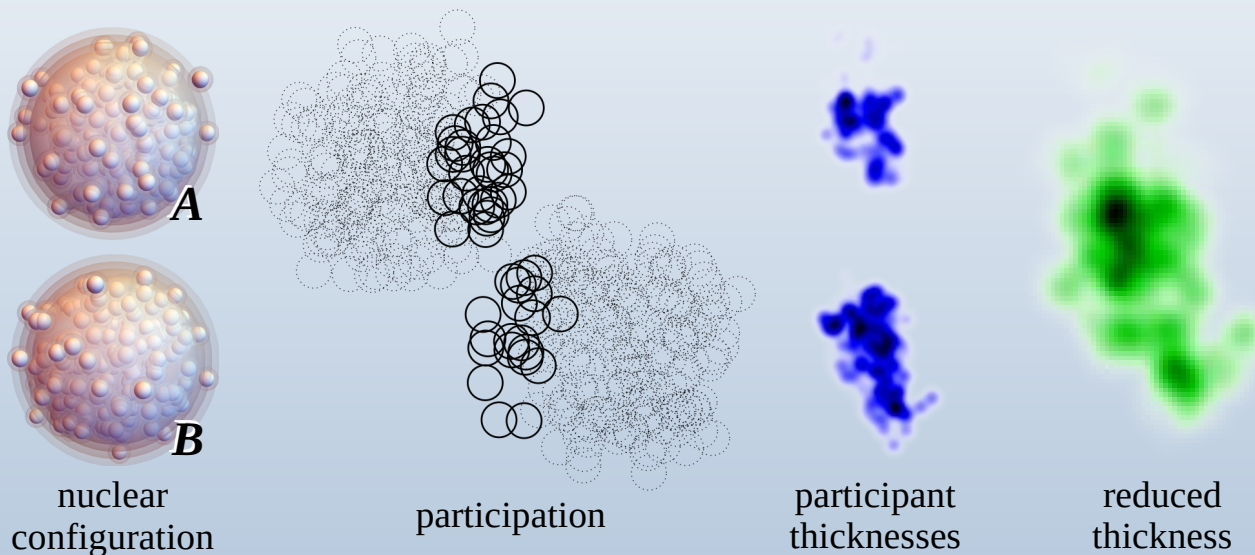


## T<sub>R</sub>ENTo (2D)



## First Results for the New T<sub>R</sub>ENTo-3D Initial-Conditions Ansatz

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We assume a **central fireball** and two **fragmentation regions**, with profile functions:

$$F_{\text{cent}}(\Delta\eta_s) = \exp\left[\left(\frac{-\Delta\eta_s^2}{2\eta_{s,\text{max}}}\right)^f\right] \times \left(1 - \left(\frac{\Delta\eta_s}{\eta_{s,\text{max}}}\right)^4\right)^4 \quad \text{and} \quad F_{\text{frag}}(x) = (-\ln x)^\alpha x^{\beta+1} \exp\left(\frac{-2k_{T,\text{min}}}{x\sqrt{s_{NN}}}\right)$$

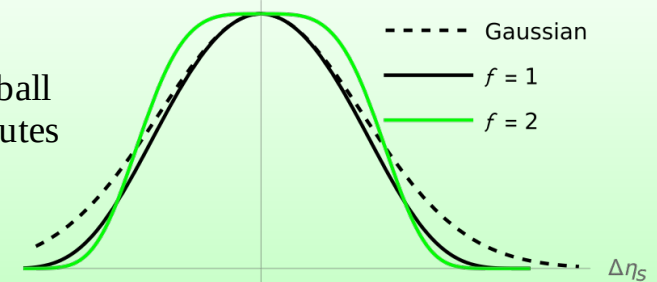
## Central Fireball

- Gaussian-like profile function with flatness parameter  $f$ :
- On average,  $N$ - $N$  participant pair deposits fraction  $\langle x_{\text{loss}} \rangle$  of energy into fireball
- Event by event, each constituent (subnucleonic degree of freedom) contributes an independent fraction  $x_{\text{loss}}$  of its energy to fireball, distributed as:

$$x_{\text{loss}} \sim \text{Beta}[\langle x_{\text{loss}} \rangle K, (1 - \langle x_{\text{loss}} \rangle) K]$$

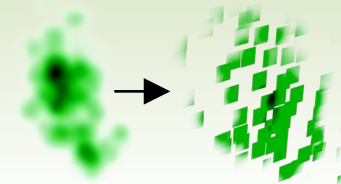
with mean  $\langle x_{\text{loss}} \rangle$  and fluctuation parameter  $K$

$F_{\text{cent}}(\Delta\eta_s) @ \sqrt{s_{NN}} = 200 \text{ GeV}$



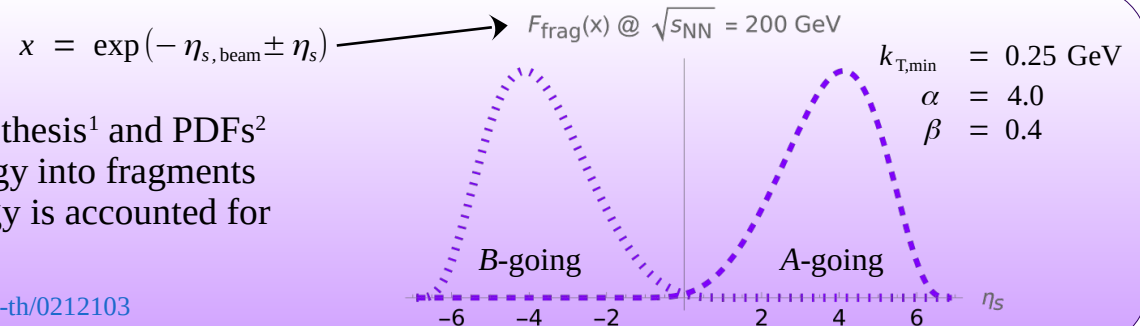
$\Delta\eta_s = \eta_s - \eta_{s,\text{cm}}$ , with center of central fireball  $\eta_{s,\text{cm}}$  varying over  $x_T$  with fireball thicknesses:

$$P_{\text{cm}}^\mu(\vec{x}_T) = \left( (T_A(\vec{x}_T) + T_B(\vec{x}_T)) \frac{\sqrt{s_{NN}}}{2}, \hat{z} (T_A(\vec{x}_T) - T_B(\vec{x}_T)) \sqrt{\frac{s_{NN}}{4} - m_p^2} \right), \quad \eta_{s,\text{cm}} = \text{asinh}(p_{\text{cm}} / \sqrt{P_{\text{cm}}^\mu P_{\mu}^{\text{cm}}})$$



## Fragmentation Regions

- Form inspired by limiting fragmentation hypothesis<sup>1</sup> and PDFs<sup>2</sup>
- Constituents deposit fraction  $(1 - x_{\text{loss}})$  of energy into fragments
- **Conservation of energy:** all participant energy is accounted for



<sup>1</sup> Benecke *et al.*, *Phys. Rev.* **188**, 2159 (1969)


<sup>2</sup> Bass, Müller, Srivastava, *Phys. Rev. Lett.* **91** (2003) 052302, [nucl-th/0212103](https://arxiv.org/abs/nucl-th/0212103)

We assume a **central fireball** and two **fragmentation regions**, with profile functions:

$$F_{\text{cent}}(\Delta\eta_s) = \exp\left[\left(\frac{-\Delta\eta_s^2}{2\eta_{s,\text{max}}}\right)^f\right] \times \left(1 - \left(\frac{\Delta\eta_s}{\eta_{s,\text{max}}}\right)^4\right)^4 \quad \text{and} \quad F_{\text{frag}}(x) = (-\ln x)^\alpha x^{\beta+1} \exp\left(\frac{-2k_{T,\text{min}}}{x\sqrt{s_{NN}}}\right)$$

$$\varepsilon(\vec{x}_T, \eta_s) = N_{\text{Trento}} \sqrt{T_A(\vec{x}_T) T_B(\vec{x}_T)} F_{\text{cent}}(\eta_s - \eta_{s,\text{cm}}(\vec{x}_T)) + \frac{k_{T,\text{min}}}{\int_{k_{T,\text{min}}/\sqrt{s_{NN}}}^1 dx F_{\text{frag}}(x)} \times$$

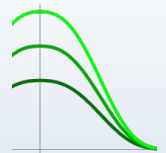
(T<sub>R</sub>ENTO  $p=0$ )  $\nearrow$   $\{F_A(\vec{x}_T) F_{\text{frag}}(\exp(-\eta_{s,\text{max}} + \eta_s)) + F_B(\vec{x}_T) F_{\text{frag}}(\exp(-\eta_{s,\text{max}} - \eta_s))\}$



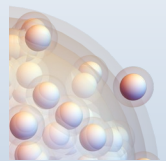
## Model parameters



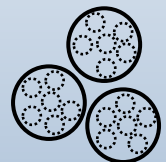
Fluctuation ( $k$ )



Fireball normalization ( $N_{\text{Trento}}$ )

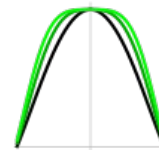


Nucleon width ( $w$ )

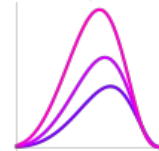


Constituent width ( $v$ )

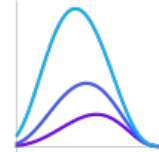
Constituent number ( $n_c$ )



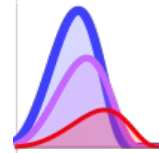
Flatness ( $f$ )



Fragmentation shape  $\alpha$

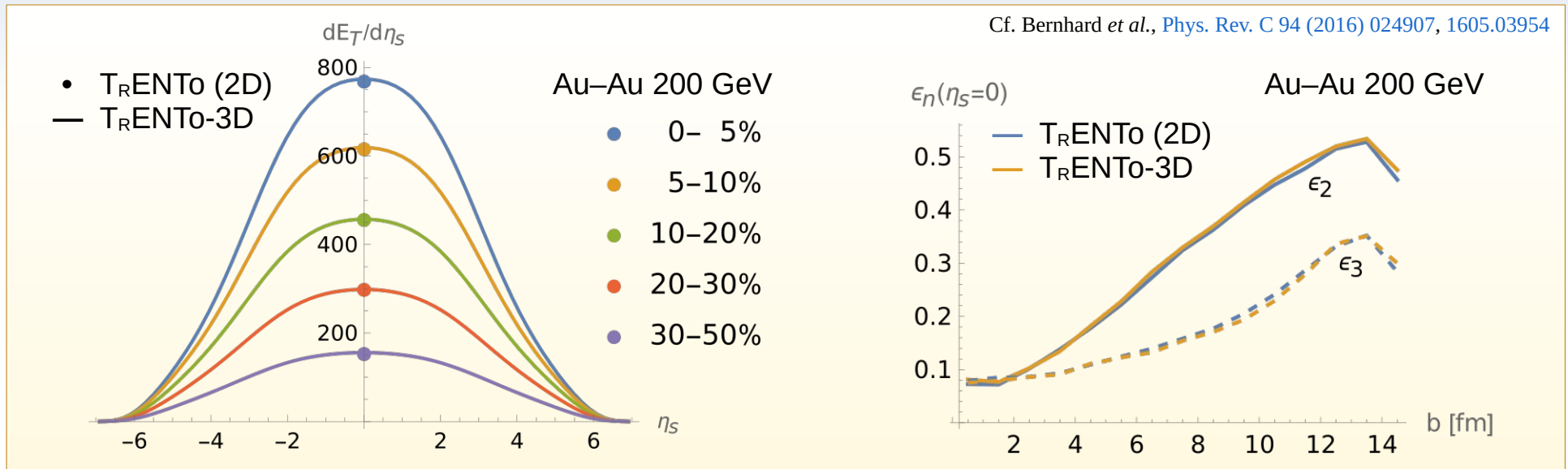
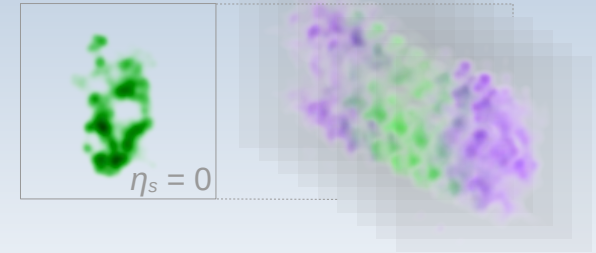


Fragmentation shape  $\beta$



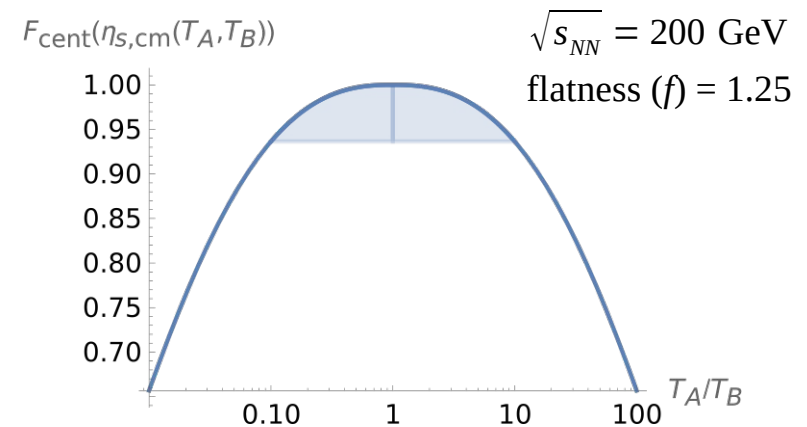
Typical parton transverse momentum scale ( $k_{T,\text{min}}$ )

Confirming that T<sub>RENT</sub>o-3D matches 2D T<sub>RENT</sub>o at midrapidity:\*



\* Some considerations:

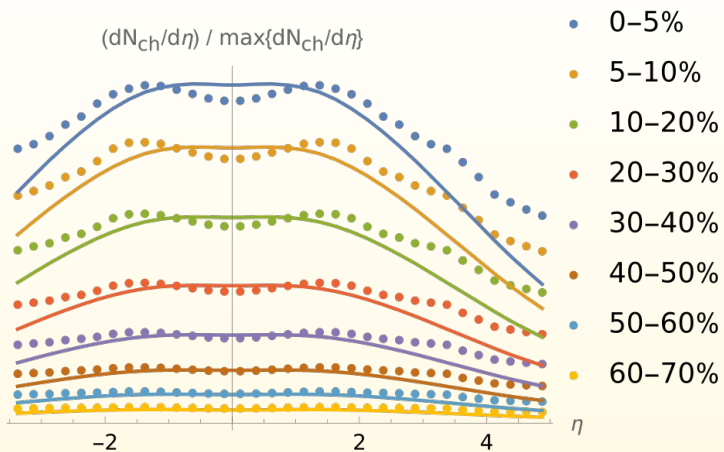
- T<sub>RENT</sub>o-3D's  $N_{\text{Trento}}$  is  $\times 1/\langle x_{\text{loss}} \rangle = 1.43$  here because its fireball–fragment fluctuation beta distribution has a **mean of  $\langle x_{\text{loss}} \rangle$** , whereas T<sub>RENT</sub>o (2D) fluctuates participants by a gamma distribution with **mean 1**
- T<sub>RENT</sub>o-3D's **central fireball profile** applies at midrapidity wherever  $\eta_{s,\text{cm}}(x_T) \neq 0$ , but in large symmetric systems it **can be neglected**:



- To connect TRENTo-3D with experiment, we use a linearized (1+1)D hydrodynamic code (“LH”) and Cooper–Frye with a transverse-Gaussian assumption (“CF”)
- No hadronic transport/decays—instead, we normalize data sets to peak at unity
- Simultaneously fitting to three experimental data sets:

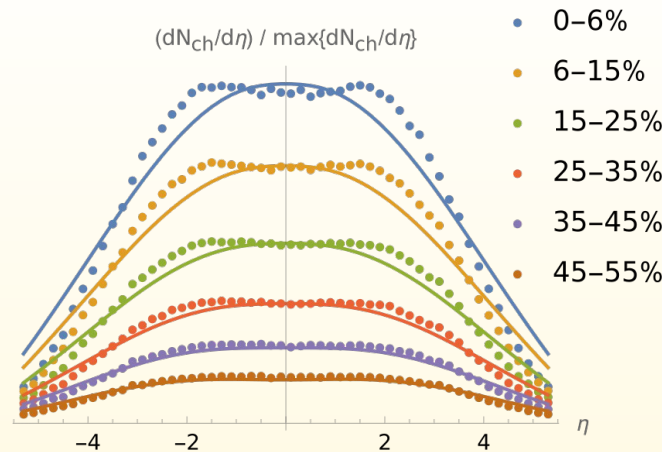
- Experiment
- T<sub>R</sub>ENTo-3D → LH → CF

## Pb–Pb 5.02 TeV



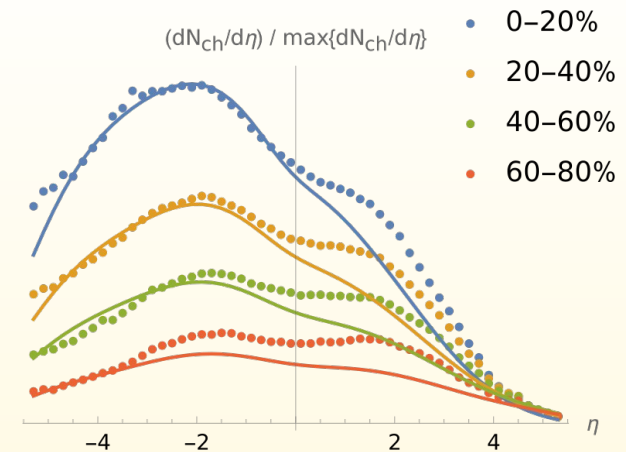
ALICE, *Phys. Lett. B* 772 (2017) 567-577,  
1612.08966

## Au–Au 200 GeV



PHOBOS, *Phys. Rev. Lett* 91 (2003) 052303,  
nucl-ex/0210015

## d–Au 200 GeV



PHOBOS, *Phys. Rev. C* 72 (2005) 031901,  
nucl-ex/0409021

- Parameter values used above:

$$\begin{aligned}
 w &= 0.41 \text{ [fm]} & v &= 0.19 \text{ [fm]} & n_c &= 9.39 & f &= 1.34 & \alpha &= 3.41 & \beta &= 0.20 \\
 N_{200 \text{ GeV}} &= 3.52 & N_{5.02 \text{ TeV}} &= 14.7 & k &= 0.11 & k_{T,\text{min}} &= 0.22 \text{ [GeV]} & \tau_0 \text{ (LH)} &= 1.48 \text{ [fm/c]}
 \end{aligned}$$

- Calibration is ongoing: more systems and energies
- Parametrization seems to challenge Bayesian parameter estimation, may need revision