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Gluodynamics

# Intermediate mass dileptons as a speedometer of QGP equilibration

(arXiv:2112.13876)

Quark Matter 2022

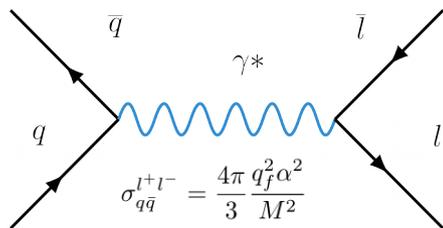
Maurice Coquet, Xiaojian Du, Jean-Yves Ollitrault, Sören Schlichting, Michael Winn

# Introduction : The ideal spectrum

- Dileptons are created throughout the history of the collision → probe entire space-time dynamics including early stages and thermalization

- At LO, production by quark-antiquark annihilation :

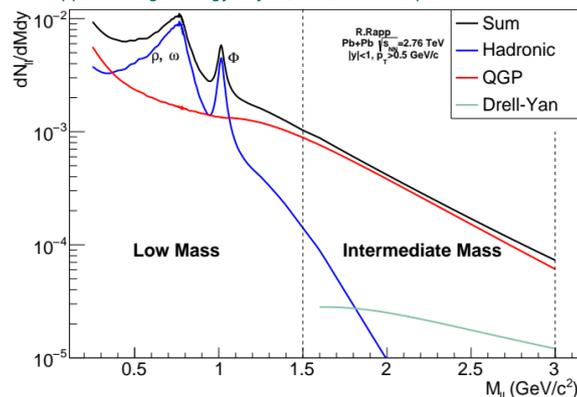
$$\frac{dN^{l^+l^-}}{d^4x d^4K} = \int \frac{d^3p_1}{(2\pi)^3 2p_1} \frac{d^3p_2}{(2\pi)^3 2p_2} f_q(x, \mathbf{p}_1) f_{\bar{q}}(x, \mathbf{p}_2) |\mathcal{A}|^2 (2\pi)^4 \delta^{(4)}(P_1 + P_2 - K),$$



- Thermal dileptons from QGP dominate invariant mass spectrum for  $M \gtrsim 1.5 \text{ GeV}/c^2$

→ **how early pre-equilibrium effects manifest themselves in the spectra of dileptons ?**

R.Rapp, Adv. High Energy Phys. (2013) 148253, private communication



Start by considering local thermal equilibrium (LTE):

- Assuming boost invariance of the expanding QGP at early times and neglecting transverse flow, production rate **depends only on transverse mass  $M_t$**  :

$$\frac{dN^{l^+l^-}}{d^4K} = C \int d\mathbf{x}_\perp \int_0^\infty \tau d\tau \int_{-\infty}^{+\infty} dy_f \exp\left(-\frac{M_t \cosh(y - y_f)}{T(\mathbf{x}_\perp, \tau)}\right)$$

- Further assuming a conformal equation of state and a uniform temperature profile in transverse plane, one finds the McLerran-Tomeila formula :

$$\left(\frac{dN^{l^+l^-}}{d^4K}\right)_{\text{ideal}} = \frac{32 N_c \alpha^2 \sum_f q_f^2 A_\perp (\tau T^3)^2}{\pi^4 M_t^6}$$

# Pre-equilibrium dynamics

Effects of pre-equilibrium dynamics :

- Quarks are underpopulated relative to gluons at early times
- Momentum distributions of quarks and gluons are anisotropic (rapid longitudinal expansion)

→ Encapsulated in parametrization of the quark distribution:

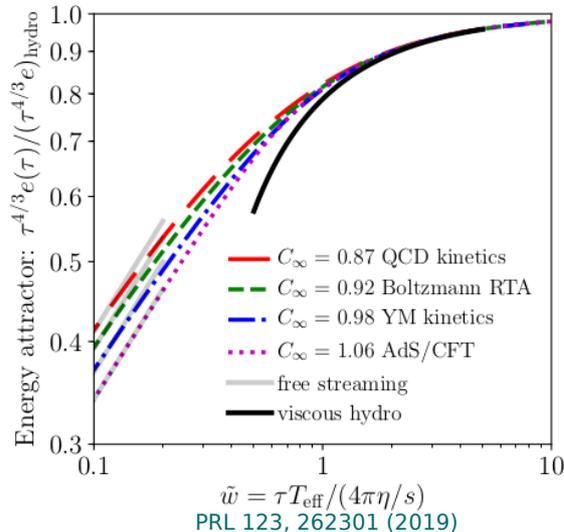
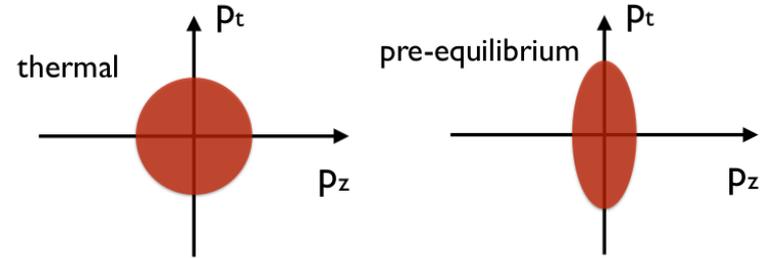
Strickland PRD 92, 025026 (2015), PRD 99 3, 034015 (2019)

$$f_q(\tau, p_T, p_L) = q_s(\tau) f_{FD} \left( -\sqrt{p_T^2 + \xi^2(\tau) p_L^2} / \Lambda(\tau) \right)$$

$q_s$  : quark suppression factor,  $\xi$  : anisotropy parameter,  $\Lambda$  : anisotropic temperature

→ Evolution of parameters computed in **QCD kinetics**

(X.Du Poster Session 1 T01) X. Du, S. Schlichting: PRD 104, 054011 (2021)



- Quark suppression implies suppression of dilepton production, which is a global factor → **preserving  $M_t$  scaling**  
Deviation from LTE calculation, more pronounced at high mass (↔ early times ↔ fewer quarks). Relative suppression scales like :

$$Re^{-1}(M_t) \equiv \frac{\eta}{s} \frac{M_t^2}{\tau T^3}$$

$\eta/s$  : early time shear to entropy density ratio → departure from local thermal equilibrium

- Fast longitudinal expansion → pressure asymmetry → momentum anisotropy which **breaks  $M_t$  scaling**, favoring small masses for a given  $M_t$  value

# Quantitative results

- Breaking of  $M_t$  scaling modest, compared to overall suppression of production yield due to quark underpopulation
- Spectra well fitted by the following formula :

$$\frac{dN^{l^+l^-}}{d^4K} \simeq \left( \frac{dN^{l^+l^-}}{d^4K} \right)_{\text{ideal}} \frac{\left(1 + a \frac{\eta}{s} M_t^2 / n\right)^{-n}}{\sqrt{1 + b \frac{\eta}{s} M_t^2}}$$

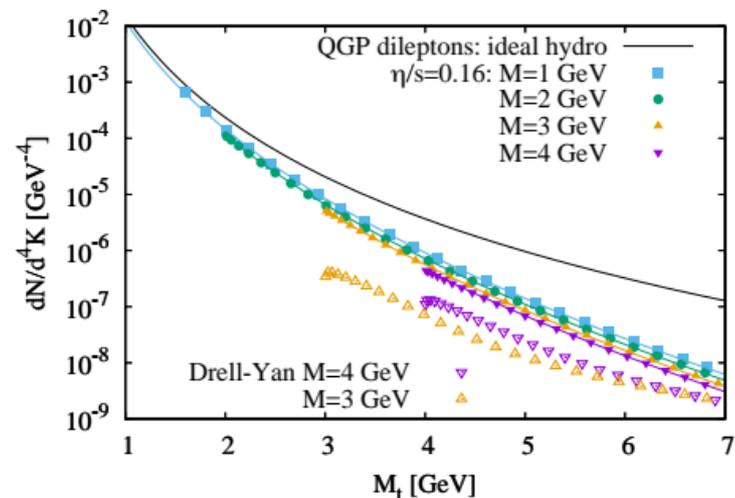
- a :  $M_t$  dependent suppression of the production, b : breaking of  $M_t$  scaling
- corrections to ideal yield linear in  $\eta/s$  for small  $\eta/s$ .
- $\sim 1/M$  for large  $\eta/s$ , showing breaking of  $M_t$  scaling.

- larger  $\eta/s$  → later thermalization → temperature starts decreasing as  $\tau^{-1/3}$  later on → lower initial temperature for fixed final energy density

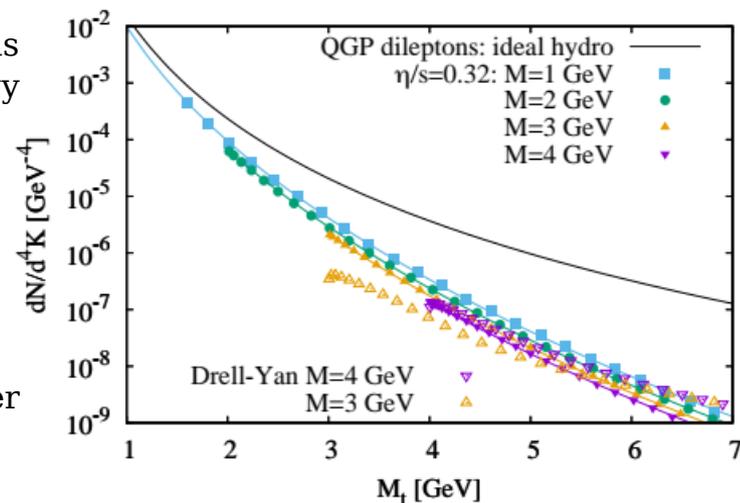
- Inverse slope of the  $M_t$  spectrum → effective temperature :

$$T_{\text{eff}}(M_t) \equiv - \left[ \frac{d}{dM_t} \ln \left( \frac{dN^{l^+l^-}}{d^4K} \right) \right]^{-1} \rightarrow T_{\text{eff}}(M_t) \simeq \frac{M_t}{6 + 2a \frac{\eta}{s} M_t^2}$$

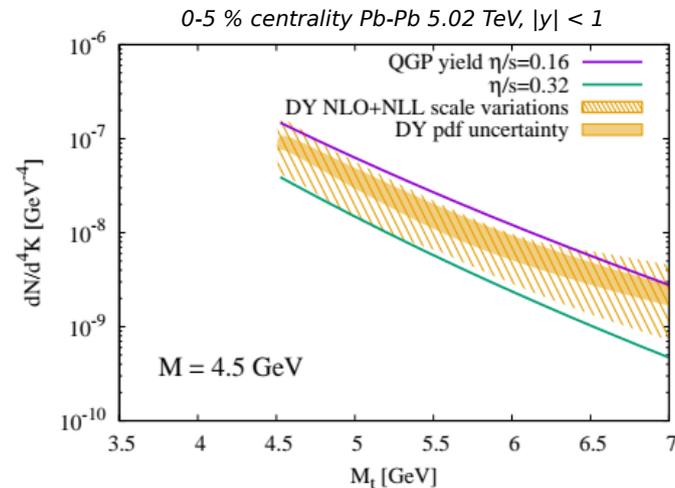
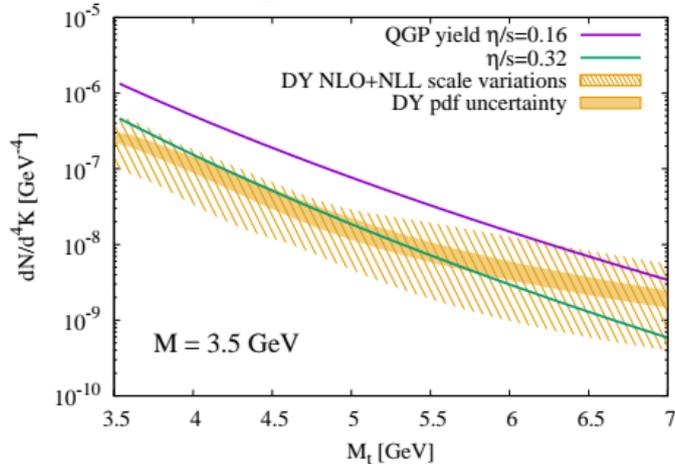
- Drell-Yan dominates spectrum at high  $M_t$  → DY enhanced for larger values of M at fixed  $M_t$ , **opposite behavior to QGP emission.**



0-5 % centrality Pb-Pb 5.02 TeV,  $|y| < 1$



# Backgrounds & scalings



- Main backgrounds in intermediate mass region :  
 → semileptonic decays of heavy flavours (rejectable based on displacement from primary vertex)  
 → Drell-Yan production in the initial state.
- Drell-Yan contribution calculated using DYTurbo software, evaluated at NLO with resummed NLL (+Sudakov form factor includes non-perturbative contribution)

## Scaling properties

- System size/centrality: Ideal spectrum scales with system size like  $(dN_{\text{ch}}/d\eta)^{4/3}$   
 → scales like space-time volume  
 → pre-equilibrium effects (a & b parameters) scale like  $(dN_{\text{ch}}/d\eta)^{-1/3}$  (up to event by event fluct.)
- Collision energy: Ideal spectrum scales with  $\sqrt{s_{NN}}$  like  $(dN_{\text{ch}}/d\eta)^2$   
 → pre-equilibrium effects scale like  $(dN_{\text{ch}}/d\eta)^{-1}$

# Conclusion

- Dilepton spectrum gives **access to  $\eta/s$**  that controls the **equilibration time** ( $\rightarrow$  can be inferred by measuring the slope), as well as early-stage chemistry
- Distinguish experimentally QGP production from Drell-Yan production by studying variation of yield as a function of  **$M$  at fixed  $M_t$**
- Our calculation : neglected transverse flow, sizeable effect for  $\tau_T > 5$  fm/c. Estimated fraction of dilepton yield emitted after  $\tau_T$  :  
for  $M_t=2$  GeV 25 % of spectrum, for  $M_t=3$  GeV only 4 % of the spectrum is late emission.

# Thank you !