

# CHARMONIUM ABUNDANCE AS A PROBE FOR REMNANTS OF CONFINEMENT

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Using a simple analytically solvable model, we argue that it makes charmonium abundance the ideal probe of remnants of confinement, expected to be present at arbitrary high temperature [1,2] but thermodynamically irrelevant, since the nearest neighbor thermal scale  $1/T$  is parametrically smaller than the confinement scale.

[1] D. Zwanziger, Phys. Rev. D **69**, 016002 (2004)

[2] G. S. Bali, J. Fingberg, U. M. Heller, F. Karsch, and K. Schilling, Phys. Rev. Lett. **71**, 3059 (1993)

The rate equations for  $J/\psi$  abundance are given by the following equations [3]:

$$\frac{dN_{J/\psi}}{d\tau} = -\Gamma(T)N_{J/\psi} + \alpha(T)\frac{N_{c\bar{c}}^2}{V} = -\frac{dN_{c\bar{c}}}{d\tau}$$

with

$$\Gamma \sim \lambda_D T \quad \text{and} \quad \alpha \sim \lambda_R \frac{1}{T^2}.$$

$\lambda_D$  and  $\lambda_R$  are dimensionless parameters (dissociation and recombination).

$\sigma_{ff \rightarrow J/\psi X} \ll \sigma_{ff \rightarrow ccX}$ ,  $N_{J/\psi} \ll N_{c\bar{c}}$  and the dependence on regeneration on  $N_{c\bar{c}}^2$ , the regeneration term will soon take over, leading to an enhancement of  $N_{J/\psi}$ .

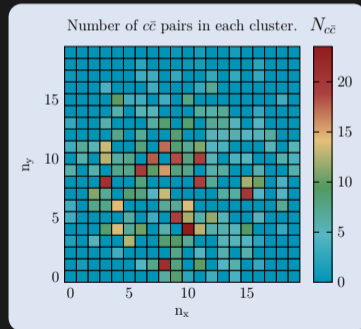
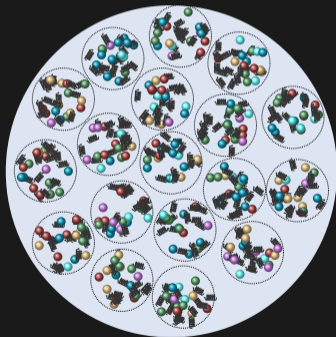
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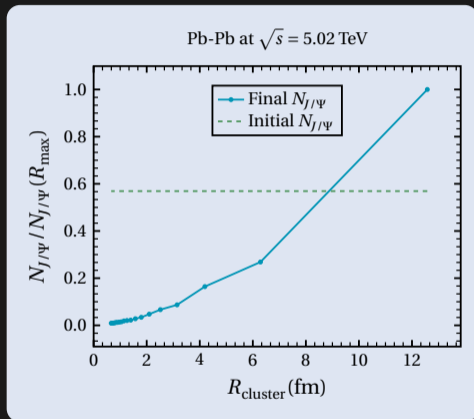
[3] R. L. Thews, M. Schroedter and J. Rafelski, Phys. Rev. C **63**, 054905 (2001)

## REMNANTS OF CONFINEMENT

Let us now implement the assumption that charm quarks cannot move beyond a volume of  $V_c = R_c^3$ .

- One has a “locally” quasi-free gas, however, all clusters of volume  $V_c$  are color neutral.
- Solving in each bin of transverse area and summing together will give the final  $N_{J/\psi}$ , and it will generally be much smaller than if  $c$  and  $\bar{c}$  are allowed to move around freely.





- Calculate  $R_{AA}$  as a function of  $R_c$ .
- Look for the  $J/\psi$  data.