

Measuring the speed of sound using cumulants of baryon number

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In collaboration with
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04/05/2022



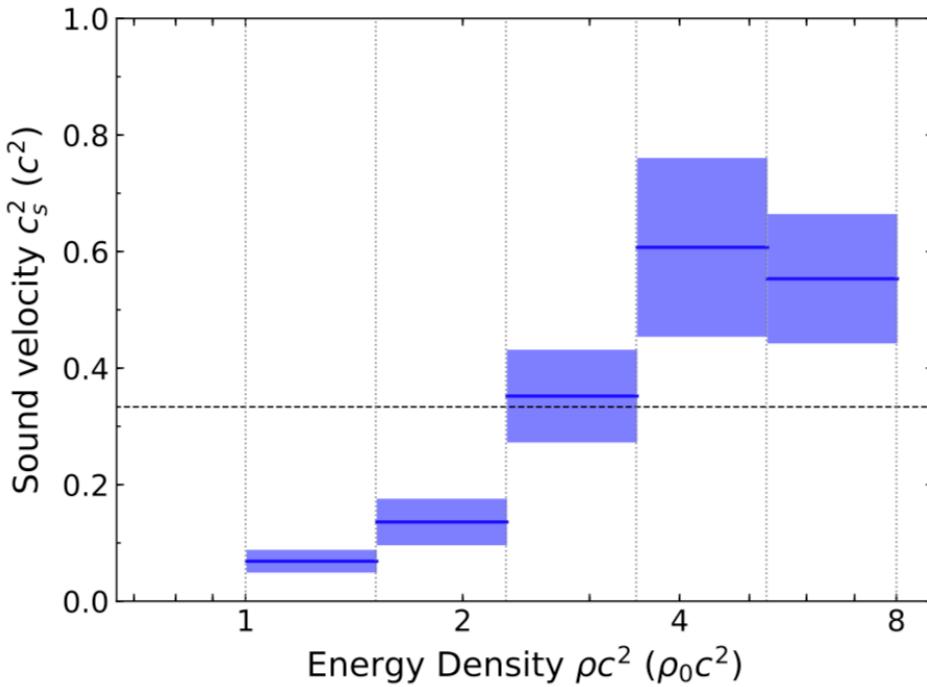
c_s^2 in neutron stars and heavy-ion collisions

In neutron stars: based on

- neutron stars with $M \gtrsim 2M_\odot$
 - knowledge of the EOS of hadronic matter at “low” densities
- c_s^2 may exceed the conformal limit of $1/3$

P. Bedaque and A. W. Steiner, Phys. Rev. Lett. **114**, no.3, 031103 (2015), arXiv: 1408.5116, Bedaque:2014sqa

I. Tews, J. Carlson, S. Gandolfi and S. Reddy, Astrophys. J. **860**, no.2, 149 (2018), arXiv:1801.01923, Tews:2018kmu

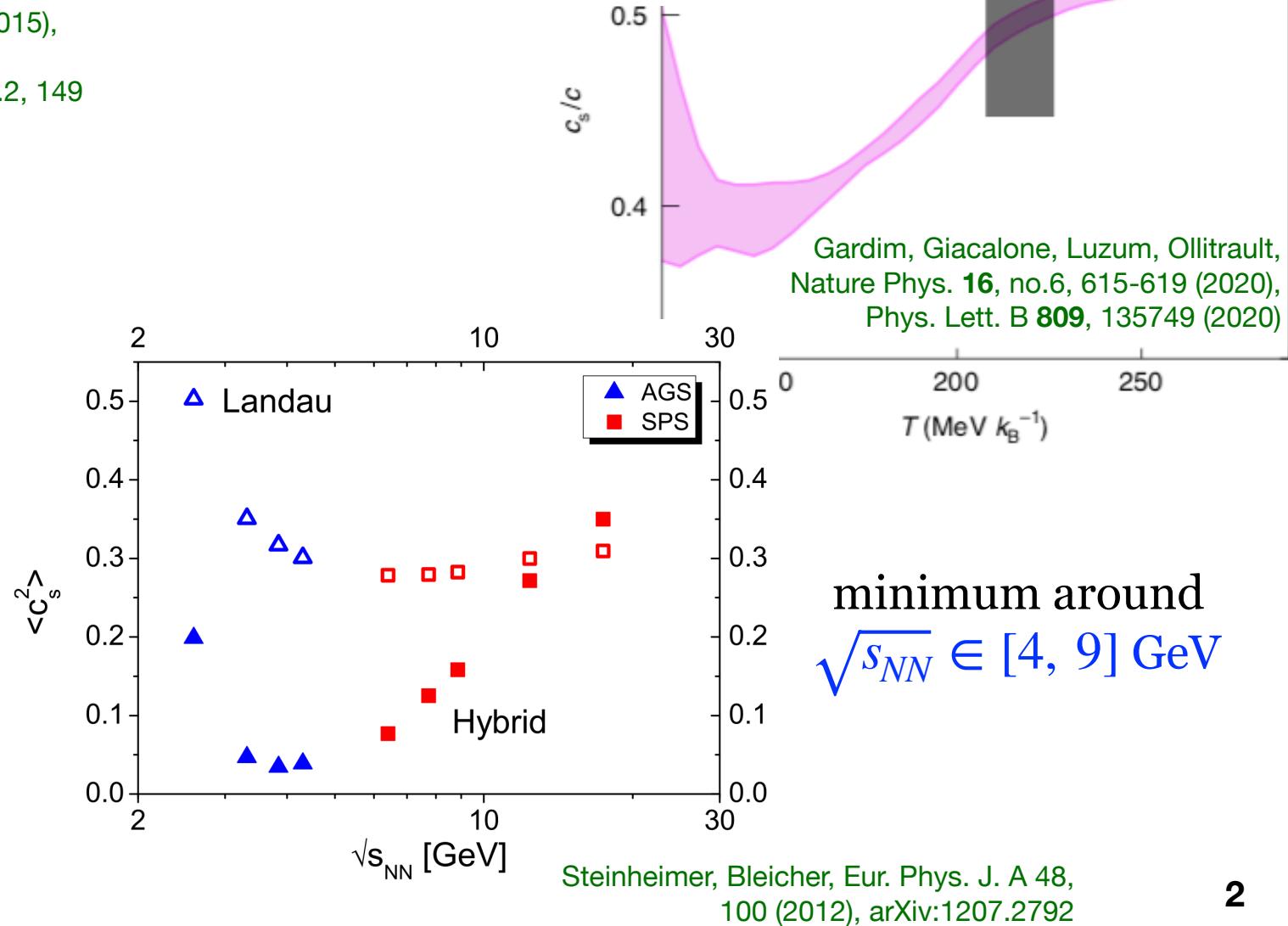


Fujimoto, Fukushima, Murase, Phys. Rev. D **101**, no.5, 054016 (2020), arXiv:1903.03400

Agnieszka Sorensen (INT)

In heavy-ion collisions:

$$c_s^2 \Big|_V \approx \frac{d \ln \langle p_T \rangle}{d \ln N_{\text{ch}}} = 0.24 \pm 0.04$$



How to measure c_s^2 in heavy-ion collisions?

At finite T , one often considers the following two expressions for the speed of sound:

the isentropic speed of sound
($\sigma = s/n_B = \text{const}$):

$$c_\sigma^2 \equiv \left(\frac{dP}{d\mathcal{E}} \right)_\sigma = \frac{\frac{s}{n_B} \left(\frac{dP}{dT} \right)_{n_B} + \left(\frac{dP}{dn_B} \right)_T \left(\frac{ds}{dT} \right)_{n_B} - \left(\frac{dP}{dT} \right)_{n_B} \left(\frac{ds}{dn_B} \right)_T}{\left(\frac{\mathcal{E} + P}{n_B} \right) \left(\frac{ds}{dT} \right)_{n_B}}$$

the isothermal speed of sound
($T = \text{const}$):

$$c_T^2 \equiv \left(\frac{dP}{d\mathcal{E}} \right)_T = \frac{\left(\frac{dP}{dn_B} \right)_T}{T \left(\frac{ds}{dn_B} \right)_T + \mu_B}$$

common $T = 0$ limit:

$$c_T^2 \Big|_{T=0} = c_\sigma^2 \Big|_{T=0} = \frac{n_B}{\mu_B} \left(\frac{d\mu_B}{dn_B} \right)_T = \frac{1}{\mu_B} \left(\frac{dP}{dn_B} \right)_T = \frac{T}{\mu_B} \frac{\kappa_1}{\kappa_2}$$

$$\kappa_1 = V n_B$$

Easy to notice: the speed of sound \sim derivatives of pressure

$$\kappa_2 = V T n_B \left(\frac{dP}{dn_B} \right)_T^{-1}$$

What “measures” derivatives of pressure? **Cumulants of baryon number!**

Can one connect the speed of sound with the cumulants?

Problematic: difficult to estimate T derivatives of cumulants from experiment

$$c_T^2 = \left[\frac{T}{\kappa_1} \left(\frac{d\kappa_1}{dT} \right)_{\mu_B} + \frac{\mu_B}{T} \frac{\kappa_2}{\kappa_1} \right]^{-1}$$

$$T/\mu_B \ll 1$$

$$c_T^2 \approx \frac{T\kappa_1}{\mu_B\kappa_2}$$

an upper limit to c_T^2

$$\left(\frac{d \ln c_T^2}{d \ln n_B} \right)_T + c_T^2 \approx 1 - \frac{\kappa_3 \kappa_1}{\kappa_2^2}$$

AS, D. Oliinychenko, V. Koch, L. McLerran, Phys. Rev. Lett. **127** (2021) 042303, arXiv:2103.07365, Sorensen:2021zme

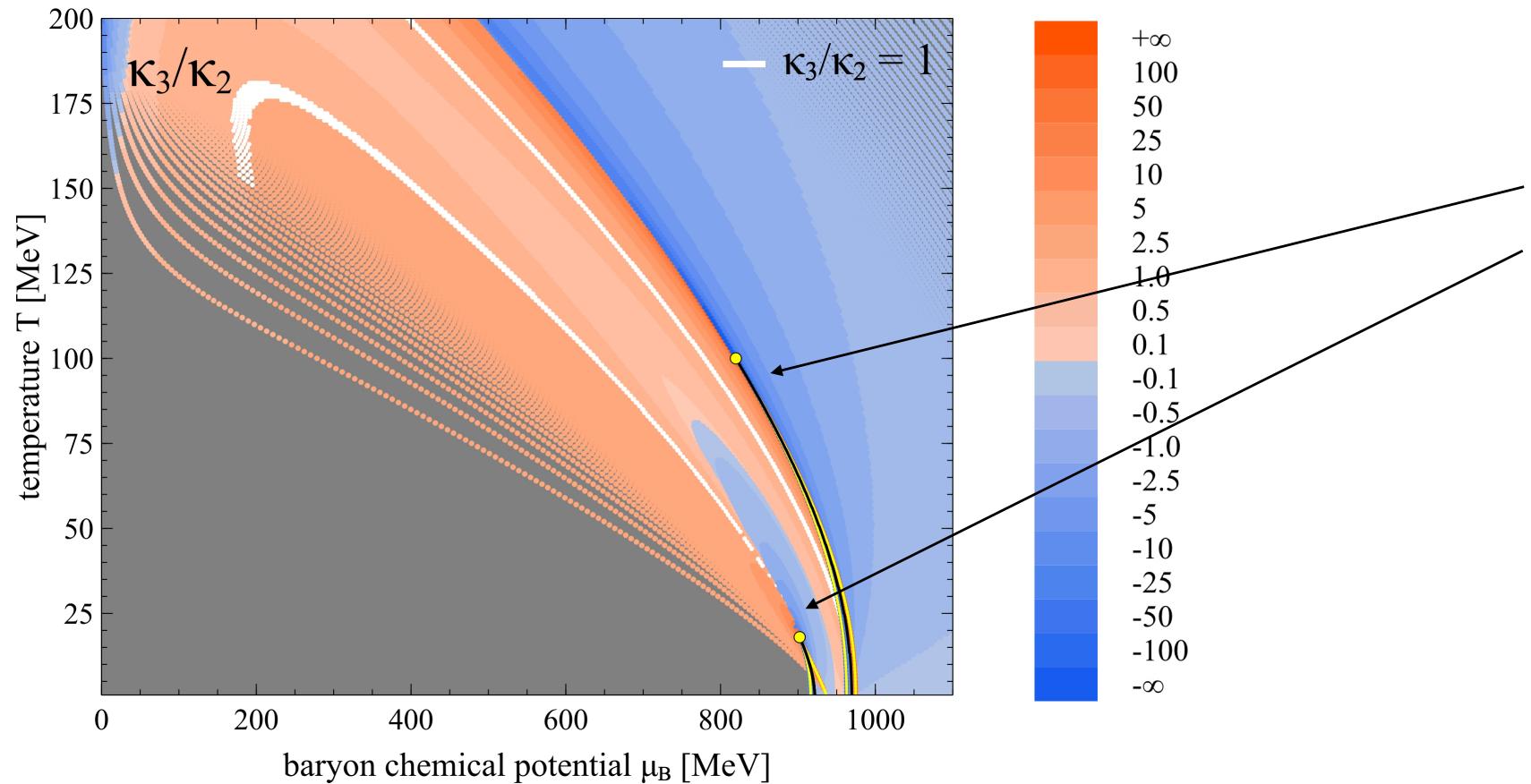
From this arose a somewhat audacious idea:
to measure c_s^2 using not only the qualitative,
but also the quantitative behavior of the cumulants

Tests in model(s)

AS, V. Koch, Phys. Rev. C **104** no. 3 (2021) 034904,
arXiv:2011.06635, Sorensen:2020ygf

the VDF (vector density functional) model:

$$P = g \int \frac{d^3 p}{(2\pi)^3} \frac{p^2}{3\sqrt{p^2 + m^2}} f_{\mathbf{p}} + \sum_{i=1}^4 C_i \frac{b_i - 1}{b_i} n_B^{b_i}$$

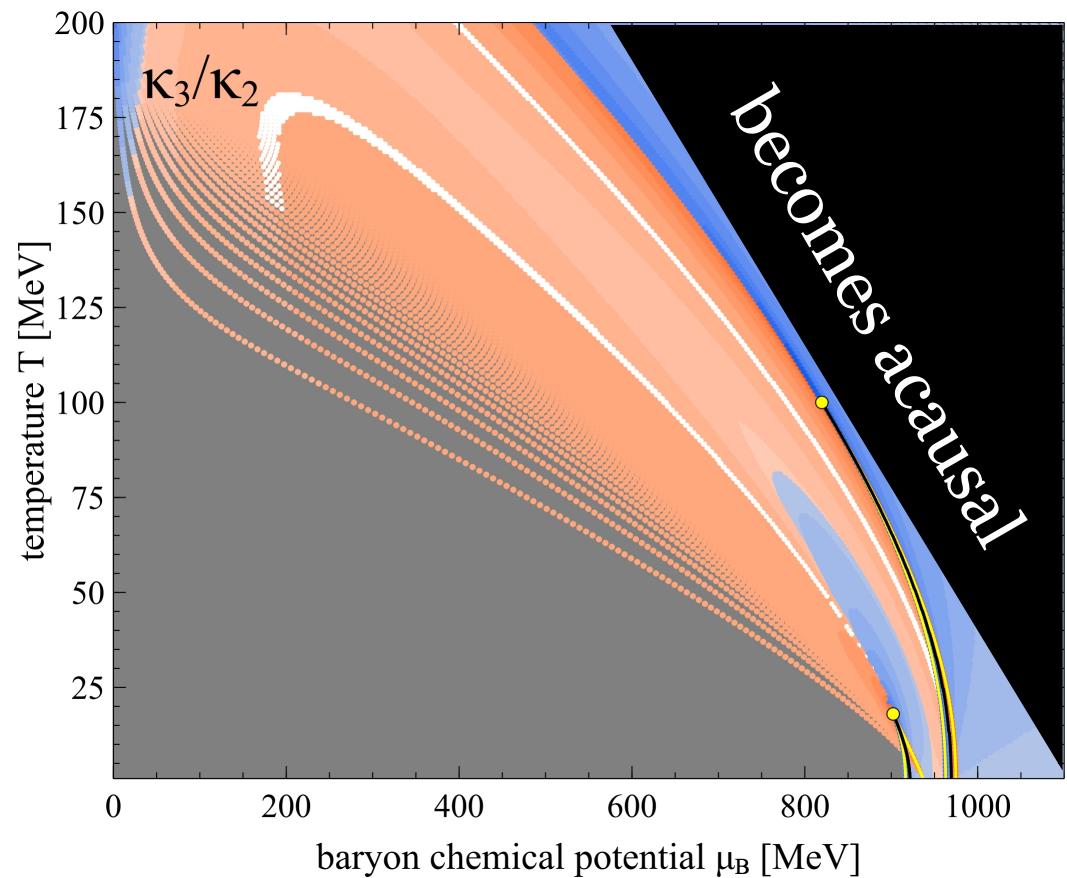


two phase transitions:
the known nuclear liquid-gas PT
and a “QGP-like” PT

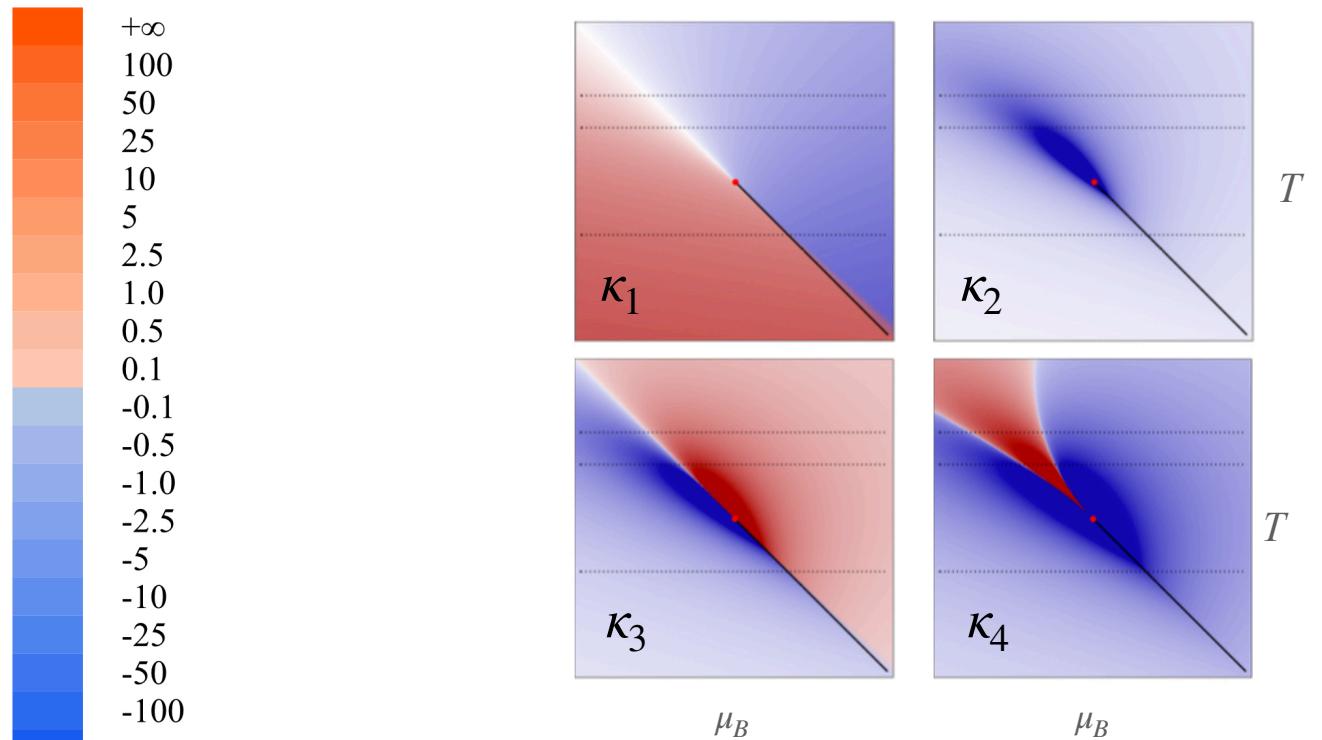
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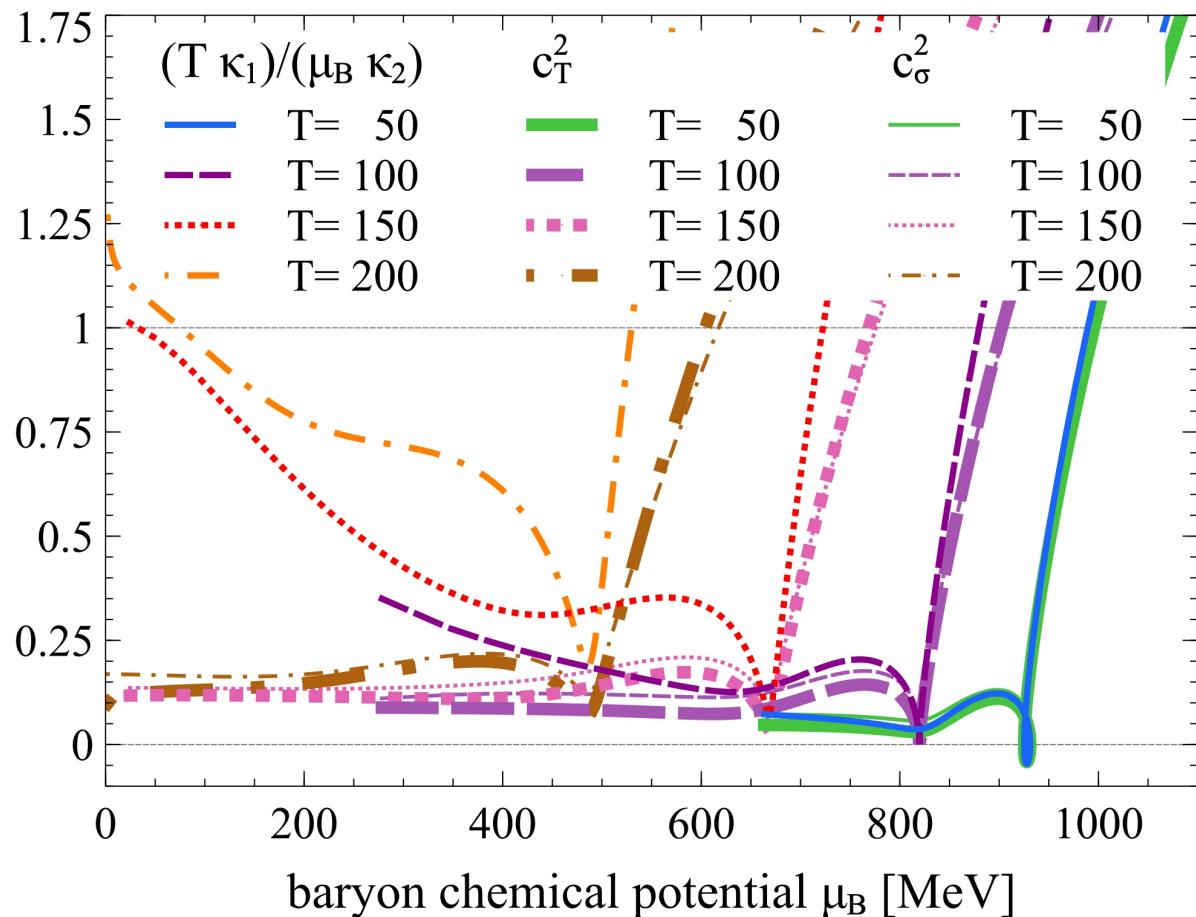
A. Bzdak *et al.*, Physics Reports **853** (2020) 1-87,
arXiv:1906.00936, Bzdak:2019pkr

orange = positive values, blue = negative values, white lines for values = 1 (Poisson limit),
yellow lines = spinodal regions, black lines for coexistence regions

Tests in the VDF model

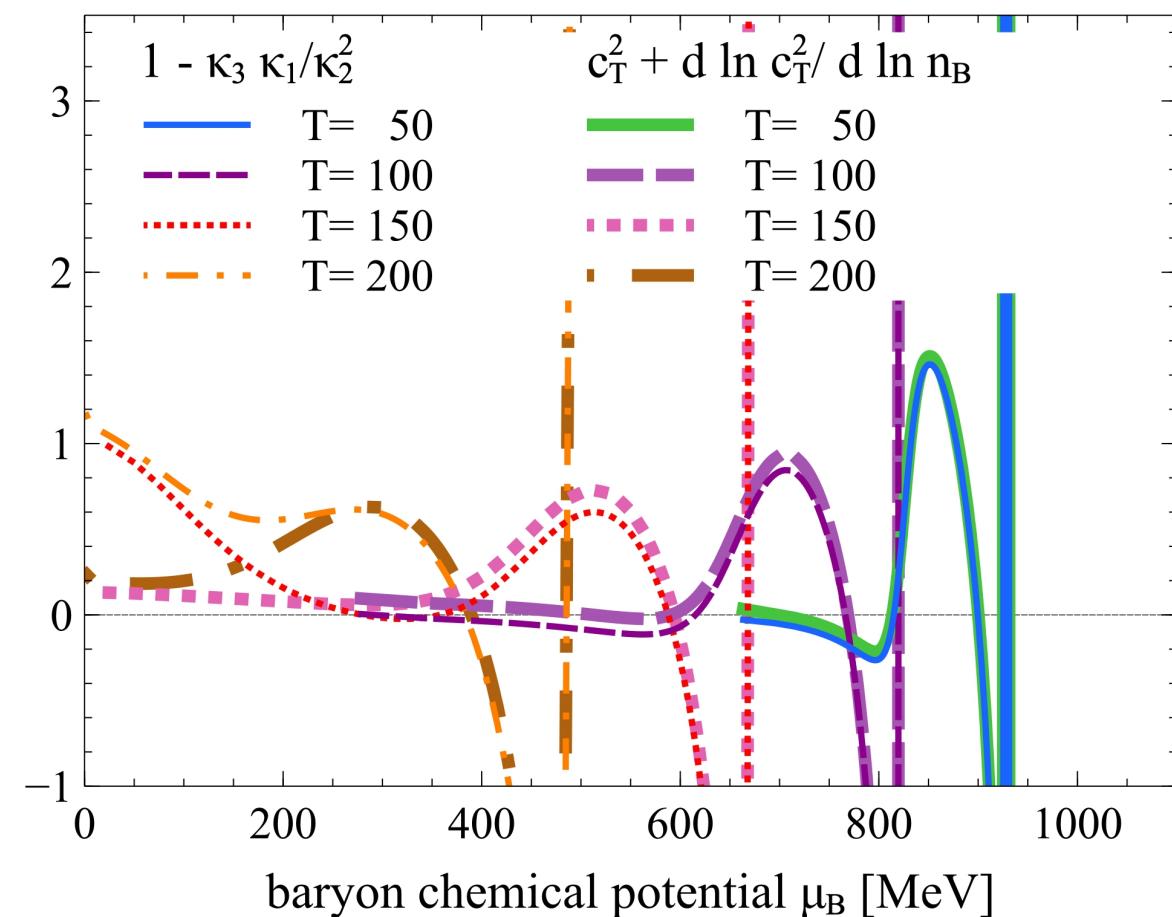
$$c_T^2 \approx \frac{T\kappa_1}{\mu_B\kappa_2}$$

works for
 $\mu_B \gtrsim 600$ MeV and $T < 150$

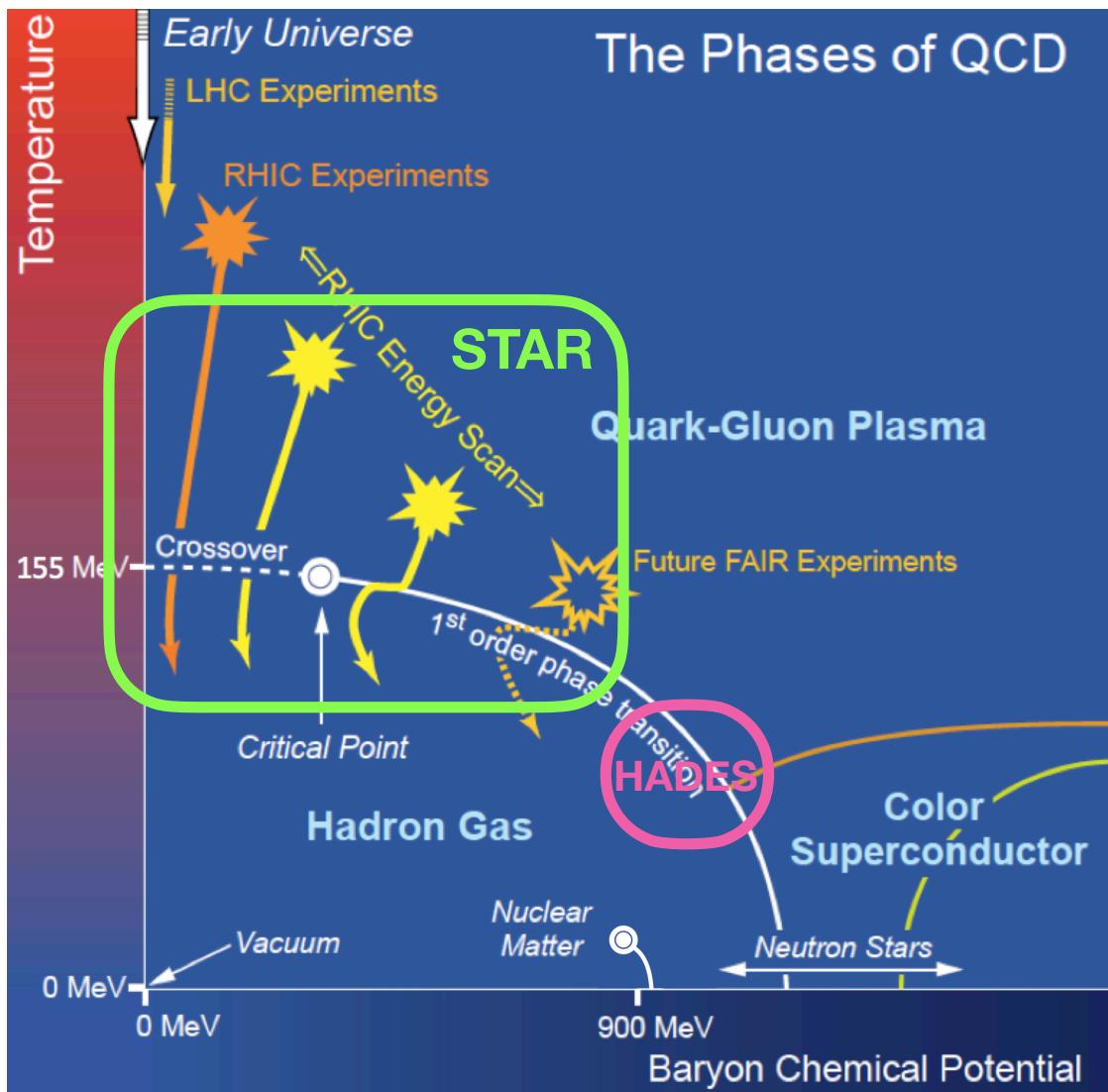


$$\left(\frac{d \ln c_T^2}{d \ln n_B} \right)_T + c_T^2 \approx 1 - \frac{\kappa_3 \kappa_1}{\kappa_2^2}$$

works for
 $\mu_B \gtrsim 200$ MeV



Experimental data



The freeze-out parameters (T_{fo}, μ_{fo}) are obtained from particle yields:

\sqrt{s} [GeV]	T_{fo} [MeV]	μ_{fo} [MeV]
200	164.3	28
62.4	160.3	70
54.4	160.0	83
39	156.4	103
27	155.0	144
19.6	153.9	188
14.5	151.6	264
11.5	149.4	287
7.7	144.3	398
2.4	65	784

Cumulants $\kappa_2, \kappa_3, \dots$ are fluctuations around the means

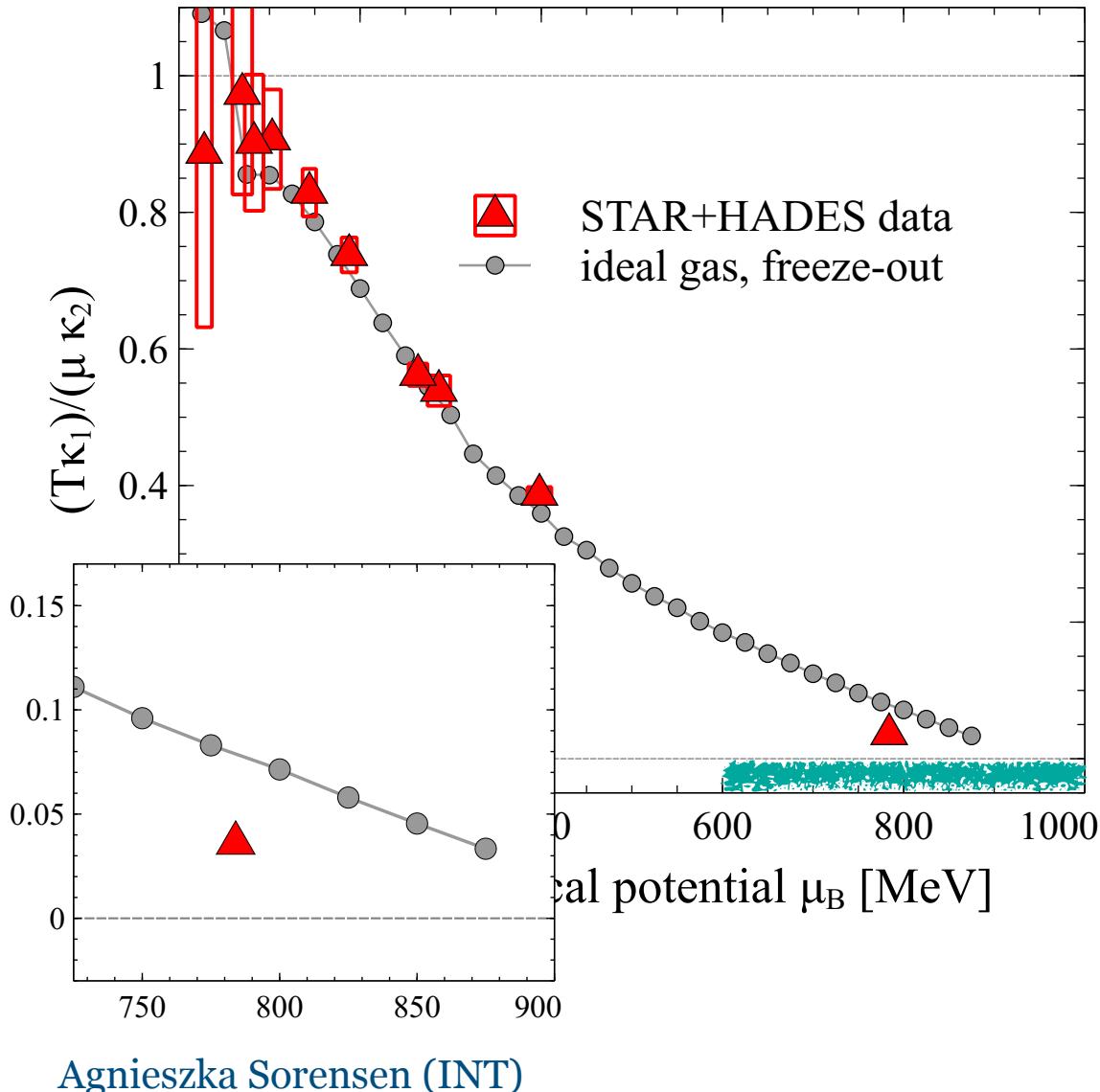
M. Abdallah *et al.* (STAR), Phys. Rev. C **104** (2021) no. 2 024902,
arXiv:2101.12413, STAR:2021iop

M. L. for the HADES collaboration (2019), 3rd EMMI Workshop

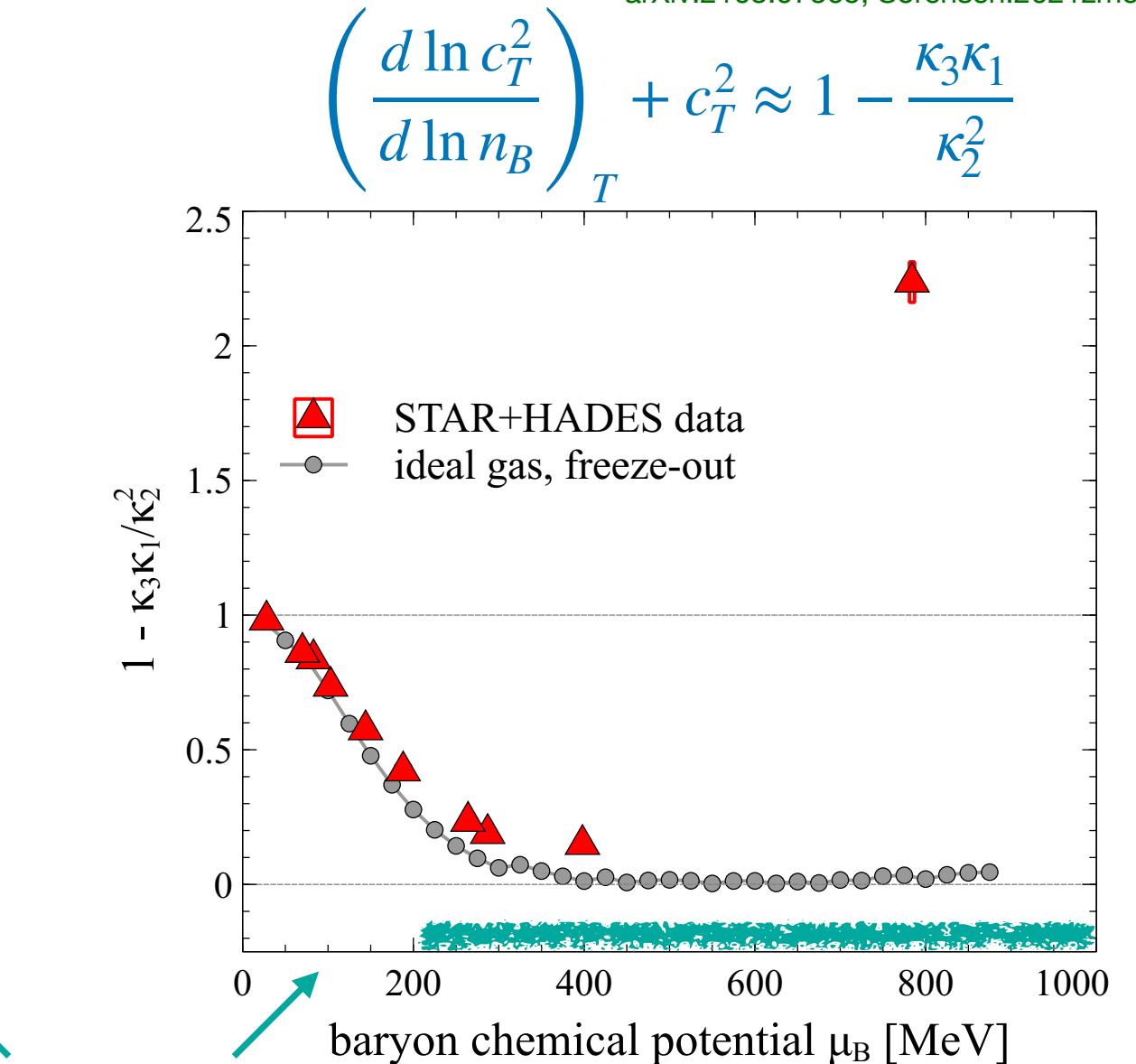
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AS, D. Oliynychenko, V. Koch, L. McLerran, Phys. Rev. Lett. **127** (2021) 042303
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$$c_T^2 \approx \frac{T\kappa_1}{\mu_B\kappa_2}$$

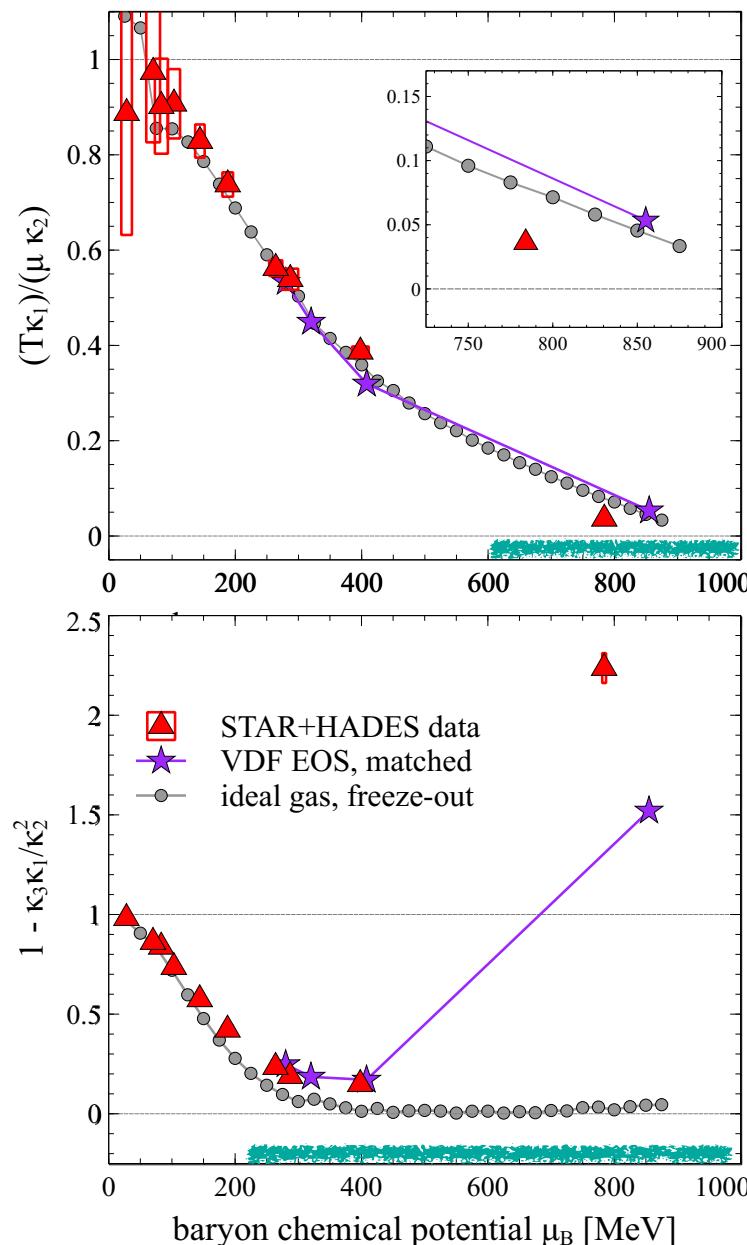


region of validity



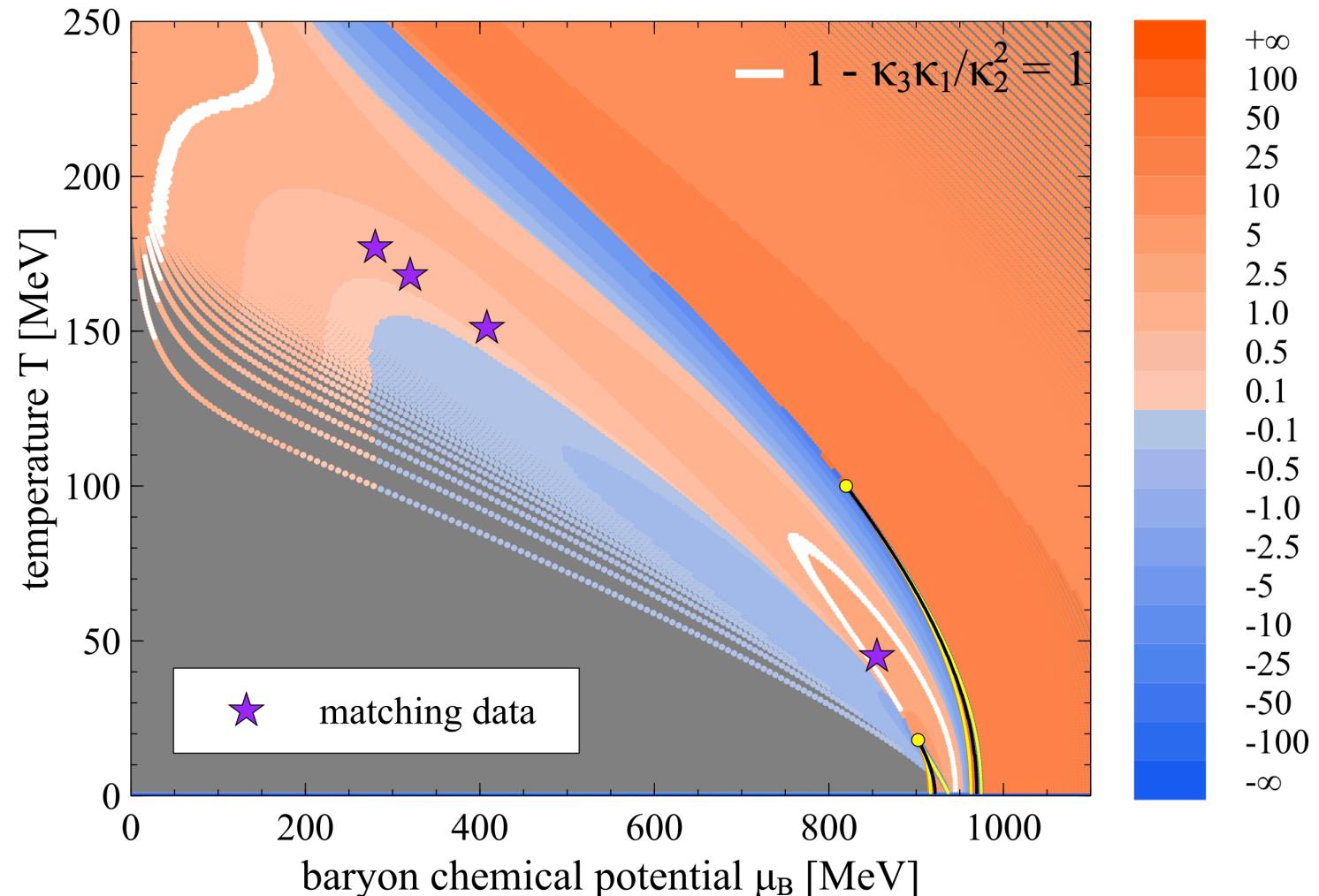
M. Abdallah et al. (STAR), Phys. Rev. C **104** (2021) 2, 024902,
arXiv:2101.12413 ,STAR:2021iop
J. Adamczewski-Musch et al. (HADES), Phys. Rev. C **102** (2020),
arXiv:2002.08701, HADES:2020wpc

Experimental data: can we understand what is happening?

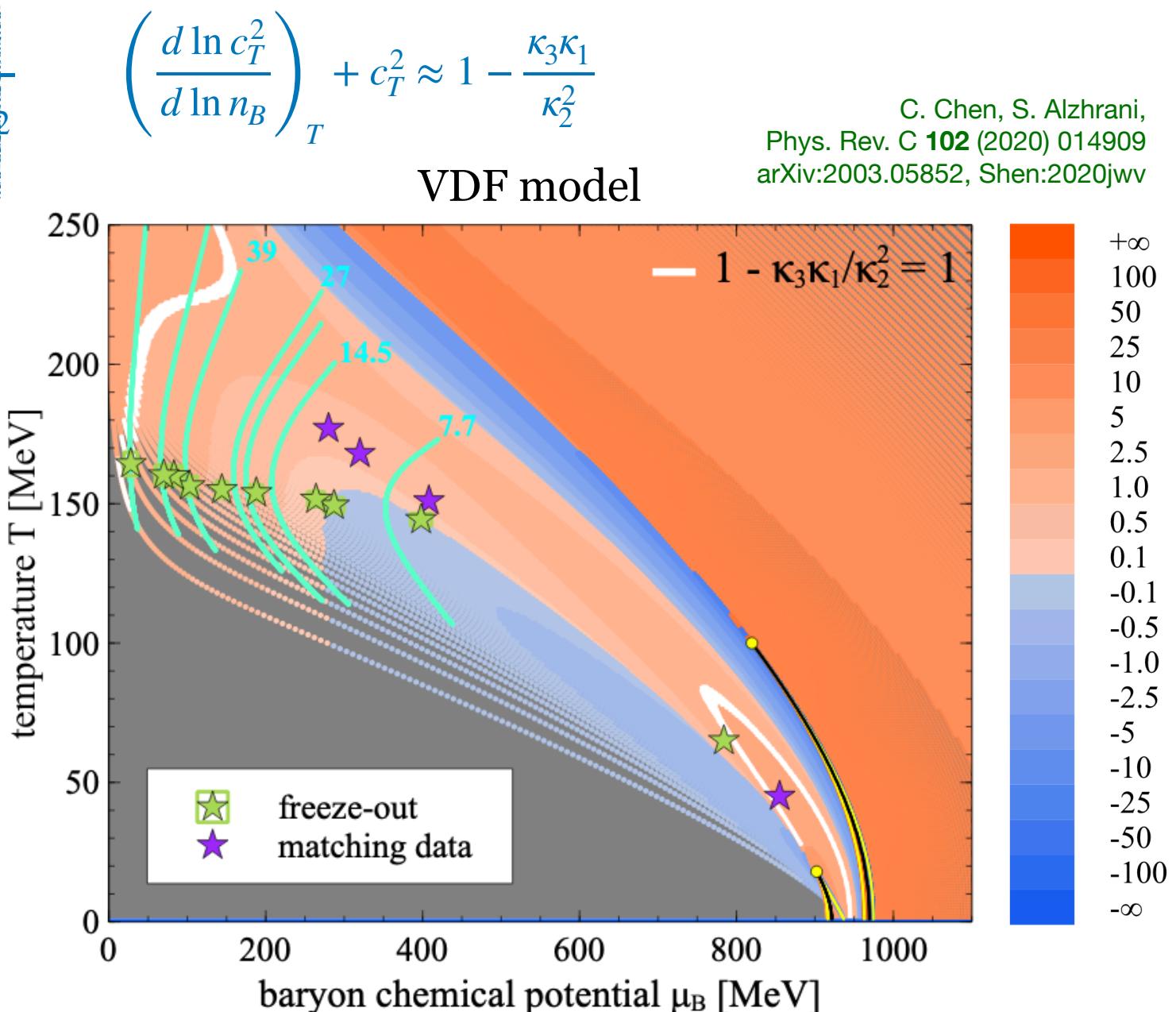
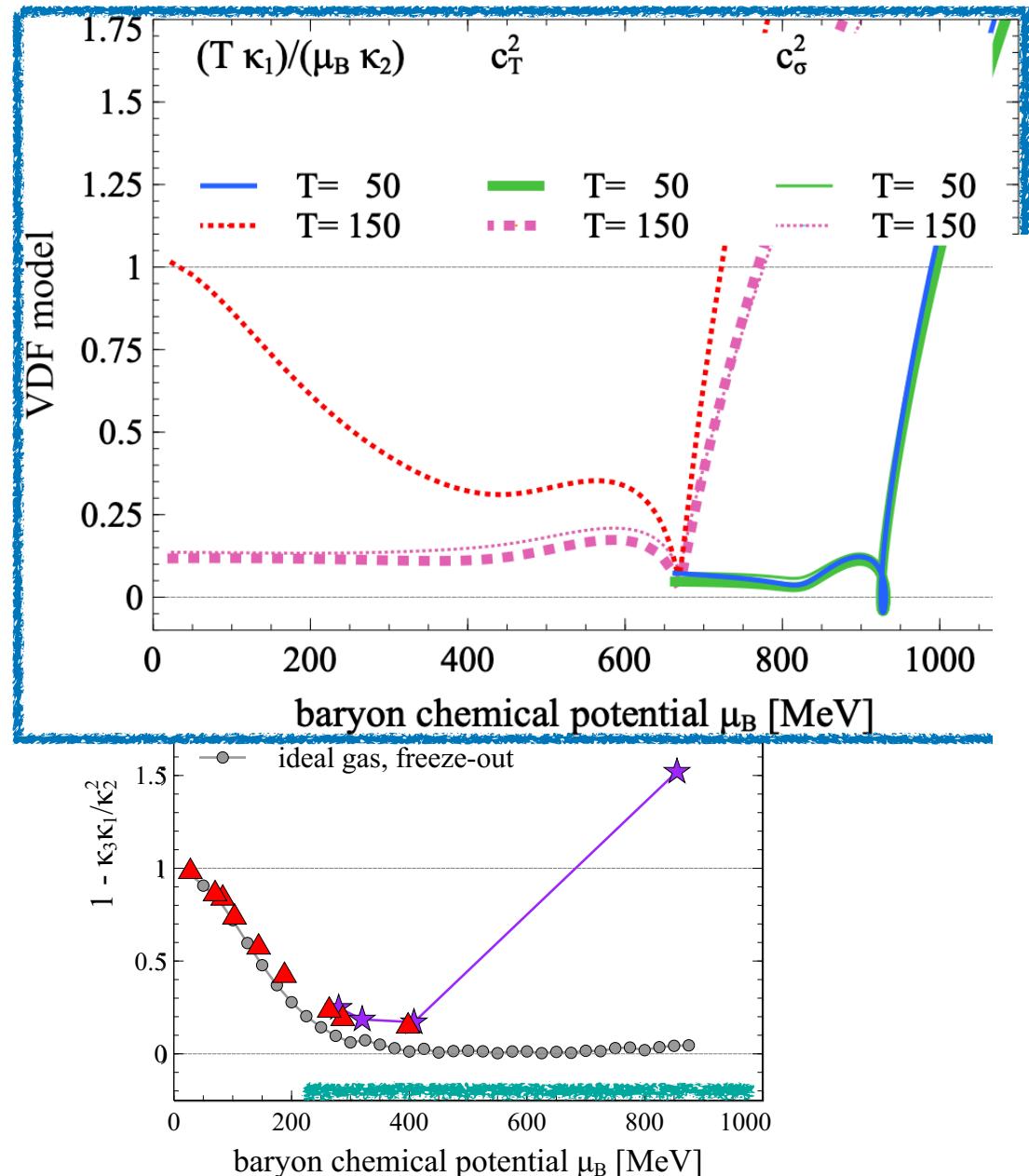


$$c_T^2 \approx \frac{T \kappa_1}{\mu_B \kappa_2} \quad \left(\frac{d \ln c_T^2}{d \ln n_B} \right)_T + c_T^2 \approx 1 - \frac{\kappa_3 \kappa_1}{\kappa_2^2}$$

VDF model

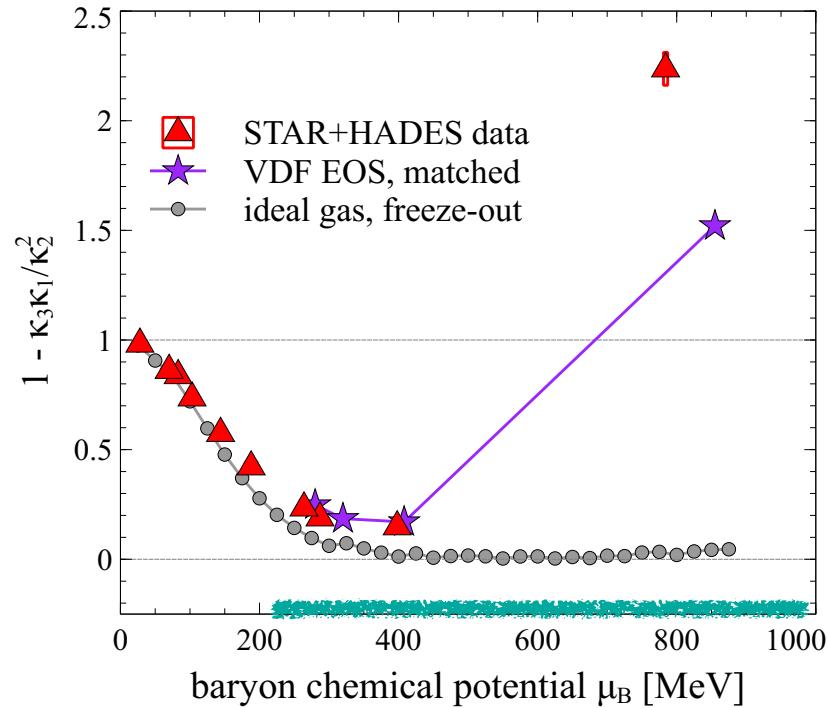
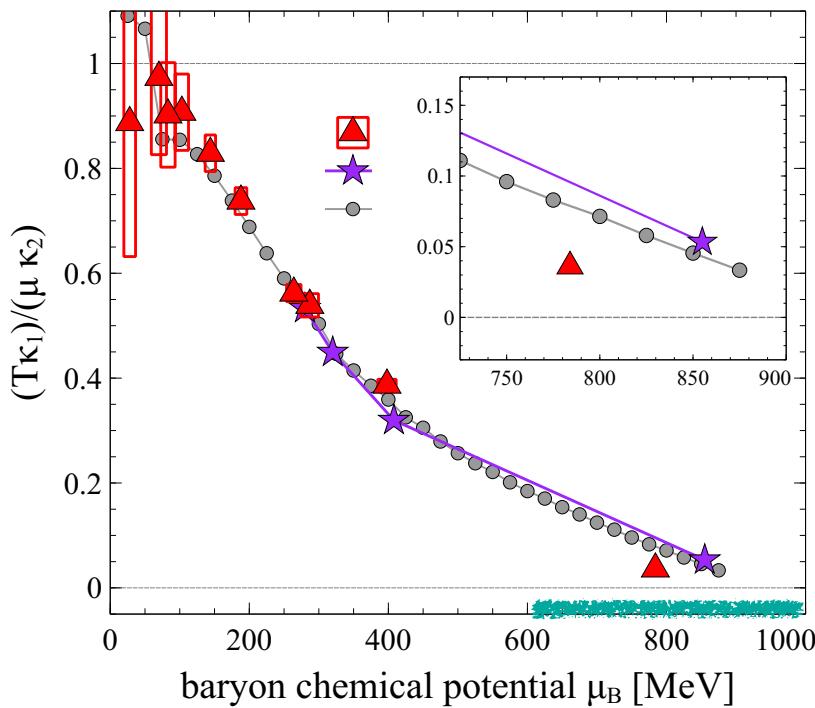


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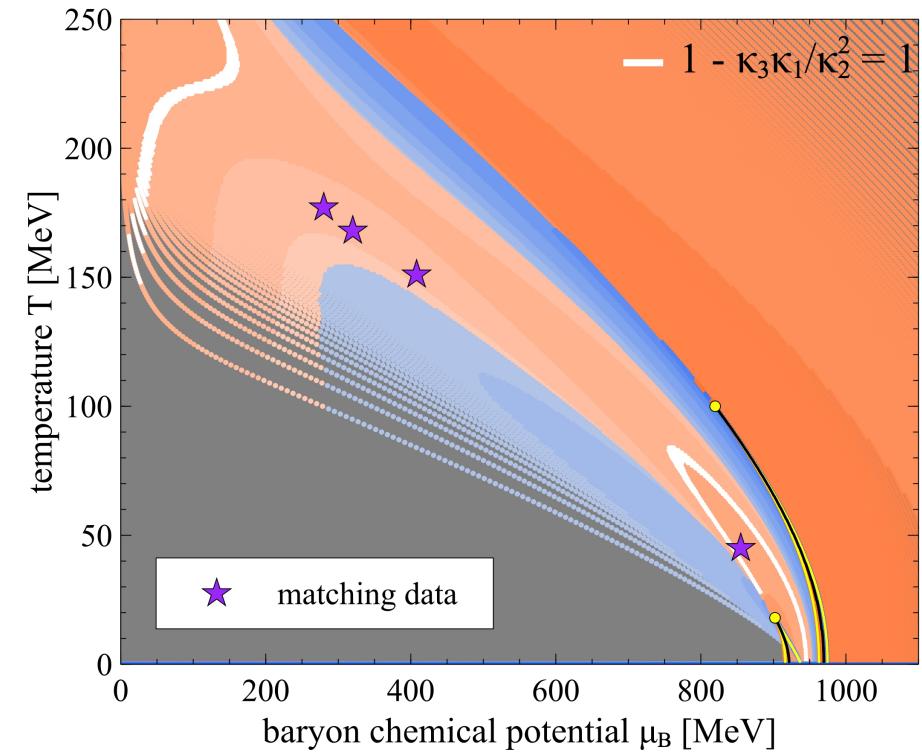


The results prompt more questions

STAR + HADES data



VDF model



- *something* significant is happening: κ_3 changes sign!
- results suggest vastly different behavior of c_s^2 at energies probed by STAR (collider) and HADES
- is behavior of the cumulants at low energies dominated by hadronic effects and the nuclear liquid-gas phase transition???

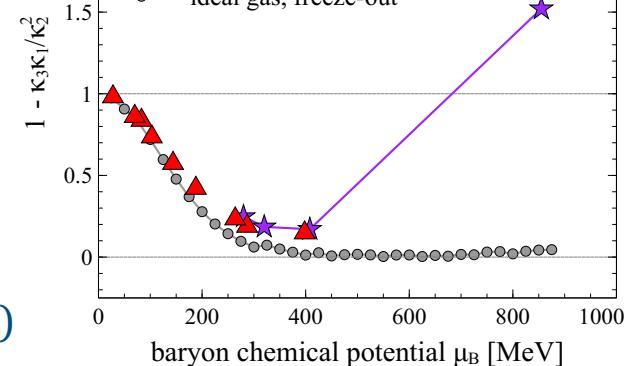
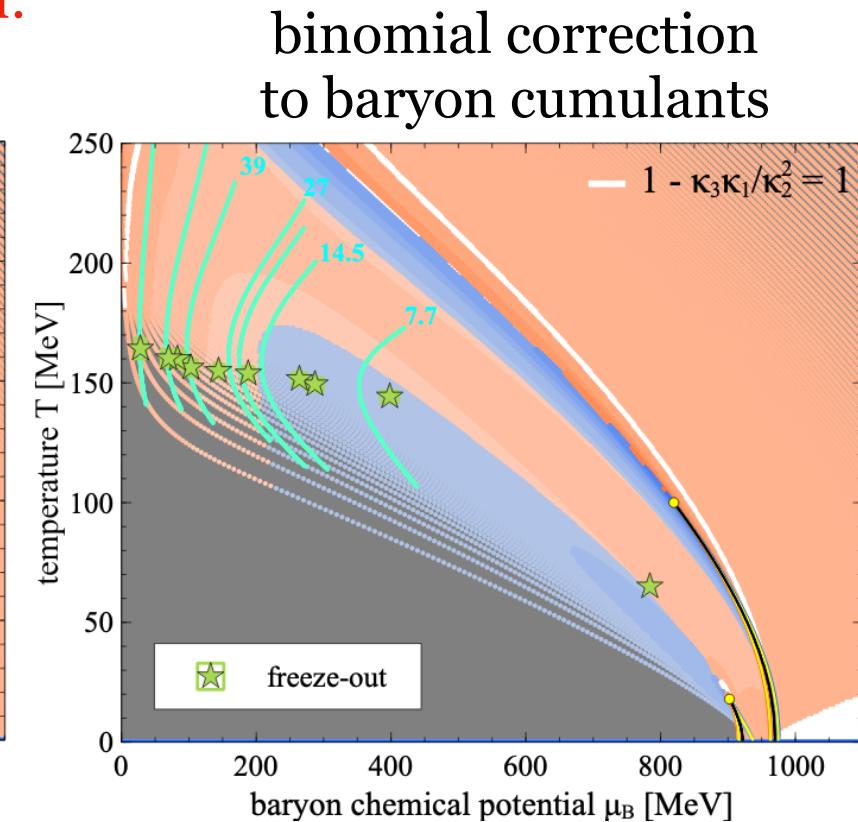
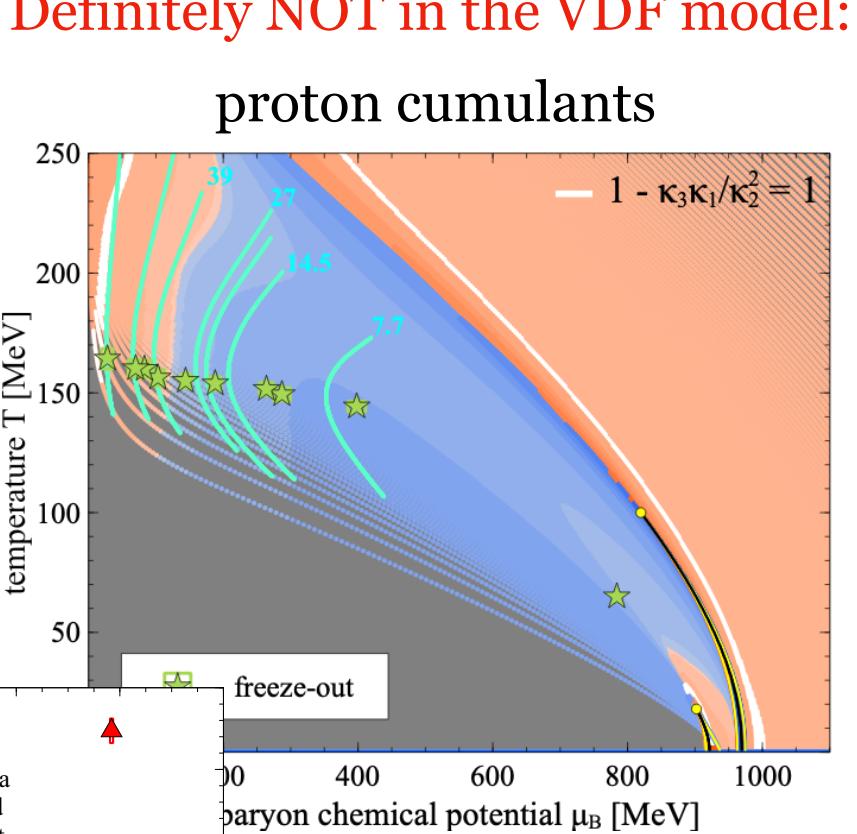
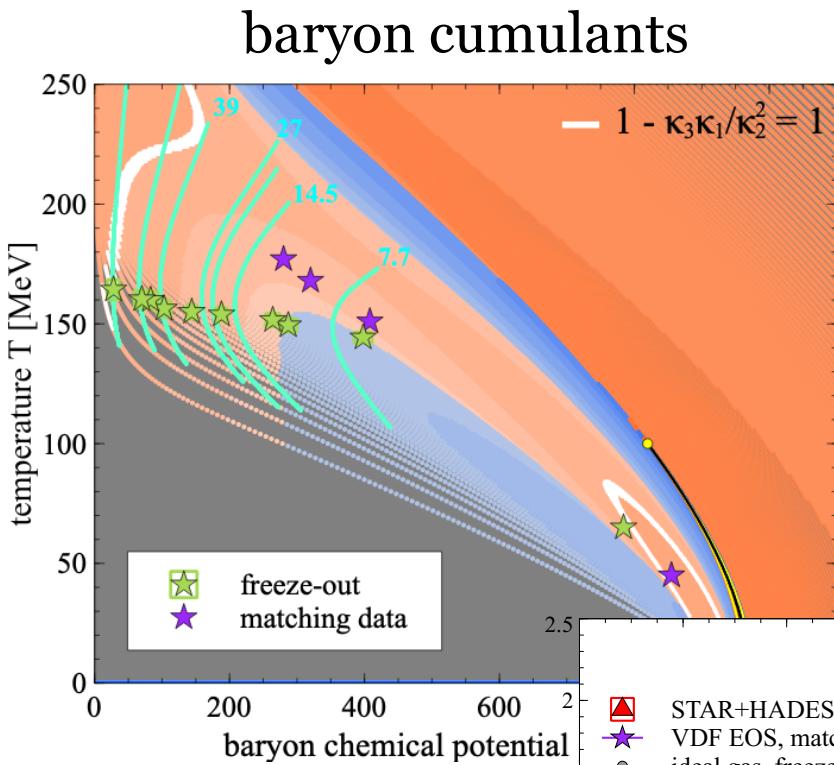
Need more experimental data and comparisons with simulations!

Issue 1: Proton vs. baryon cumulants

Y. Hatta, Y. and M. A. Stephanov, Phys. Rev. Lett. **91** (2003) 102003,
arXiv:0302002, Hatta:2003wn

Proton cumulants said to be a good proxy for baryon cumulants.

Are they?



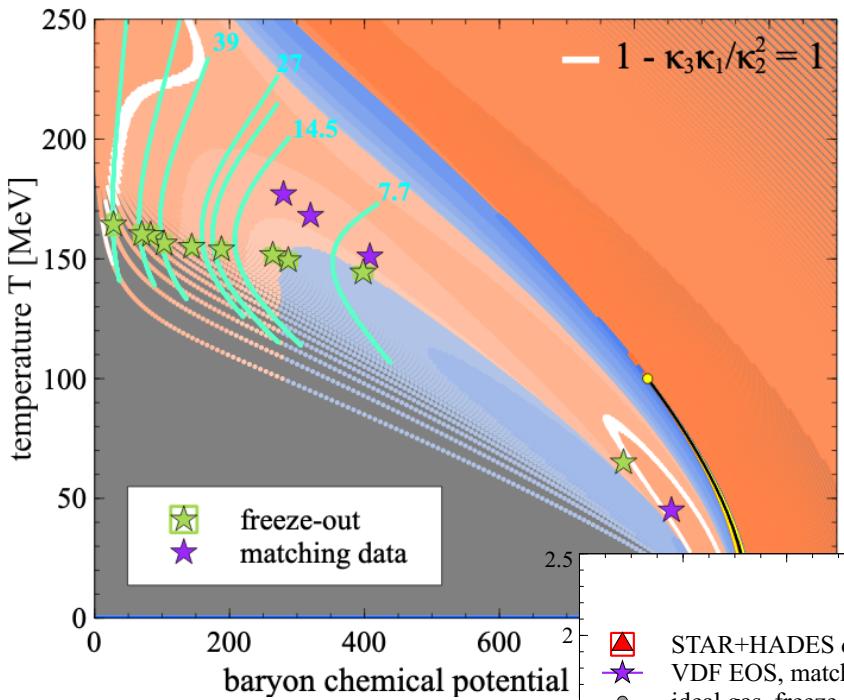
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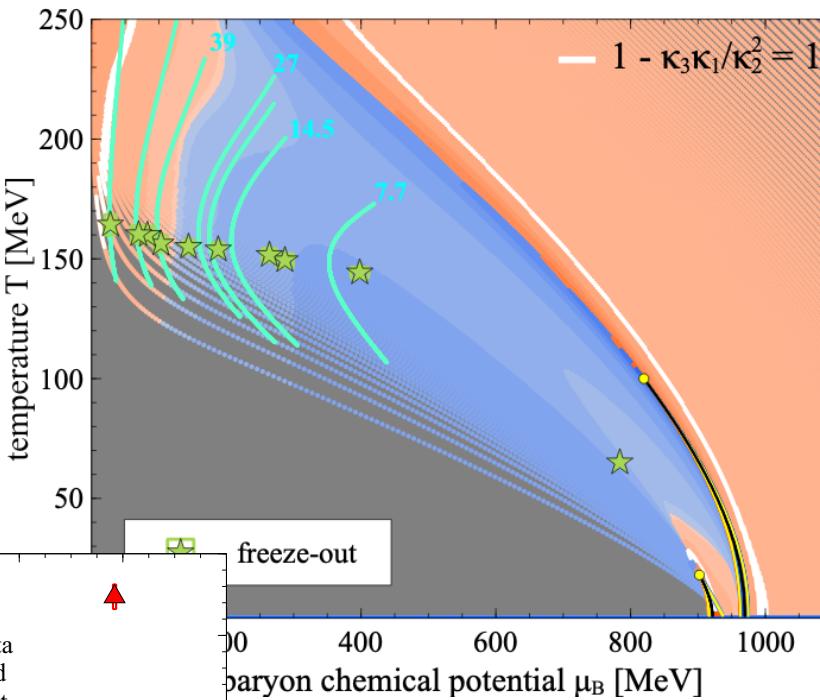
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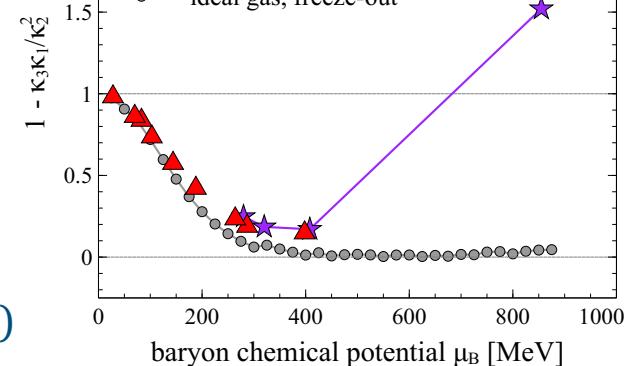
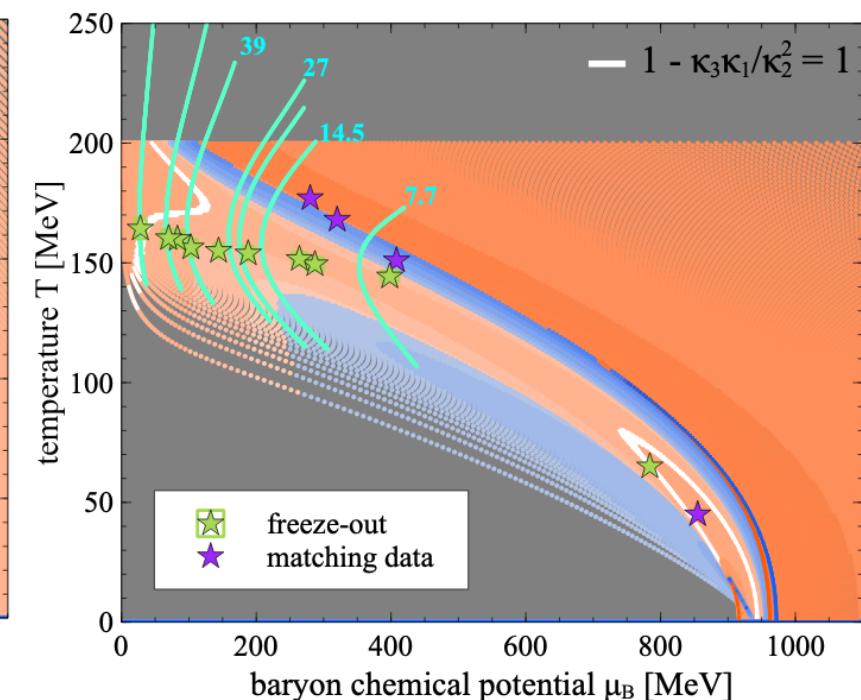
baryon cumulants



proton cumulants



baryon cumulants
(63 baryon species)

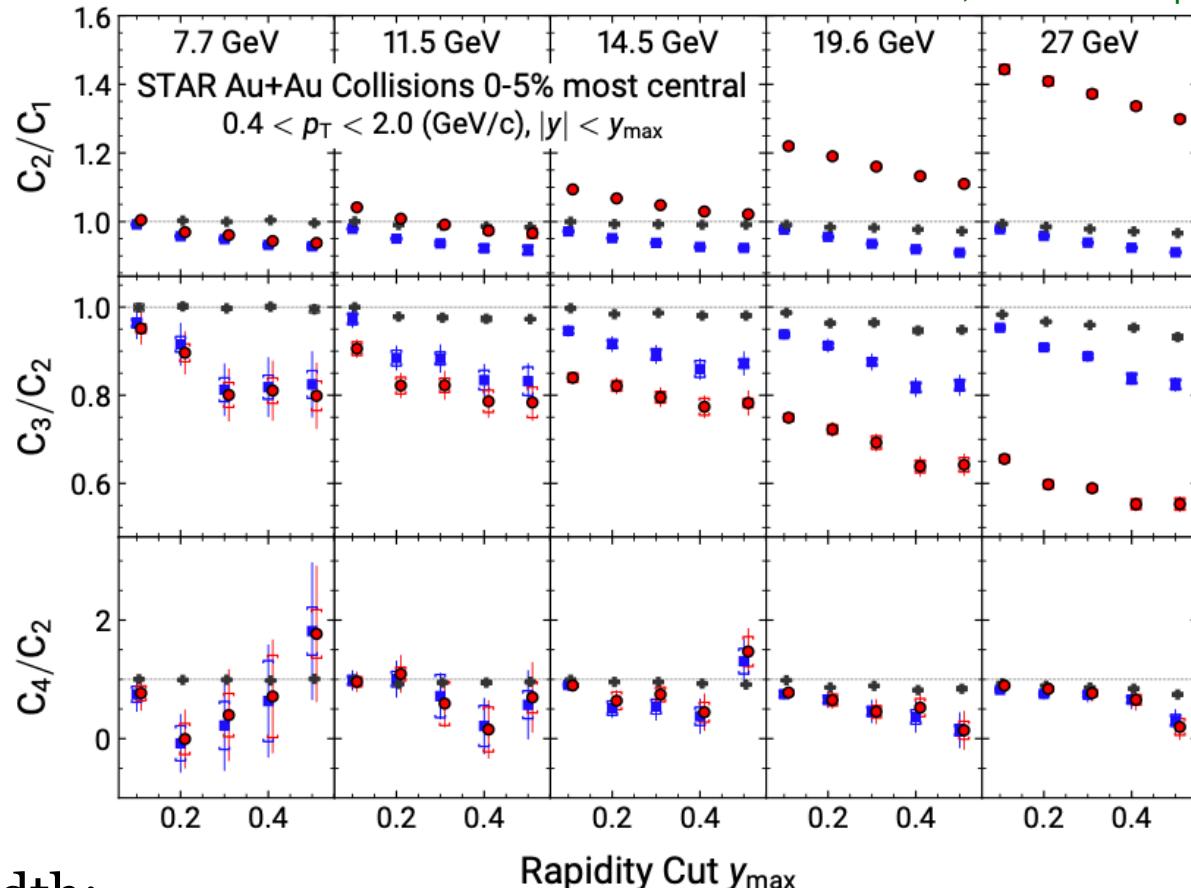


Also: momentum space vs coordinate space cumulants

V. Vovchenko and V. Koch,
“Thermodynamic approach to proton number fluctuations in baryon-rich heavy-ion matter created at moderate collision energies”,
arXiv:2204.00137, Vovchenko:2022szk

Issue 2: Rapidity window dependence vs. theory calculations

STAR, Phys. Rev. C **104** (2021) no. 2, 024902,
arXiv:2101.12413, STAR:2021iop



Changing the rapidity width:

- ⇒ changes the probed scale (what bin width will “capture” the correlations?)
- ⇒ can increase/reduce baryon number conservation effects

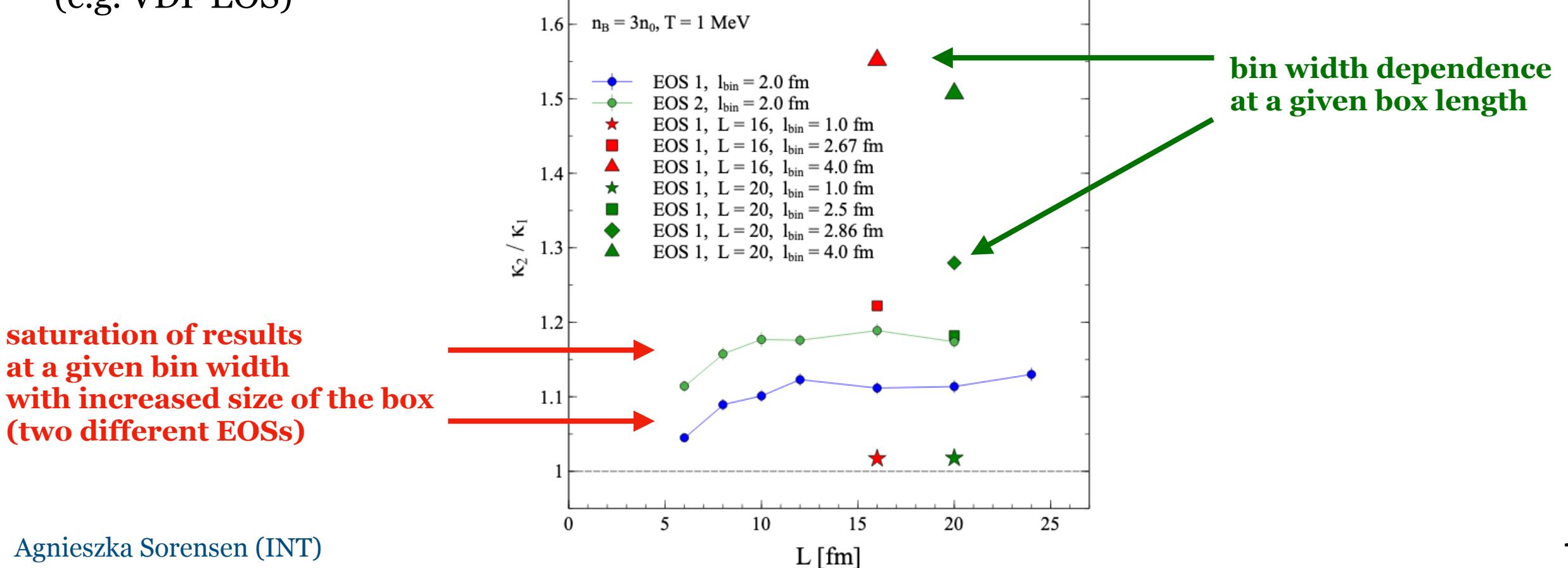
V. Vovchenko, O. Savchuk, R. Poberezhnyuk, M. Gorenstein, V. Koch,
Phys. Lett. B **811** (2020) 135868, arXiv:2003.13905, Vovchenko:2020tsr

Results at which bin width should be compared with theory?

Issue 3: Comparison with hadronic transport simulations

Heavy-ion collisions are dynamical, messy phenomena occurring in a finite volume and within a finite time span. At low energies hadronic transport is the appropriate tool to use.

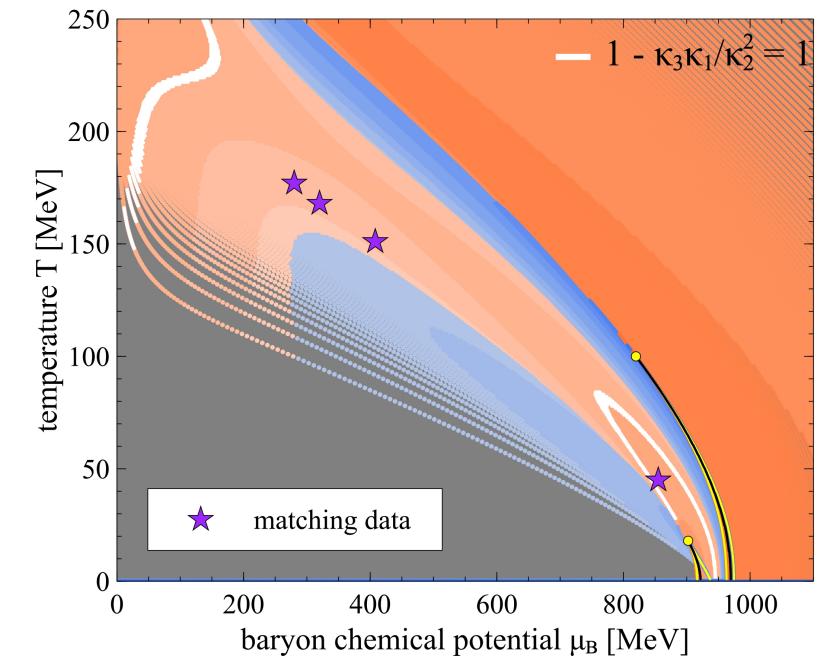
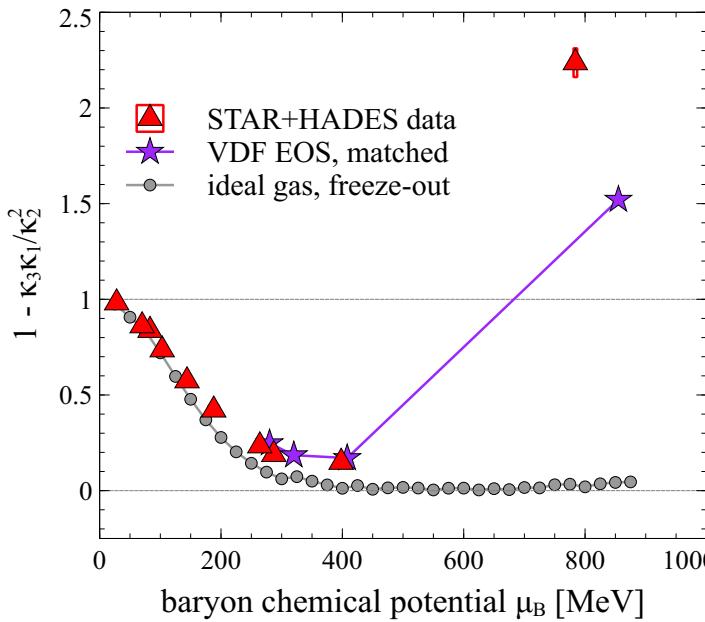
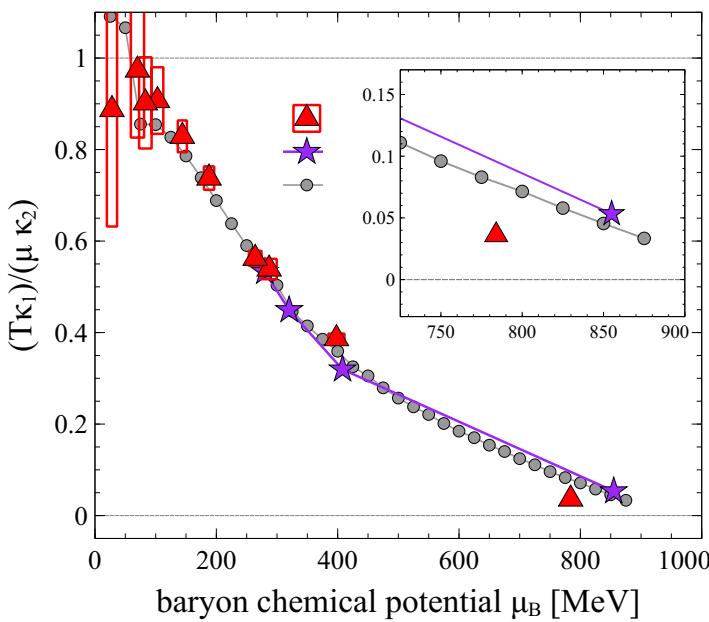
- Representing a single baryon with many ($N_T \gg 1$) test particles with fractional baryon number ($Q_B = 1/N_T$) to calculate mean-fields: has to be accounted for when comparing to experiment
- Connecting values of cumulants from simulations (e.g. SMASH) to values from the theory (e.g. VDF EOS)



Summary

- Cumulants reimagined: if we understand them quantitatively, they can reveal the speed of sound
- Independently of the model interpretation, the data points towards interesting behavior
- The path forward is to understand finite number, binning, conservation, ... effects in fluctuation observables and connect results from measurements to theory

Thank you for your attention



This material is based upon work supported by the U.S. Department of Energy, Office of Science, Office of Nuclear Physics, under contract number DE-AC02-05CH11231.