



29TH INTERNATIONAL  
CONFERENCE ON ULTRARELATIVISTIC  
NUCLEUS - NUCLEUS COLLISIONS  
APRIL 4-10, 2022  
KRAKÓW, POLAND

# Multiple Parton Scattering and Gluon Saturation in Dijet Production at EIC



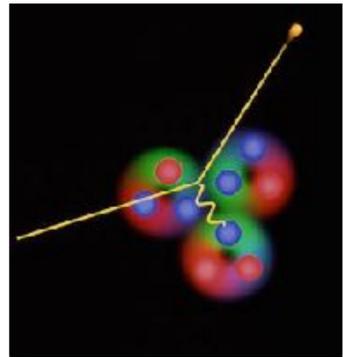
***Phys.Rev.D 105 (2022) 3, 034015***

**Yuan-Yuan Zhang (CUHKSZ)**

**Collaborator : Xin-Nian Wang (LBL)**

# Probe $\hat{q}$ (jet quenching) in Cold Nuclei

Electron Ion Collision

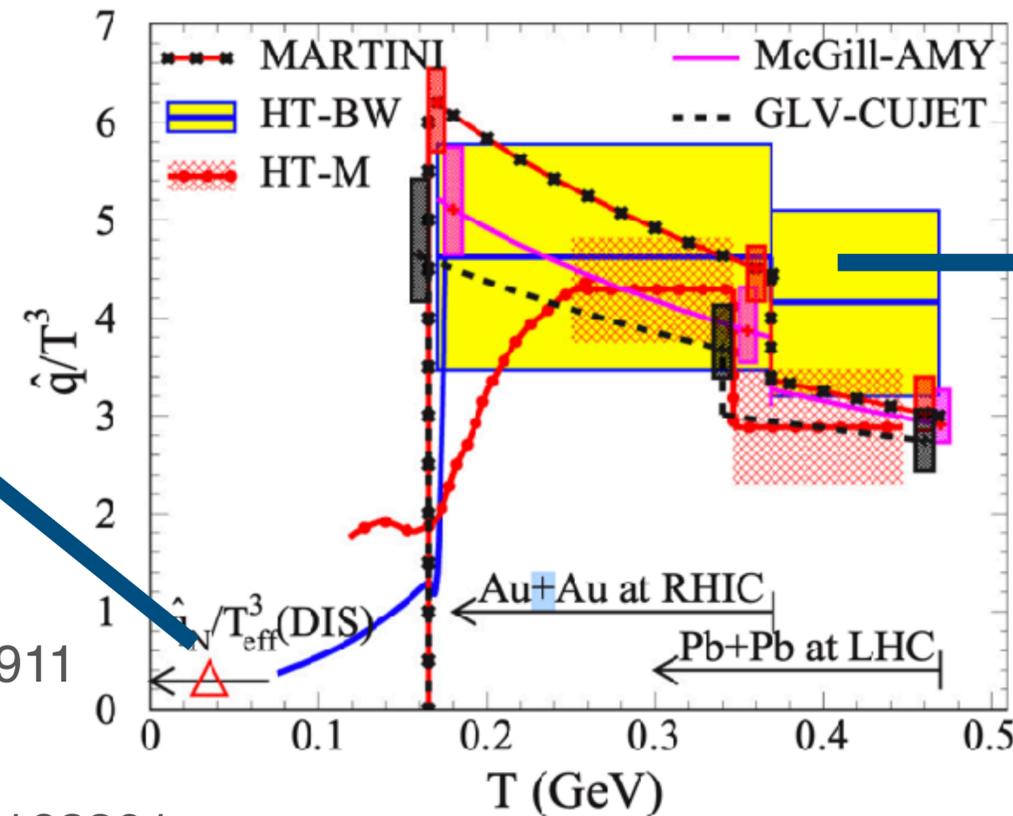


$$\hat{q} \approx 0.015 \text{ GeV}^2/\text{fm}$$

NB Chang et al, PRC **89.3** (2014): 034911

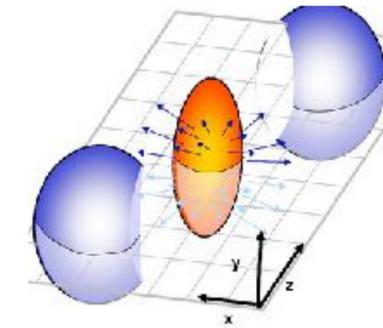
P. Ru et al, arXiv:2004.00027

E Wang, XN Wang, PRL **89.16** (2002): 162301



JET Collaboration, PRC **90**, 014909

Heavy Ion Collision



$$\hat{q} \approx 1.2 \pm 0.3 \text{ GeV}^2/\text{fm} \quad \text{RHIC}$$

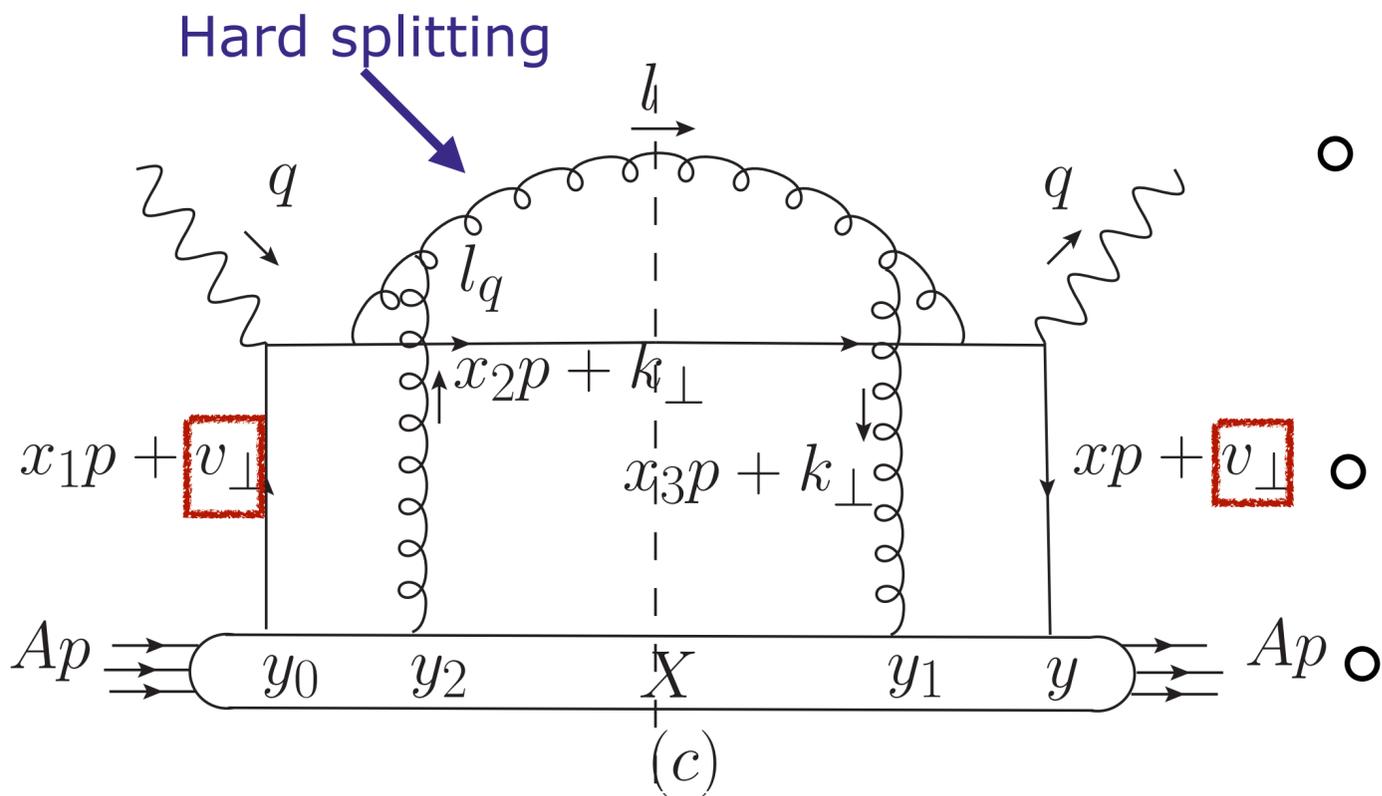
$$\hat{q} \approx 1.9 \pm 0.7 \text{ GeV}^2/\text{fm} \quad \text{LHC}$$

JET Collaboration, PRC **90**, 014909

○  $\hat{q}$  average  $p_T$  transfer squared per unit length

○  $\hat{q}$  related to gluon TMD PDF  $\hat{q} = \int d^2k_{\perp} \hat{q}(k_{\perp}) = C \int d^2k_{\perp} \rho \phi(x_g, k_{\perp})$

# Use Dijet to probe gluon TMD PDF or $\hat{q}$



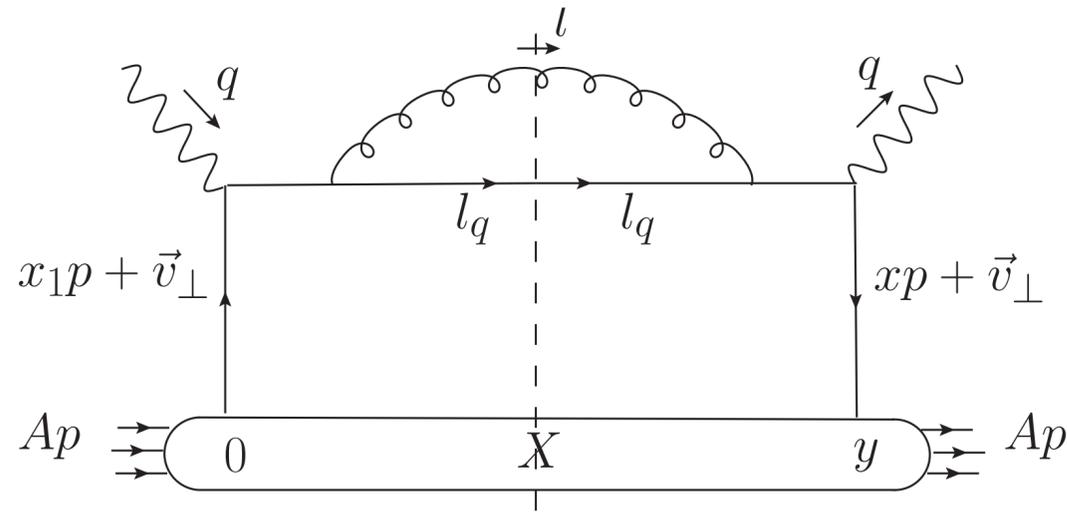
- Based on Generalized High Twist Approach
- Relax  $l_{\perp}, l_{q\perp} \gg k_{\perp}$ , no  $k_{\perp}$  /twist expansion
- The second scattering probe gluon TMD pdf
- Include initial quark  $v_{\perp}$  and pT broadening

Include **saturation** in gluon TMD PDF , determine saturation scale self-consistently

$$\phi_N(x_G, k_{\perp}, \mu^2) = \begin{cases} \phi_N^0 \text{ at } Q_s, & k_{\perp} < Q_s; \\ \phi_N^0(\mu^2 = k_{\perp}^2), & k_{\perp} > Q_s, \end{cases} \quad Q_s^2(x_B, Q^2, b_{\perp}) = C \int dy_0^- \rho(y_0^-, b_{\perp}) \int \frac{d^2 k_{\perp}}{(2\pi)^2} \alpha_s(\mu) \phi_N(x_G, k_{\perp}, \mu^2)$$

**TMDlib Package** F Hautmann et al., EPJC 74 (2014), 3220

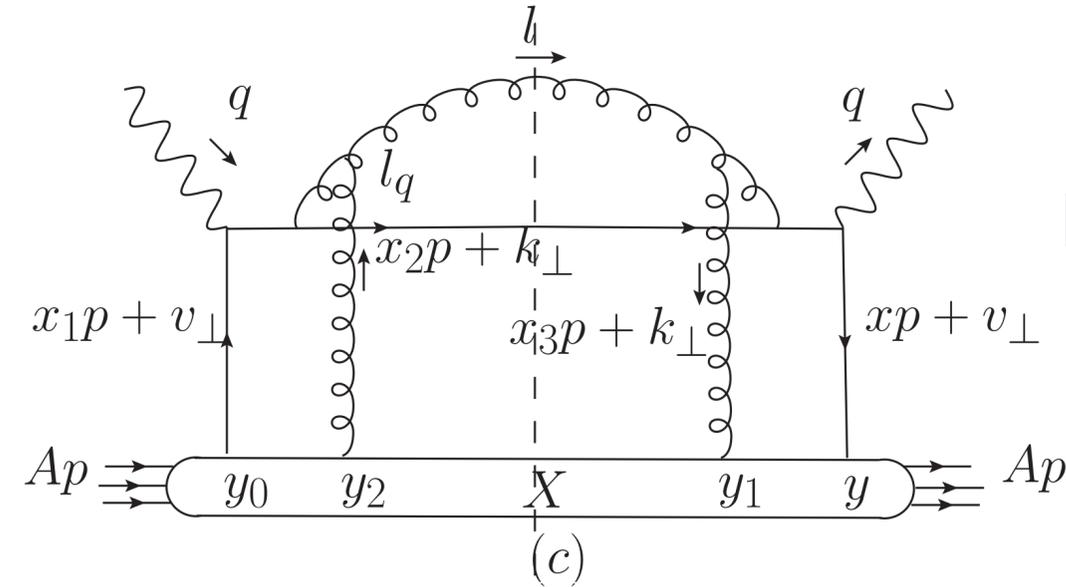
# Dijet in e+A: Single scattering & Double scattering



$$\frac{d\hat{\sigma}_{eA}^S}{dx_B dQ^2 dz d^2l_\perp d^2l_{q\perp}} \propto q_A(x_B, \vec{v}_\perp)$$

$$q_A = \int dy_0^- d^2b_\perp \rho q_N$$

nuclear modification :  
initial quark pT broadening



LPM effect

$$\frac{d\hat{\sigma}_{eA}^D}{dx_B dQ^2 dz d^2l_\perp d^2l_{q\perp}} \propto q_N(x_B, \vec{v}_\perp, \vec{b}_\perp) \otimes \phi_N(x_G, \vec{k}_\perp)$$

$$\otimes [\mathcal{N}_g^{\text{nonLPM}} + \mathcal{N}_g^{\text{qLPM}} + \mathcal{N}_g^{\text{gLPM}}]$$

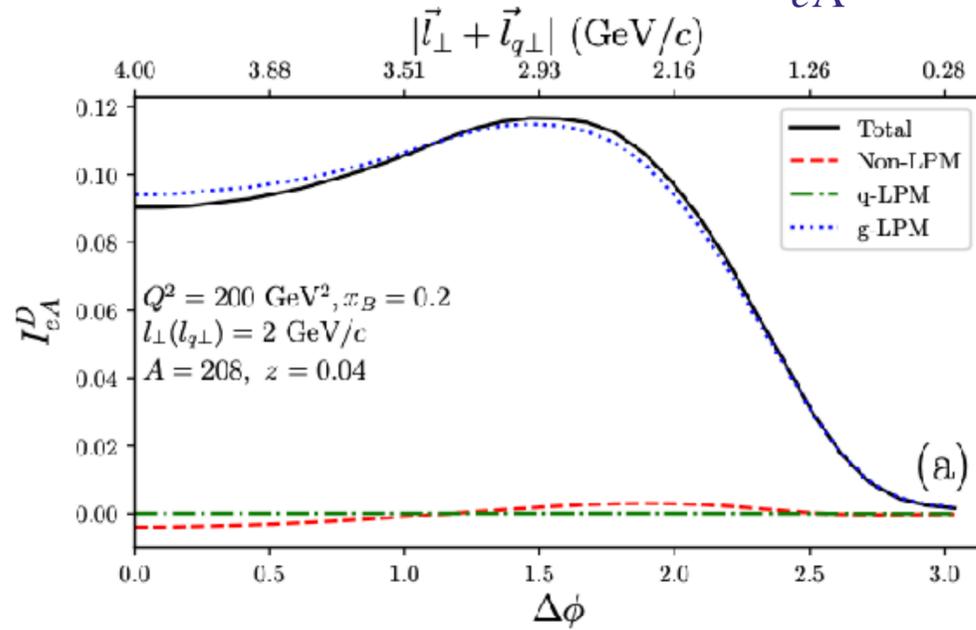
nuclear modification : initial quark pT broadening  
+ second scattering

Liang, Z. T., Wang, X. N., & Zhou, J. (2008). *PRD*, 77(12), 125010.

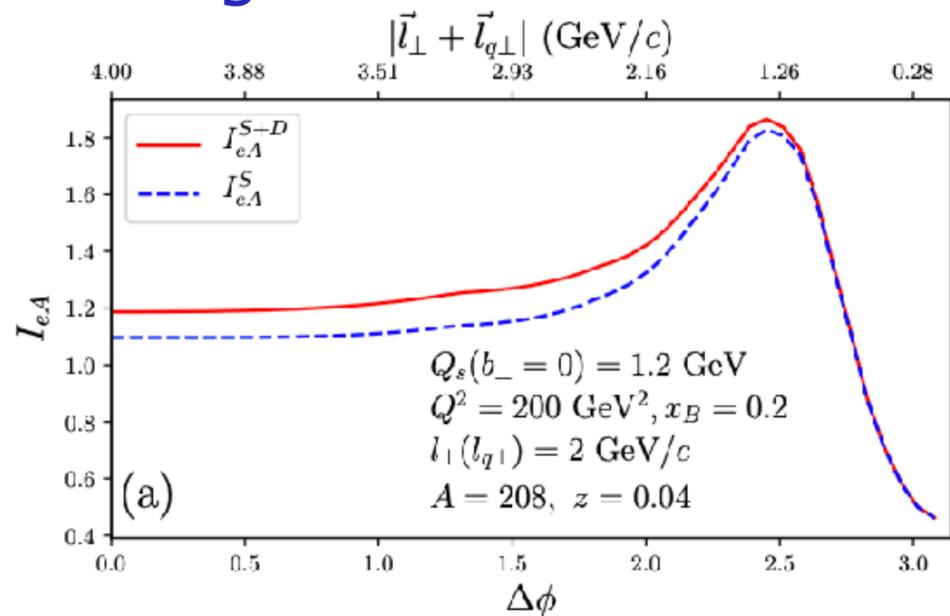
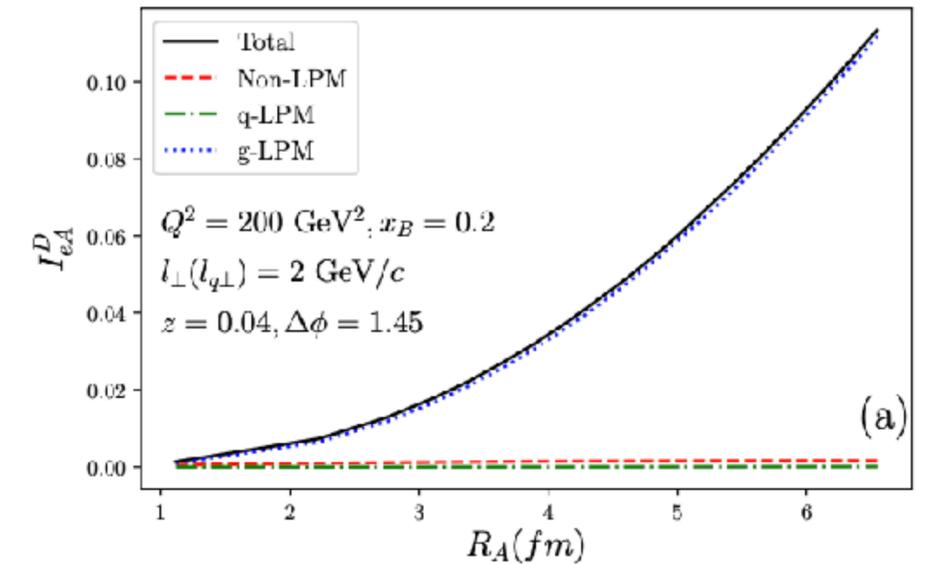
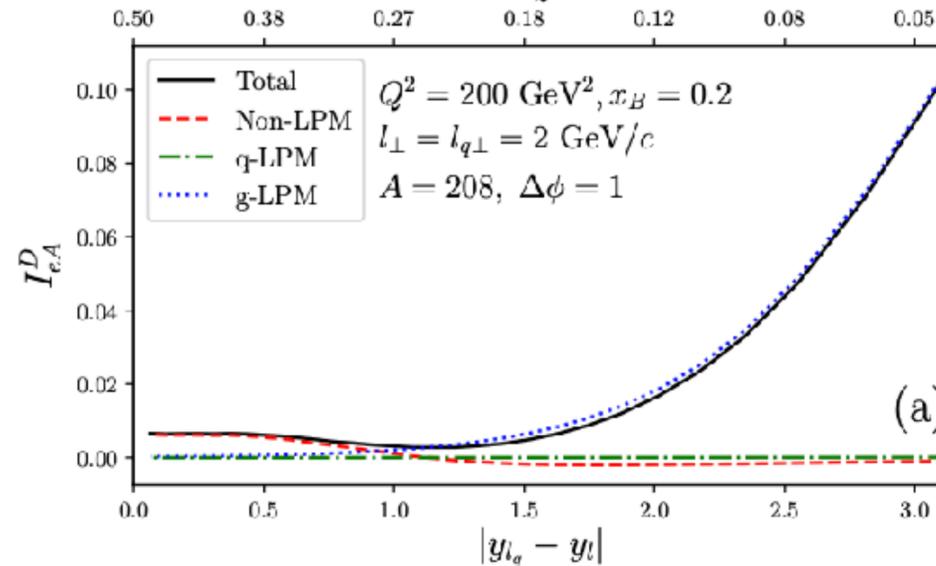
- Embed pT broadening (Gaussian, width  $\Delta = \int dy_0^- \hat{q}_F$ ) in nucleon quark pdf  $q_N(x, \vec{v}_\perp, \vec{b}_\perp)$
- Gluon TMD pdf related to  $\hat{q}$  :  $\hat{q} = \int d^2k_\perp \hat{q}(k_\perp) = C \int d^2k_\perp \rho \phi(x_g, k_\perp)$

# Nuclear Modification : $\Delta\phi$ , $|y_{l_q} - y_l|$ , $R_A$ dependence

$$I_{eA}^{S(D)}(l_{\perp}, l_{q\perp}, \Delta\phi, z) = \frac{d\hat{\sigma}_{eA}^{S(D)}}{d\mathcal{P}} / A \frac{d\hat{\sigma}_{ep}}{d\mathcal{P}} \quad d\mathcal{P} \equiv dx_B dQ^2 dz d^2l_{\perp} d^2l_{q\perp}$$

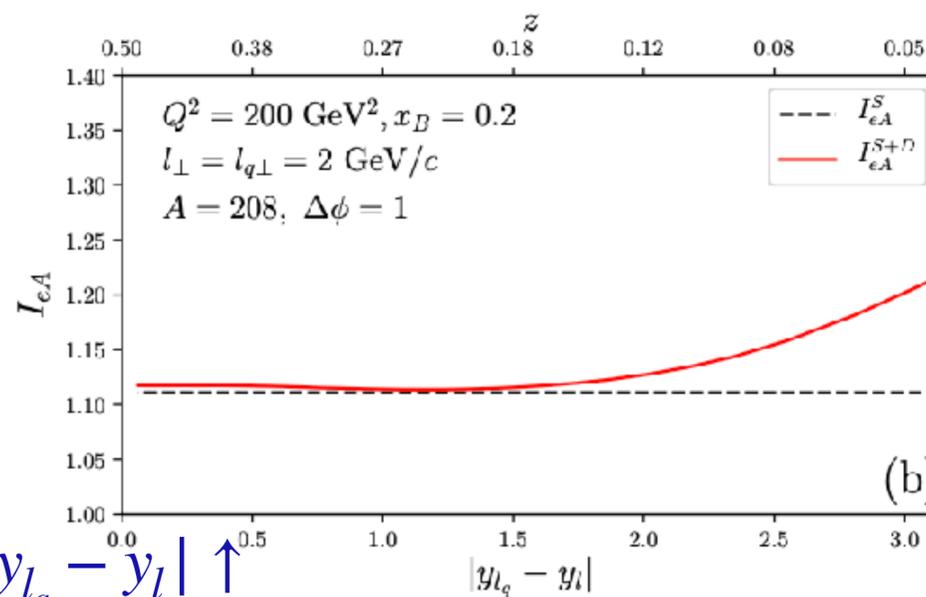


gLPM: dominant



Single scattering dominant

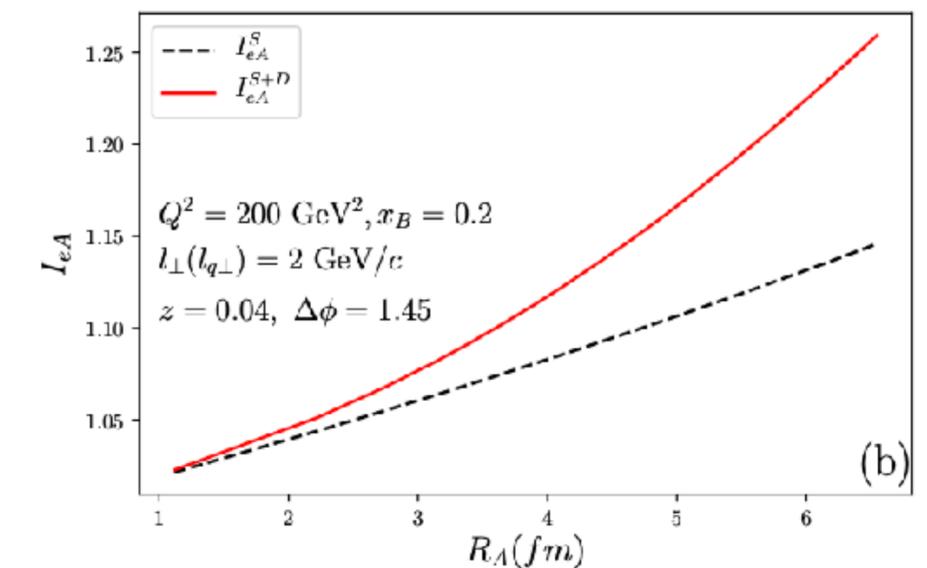
Yuan-Yuan Zhang



$|y_{l_q} - y_l| \uparrow$

LPM disappear, incoherent contri  $\uparrow$

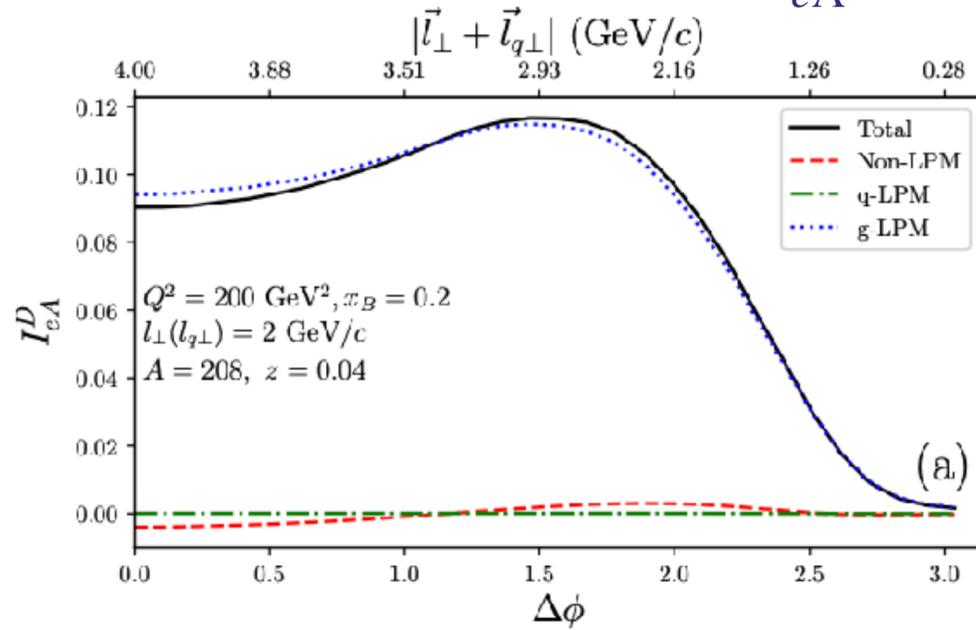
QM2022



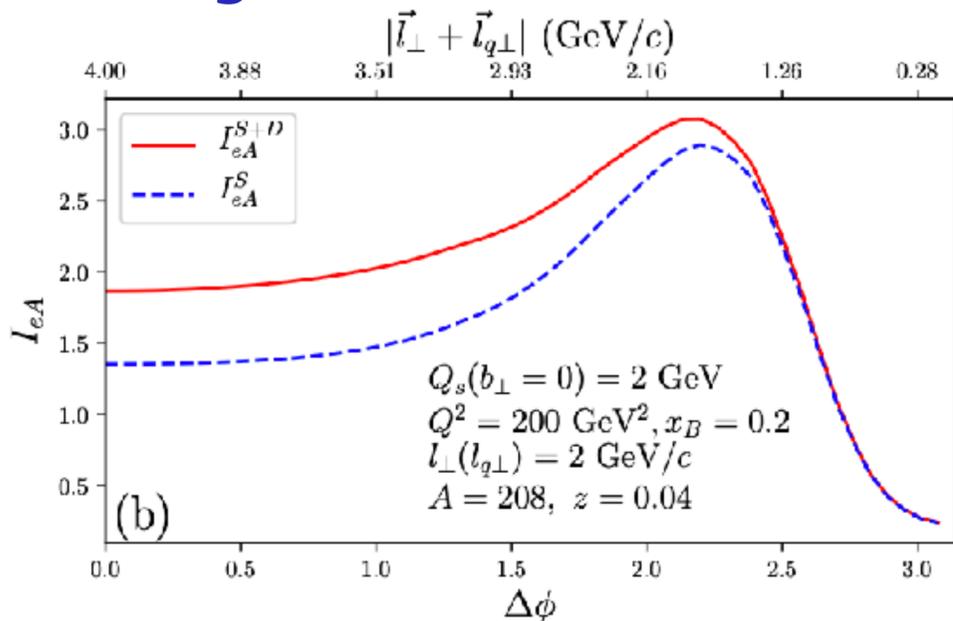
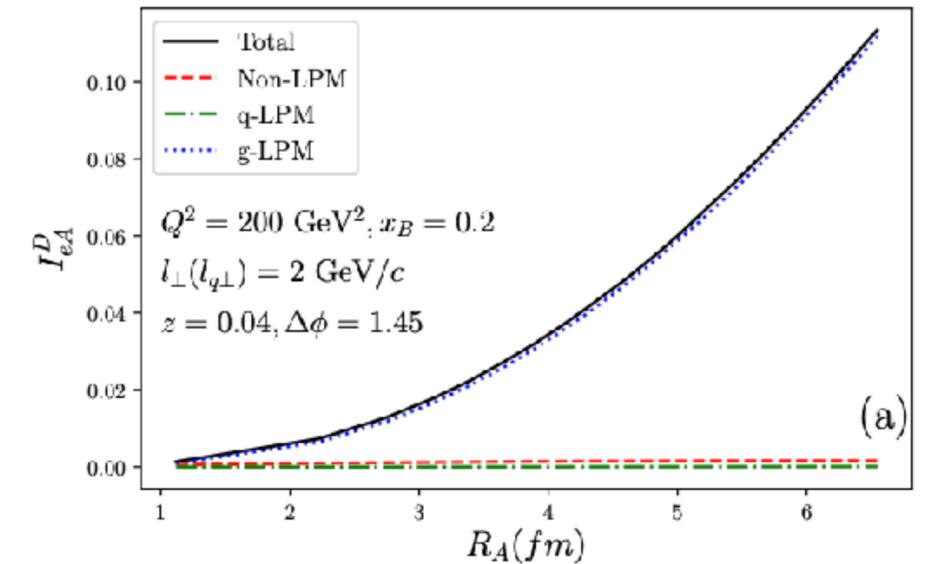
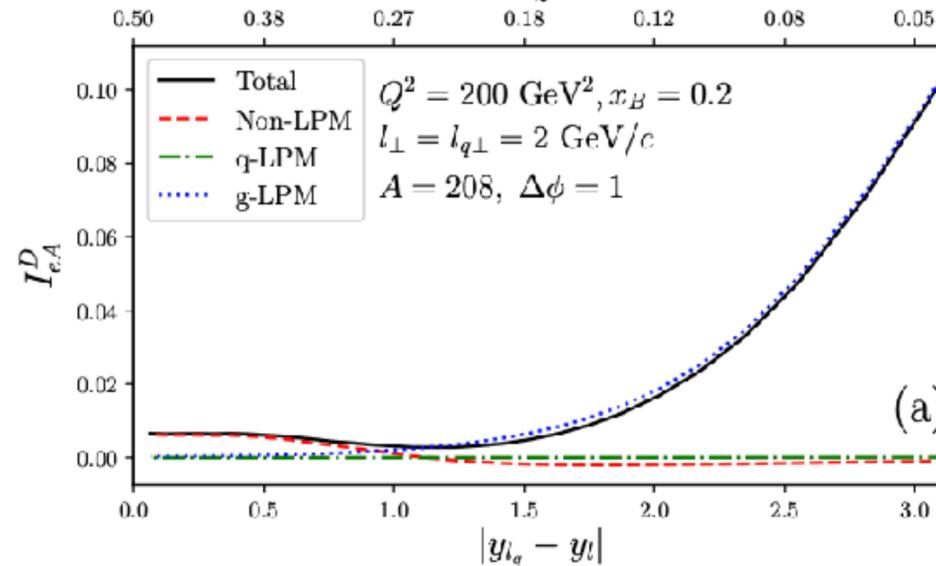
LPM  $\rightarrow$  Non-linear  $R_A$  dependence

# Nuclear Modification : $\Delta\phi$ , $|y_{l_q} - y_l|$ , $R_A$ dependence

$$I_{eA}^{S(D)}(l_{\perp}, l_{q\perp}, \Delta\phi, z) = \frac{d\hat{\sigma}_{eA}^{S(D)}}{d\mathcal{P}} / A \frac{d\hat{\sigma}_{ep}}{d\mathcal{P}} \quad d\mathcal{P} \equiv dx_B dQ^2 dz d^2l_{\perp} d^2l_{q\perp}$$

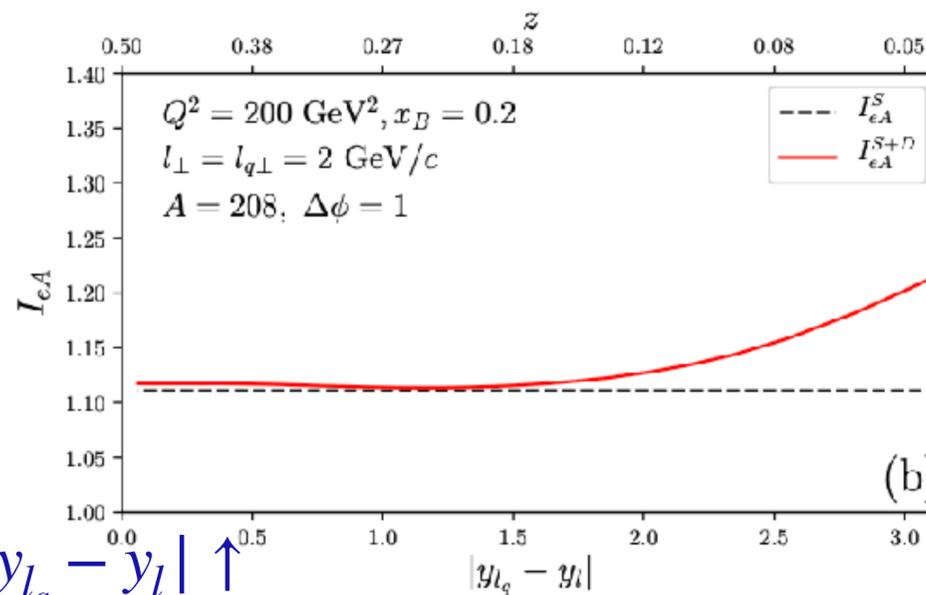


gLPM: dominant



Single scattering dominant

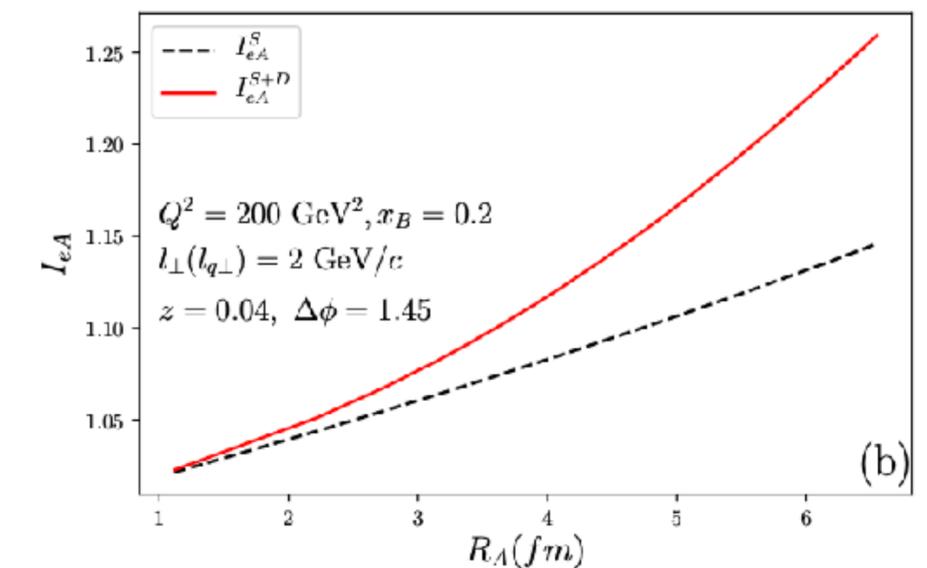
Yuan-Yuan Zhang



$|y_{l_q} - y_l| \uparrow$

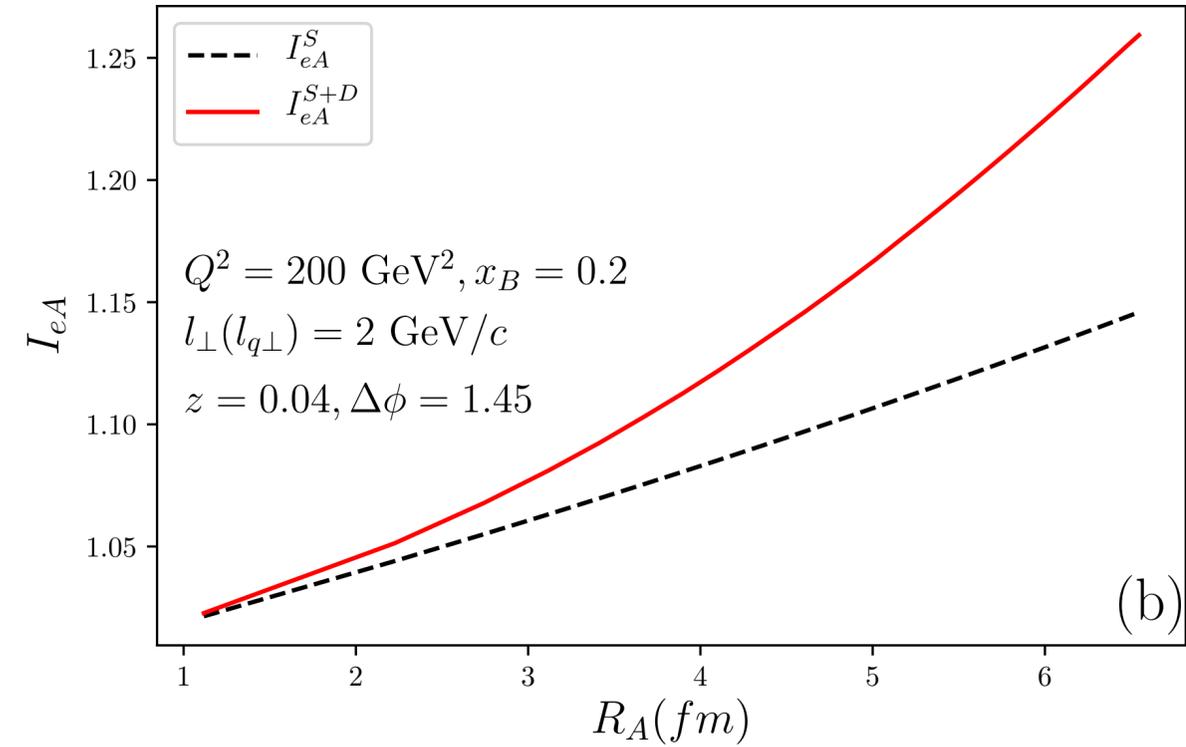
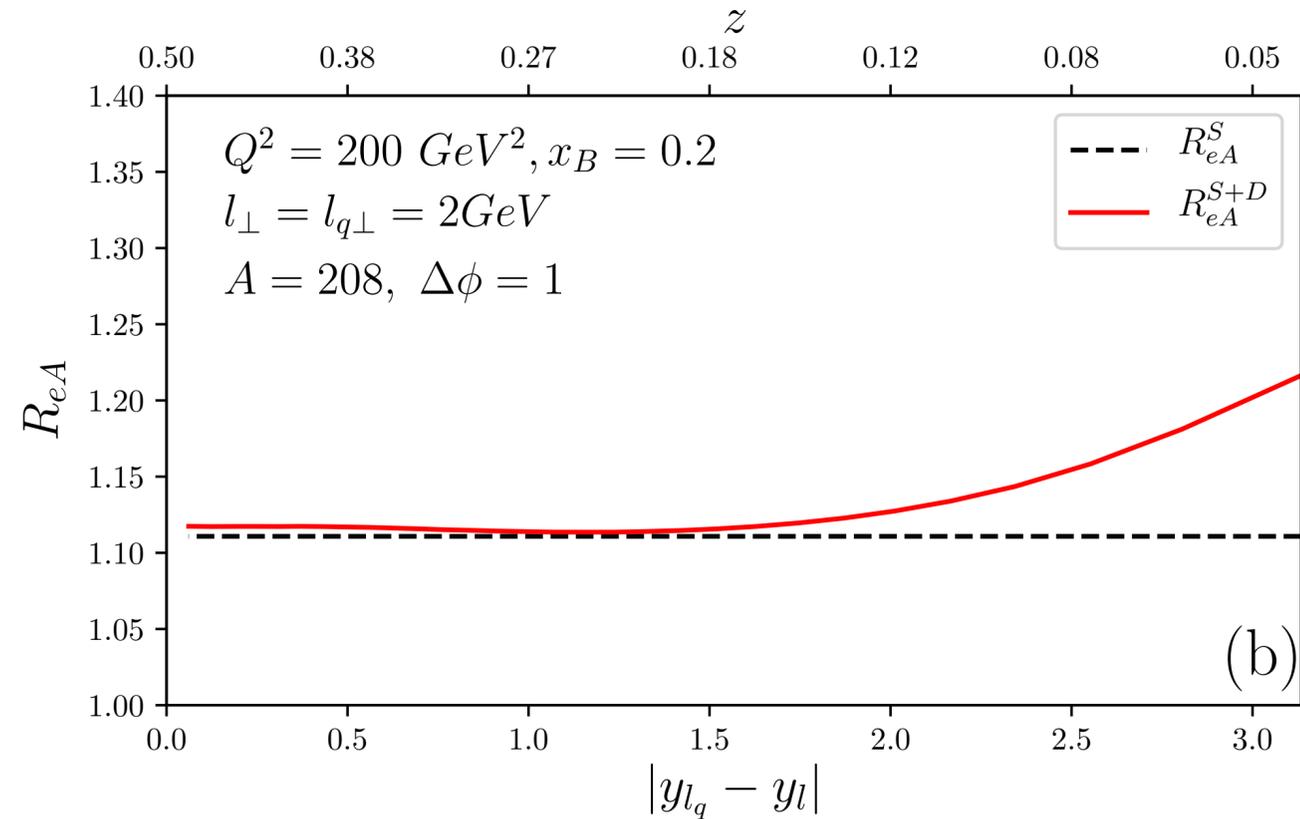
LPM disappear, incoherent contri  $\uparrow$

QM2022



LPM  $\rightarrow$  Non-linear  $R_A$  dependence

# Summary



- Use dijet correlation in e+A to probe gluon TMD pdf  $\phi(x_G, k_\perp)$  or TMD  $\hat{q}(k_\perp)$
- Nuclear modification in single and double scattering depend on  $\phi(x_G, k_\perp)$
- LPM effect and gluon saturation embedded in  $\phi(x_G, k_\perp)$  bring unique features in  $\Delta\phi, |y_{l_q} - y_l|, R_A$  dependence of nuclear modification ratio



29TH INTERNATIONAL  
CONFERENCE ON ULTRARELATIVISTIC  
NUCLEUS - NUCLEUS COLLISIONS

APRIL 4-10, 2022

KRAKÓW, POLAND

Thanks !

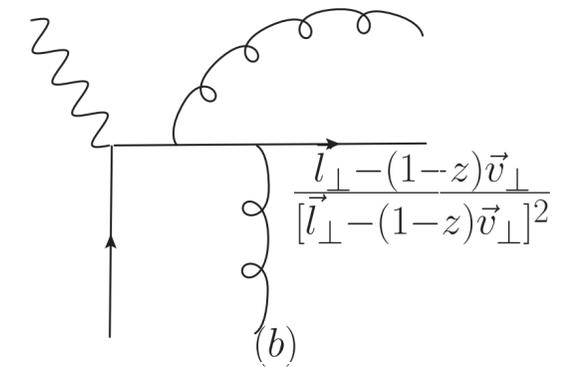


# Backup: Different contribution in double scattering

Contribution divided by how gluon radiated, understand from central cut diagrams

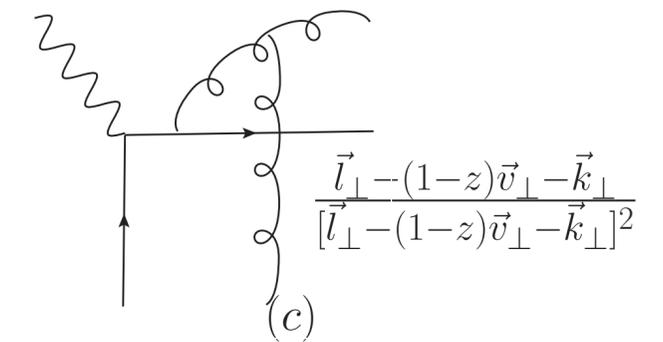
$$\mathcal{N}_g^{\text{qLPM}} = \frac{1}{N_c} (\dots) \left[ 1 - \cos\left(\frac{y_1^- - y_0^-}{\tau_{gf}}\right) \right]$$

$$\tau_{qf} = \frac{2q^-z(1-z)}{[\vec{l}_\perp - (1-z)\vec{v}_\perp]^2}$$



$$\mathcal{N}_g^{\text{gLPM}} = C_A (\dots) \left[ 1 - \cos\left(\frac{y_1^- - y_0^-}{\tau_{gf}}\right) \right]$$

$$\tau_{gf} = \frac{2q^-z(1-z)}{[\vec{l}_\perp - (1-z)\vec{v}_\perp - \vec{k}_\perp]^2}$$



$$\mathcal{N}_g^{\text{nonLPM}} = C_F (\dots)$$

no LPM interference

