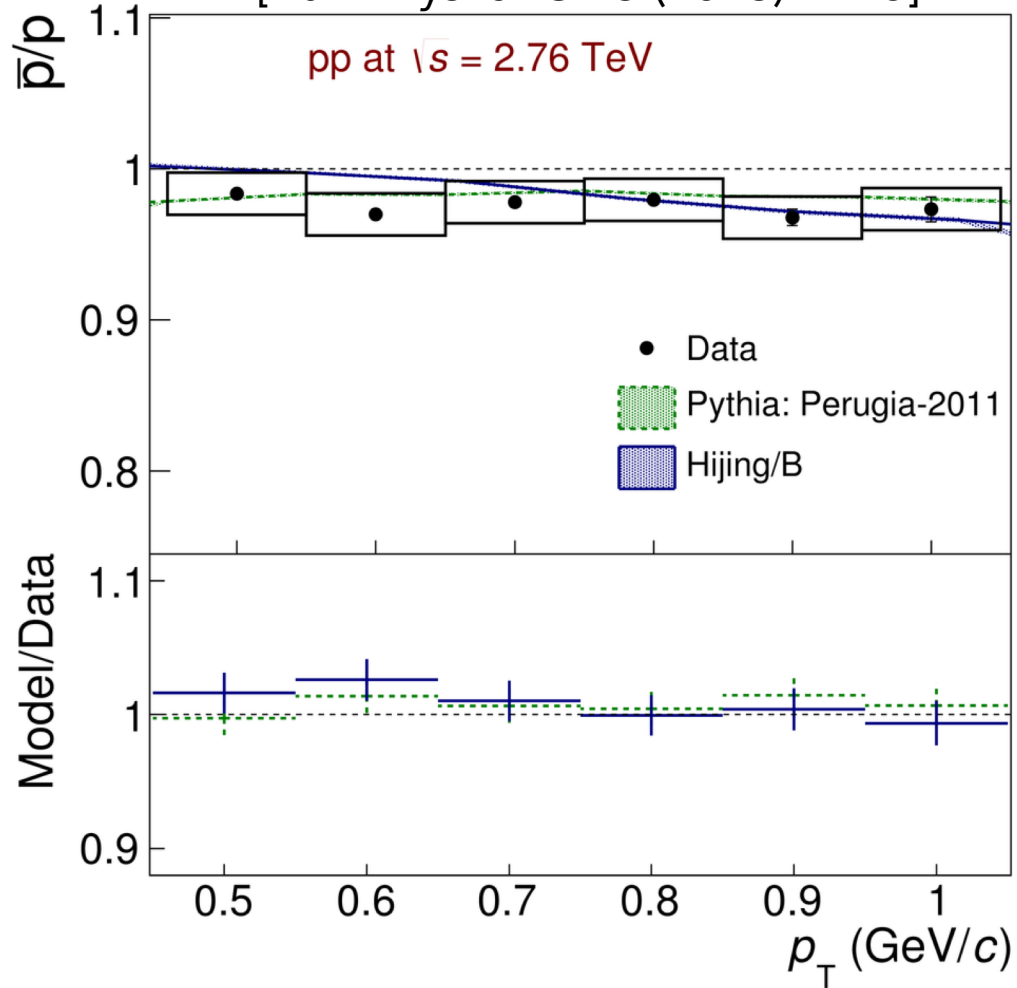


Search for dark  
matter in space

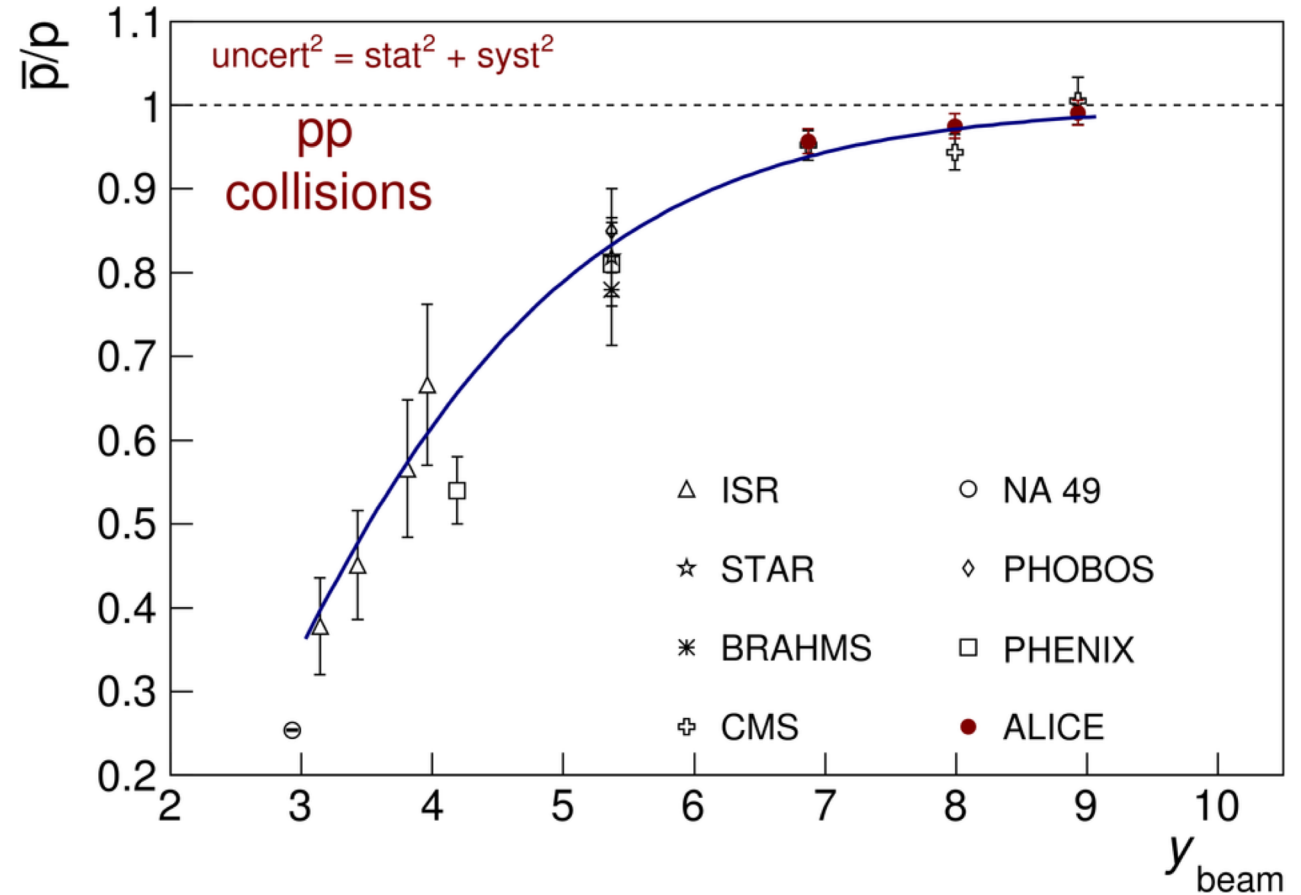
Hyperon-nucleon  
interactions

# Particles and anti-particles

[Eur. Phys. J. C 73 (2013) 2496]

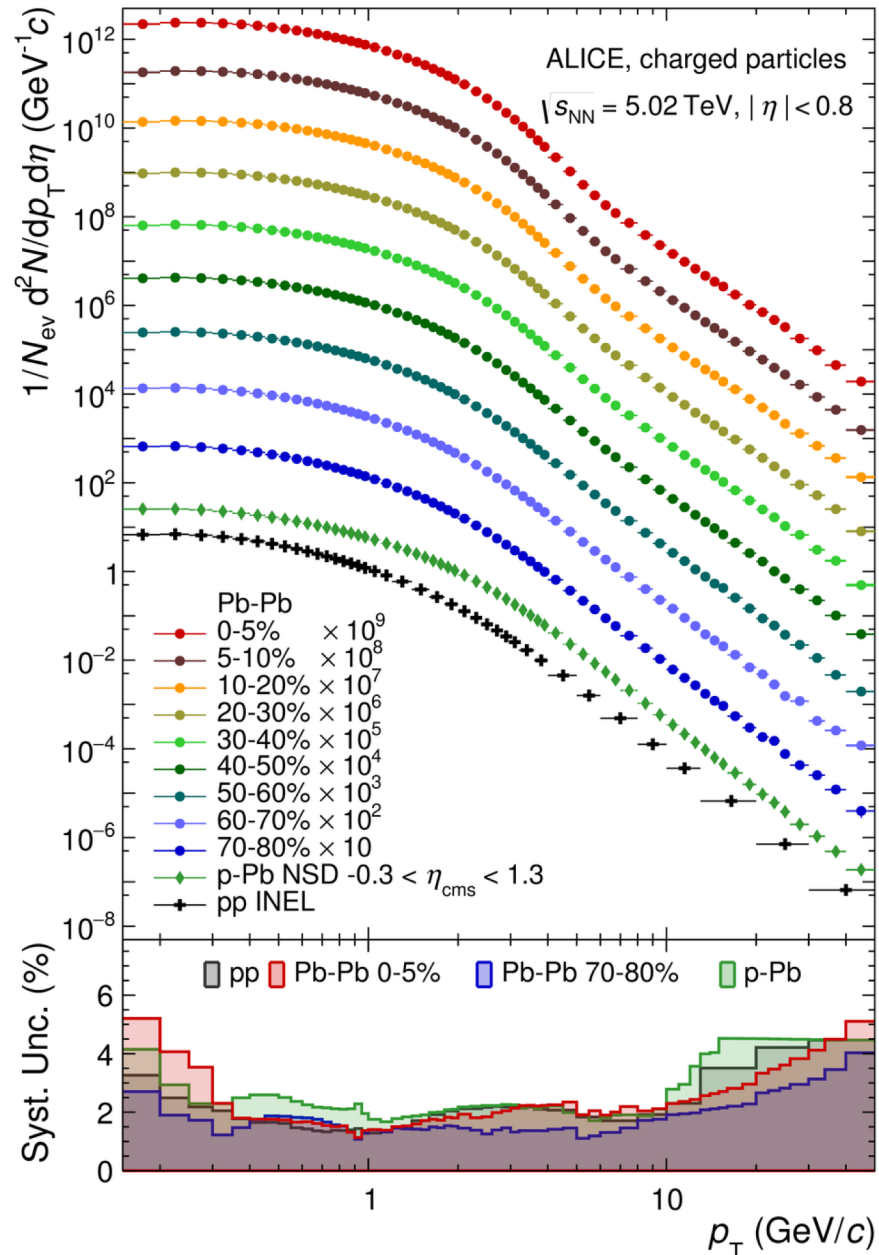


[Eur. Phys. J. C 73 (2013) 2496]



At LHC energies, particles and anti-particles are produced in equal abundance at mid-rapidity.

# Particle production at LHC energies (1)

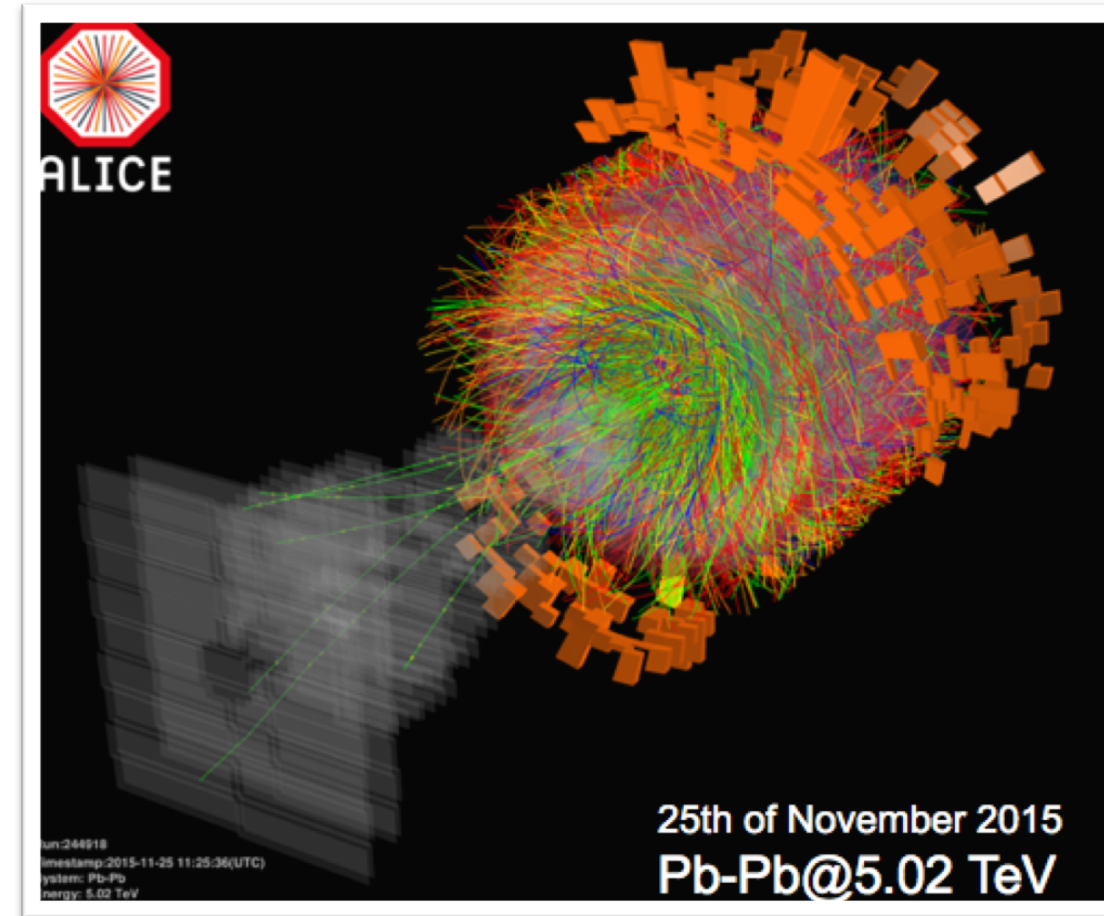
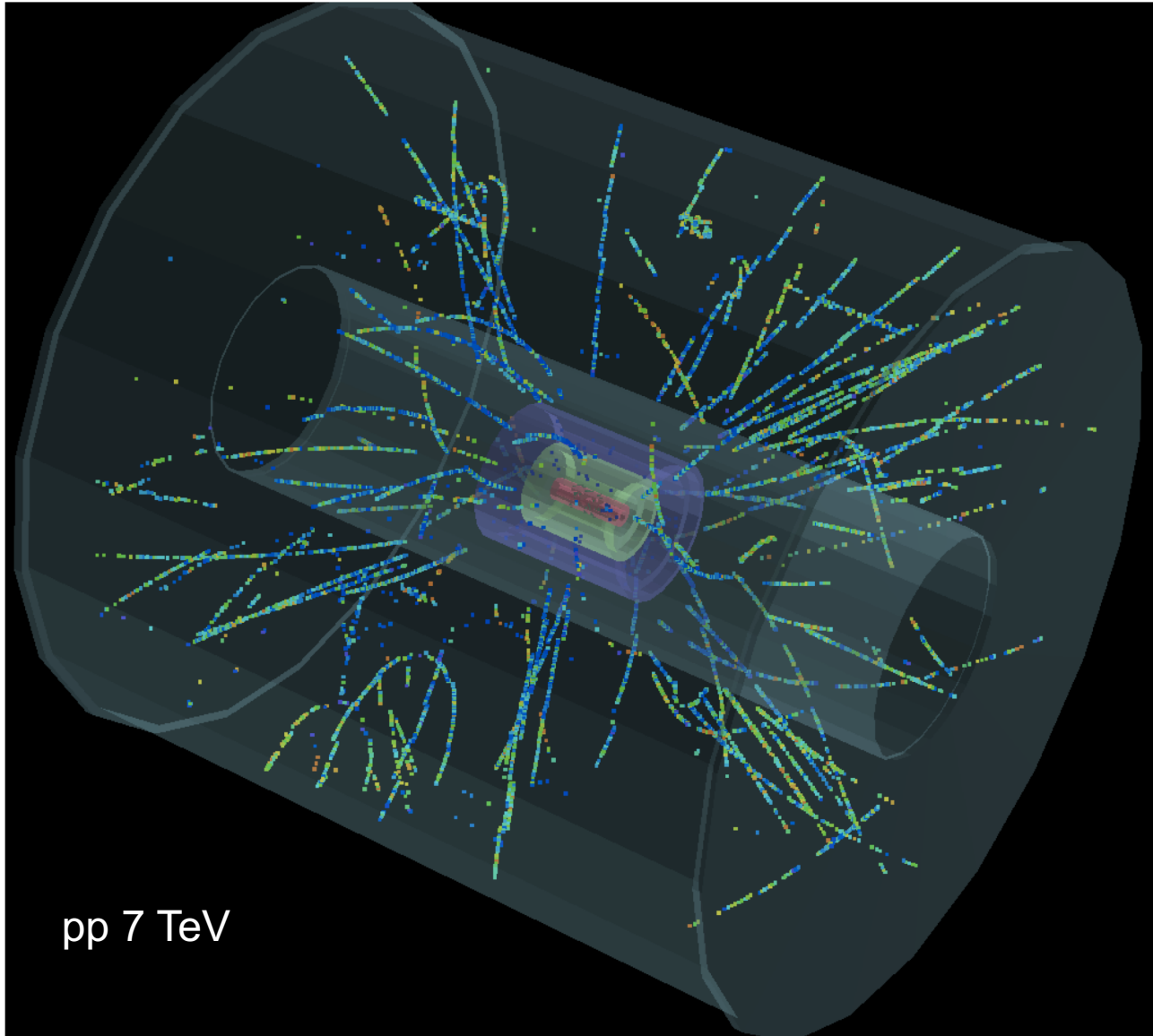


→ Even at LHC energies, 95% of all particles are produced with  $p_T < 2 \text{ GeV}/c$  in pp and Pb-Pb collisions.

→ Bulk particle production and the study of collective phenomena are associated with **"soft" physics** in the non-perturbative regime of QCD.

[ALICE, JHEP 1811 (2018) 013]

# Particle production at LHC energies (2)

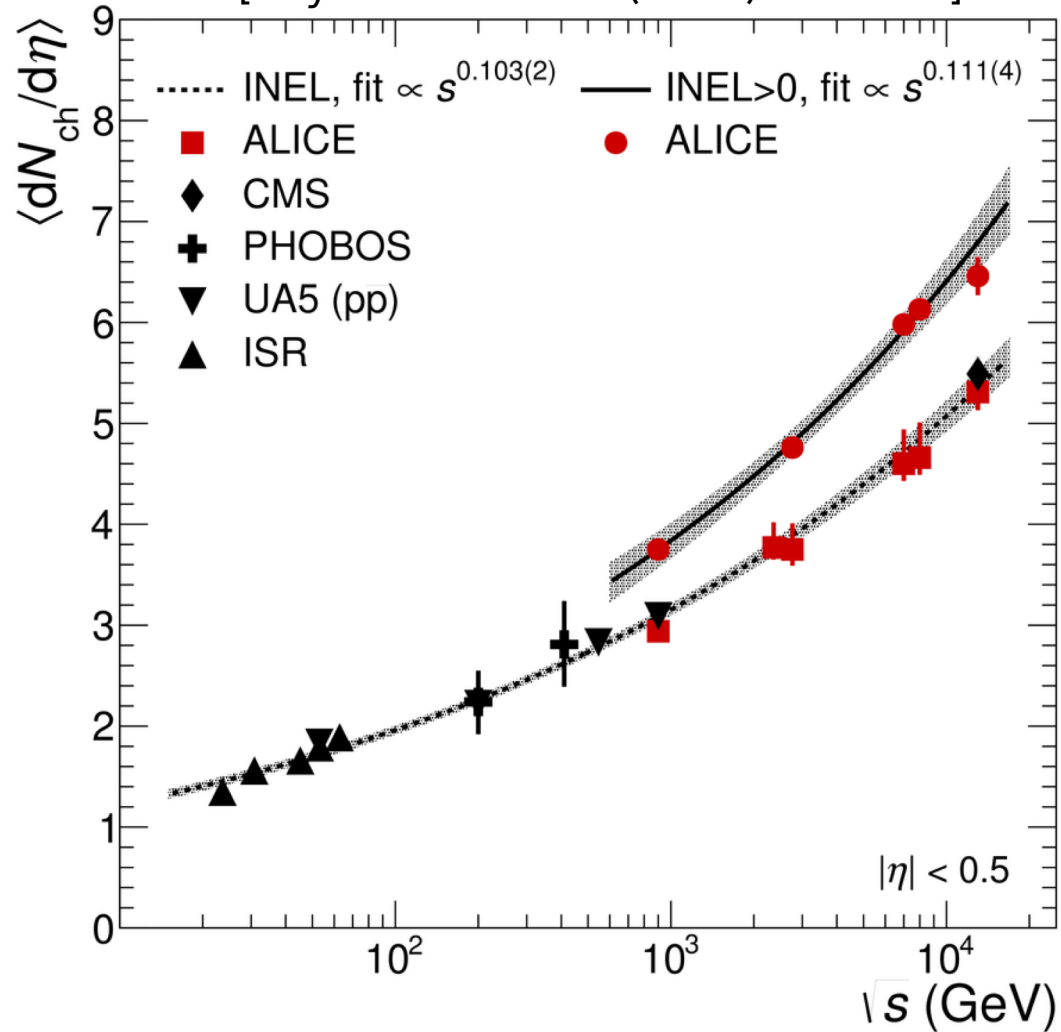


→ How many particles are produced in such a collision?



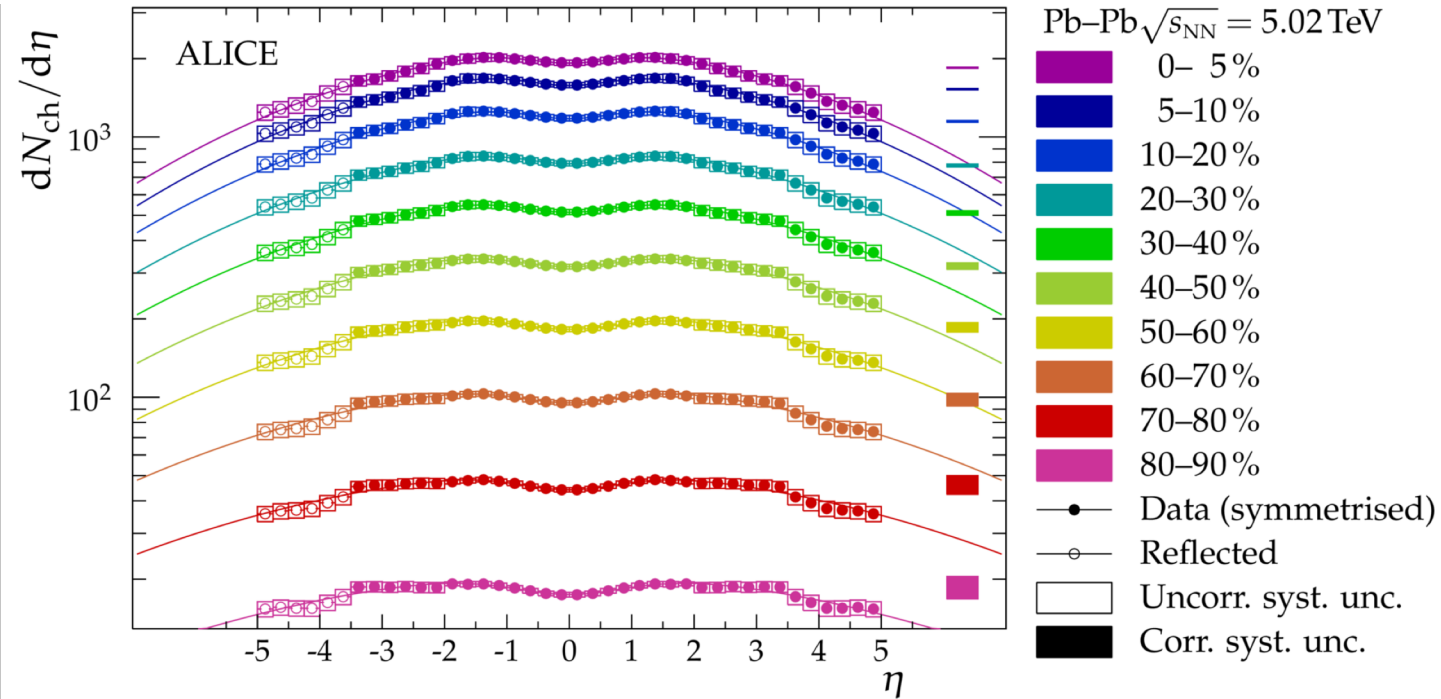
# Particle production at LHC energies (3)

[Phys. Lett. B 753 (2016) 319-329]



$dN_{ch}/d\eta \approx 4-7$  in pp collisions at LHC energies

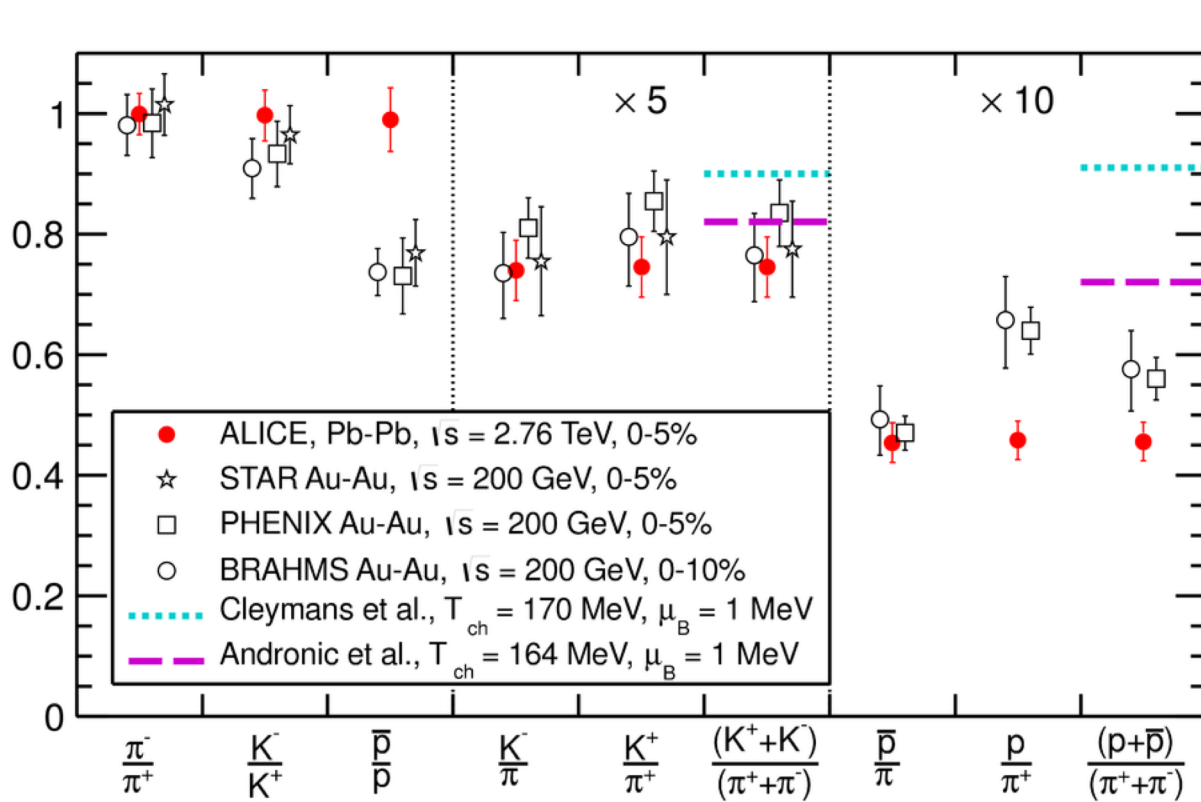
[Phys.Lett. B 772 (2017) 567-577]



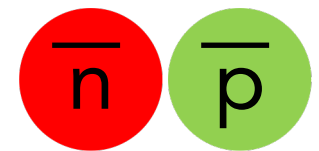
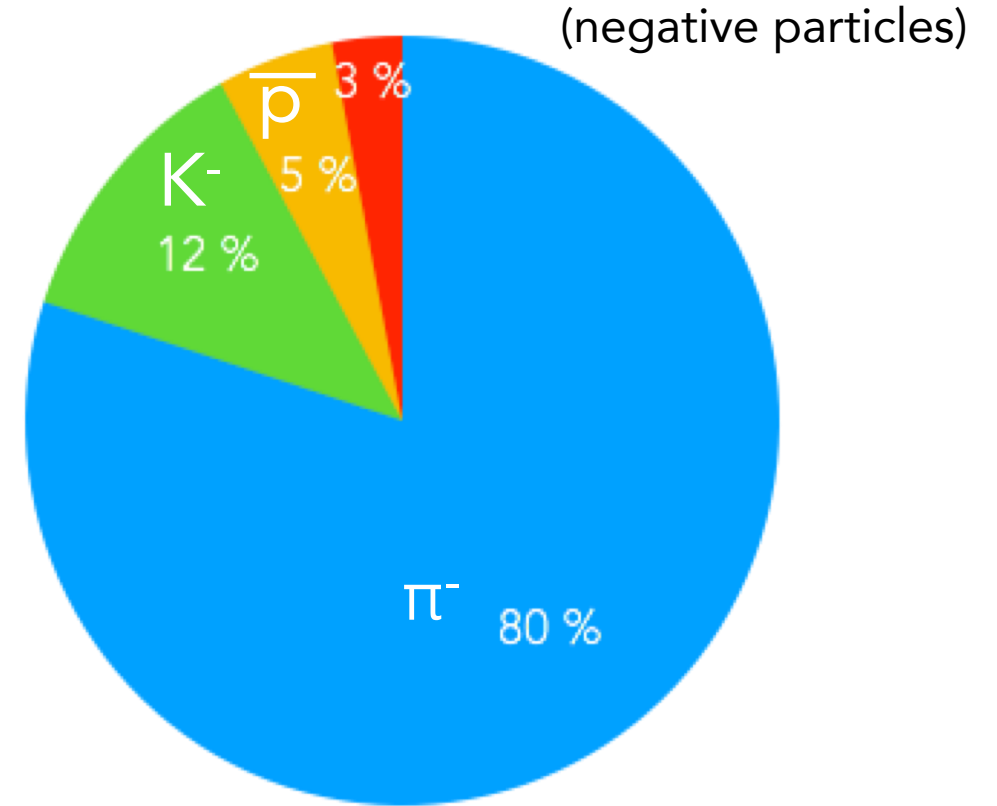
ALI-PUB-115086

$dN_{ch}/d\eta \approx 1943 \pm 54$  at midrapidity in Pb-Pb.

# Particle production at LHC energies (4)



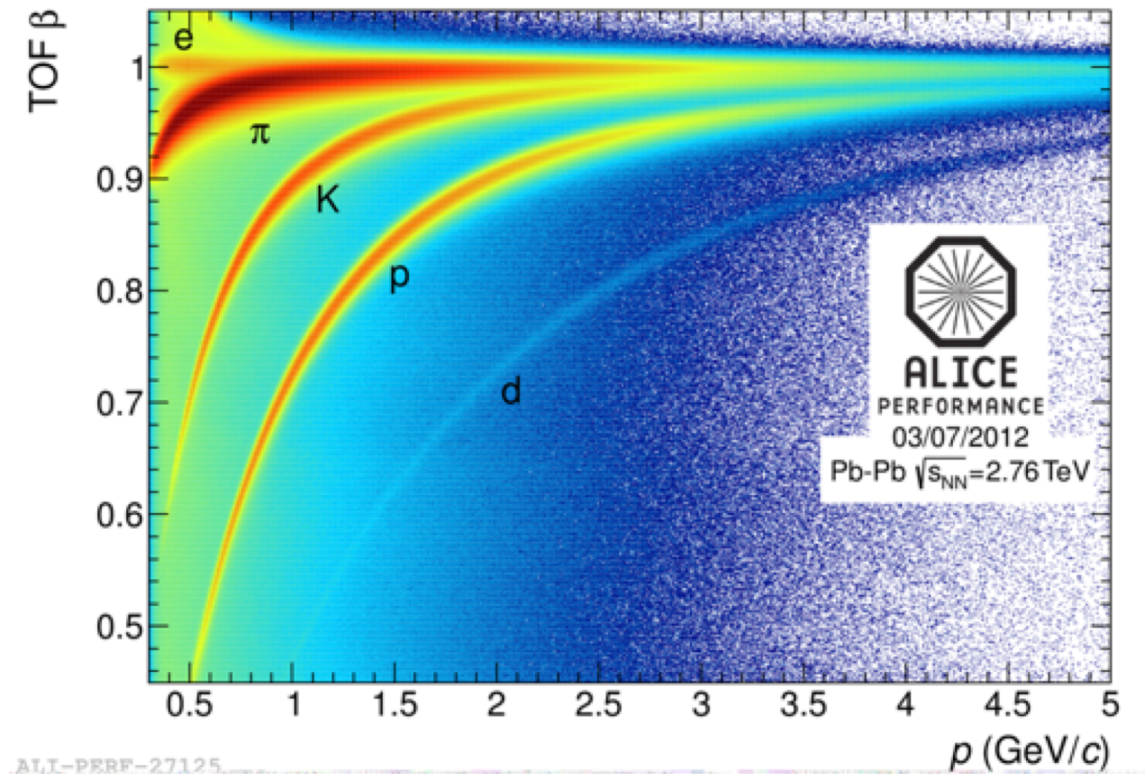
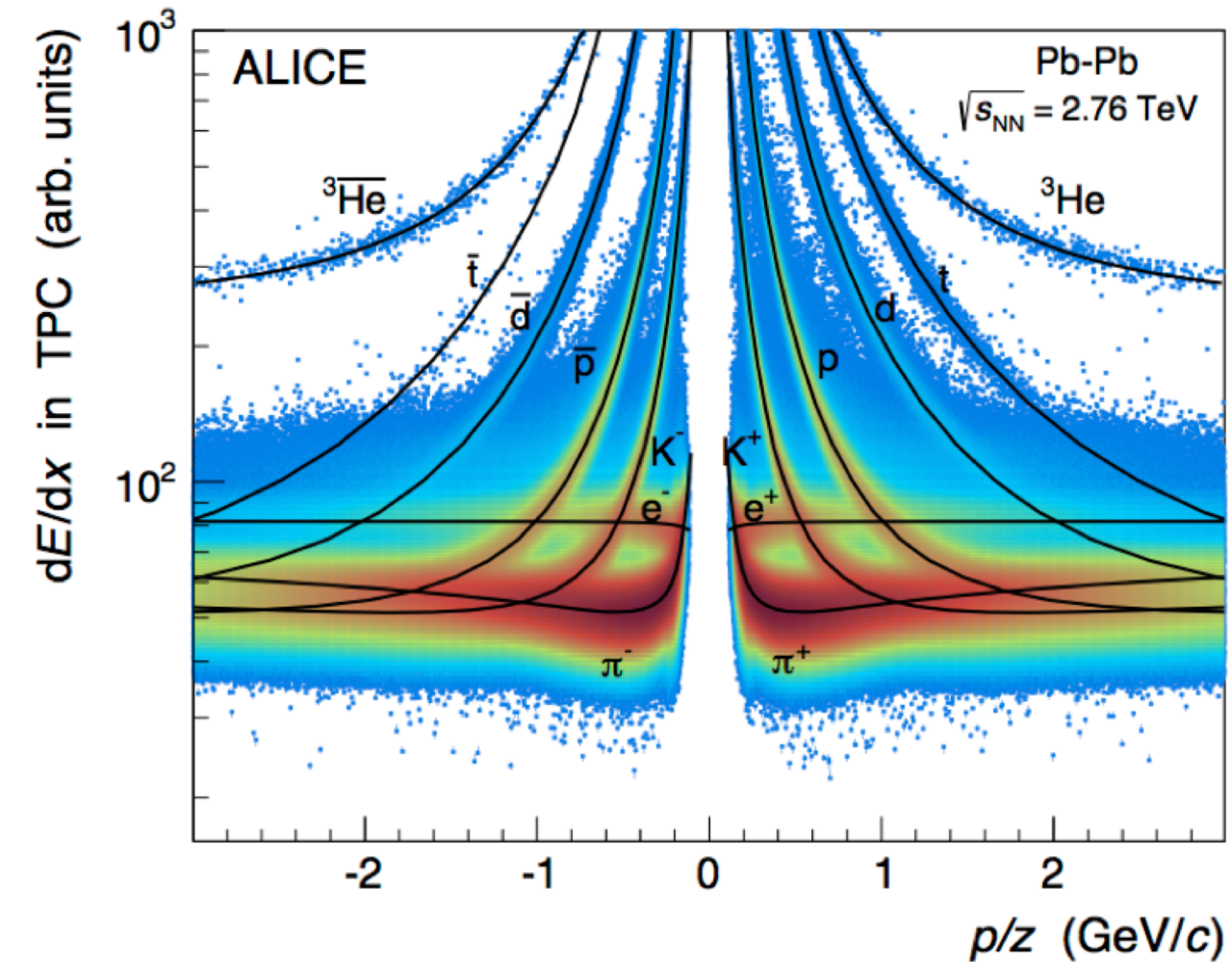
[Phys. Rev. Lett. 109 (2012) 252301]



anti-deuteron

- Only ~5% of all negative particles are anti-protons (the “lightest anti-nucleus”).
- The production of composite anti-particles is **very rare**: ~ **0.005%** are anti-deuterons. → **Clean PID needed!**

# Anti-nuclei identification in ALICE



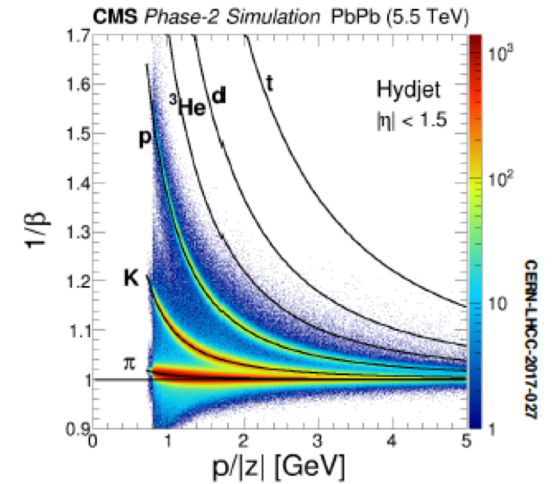
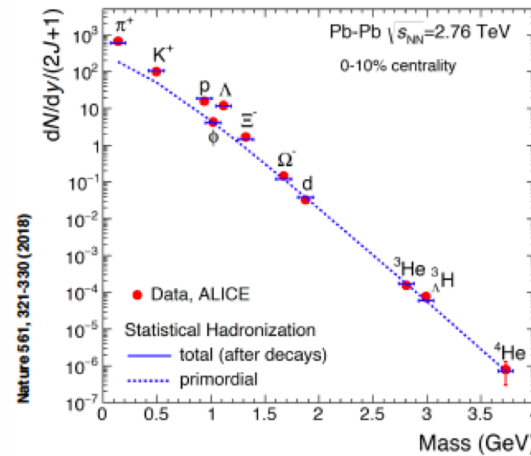
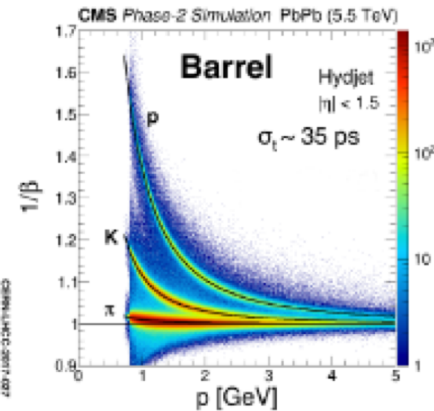
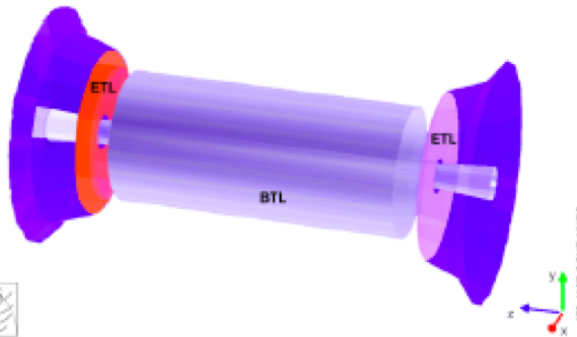
[Phys. Rev. C 93 (2015) 024917]

# Prospects for (anti-)(Hyper-)nuclei measurements in CMS

[QM talk,  
A. Govinda]

## MIP Timing Detector (2026+)

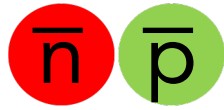
- Time-of-flight:  $\pi/K/p$  identification!
- Enhanced heavy flavour program... and more!



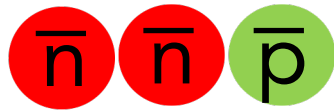
- Light ( $d$ ,  ${}^3\text{He}$ ,  $t$ ) and hyper ( ${}^3_\Lambda\text{He}$ ,  ${}^3_\Lambda\text{H}$ ) nuclei can be identified over a wide kinematic range via TOF using MTD.
- Also provide insights for dark matter searches and astrophysics.



# Light (anti-)nuclei



anti-deuteron



anti-triton



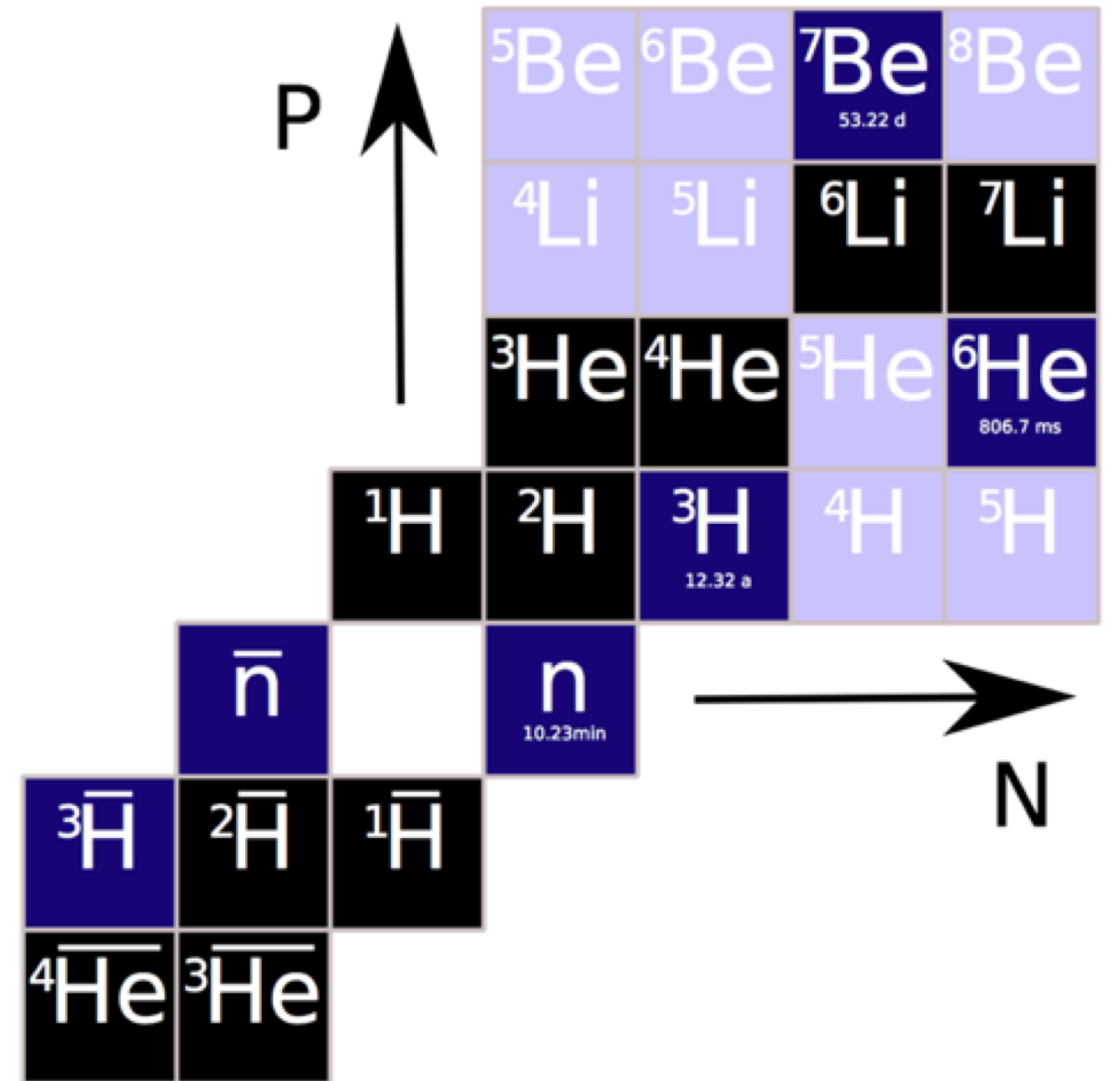
anti-hyper-triton



anti-helium3

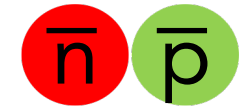
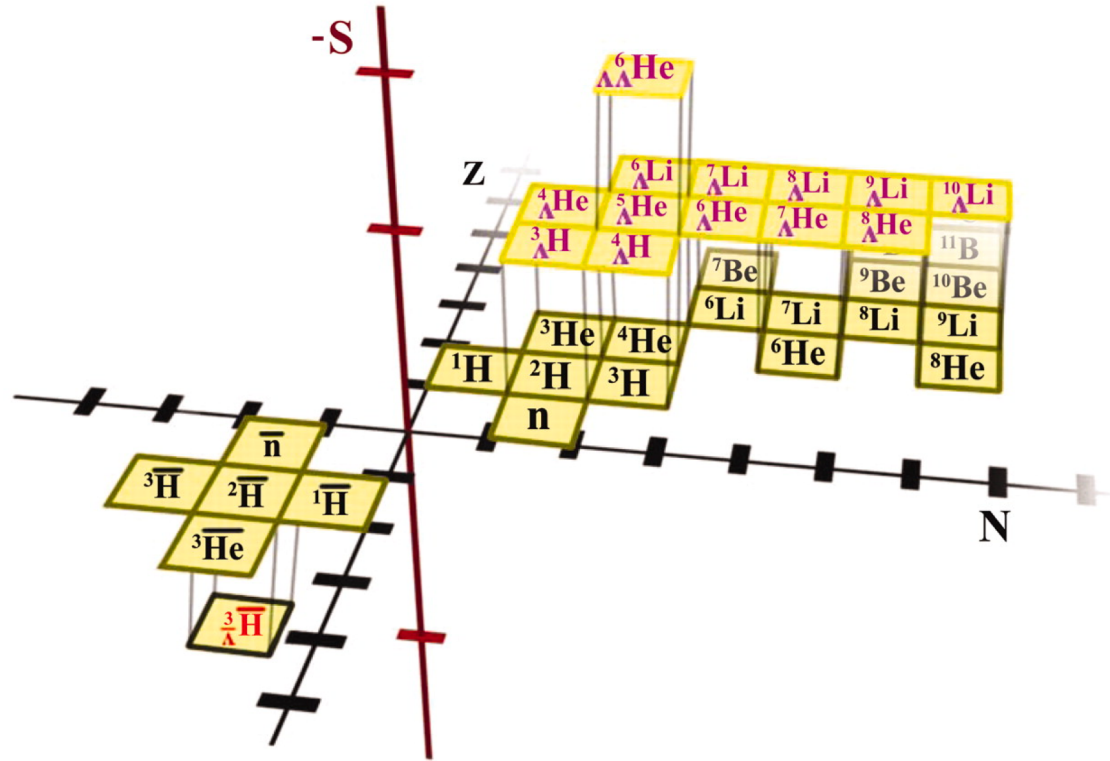


anti-alpha



Anti-(hyper)-nuclei up to  $A=4$  are currently in reach at accelerators. The anti-alpha is the heaviest observed so far and was first seen by the STAR experiment in 2011.

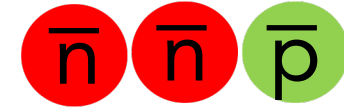
# Light (anti-)nuclei



anti-deuteron



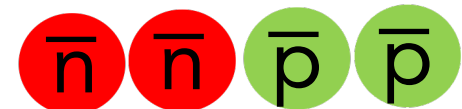
anti-hyper-triton



anti-triton



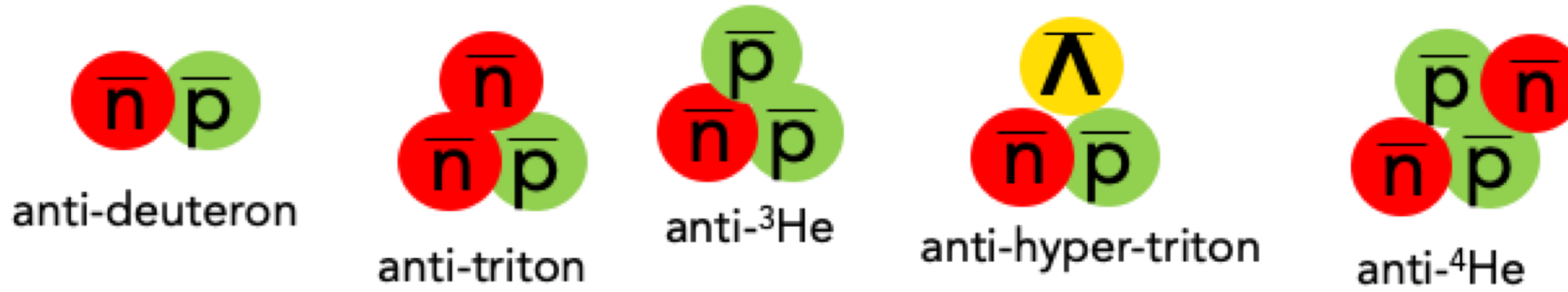
anti-helium3



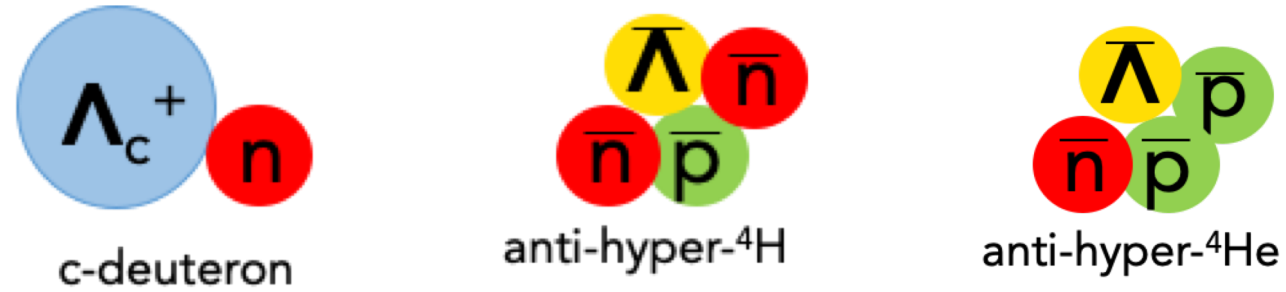
anti-alpha

# Zoo of exotic QCD bound states

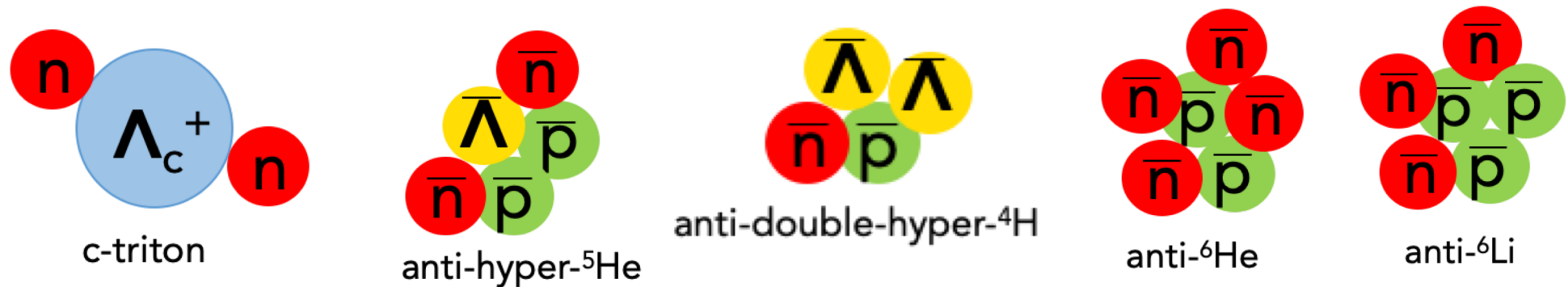
Run 1 & 2



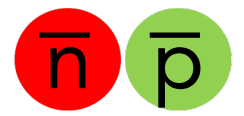
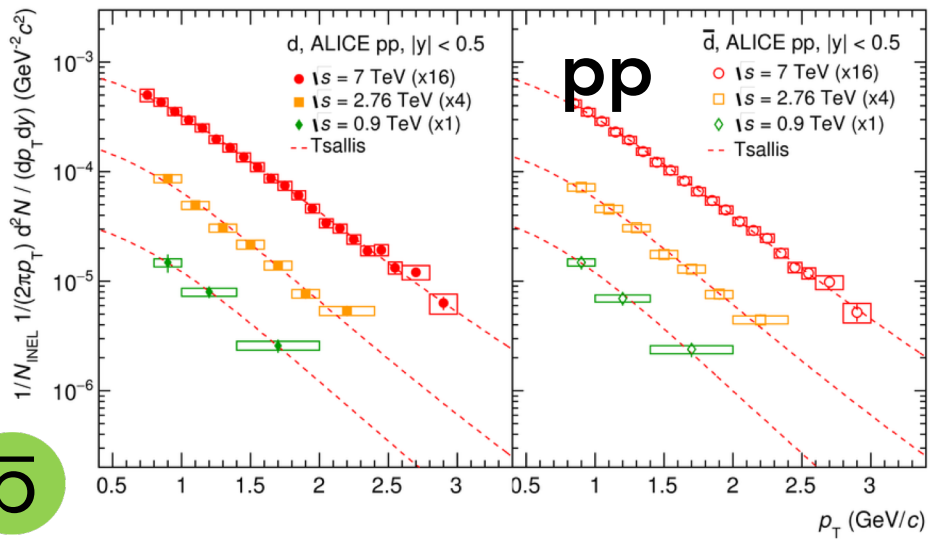
Run 3 & 4



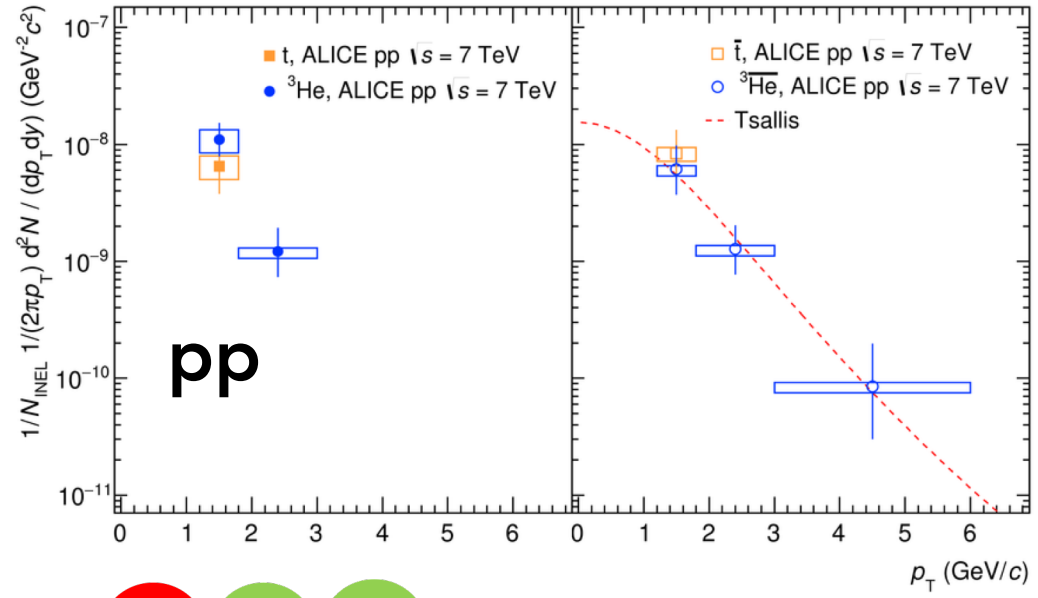
Run 5 & 6



# A wealth of data: $p_T$ spectra



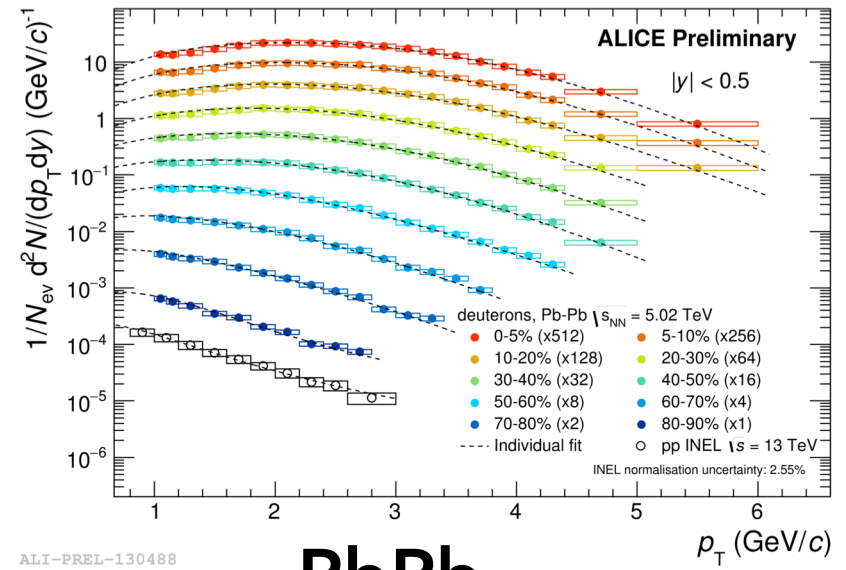
anti-deuteron



anti-helium3

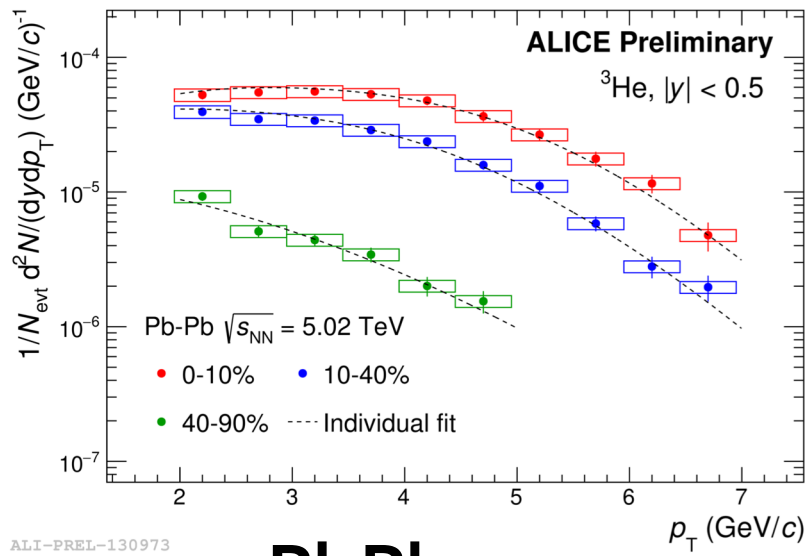


anti-hyper-triton



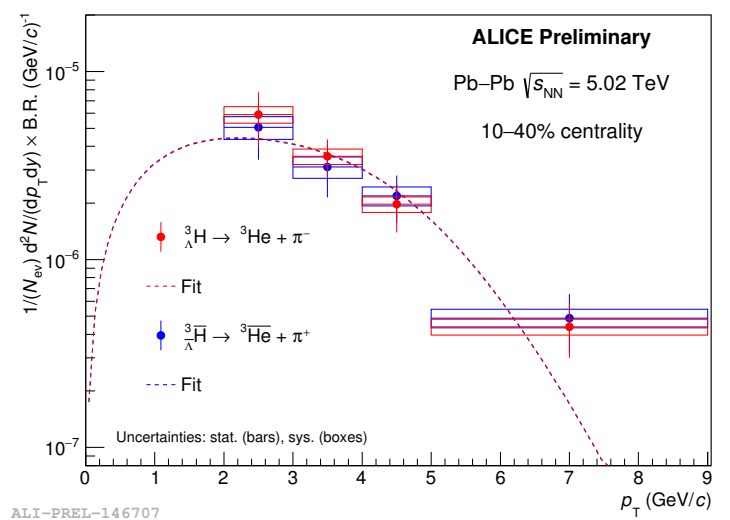
PbPb

ALI-PREL-130488



PbPb

ALI-PREL-130973

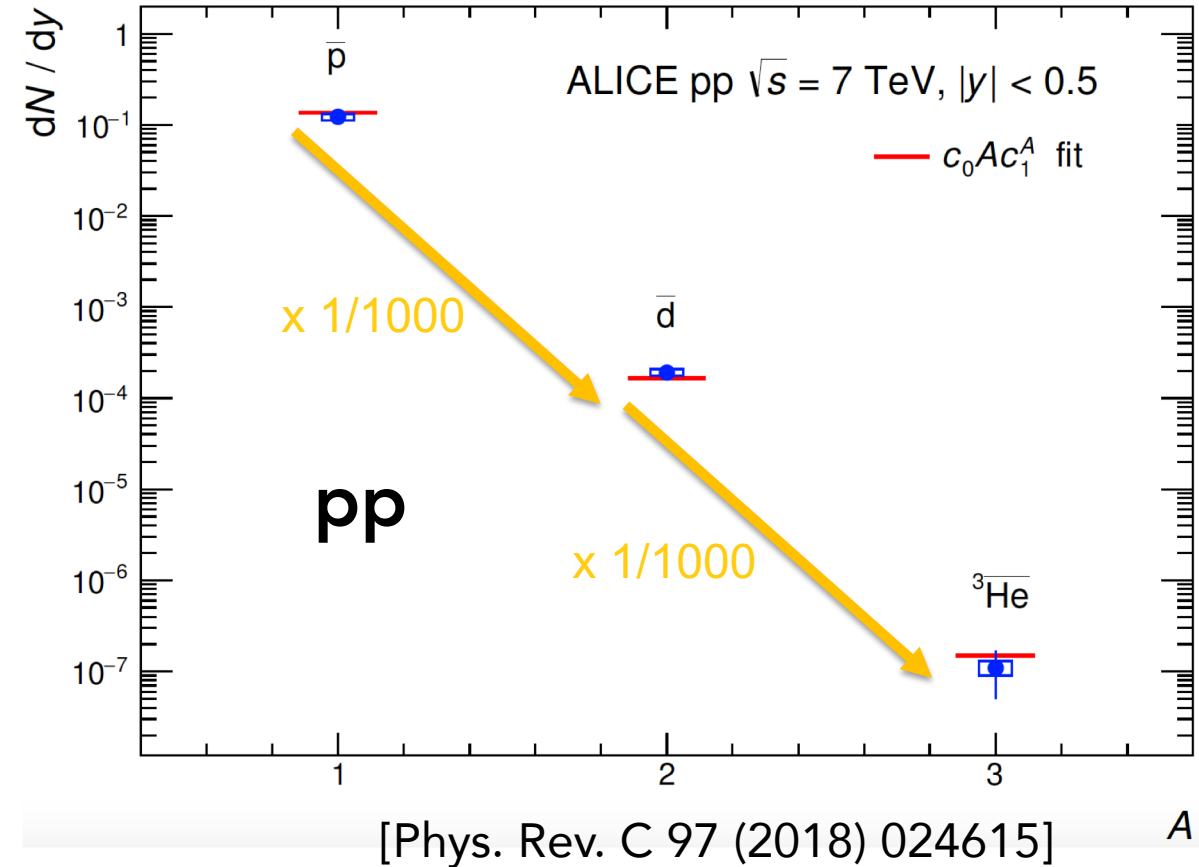
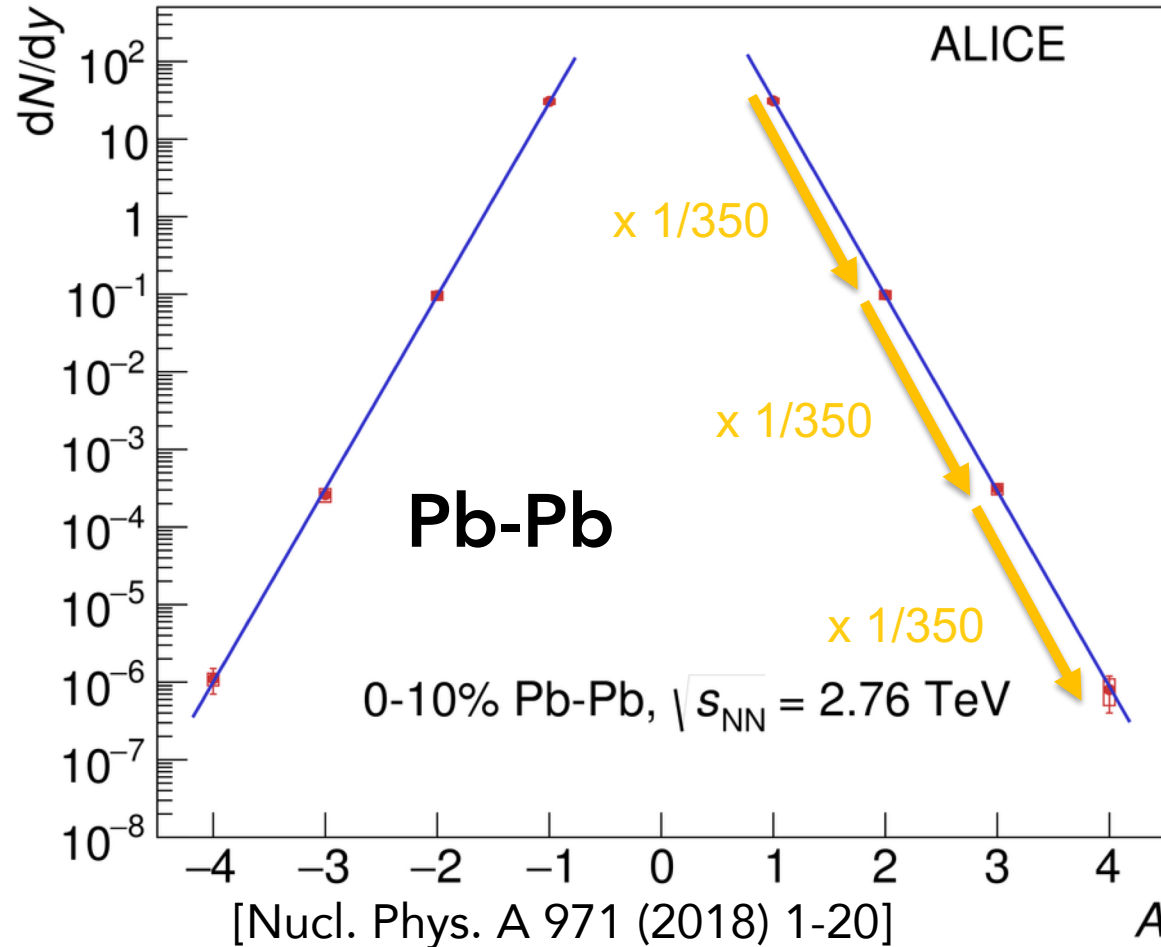


PbPb

ALI-PREL-146707



# Penalty factor at the LHC



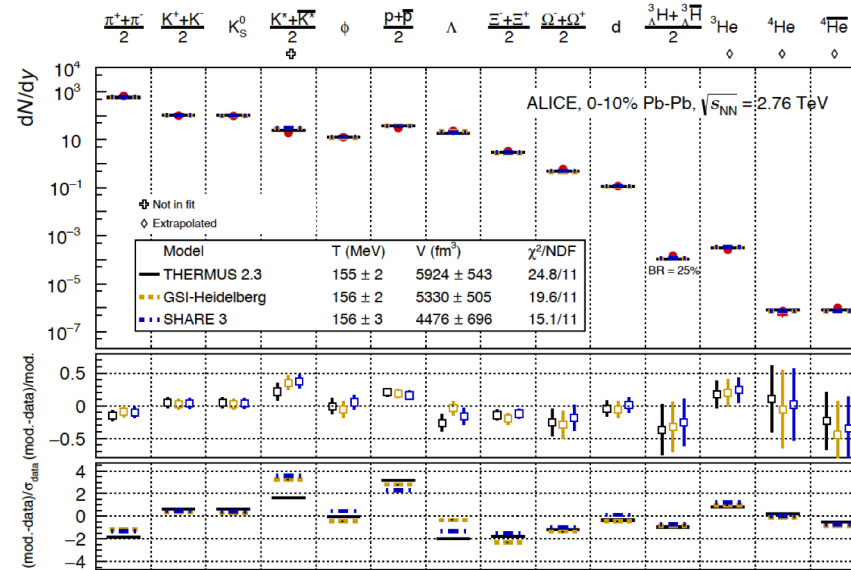
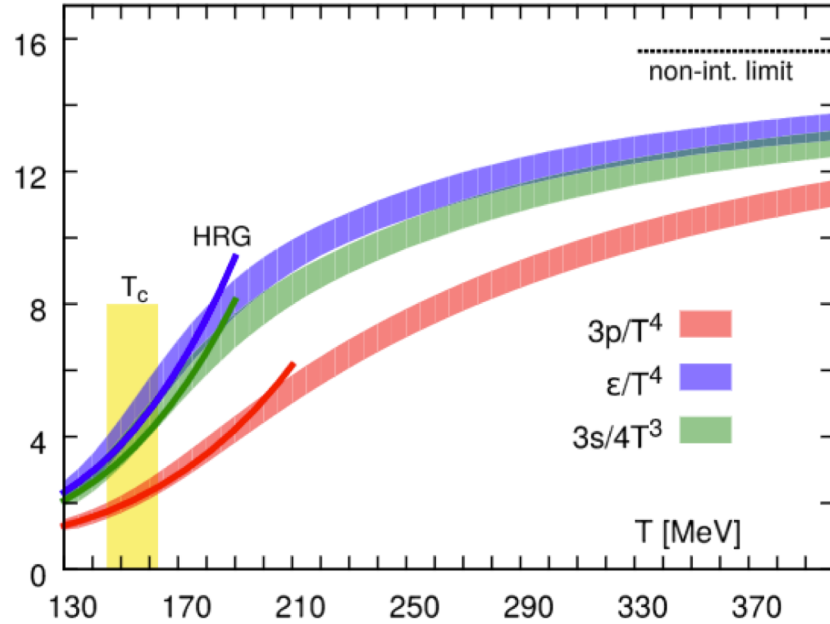
The production yield of (anti)-nuclei decreases by a factor of about  $\sim 350$  for each additional nucleon in Pb-Pb ( $\sim 1000$  in pp).

# An introduction to the “anti-nuclei puzzle”

# A problem of two scales

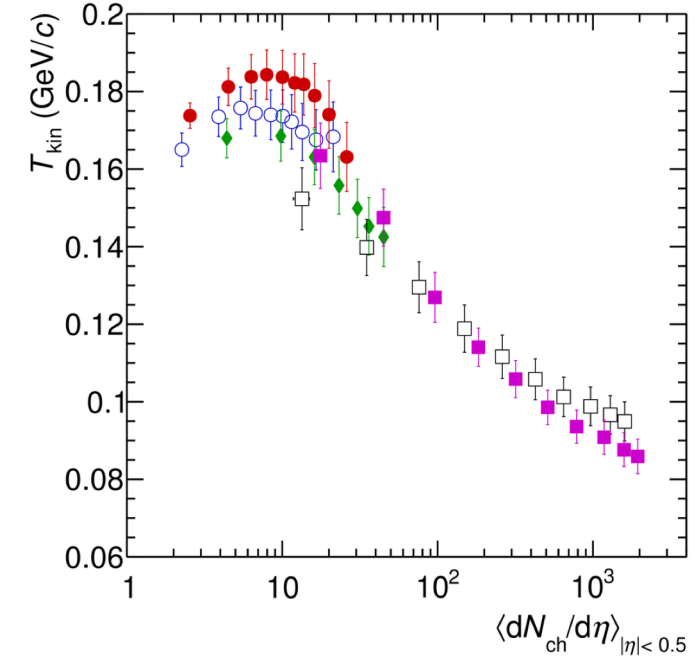
1. "Temperatures", mean-transverse momenta etc. are of the 100-200 MeV scale:

[Lattice QCD, PRD 90 094503 (2014)]



[Nucl. Phys. A 971 (2018) 1-20]

[arXiv:2003.02394]



2. Binding energies of the composite objects are of the order of  $\sim 1-10$  MeV:

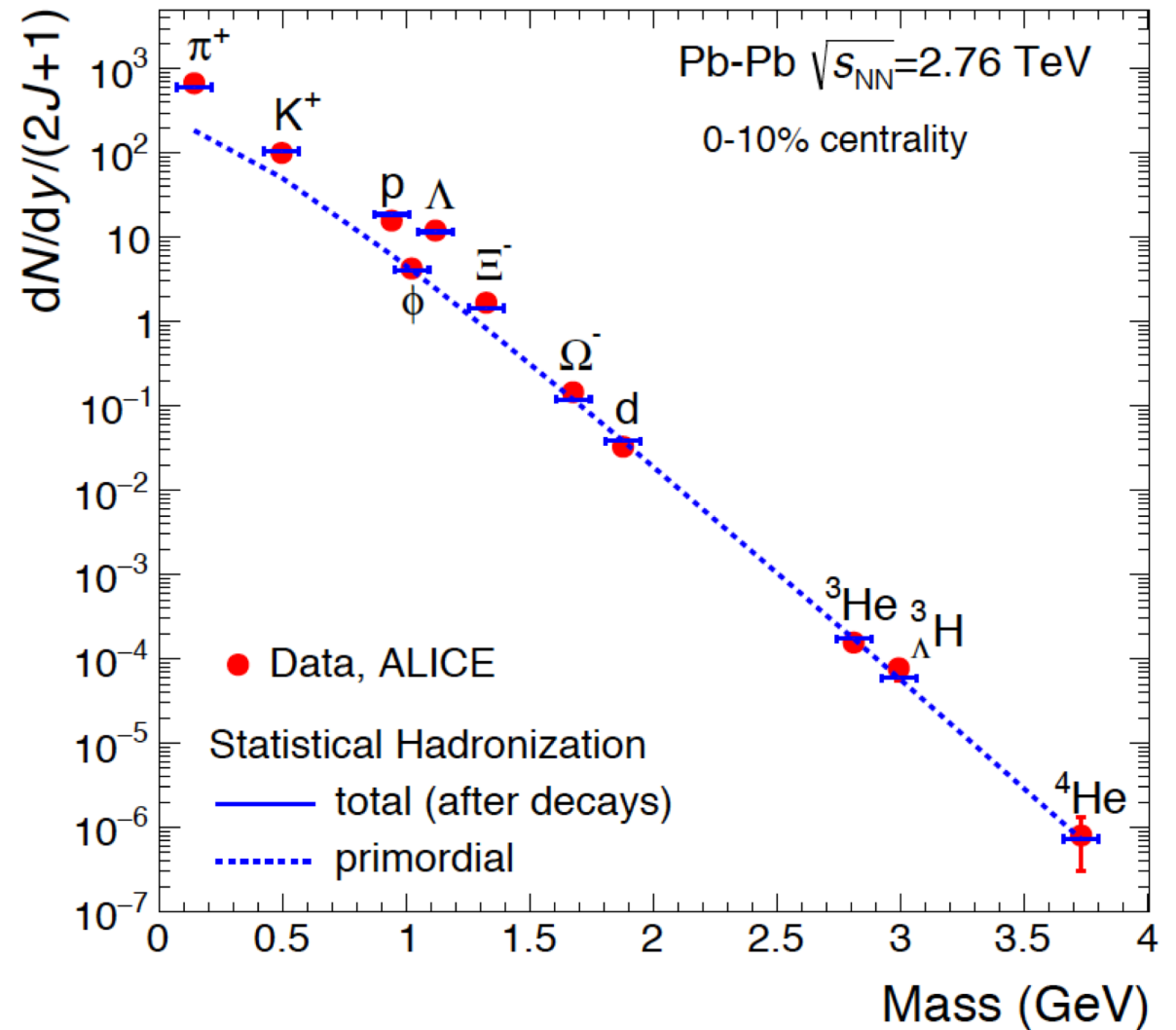
Mass number	Nucleus	Composition	Binding energy (MeV)	Spin
$A = 2$	d	pn	2.224575 (9)	1
$A = 3$	$^3\text{H}$	pnn	8.4817986 (20)	1/2
	$^3\text{He}$	ppn	7.7180428 (23)	1/2

# Chemical equilibrium at the LHC: anti-nuclei puzzle

Production yields of light flavour hadrons from a chemically equilibrated fireball can be described by statistical-thermal models (roughly  $dN/dy \sim \exp\{-m/T_{ch}\}$ , in detail derived from partition function)

→ In Pb-Pb collisions, particle yields of light flavor hadrons are described over 7 orders of magnitude with a **common** chemical freeze-out temperature of  $T_{ch} \approx 156 \text{ MeV}$ .

→ Light (anti-)nuclei are also well described despite their low binding energy ( $E_{b,d} = 2.2 \text{ MeV} \ll T_{ch}$ ).

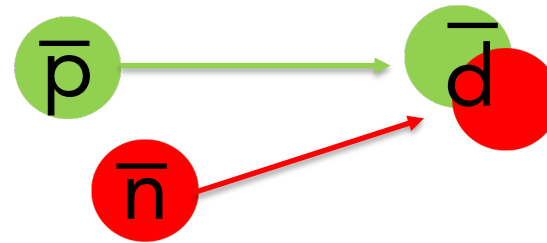


[A. Andronic *et al.*, arXiv:1710.09425]

# Coalescence parameters $B_A$

- A production by coalescence could explain this puzzle.
- (anti-)nuclei production by coalescence of (anti-)protons and (anti-)neutrons which are close by in momentum and configuration space. Roughly speaking: *"deuteron  $\propto$  proton  $\times$  neutron  $\Rightarrow$  deuteron  $\propto$  proton<sup>2</sup>"*

$$E_d \frac{d^3 N_d}{dp_d^3} = B_2 \left( E_p \frac{d^3 N_p}{dp_p^3} \right)^2$$



- Spherical approximation: maximum momentum difference (coalescence momentum  $p_0$ ) is approx. 100 MeV (5.3 MeV kinetic energy of a nucleon in the rest frame of the other).

# Fermi momentum and uncertainty principle

- “Close in phase-space”  $\leftrightarrow$  “Close in configuration and momentum space” sounds nice, but is quantum-mechanically ill defined due to the uncertainty principle!
- Imagine a point-like emission source (pp): phase-space reduces to momentum space and the **coalescence momentum**, becomes equal to the momentum of the nucleons in the bound nucleus (the “Fermi momentum” in large nuclei).
- Quantum-mechanically correct treatment: overlap of the source function with the **Wigner-function** of the nucleus.

$$P(x, p) = \frac{1}{\pi\hbar} \int_{-\infty}^{\infty} dy \psi^*(x + y)\psi(x - y)e^{2ipy/\hbar}$$

- A point for the discussion: *Should we still call it “coalescence” if there is no (or negligible) dependence on the internal structure of the formed object?*

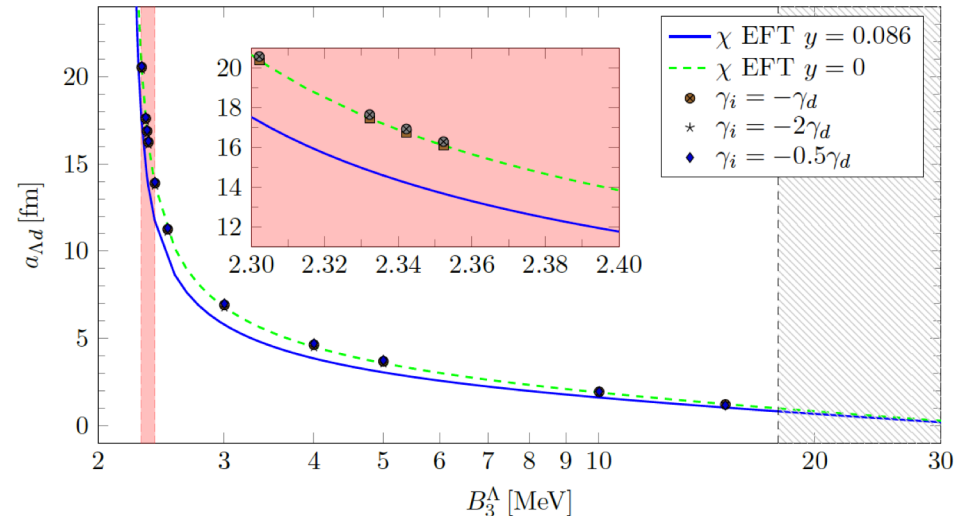
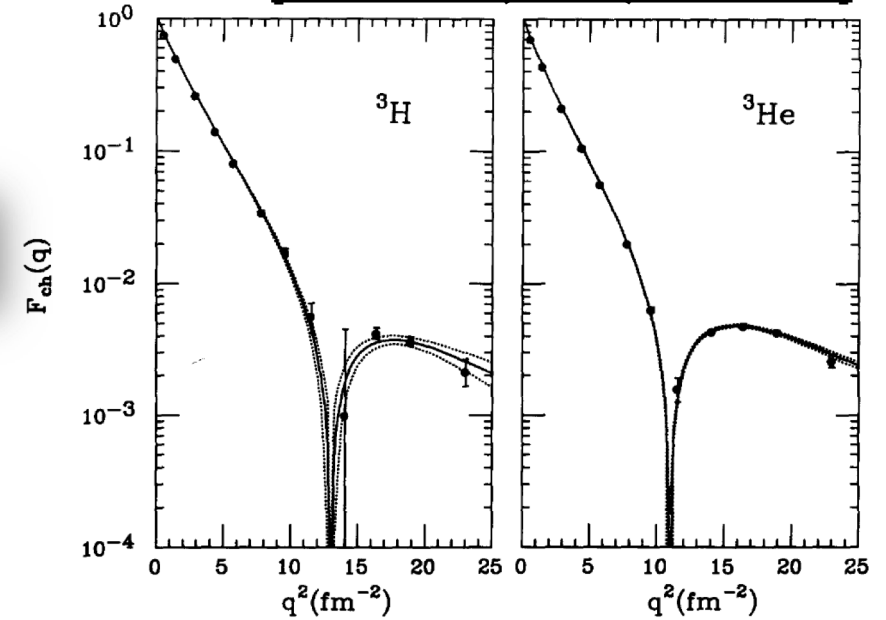
# Measuring the wave-function of nuclei

- Data from (old) electron scattering experiments for deuterons, tritons,  $^3\text{He}$ ,  $^4\text{He}$ .

The design of a tritium target containing of the order of 370 TBq (10 000 curies) of  $^3\text{H}$  required a thorough study of options and safety considerations.

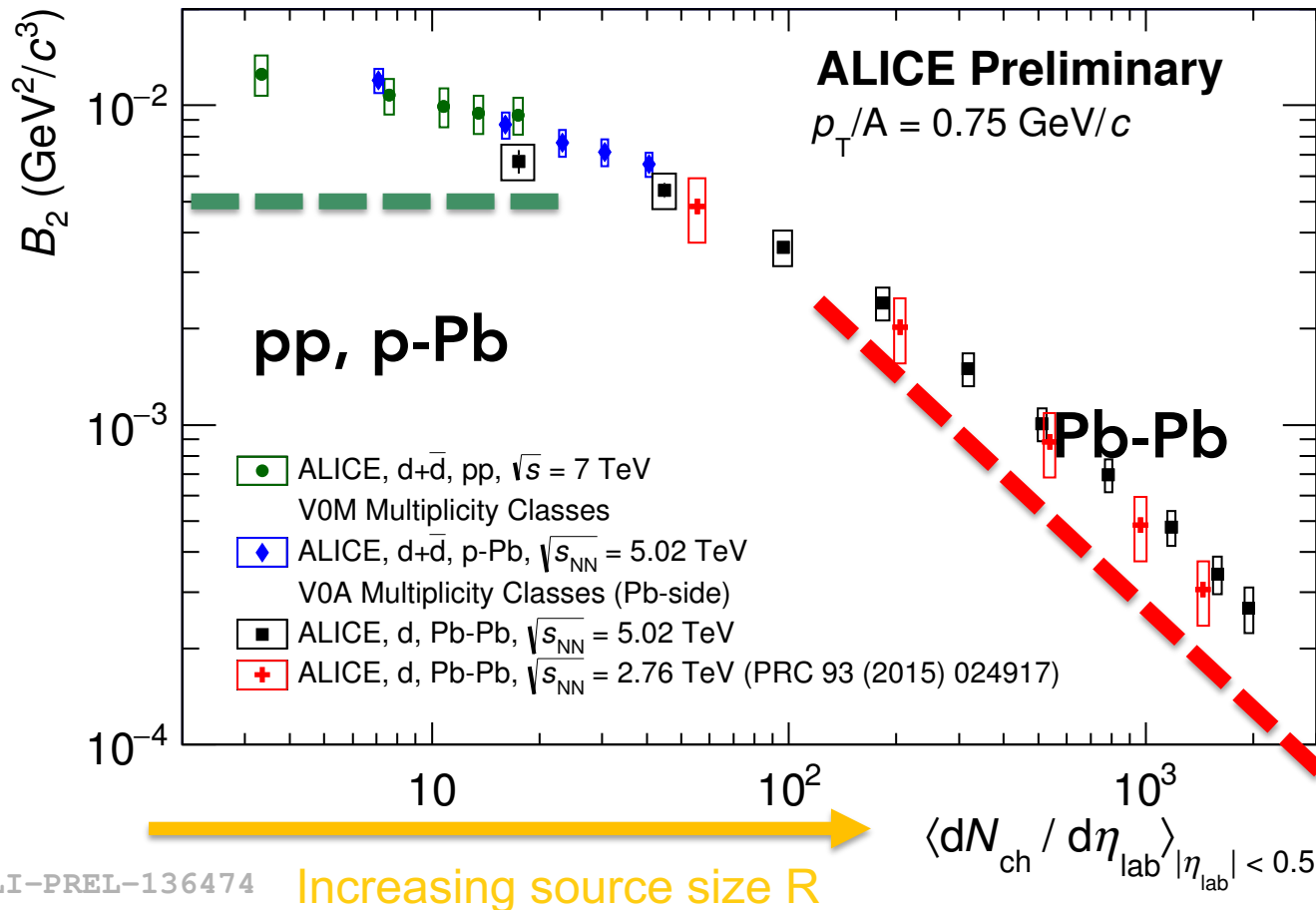
- For hyper-triton, the X(3872), and other QCD exotics, we have to rely on theory.
- A win-win situation?
  - If the coalescence model will turn out to be the correct description in small systems, can we then infer the wave-function of any QCD bound state by relating its production to the production of its constituents?
  - If thermal-statistical production turns out to be the correct scenario, we have to rethink the nature of QCD bound states.

[NPA 579 (1994) 596-626]



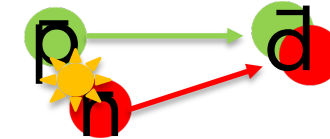
[Phys. Rev. C 100, 034002 (2019)] 25

# Coalescence models in heavy-ion collisions

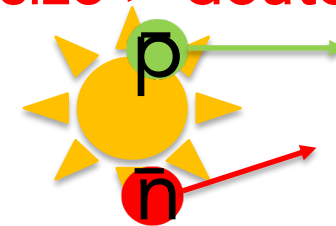


→ Two production regimes observed:

(a.) system size < deuteron size



(b.) system size > deuteron size



The trend with multiplicity is explained as an increase in the source size  $R$  in coalescence models

(e.g. *Scheibl, Heinz PRC 59 (1999) 1585*).

→ Strong dependence of  $B_2$  on collision geometry.

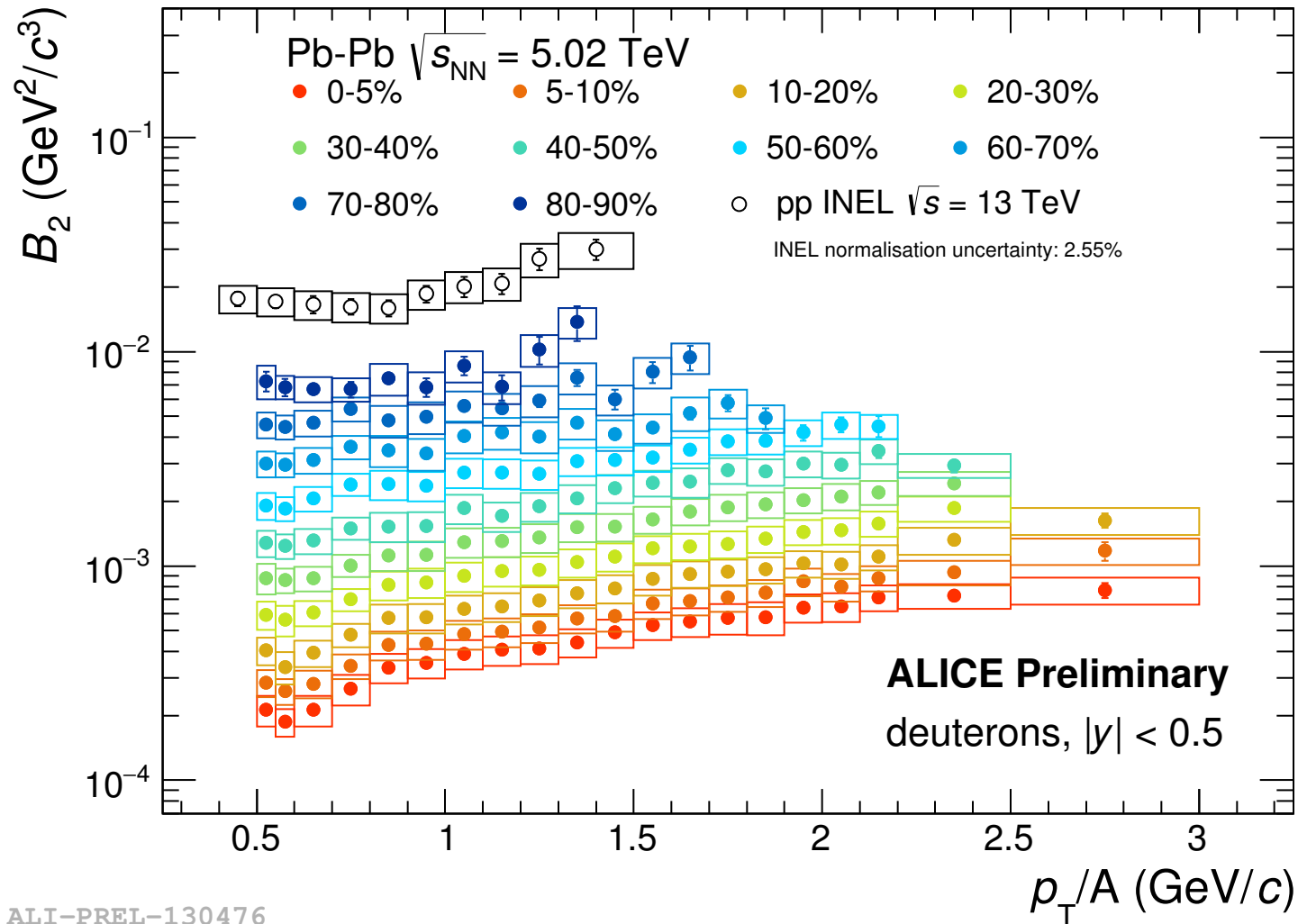
$$N_A \sim (R^2 + r_A^2/4)^{-3/2}$$

$R$  – size of the system

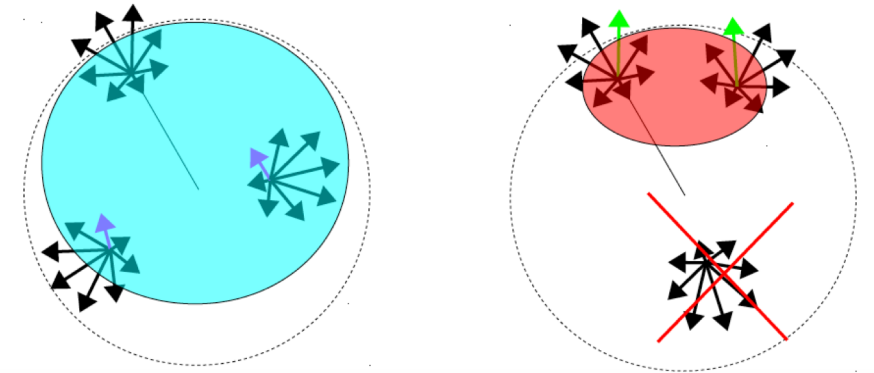
$r_A$  – size of nucleus



# $p_T$ dependence of $B_2$ in Pb-Pb



- A particle emitted from a medium will have a collective velocity  $\beta_f$  and a thermal (random) one  $\beta_t$
- As observed  $p_T$  grows, the region from where pairs with small relative momentum can be emitted gets smaller and shifted to the outside of the source



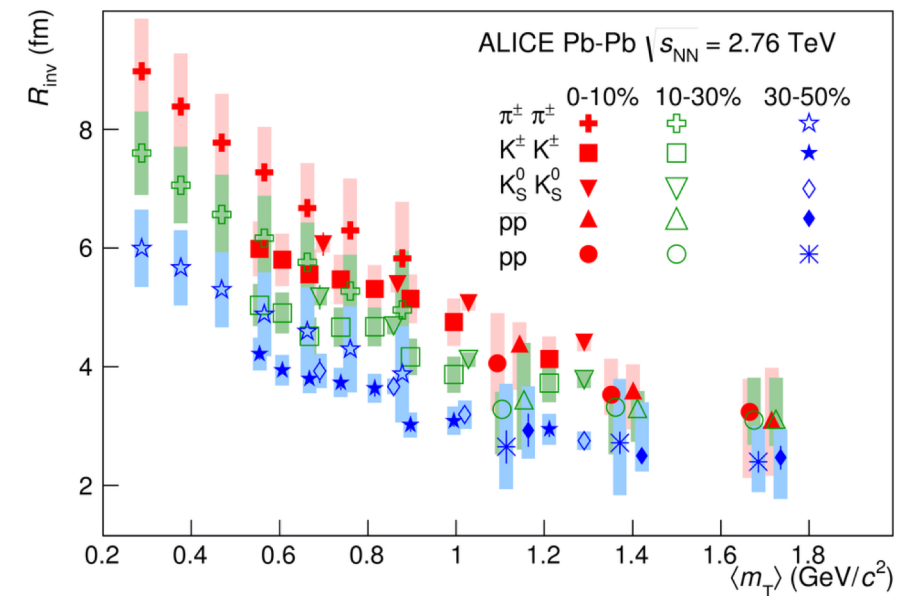
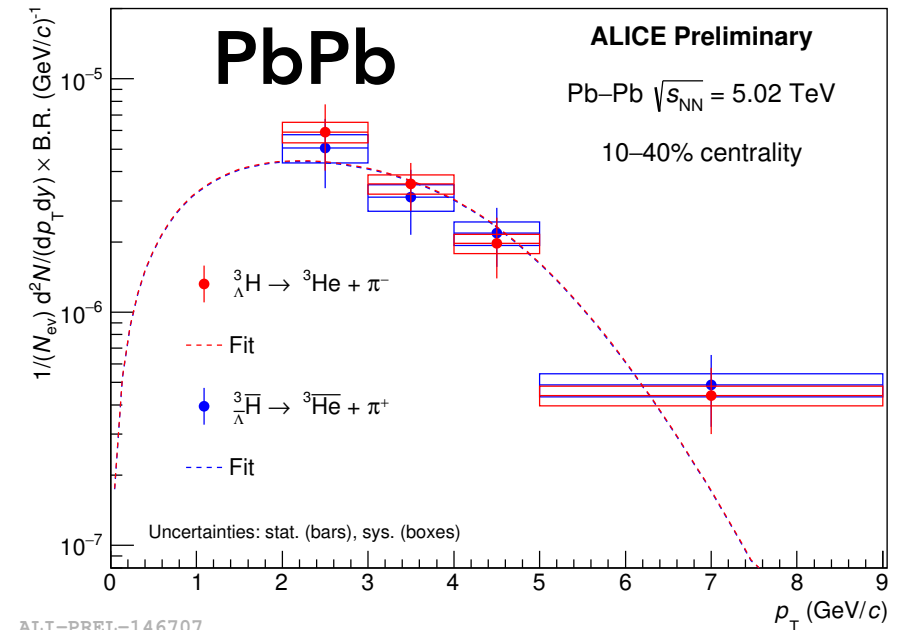
A. Kisiel, Hirscheegg 2019

High momenta means  $\Rightarrow$  thermal (random) and collective component must align  
 $\Rightarrow$  more likely that they come from the same volume element  
 $\Rightarrow$  higher coalescence probability

# Typical momenta and HBT radii

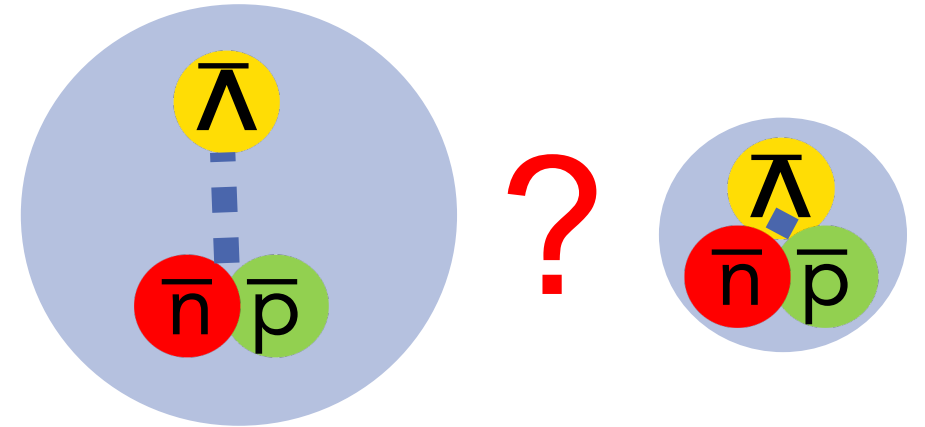
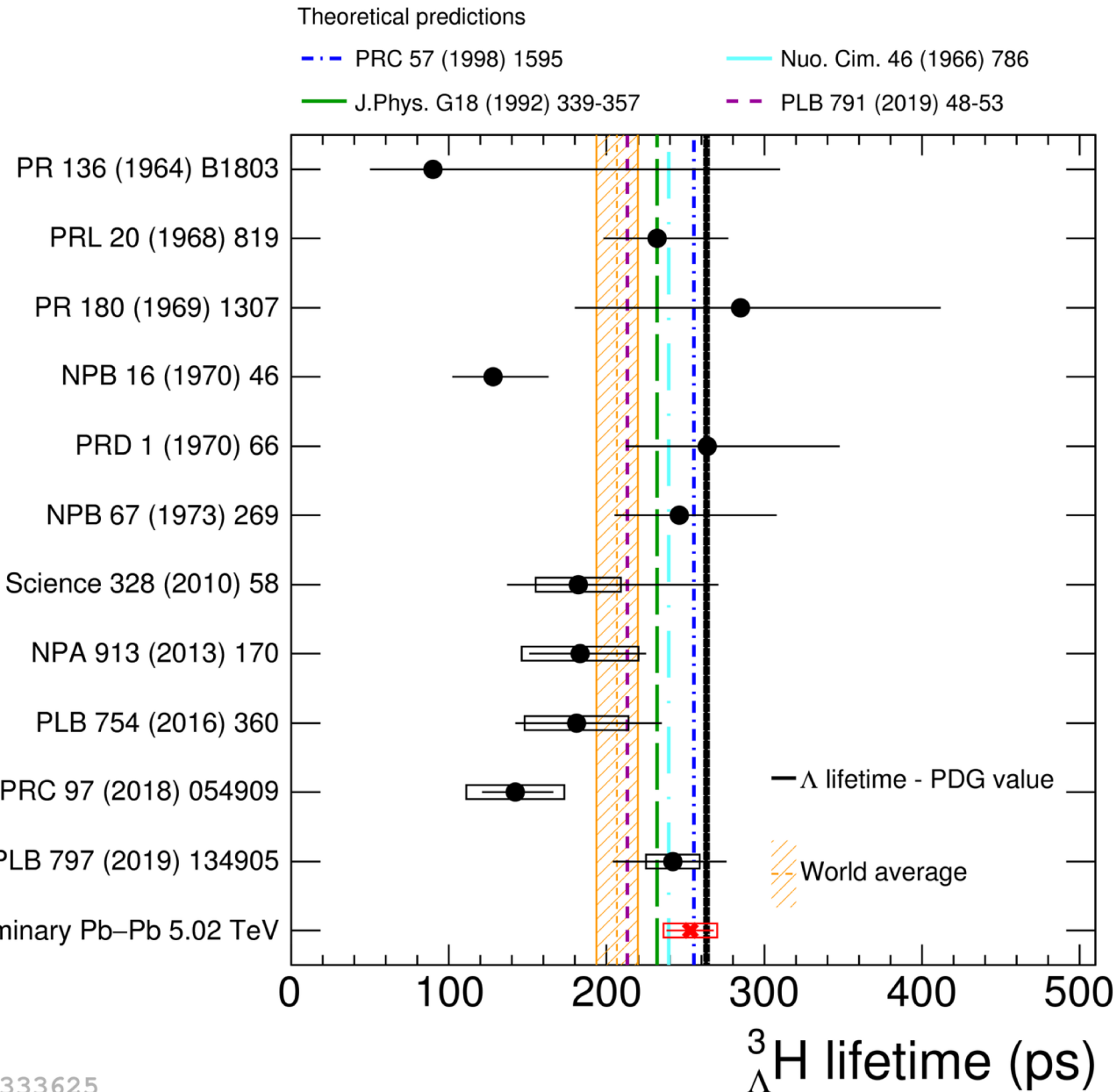
- Roughly speaking, the maximum of the production probability in a range around 0.5 to 2.0 GeV/c per nucleon (system size dependent!).
- This corresponds to an  $m_T$  range around 1 to 2 GeV/c and thus to radii around 2-6 fm.

→ Important to do such "rough" studies precisely in future and it is one of the goals of the workshop to establish a roadmap for these projects.



Towards more exotic QCD bound states:  
The  $X(3872)$  and the anti-hyper-triton

# (Anti-)Hypertriton lifetime measurement from ALICE

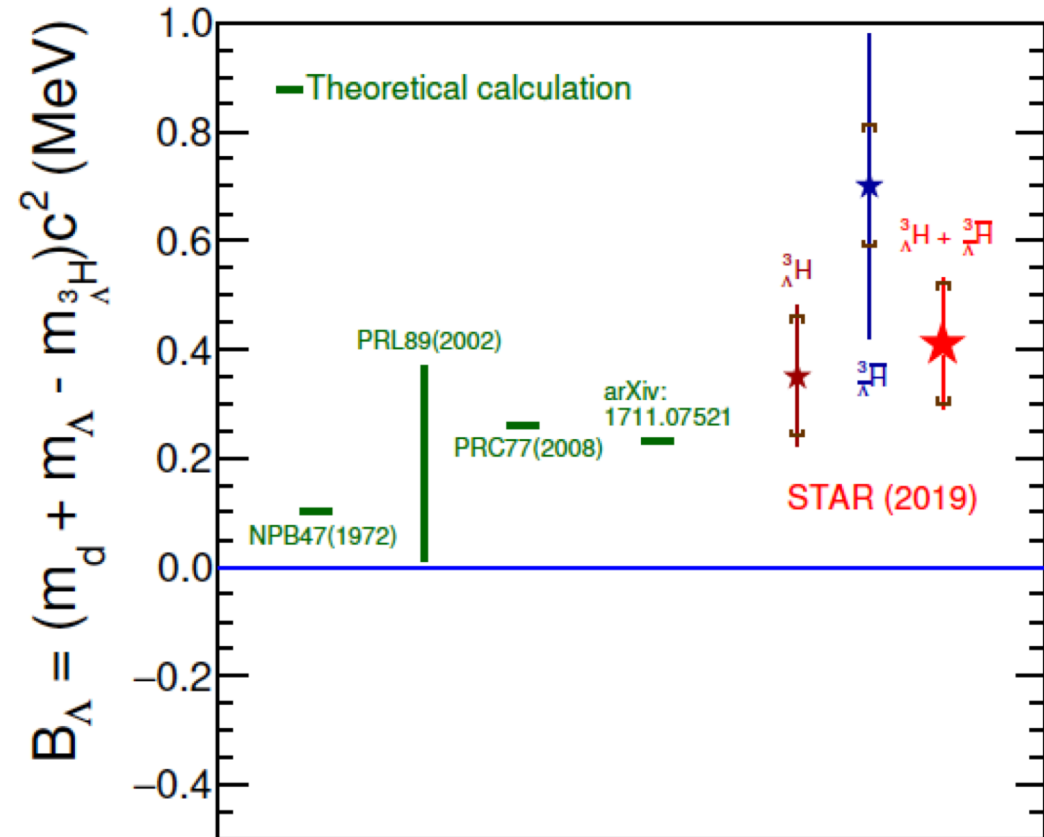
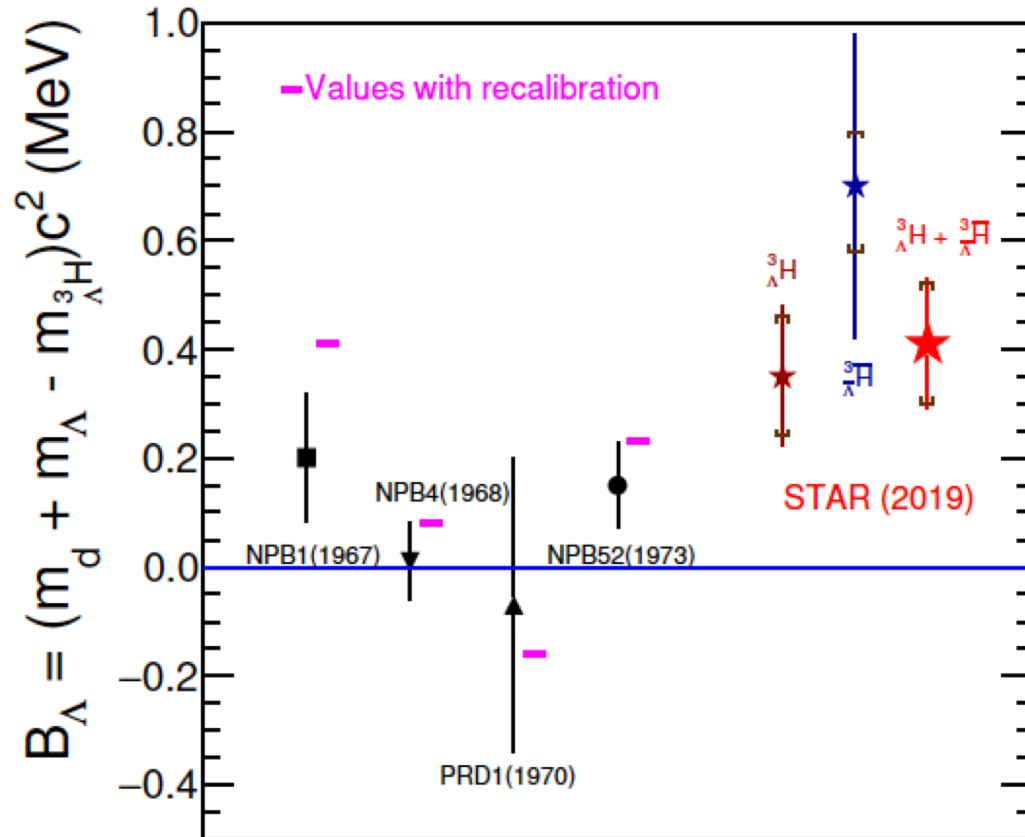


$m = 2.991 \text{ GeV}/c^2$ ,  $B_\Lambda = 130 \text{ keV}$   
 $\rightarrow$  rms-radius = 10.3 fm

- Exclude large deviations from free  $\Lambda$  life time
- Test of different models with different Hypertriton structure and final state interaction

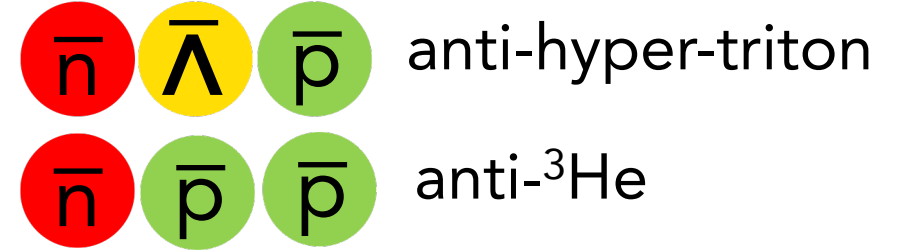
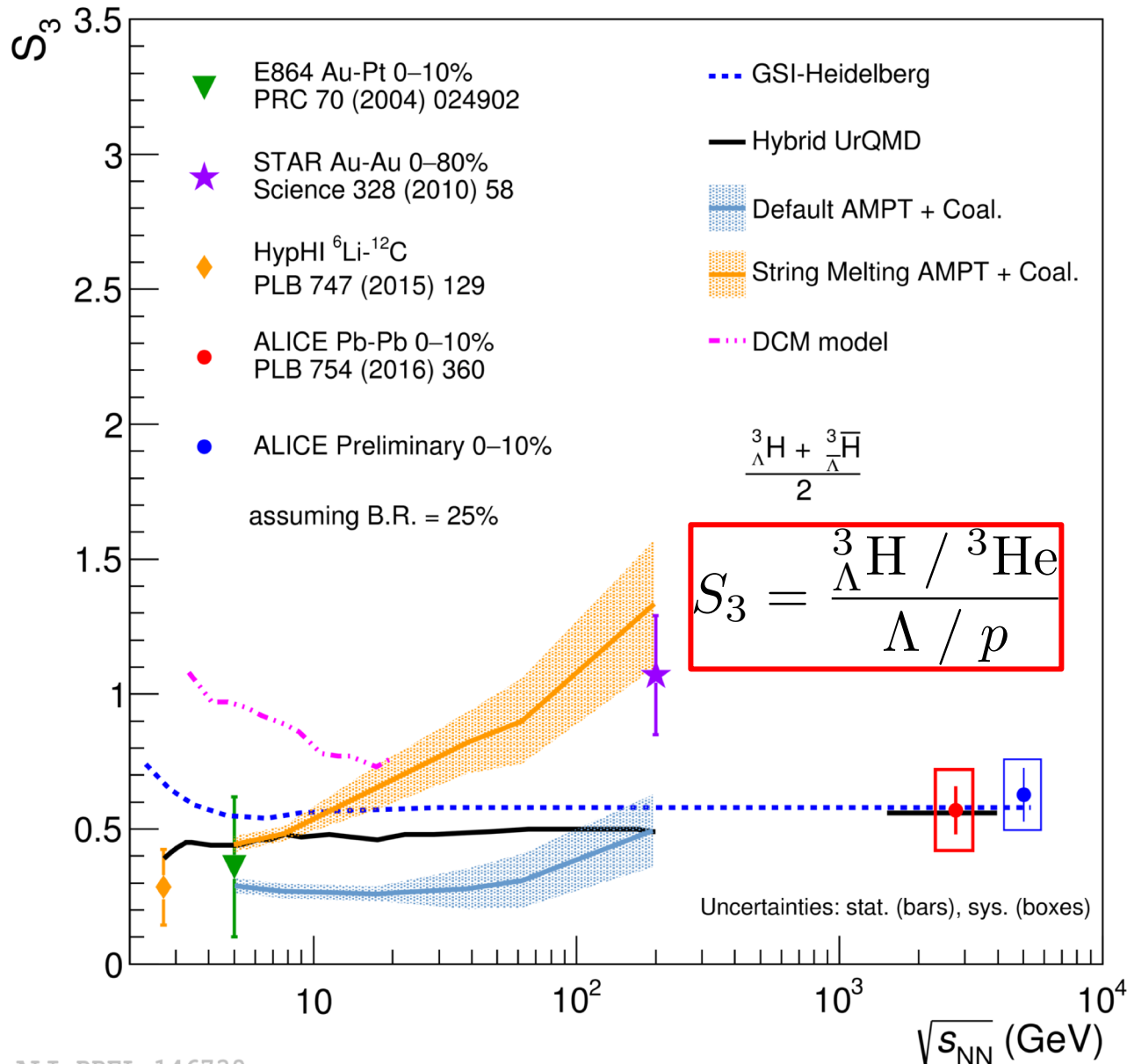
# (anti-)Hyper-triton measurements from STAR

[STAR, *Nature Phys.* 16 (2020) 4]



How precisely do we really know the Lambda separation energy in the hyper-triton?  
 → Homework for the experimentalists..

# (anti-)hyper-triton in Pb-Pb collisions at 5.02 TeV



→ Yields of heavy and fragile objects such as (anti-)(hyper-)nuclei in agreement with thermal-statistical model predictions at *chemical* freeze-out.

→ No re-scattering of anti-nuclei in hadronic phase despite large dissociation cross-section.

→ Final-state coalescence after kinetic freeze-out requires more detailed modeling: *naive coalescence* ( $S_3 \approx 1$ ) does not describe data.

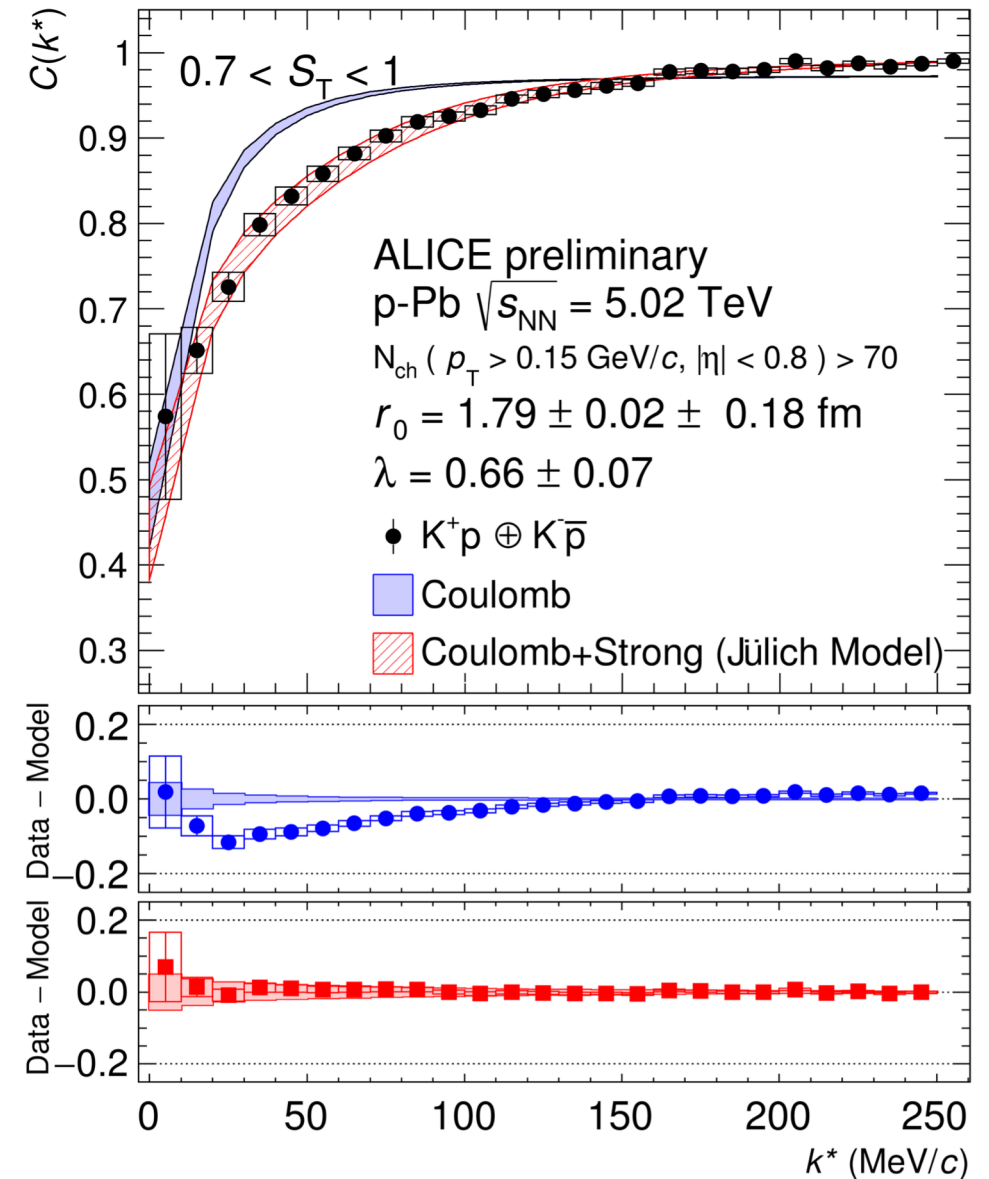
# Correlation functions

Already now, ALICE is investigating anti-nuclei correlations (e.g. with kaons) and these studies can be taken to a next level.

→ The formation of any exotic bound QCD object can be studied in

**unprecedented detail** by studying correlation functions of for instance

$(d, \Lambda) \rightarrow {}^3_{\Lambda}\text{H}$  or  
 $({}^3\text{He}, \Lambda) \rightarrow {}^4_{\Lambda}\text{He}$ .

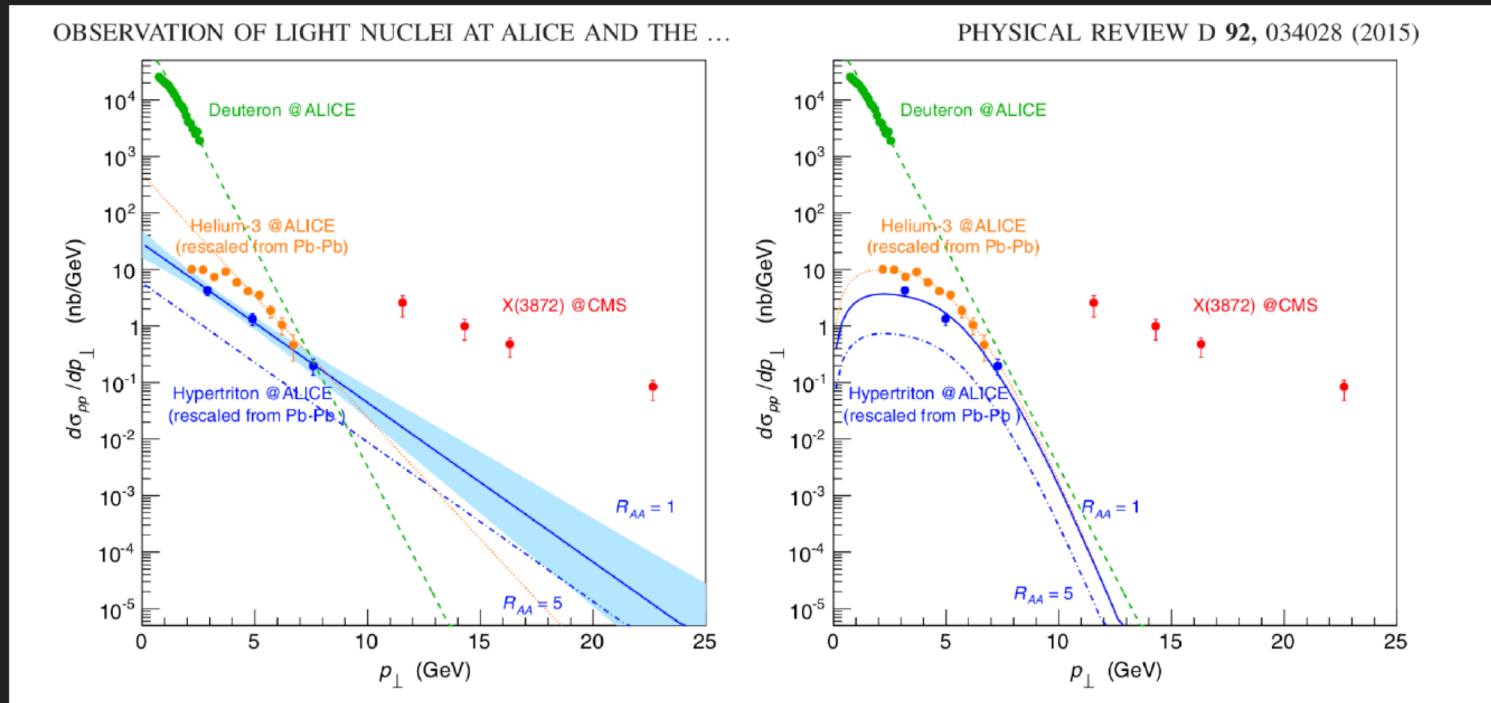


ALI-PREL-316331

# The X(3872)

The X(3872) sort of anomalous charmonium with  $1^{++}$  quantum numbers  
**right at**  $DD^*$  threshold and rather close to  $J/\psi + \rho$ .

[Talk of A. Polosa at  
the charm hadronization  
workshop]



Esposito et al. PRD92 (2015) 034028

Bignamini, Grinstein, Piccinini, ADP, Sabelli, PRL103 (2009) 162001

Esposito, Grinstein, Maiani, Piccinini, Pilloni, ADP, Riquer, 1709.09631

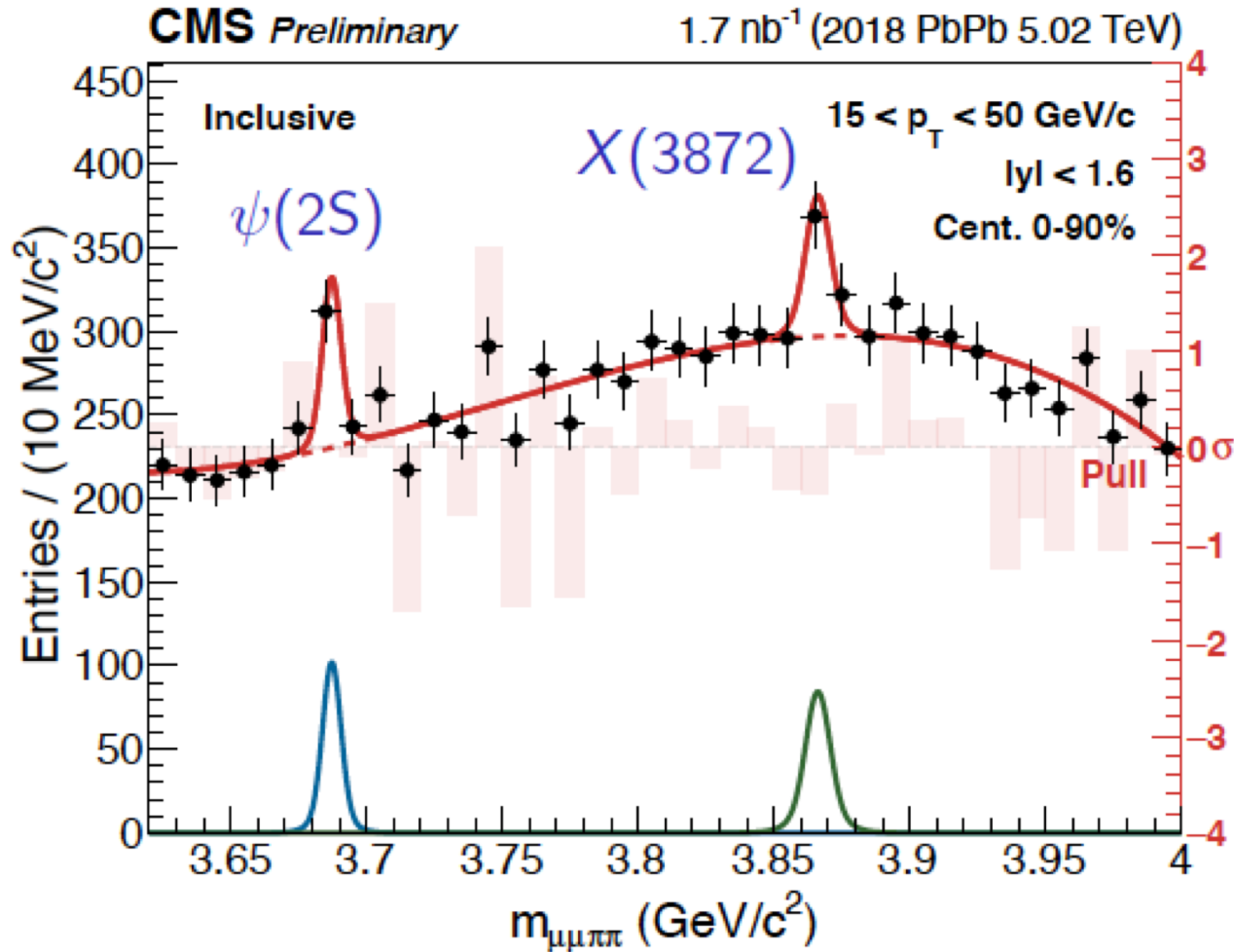


# Loosely bound objects in heavy-ion collisions

[Y. Jie Lee]

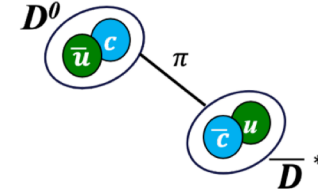
[J. Durham]

→ The X(3872) will not be suppressed if its behavior is like a hyper-triton!



• Mass is consistent with sum of  $D^0$  and  $\bar{D}^{*0}$  masses:  
 $M_{\chi_{c1}(3872)} - (M_{D^0} + M_{\bar{D}^{*0}}) = 0.01 \pm 0.27 \text{ MeV}$

*$D^0 \bar{D}^*$  Molecule*

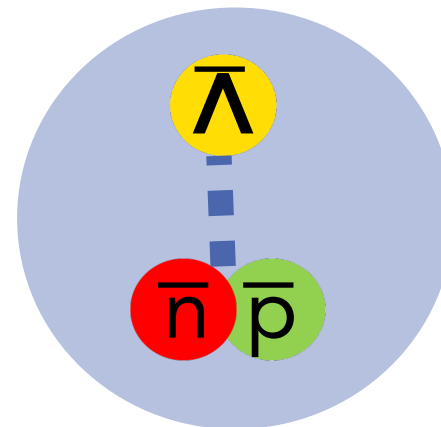


*VERY small binding energy  
VERY large radius, ~7 fm*

*Compact tetraquark*



*Tightly bound via color exchange between diquarks  
Small radius, ~1 fm*



$m = 2.991 \text{ GeV}/c^2$ ,  $B_\Lambda = 130 \text{ keV}$   
 → rms-radius = 10.3 fm

# Another snowball in hell?

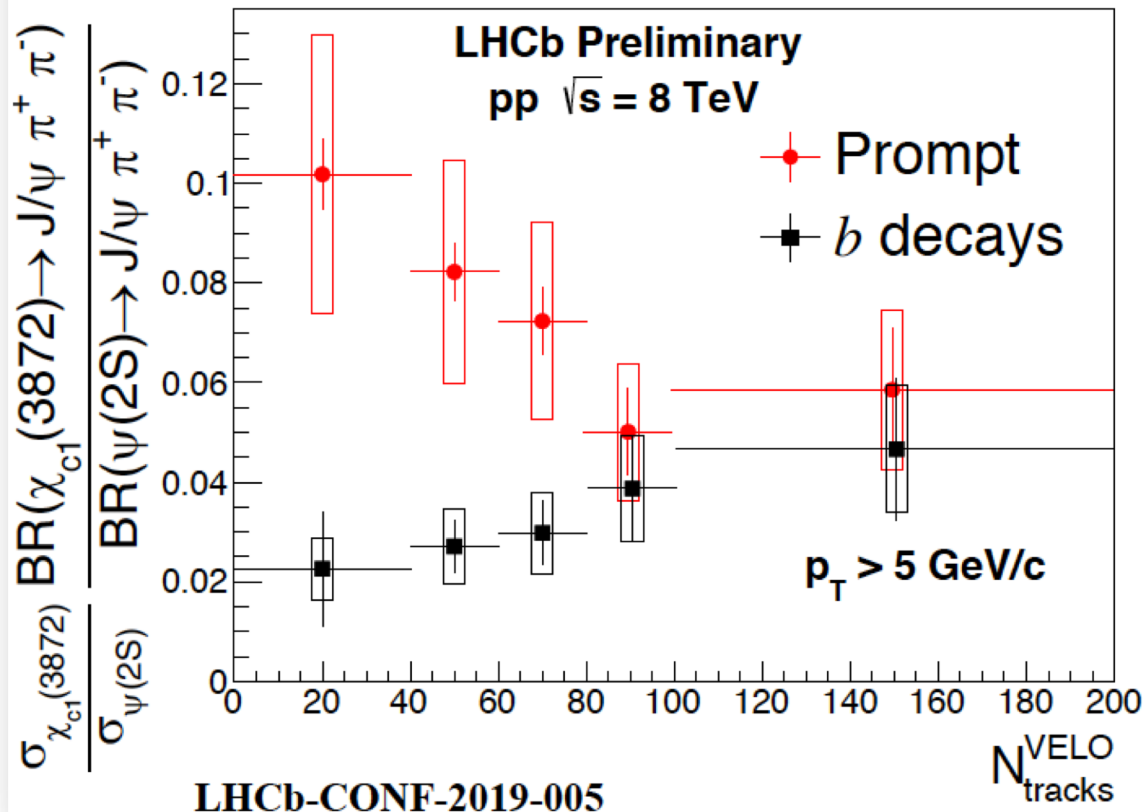
→ The X(3872) will not be suppressed if its behavior is like a hyper-triton!



## Ratio of cross sections



$$\frac{\sigma_{\chi_{c1}(3872)}}{\sigma_{\psi(2S)}} \times \frac{\mathcal{B}[\chi_{c1}(3872) \rightarrow J/\psi \pi^+ \pi^-]}{\mathcal{B}[\psi(2S) \rightarrow J/\psi \pi^+ \pi^-]} = \frac{N_{\chi_{c1}(3872)} f_{\text{prompt}}^{\chi_{c1}(3872)}}{N_{\psi(2S)} f_{\text{prompt}}^{\psi(2S)}} \times \frac{\epsilon_{\psi(2S)}}{\epsilon_{\chi_{c1}(3872)}}$$



**Prompt Component:**  
 Increasing suppression of **X(3872)** production relative to  **$\psi(2S)$**  as event activity increases

***b*-decay component:**  
 No significant change in relative production, as expected for decays in vacuum. Ratio is set by ***b*** decay branching fractions.

Consistent with ATLAS measurement  
 $R = 0.0395 \pm 0.0032 \pm 0.0008$  ( $p_T > 10$  GeV/c)

JHEP 2017:117 (2017)

# Summary and conclusions

- A lot of and fast evolving experimental and phenomenological activity in this research area.
- After the reporting of the initial findings, it is now time to enter an era of increased precision in comparisons of data and theory:
  1. Remove approximations, e.g. by measuring HBT radii and production yields in the same events
  2. If approximations cannot be removed, they need to be quantified, e.g. uncertainties in the hyper-triton wave-function.

Thank you and enjoy the workshop!