

# Modelling particle polarization with spin hydrodynamics

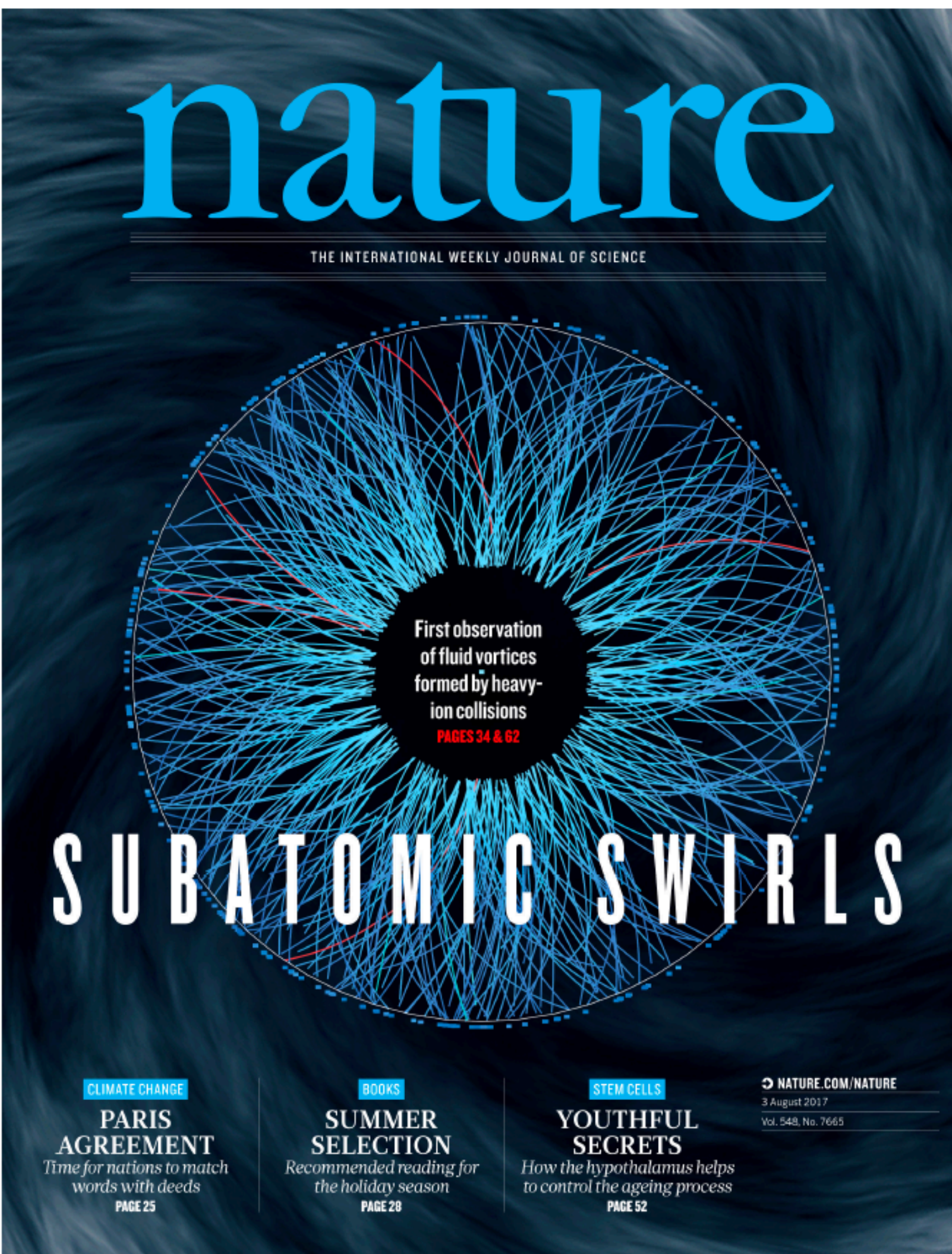
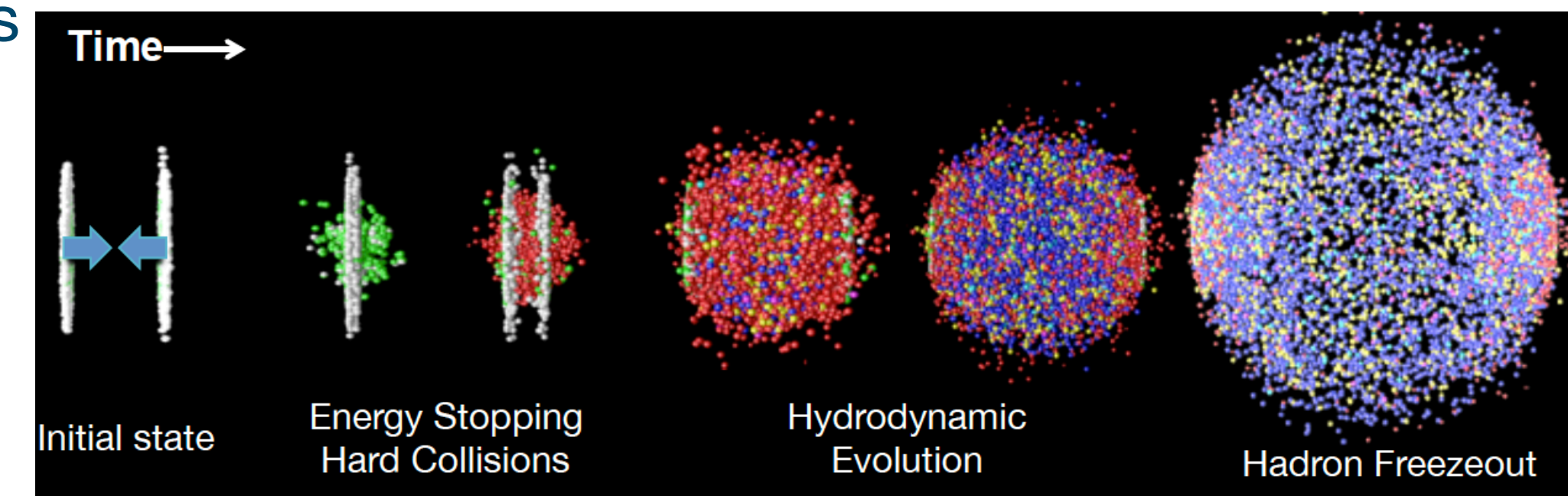
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Relativistic fluid dynamics forms  
the basis of HIC models

Primary reference: 2112.01856

**Collaborators:** W. Florkowski (UJ), R. Ryblewski (IFJ), G. Sophys (IFJ)



**Quark Matter 2022**

**6 Apr 2022**

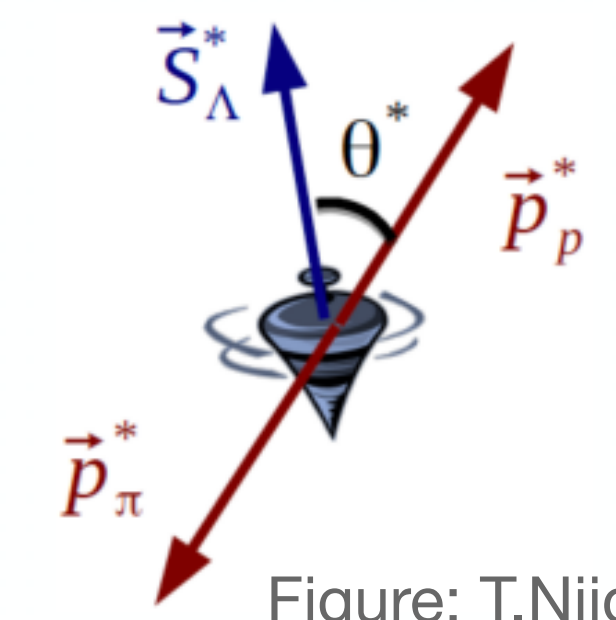


Figure: T.Niida



# Spin polarization in heavy-ion collisions: A new sensitive probe!

Non-central heavy-ion collisions create fireballs with large global orbital angular momenta

F. Becattini, F. Piccinini, J. Rizzo, PRC 77 (2008) 024906

$$\mathbf{L}_{\text{init}} \sim 10^5 \hbar$$

Part of the angular momentum can be transferred from the orbital to the spin part

$$\mathbf{J}_{\text{init}} = \mathbf{L}_{\text{init}} = \mathbf{L}_{\text{final}} + \mathbf{S}_{\text{final}}$$

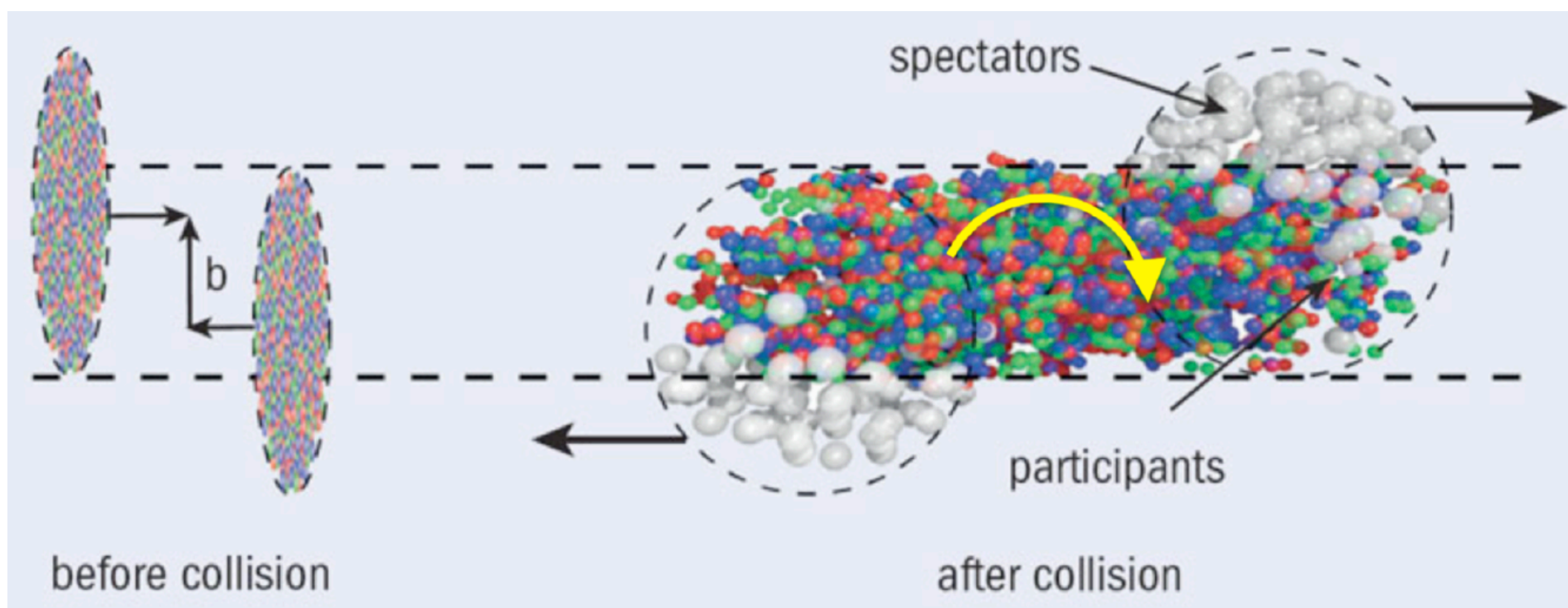


Figure: M. Lisa, talk @ "Strangeness in Quark Matter 2016"

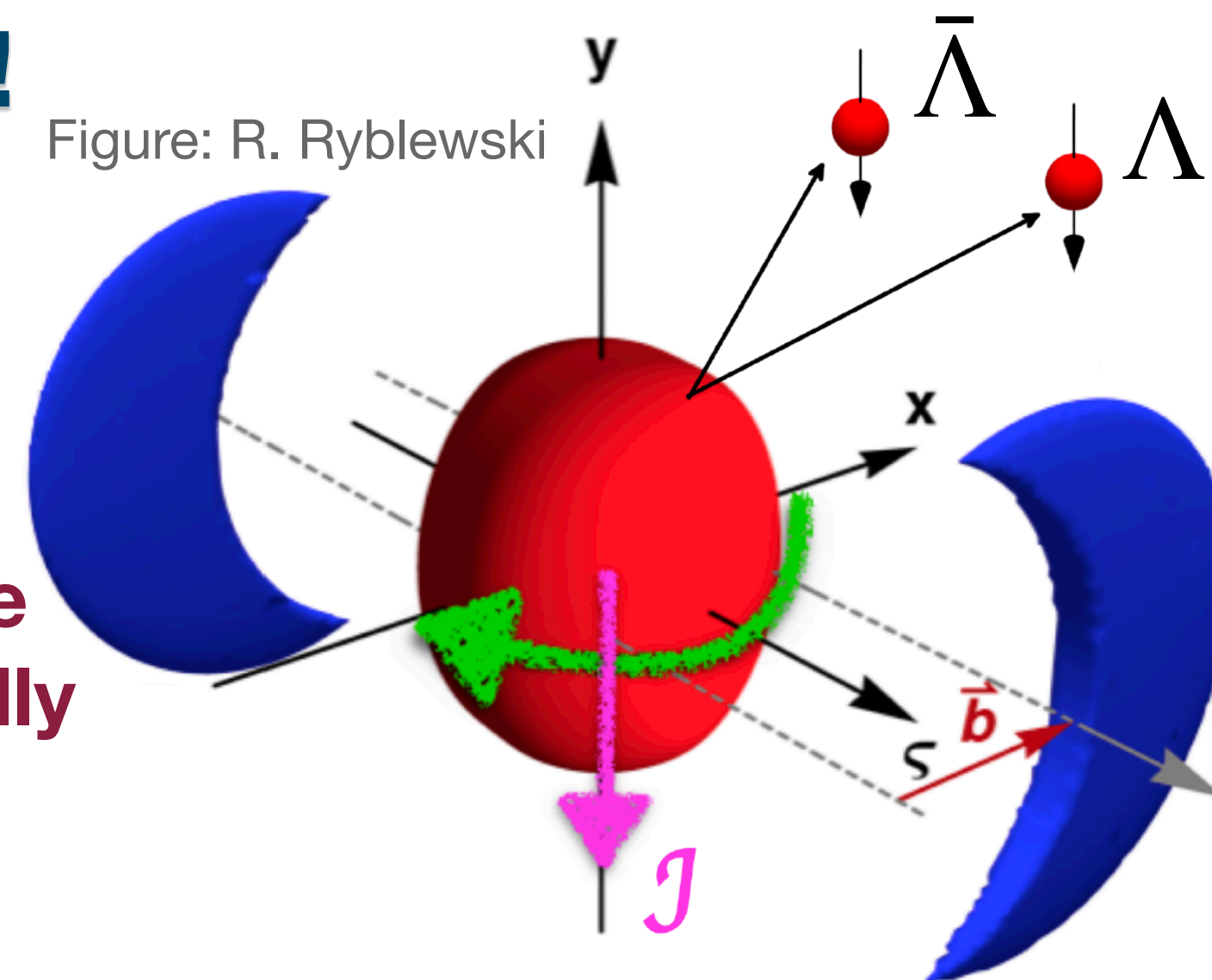
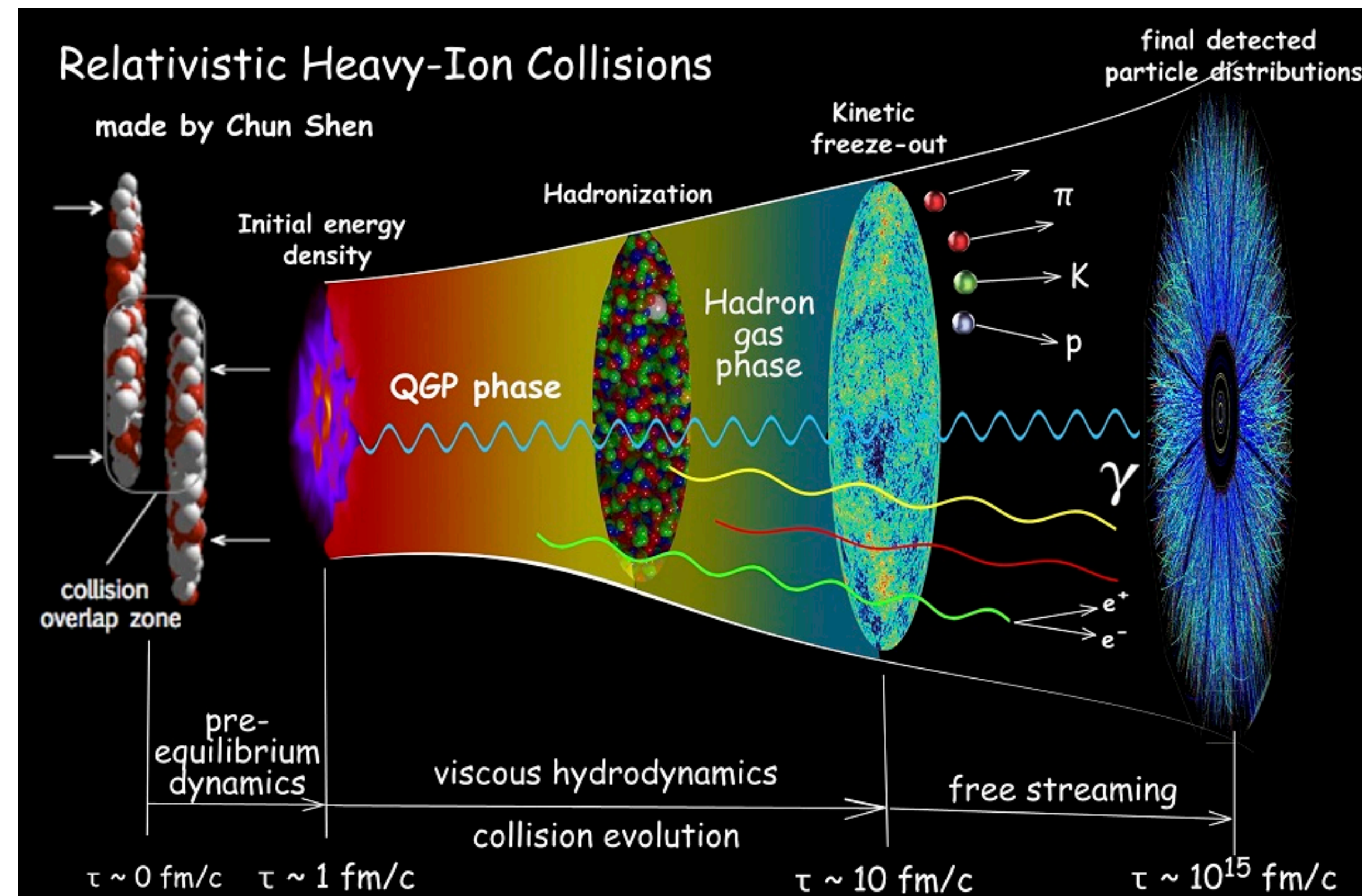


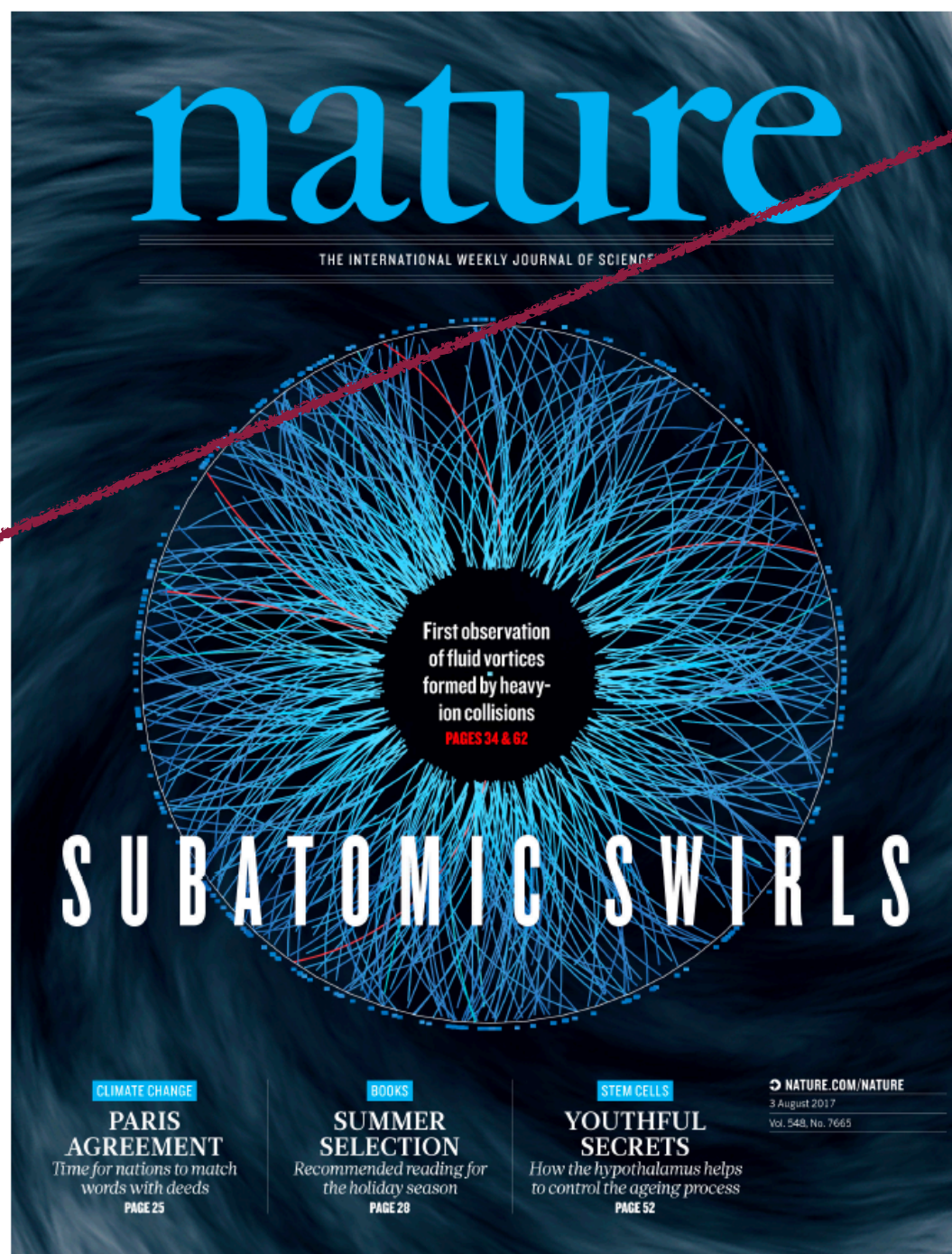
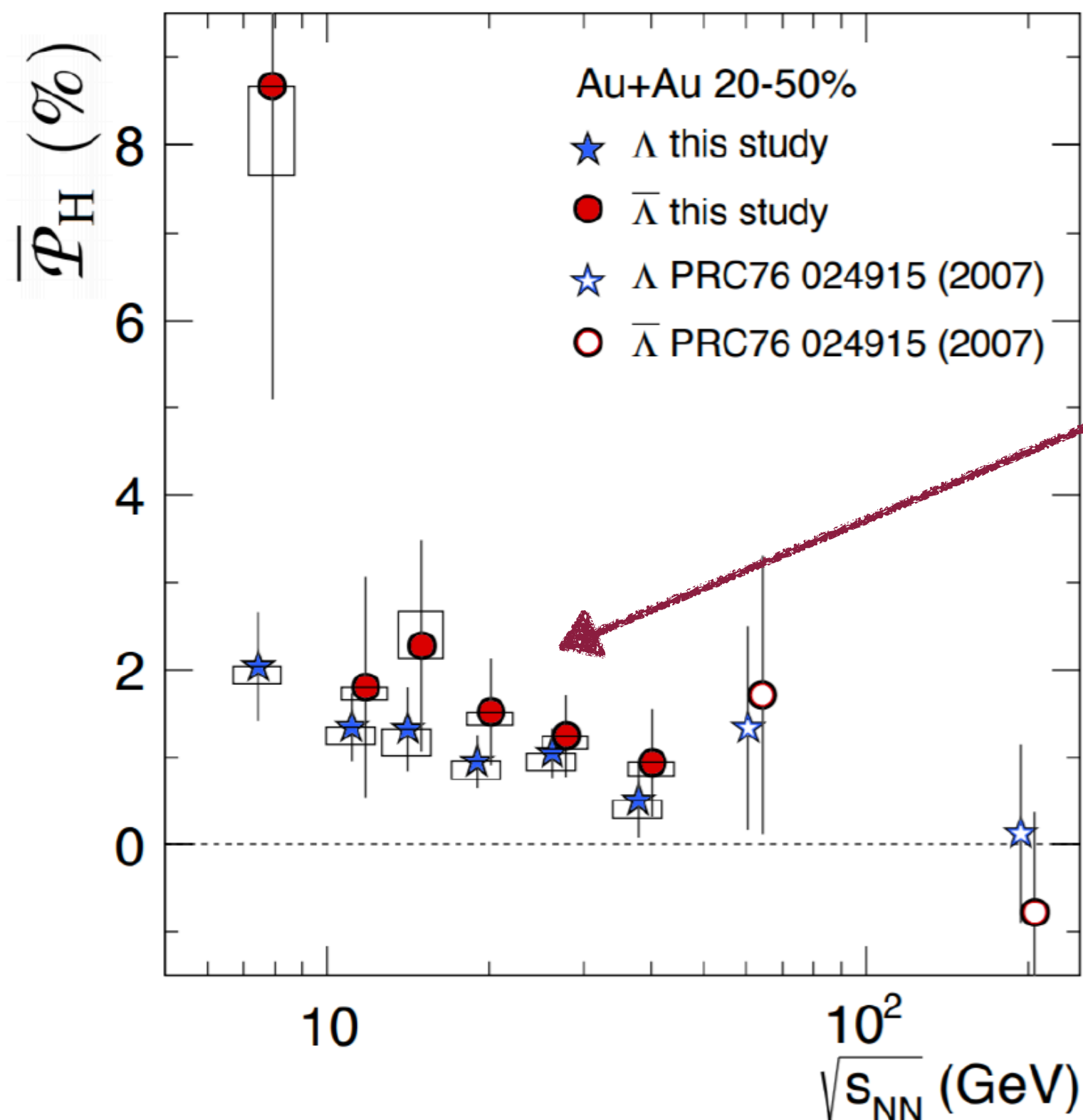
Figure: R. Ryblewski

Emitted particles are expected to be globally polarized along the system's angular momentum



# Experimental measurement of $\Lambda(\bar{\Lambda})$ spin polarization in heavy-ion collisions

L. Adamczyk et al. (STAR) (2017), Nature 548 (2017) 62-65



~2% - small but measurable effect

Self-analysing parity-violating hyperon weak decay allows to measure polarization of  $\Lambda$

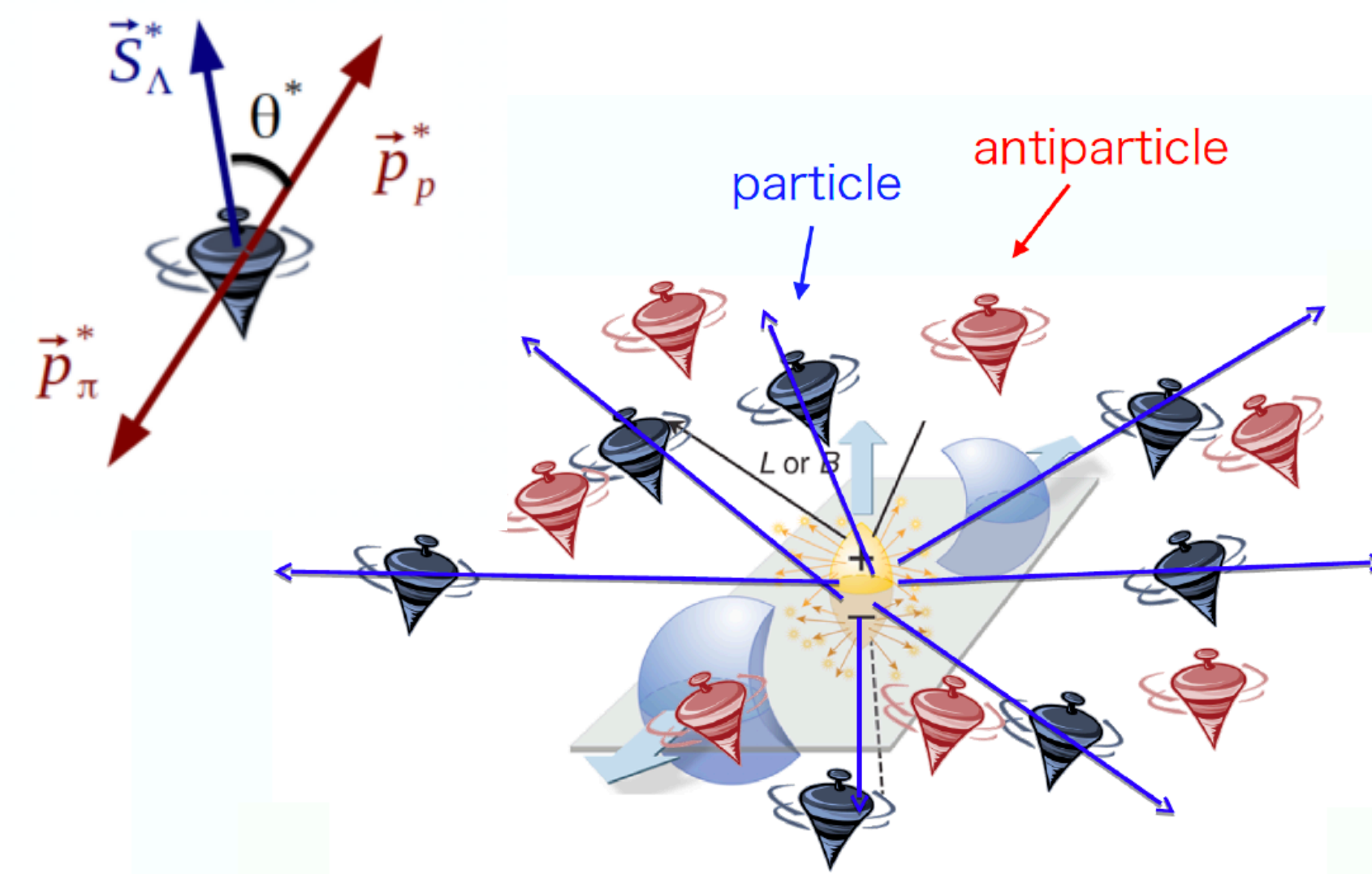


Figure: T.Niida

**QGP** is the *hottest, least viscous, and most vortical* fluid ever produced

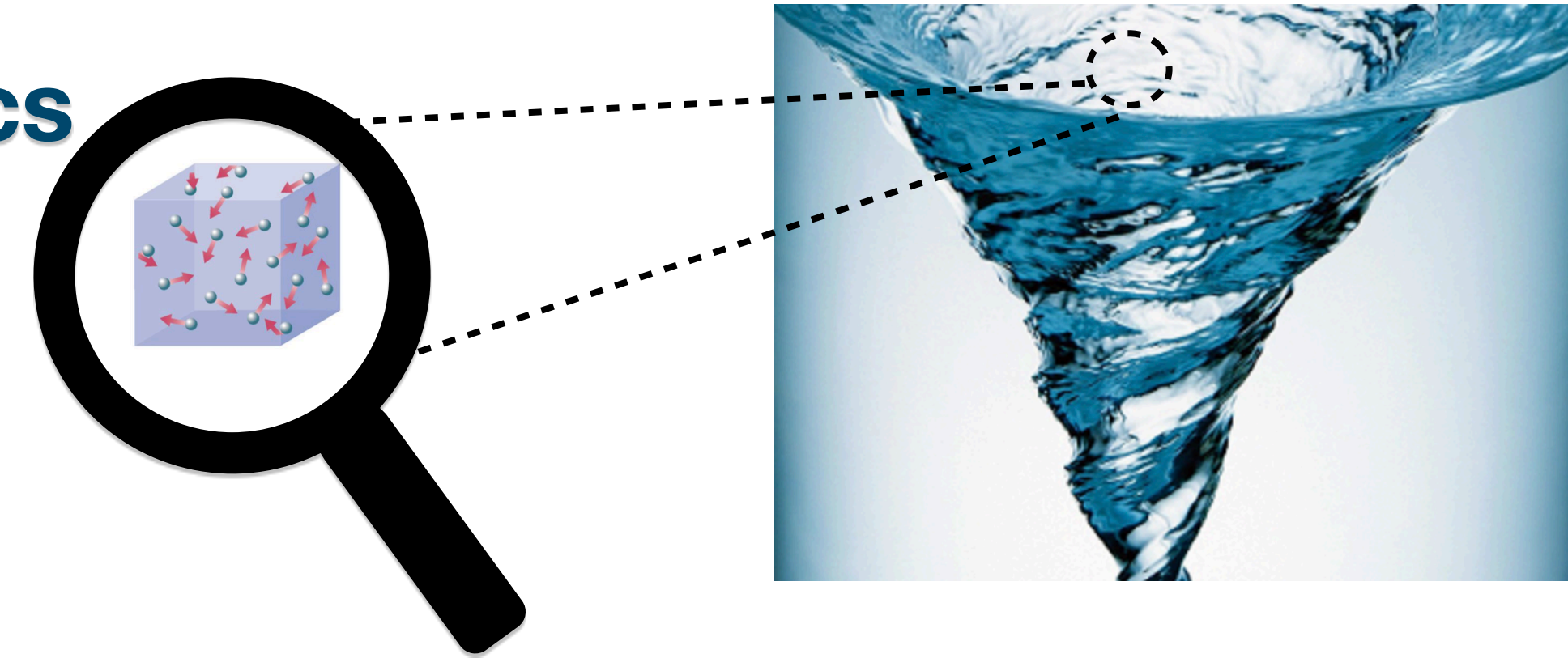
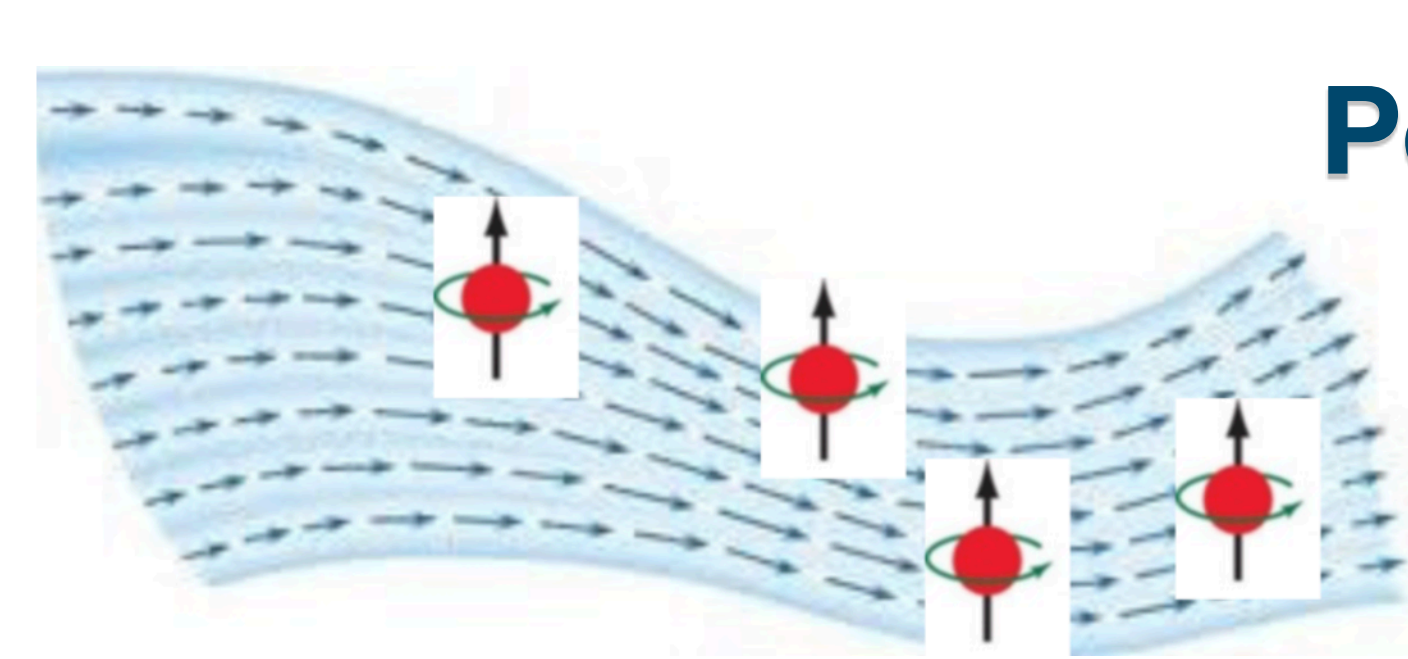
$$\omega = (P_\Lambda + P_{\bar{\Lambda}})k_B T / \hbar \sim 0.6 - 2.7 \times 10^{22} \text{ s}^{-1}$$

$$P_\Lambda \approx \frac{1}{2} \frac{\omega}{T} + \frac{\mu_\Lambda B}{T} \quad P_{\bar{\Lambda}} \approx \frac{1}{2} \frac{\omega}{T} - \frac{\mu_\Lambda B}{T}$$

$$\frac{dN}{d\Omega^*} = \frac{1}{4\pi} (1 + \alpha_H \mathbf{P}_H \cdot \mathbf{p}_p^*)$$

$P_\Lambda \approx P_{\bar{\Lambda}}$  → first direct observation of spin

# Perfect-fluid spin hydrodynamics



Relativistic kinetic theory  
formulation of ideal fluid

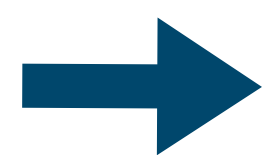


For dilute systems, the derivation of fluid  
dynamics can be done starting from the  
underlying kinetic theory

Quantum RKT

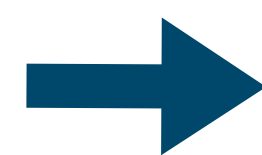
$$\left( \gamma_\mu K^\mu - m \right) \mathcal{W}(x, k) = C[\mathcal{W}(x, k)]$$

$$K^\mu = k^\mu + \frac{i}{2} (\hbar \partial^\mu)$$



$$k^\mu \partial_\mu \mathcal{F}_{\text{eq}}(x, k) = 0$$

$$k^\mu \partial_\mu \mathcal{A}_{\text{eq}}^\nu(x, k) = 0$$



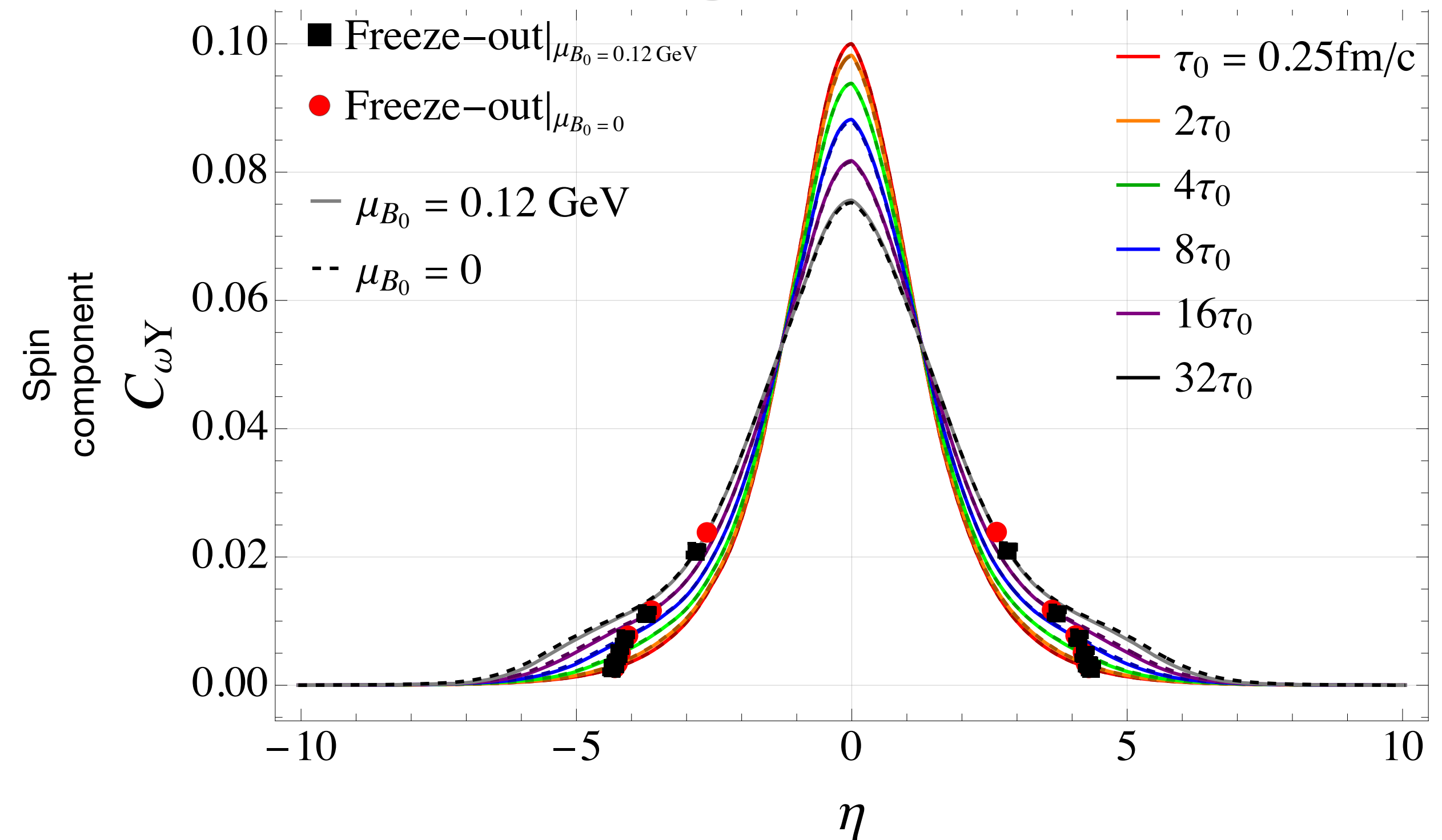
$$\partial_\mu N^\mu = 0$$

$$\partial_\nu T^{\mu\nu} = 0$$

$$\partial_\lambda S^{\lambda, \mu\nu} = 0$$

semi-classical expansion

moments method



# Classical spin treatment - perfect fluid

W. Florkowski, R. Ryblewski, A. Kumar, Prog. Part. Nucl. Phys. 108 (2019) 103709  
 J.-W. Chen, J.-y. Pang, S. Pu, Q. Wang, PRD 89 (9) (2014) 094003

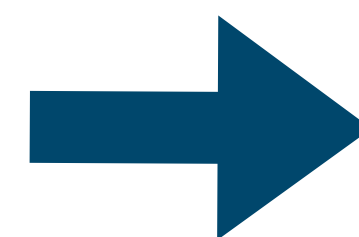
$$f_{\text{eq}}^{\pm}(x, p, s) = \exp \left( -p \cdot \beta(x) \pm \xi(x) + \frac{1}{2} \omega_{\alpha\beta}(x) s^{\alpha\beta} \right)$$

$$\int dS \dots = \frac{m}{\pi \mathfrak{B}} \int d^4 s \delta(s \cdot s + \mathfrak{B}^2) \delta(p \cdot s) \dots$$

$$N_{\text{eq}}^{\mu} = \int dP \int dS p^{\mu} [f_{\text{eq}}^{+}(x, p, s) - f_{\text{eq}}^{-}(x, p, s)]$$

$$T_{\text{eq}}^{\mu\nu} = \int dP \int dS p^{\mu} p^{\nu} [f_{\text{eq}}^{+}(x, p, s) + f_{\text{eq}}^{-}(x, p, s)]$$

$$S_{\text{eq}}^{\lambda\mu\nu} = \int dP \int dS p^{\lambda} s^{\mu\nu} [f_{\text{eq}}^{+}(x, p, s) + f_{\text{eq}}^{-}(x, p, s)]$$



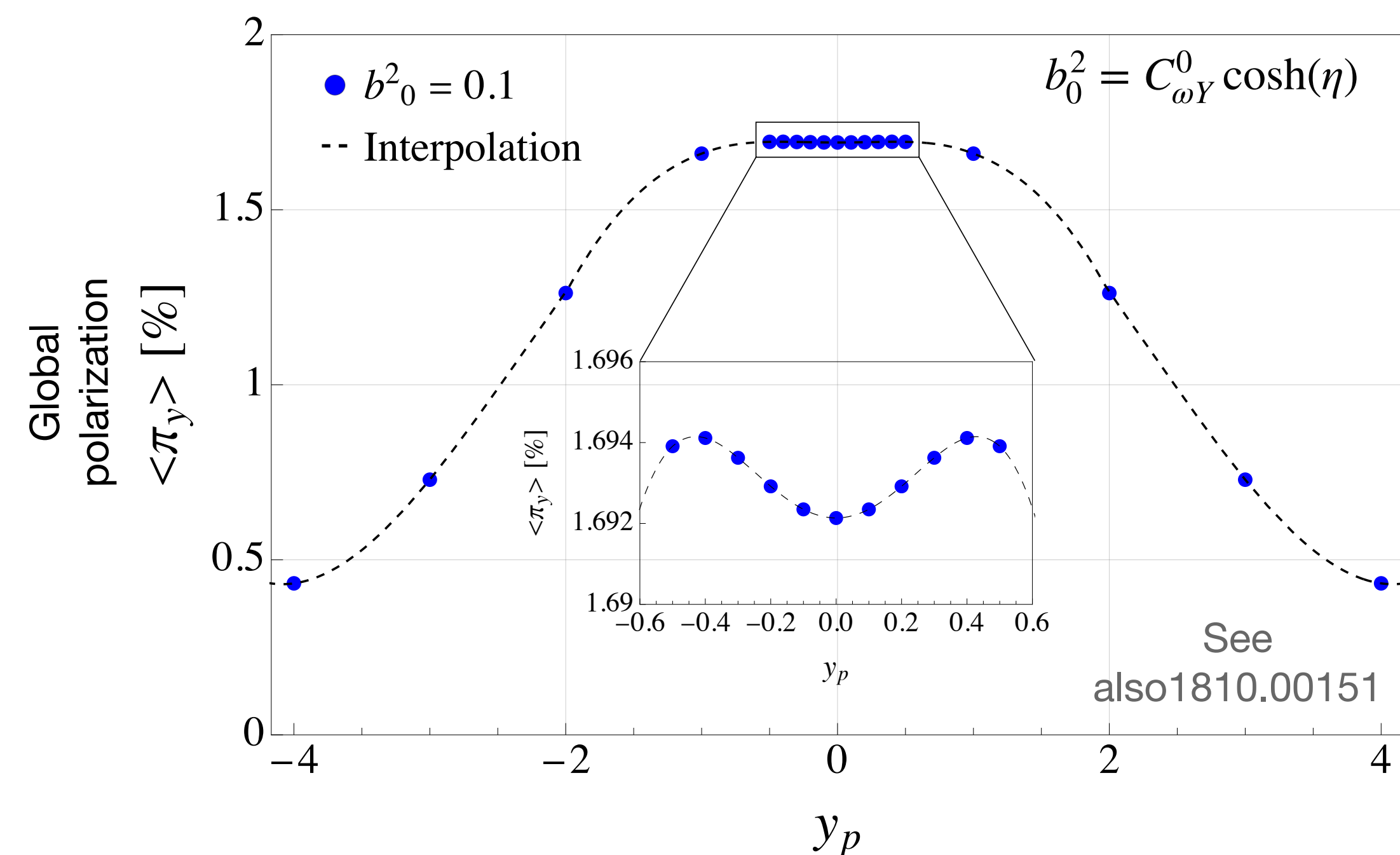
## Explicit constitutive relations

$$N_{\text{eq}}^{\alpha} = n u^{\alpha}$$

$$T_{\text{eq}}^{\alpha\beta}(x) = \varepsilon u^{\alpha} u^{\beta} - P \Delta^{\alpha\beta}$$

$$S_{\text{eq}}^{\lambda,\mu\nu} = S_{\text{GLW}}^{\lambda,\mu\nu} = C \left( n_0(T) u^{\lambda} \omega^{\mu\nu} + S_{\Delta\text{GLW}}^{\lambda,\mu\nu} \right)$$

$$S_{\Delta\text{GLW}}^{\alpha,\beta\gamma} = \mathcal{A}_0 u^{\alpha} u^{\delta} u^{[\beta} \omega_{\delta}^{\gamma]} + \mathcal{B}_0 \left( u^{[\beta} \Delta^{\alpha\delta} \omega_{\delta}^{\gamma]} + u^{\alpha} \Delta^{\delta[\beta} \omega_{\delta}^{\gamma]} + u^{\delta} \Delta^{\alpha[\beta} \omega_{\delta}^{\gamma]} \right)$$



W. Florkowski, A. Kumar, R. Ryblewski, R. Singh, Phys.Rev.C 99 (2019) 4, 044910  
 R. Singh, G. Sophys, R. Ryblewski, Phys.Rev.D 103 (2021) 7, 074024  
 R. Singh, M. Shokri, R. Ryblewski, Phys.Rev.D 103 (2021) 9, 094034  
 W. Florkowski, R. Ryblewski, R. Singh, G. Sophys, Phys.Rev.D 105 (2022) 5, 054007

**For  $|\omega_{\mu\nu}| < 1$  one obtains the formalism that agrees with that based on the quantum description of spin (in the GLW version).**

$$\langle \pi_{\mu} \rangle = \frac{\int dP \langle \pi_{\mu} \rangle_p E_p \frac{d\mathcal{N}(p)}{d^3p}}{\int dP E_p \frac{d\mathcal{N}(p)}{d^3p}} \equiv \frac{\int d^3p \frac{d\Pi_{\mu}^*(p)}{d^3p}}{\int d^3p \frac{d\mathcal{N}(p)}{d^3p}}$$