

# $Y(nS)$ meson production in Pb+Pb and $pp$ collisions

*Alexandre Lebedev for the ATLAS Collaboration*

ISU<sup>®</sup>



29<sup>TH</sup> INTERNATIONAL CONFERENCE ON ULTRARELATIVISTIC  
NUCLEUS-NUCLEUS COLLISIONS, **APRIL 4-10 2022**, KRAKÓW, POLAND

# Physics motivation

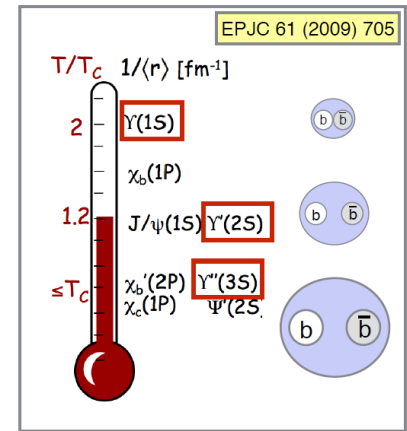
## Upsilon can serve as an important tool for studying QGP

In nucleus-nucleus collisions:

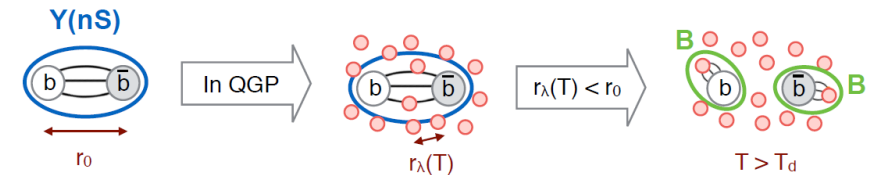
- The three Y states have similar kinematics, but different binding energies.  
-> QGP “thermometer” (sequential melting).
- Very different non-prompt fraction and regeneration compared to charmonia.

QGP-like signatures have been observed in small systems, even in pp

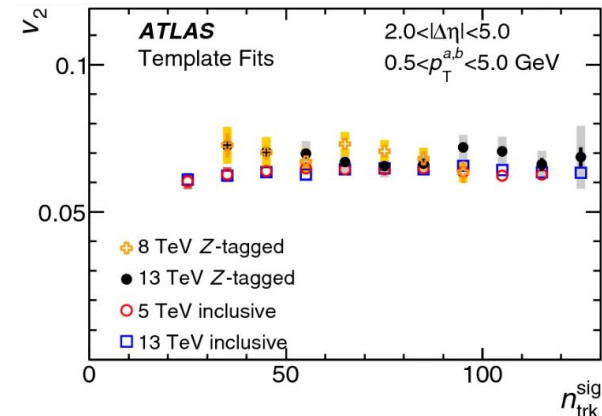
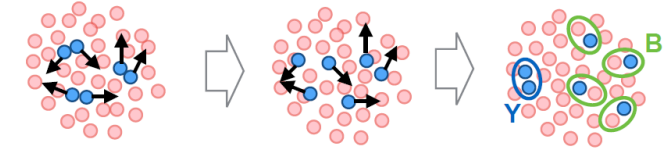
- Event multiplicity dependencies



[Color screening]



[Regeneration]



# Signal Extraction in $pp$ and Pb+Pb collisions

Upsilon mesons measured in di-muon channel at midrapidity

Y kinematics:

$p_T < 30$  GeV

$|y| < 1.5$

Centrality: 0-80%

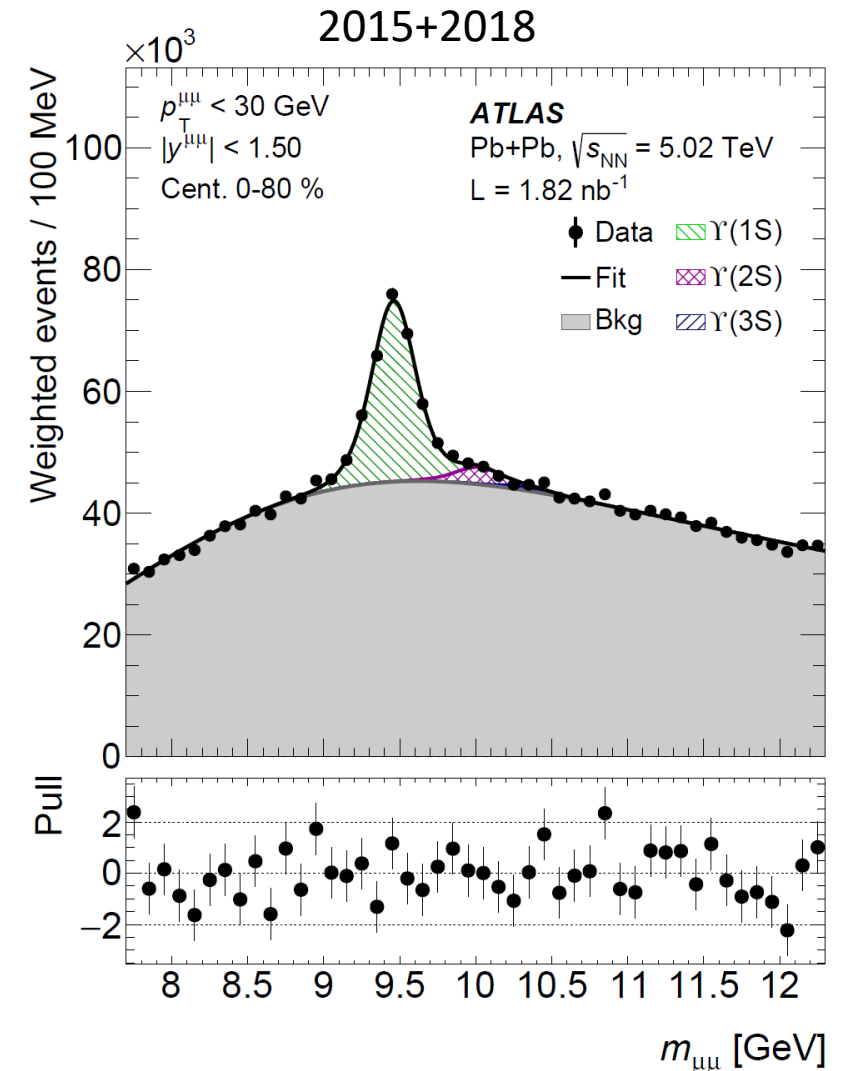
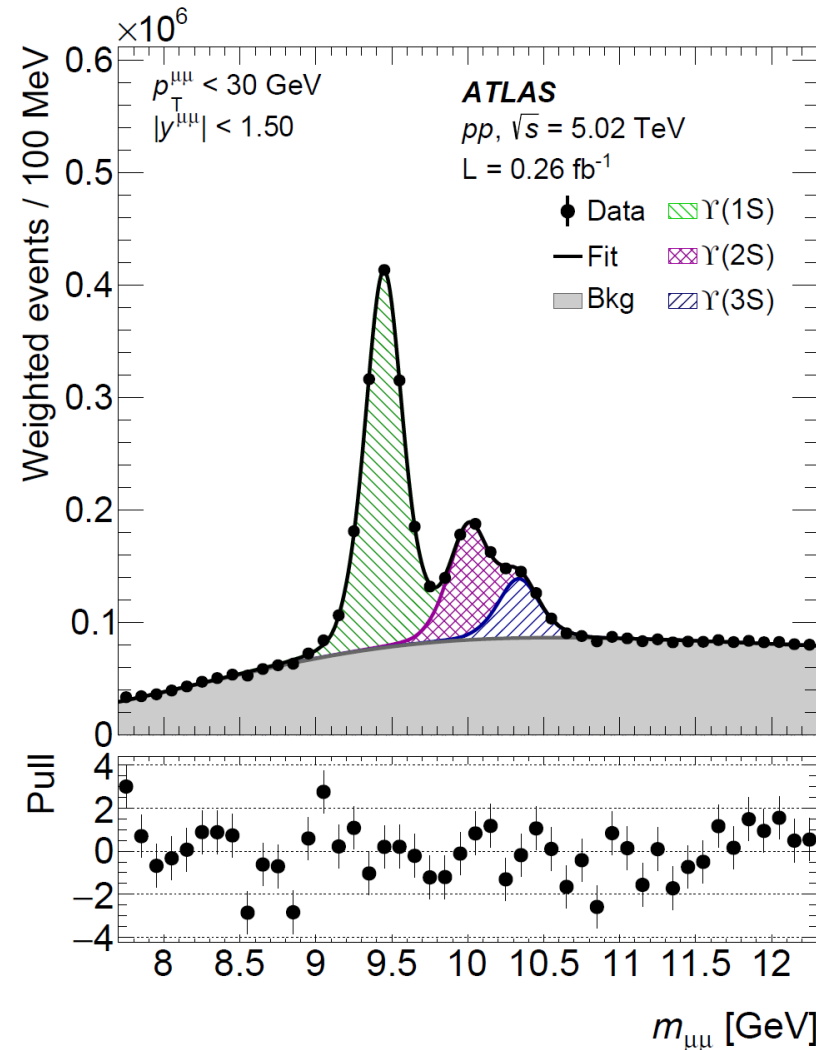
Signal:

Crystal Ball + Gauss

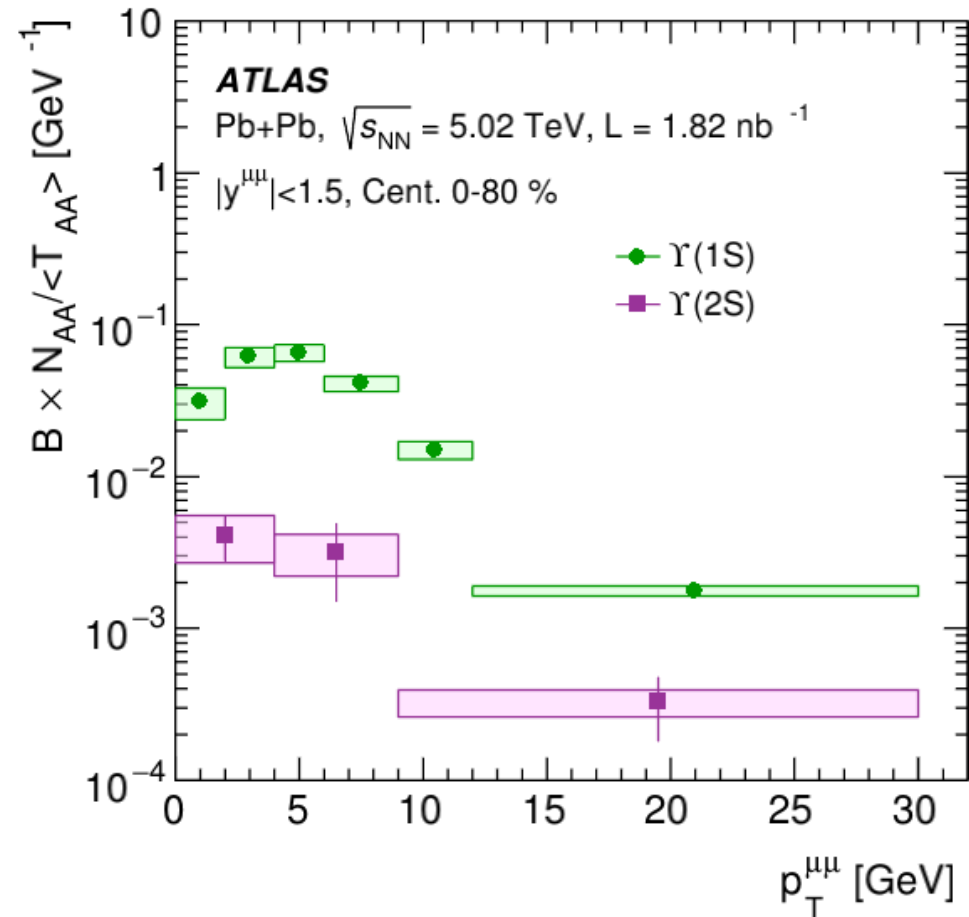
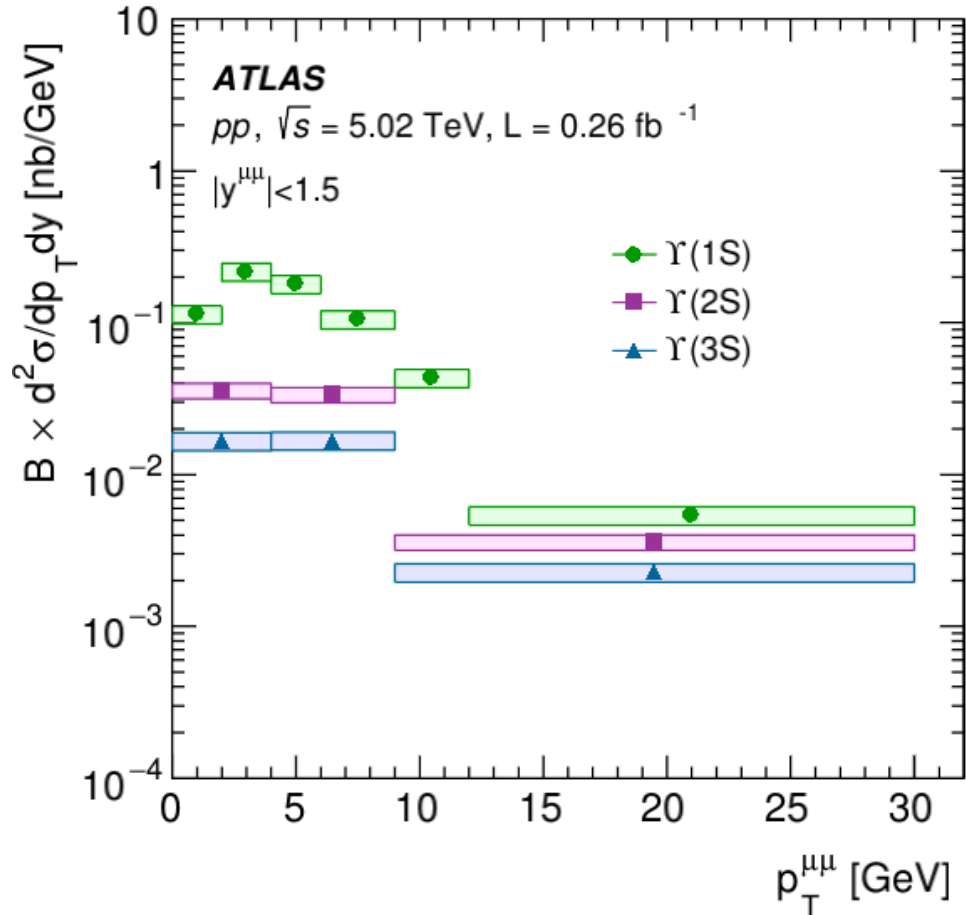
Background:

2<sup>nd</sup> order polynomial

or Erf\*Exp



# Cross-section in $pp$ and Pb+Pb



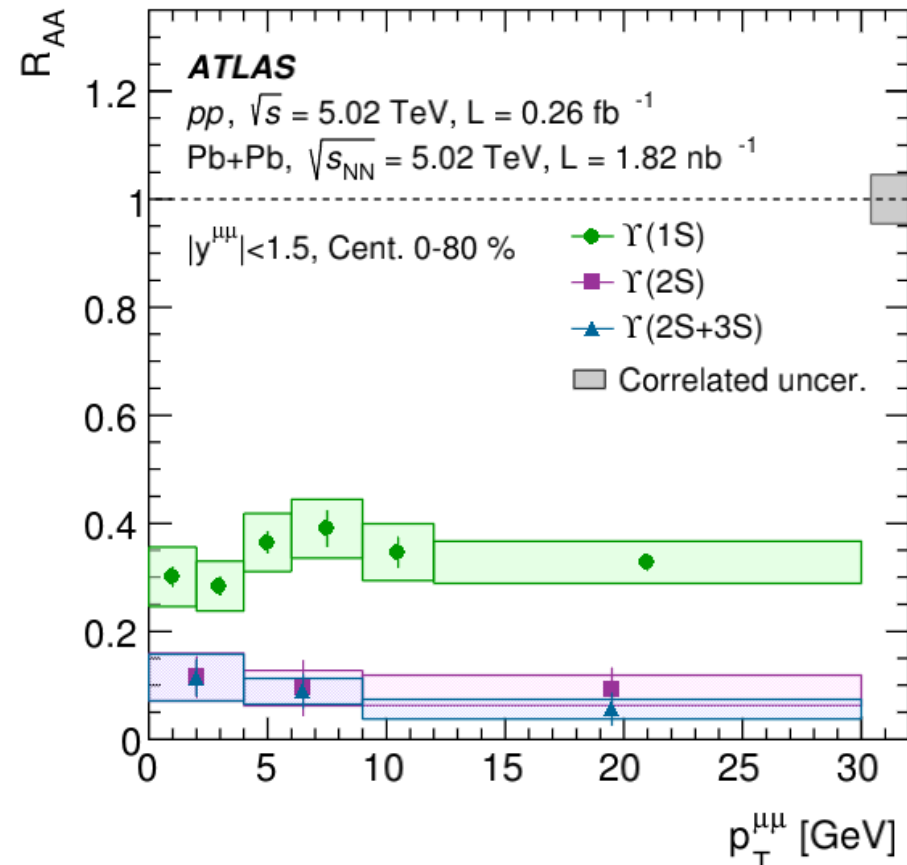
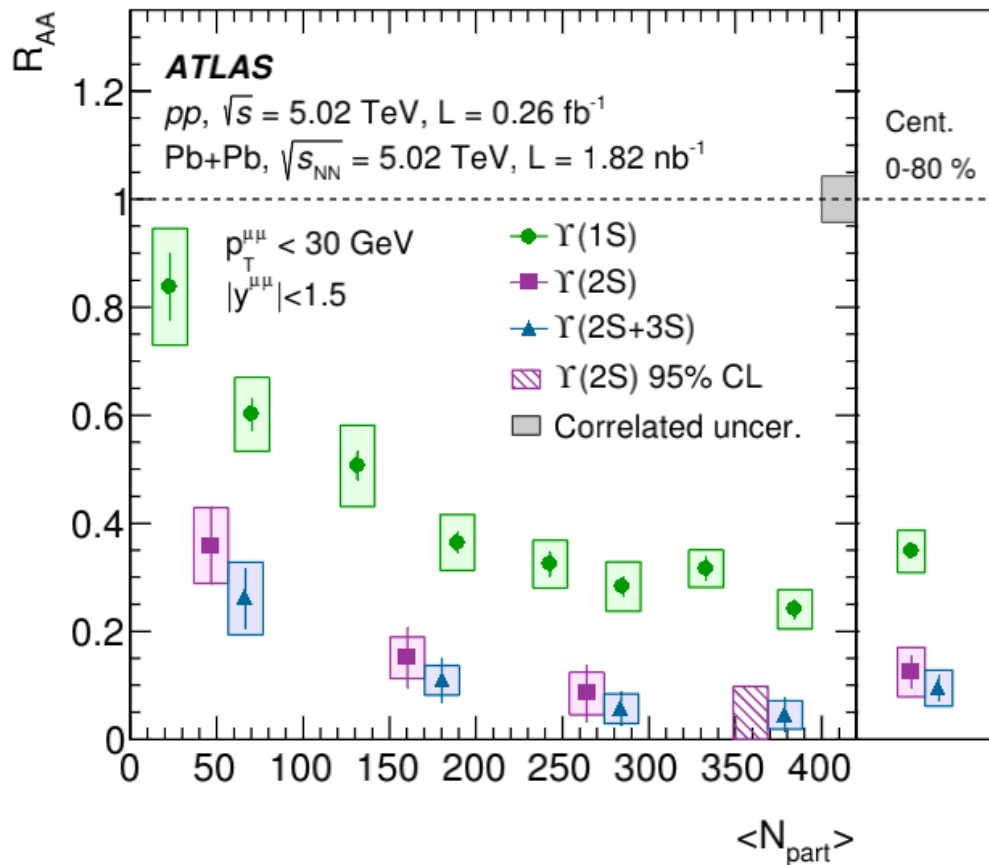
$\Upsilon(3S)$  not shown since peaks are not statistically significant.

# Nuclear modification factor $R_{AA}$

Nuclear modification factor

$$R_{AA} = \frac{N_{AA}}{\langle T_{AA} \rangle \times \sigma^{pp}}$$

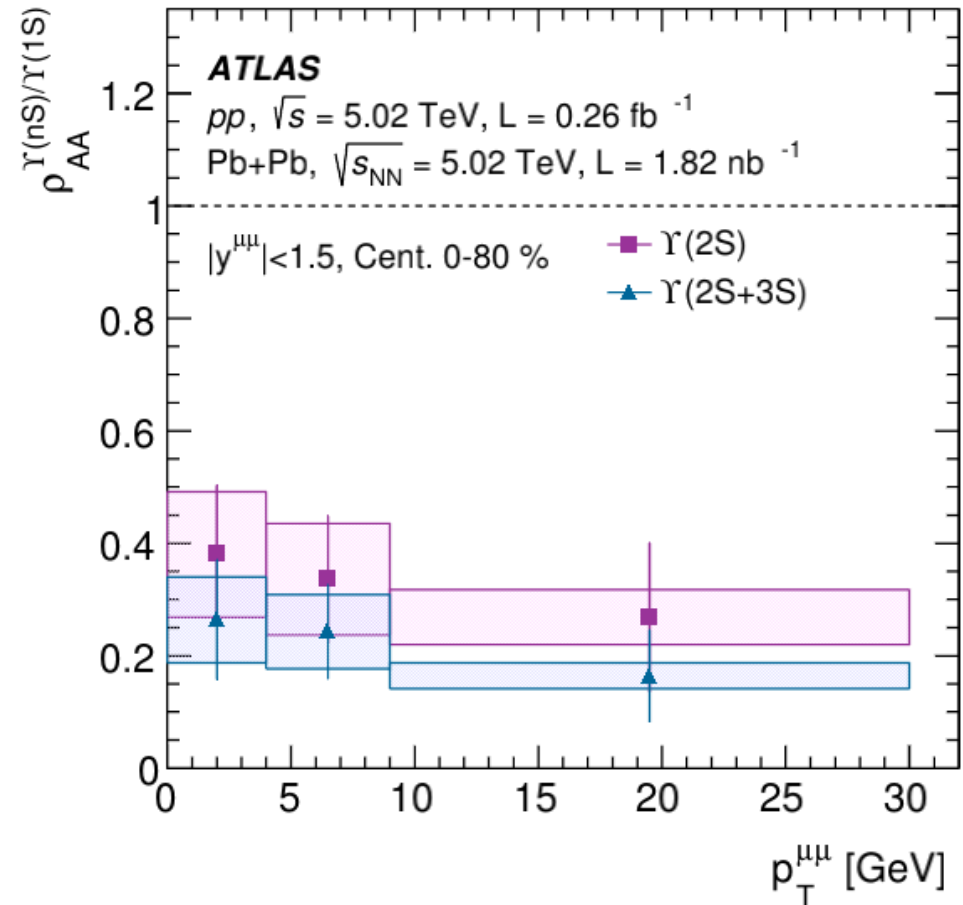
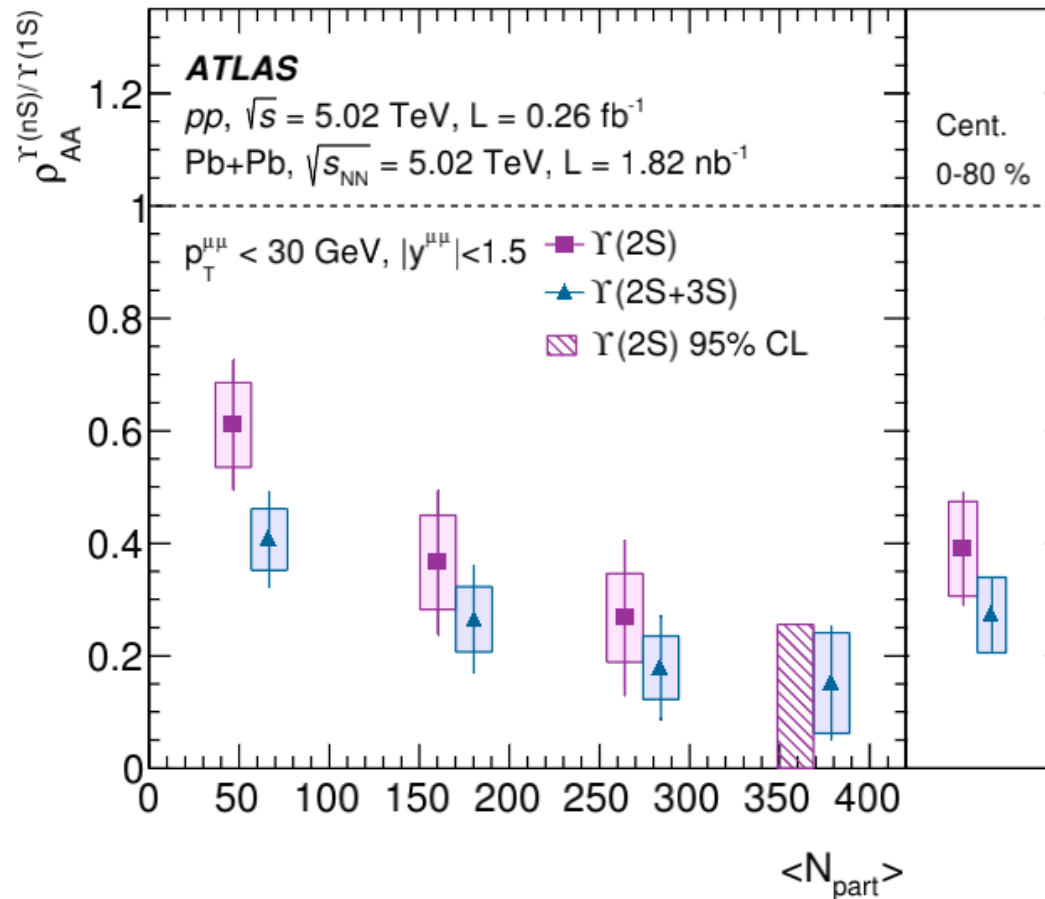
- Ordering in  $R_{AA}$ :  $Y(1S) > Y(2S) > Y(2S+3S)$
- More suppression in more central collisions
- No strong  $p_T$  dependence.
- $Y(2S+3S)$  is shown instead of  $Y(3S)$



# Double ratios

$$\rho_{AA}^{\Upsilon(nS)/\Upsilon(1S)} = \frac{\sigma_{AA}^{\Upsilon(nS)} / \sigma_{AA}^{\Upsilon(1S)}}{\sigma_{pp}^{\Upsilon(nS)} / \sigma_{pp}^{\Upsilon(1S)}} = \frac{R_{AA}(\Upsilon(nS))}{R_{AA}(\Upsilon(1S))}$$

- Luminosity and  $\langle T_{AA} \rangle$  corrections cancel
- Acceptance and efficiency corrections partially cancel
- Consistent with sequential melting
- $\Upsilon(2S+3S)$  systematically lower than  $\Upsilon(2S)$



# Comparison with models ( $R_{AA}$ )

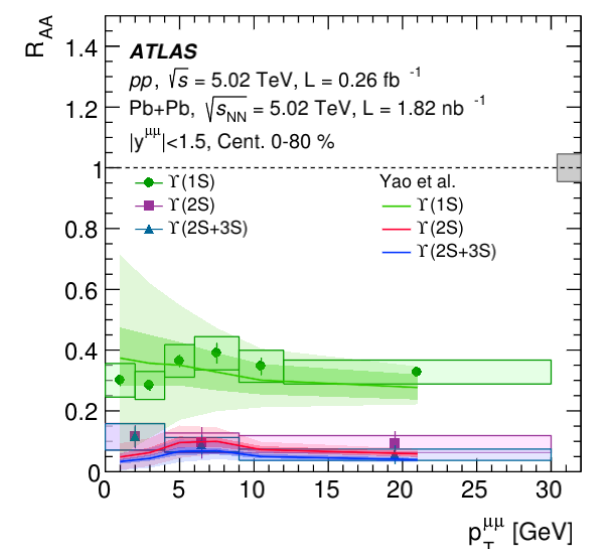
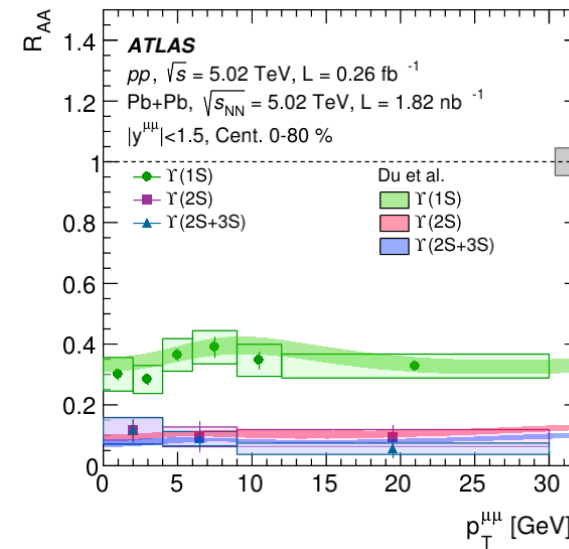
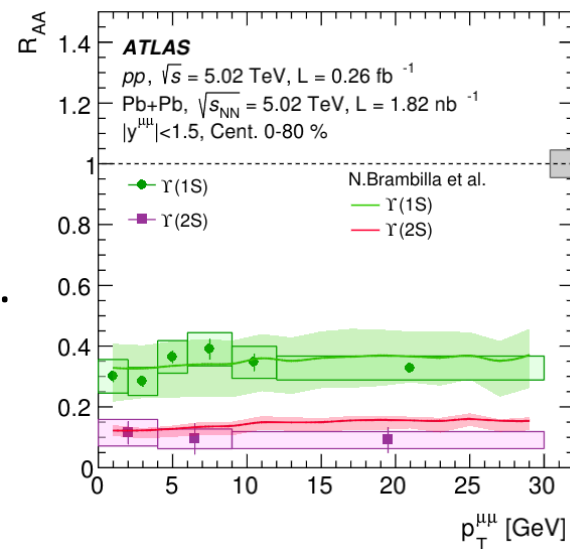
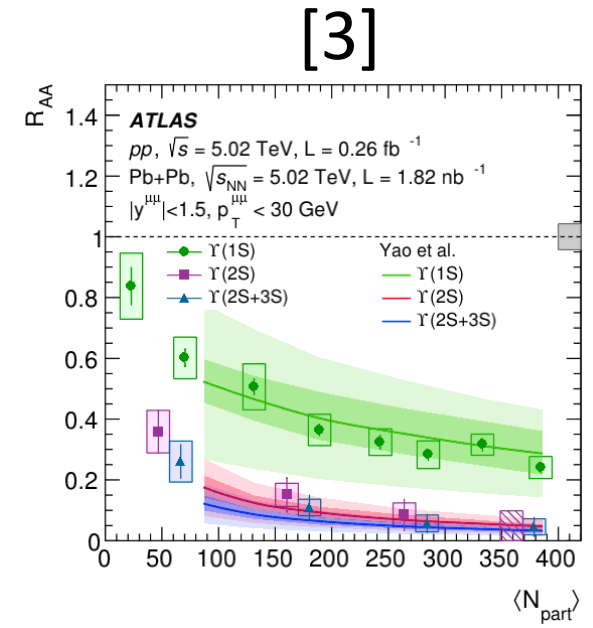
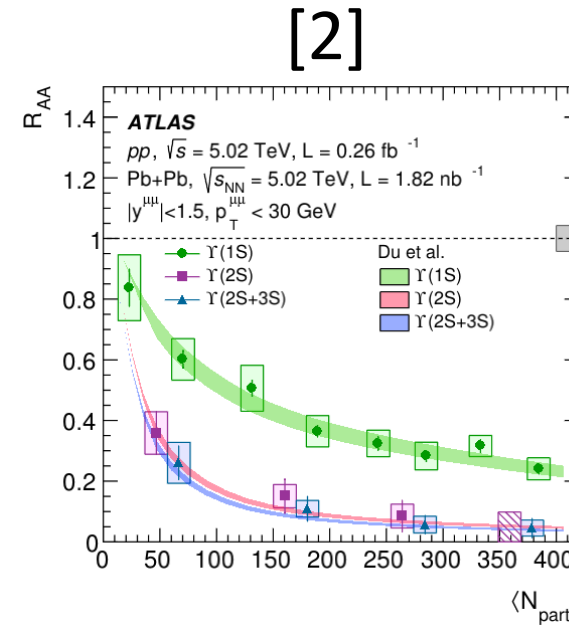
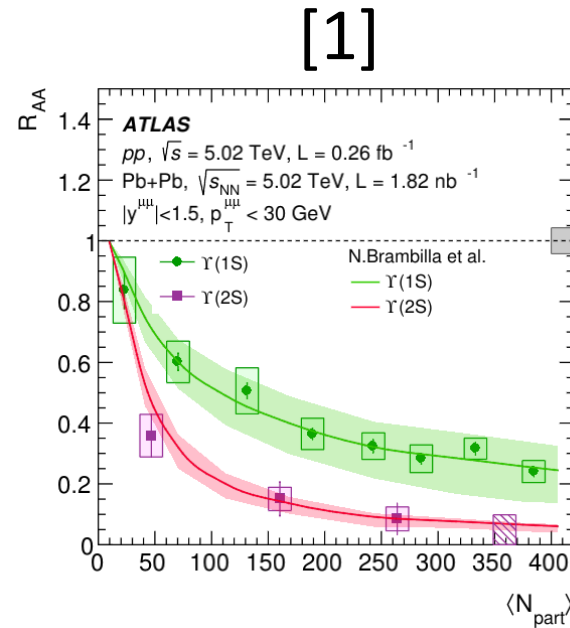
[1] N.Brambilla et al.,  
Phys. Rev. D 104 (2021) 094049

[2] M. H. X.Du and R. Rapp,  
Phys. Rev. C 96 (2017) 054901

[3] X. Yao et al.,  
JHEP 2021 (2021) 46

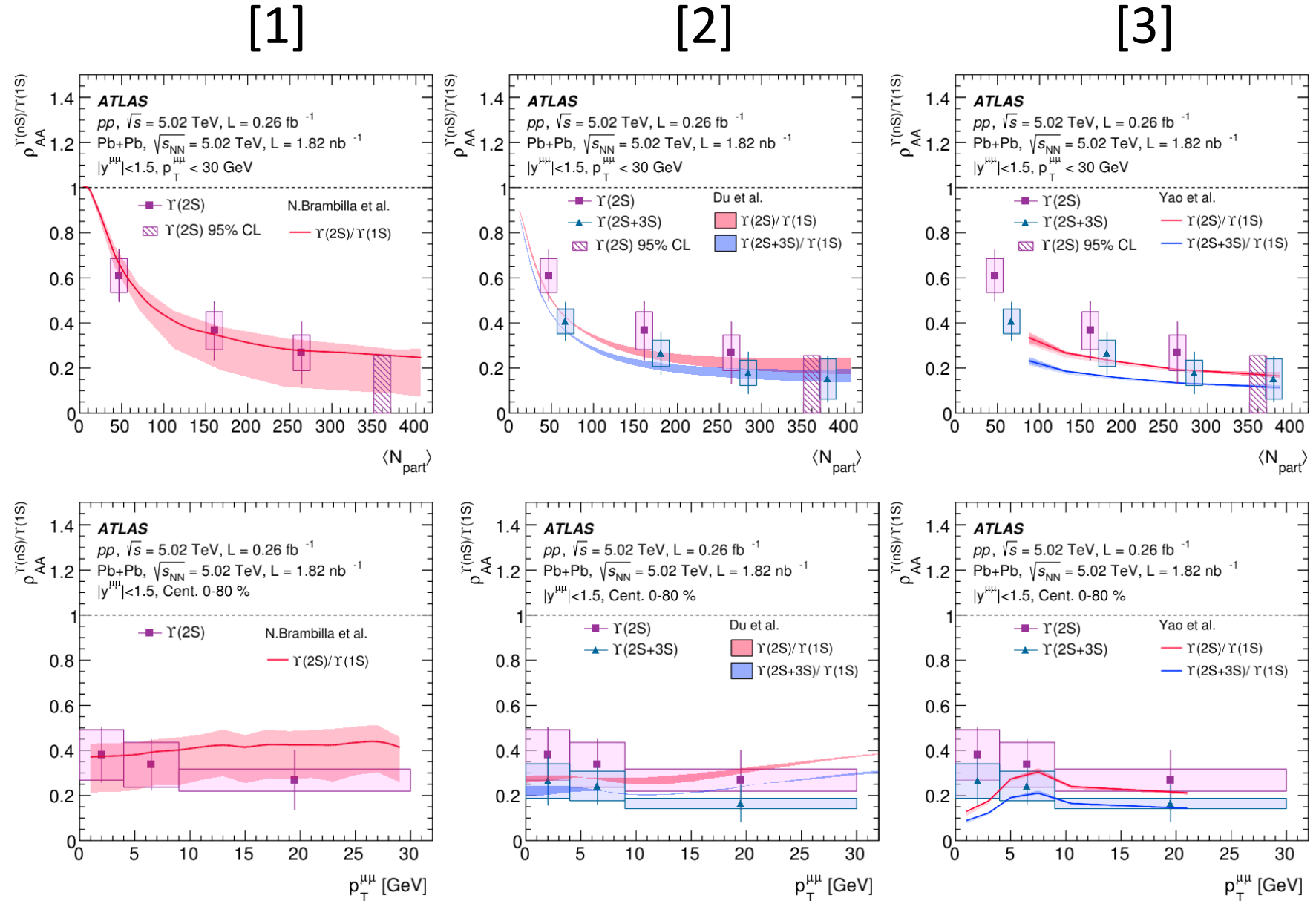
- Models use different approach to Y suppression, but include deconfinement as key ingredient.

- Good agreement with the data. Previous Y suppression data available to authors.



# Comparison with models (double ratios)

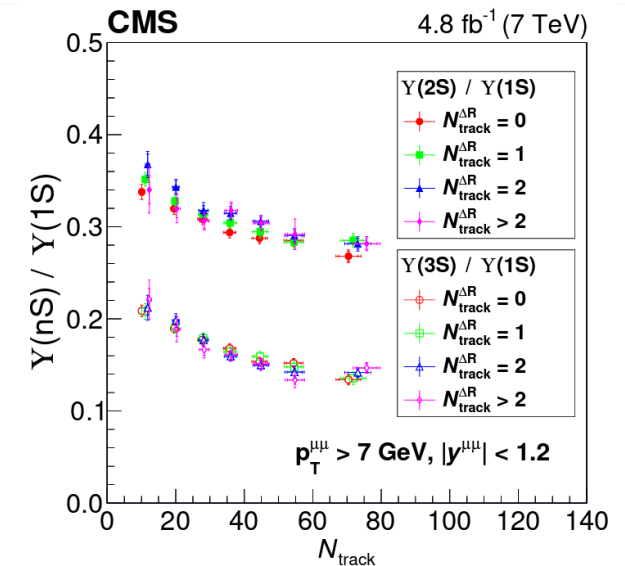
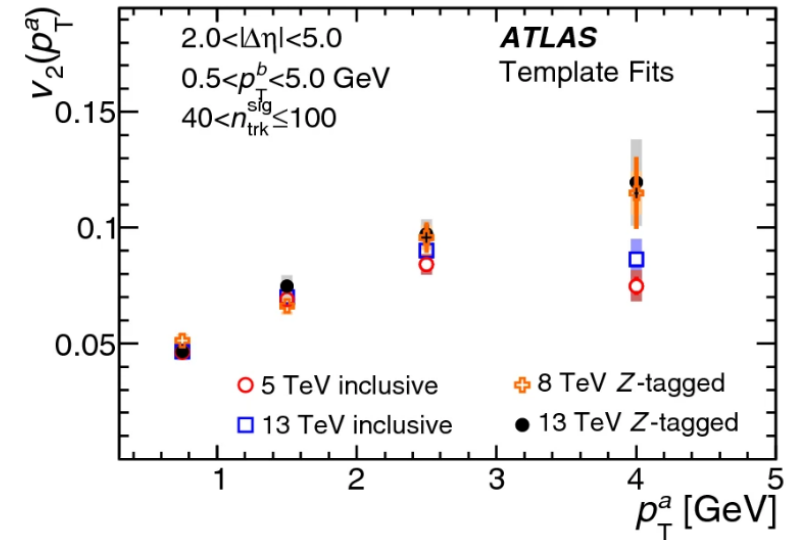
- Good agreement with the data.
- $Y(2S+3S)$  suppression relative to  $Y(2S)$  in models consistent with data.
- Many model uncertainties cancel in double ratio.





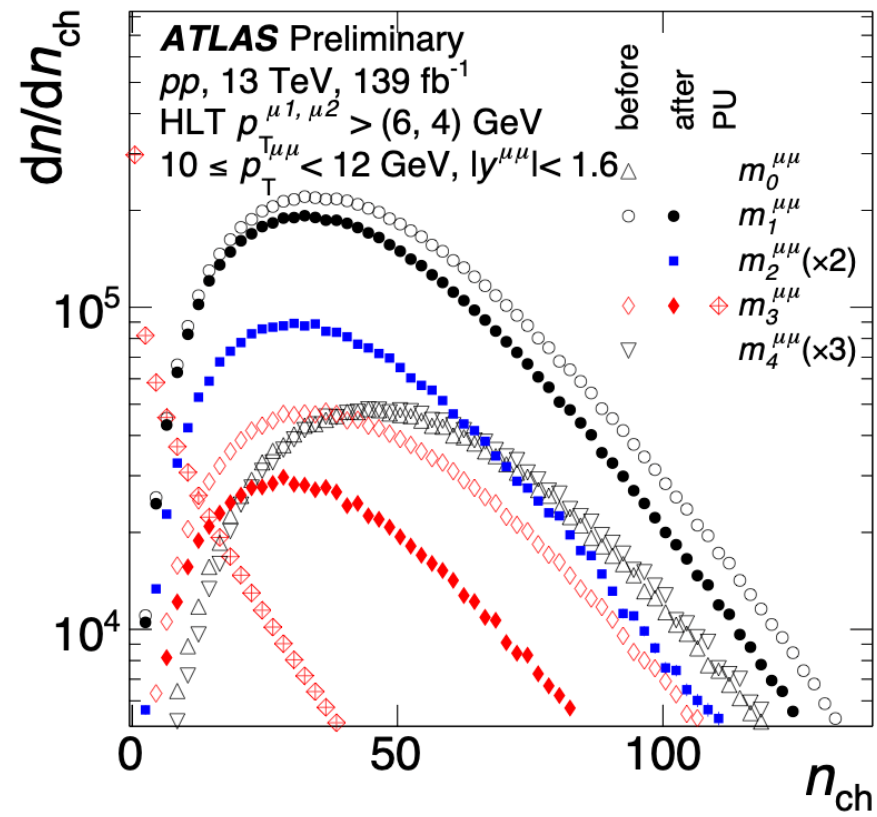
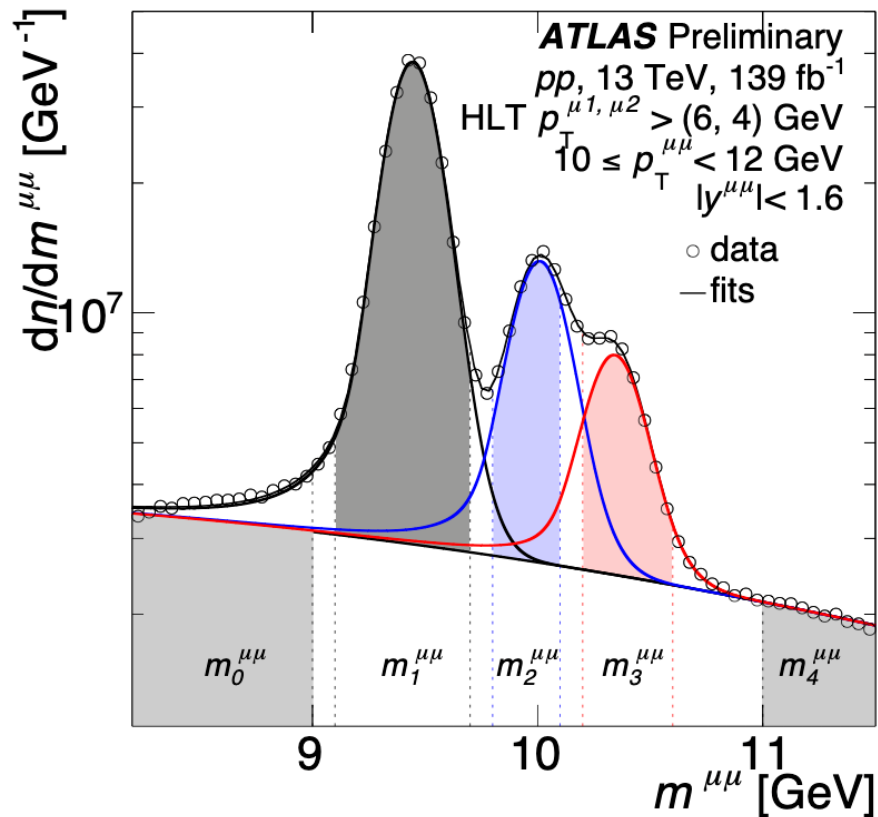
# Correlation of $Y$ meson production with the underlying event in pp collisions

- Heavy ion features have been observed in pp
  - *Flow in Z tagged pp events by ATLAS*  
[*Eur. Phys. J* **C80** (2020) 64],
  - *Y double ratios decreasing with multiplicity pp events by CMS* [*JHEP* **2020** (2020) 001]
- We search for QGP-like signatures by looking at how underlying event changes in events with Upsilon
  - *Measure charged particle multiplicity,  $p_T$  and  $\Delta\phi$  distributions in events with different Y states.*
- For details see poster by Iakov Aizenberg



# Analysis details

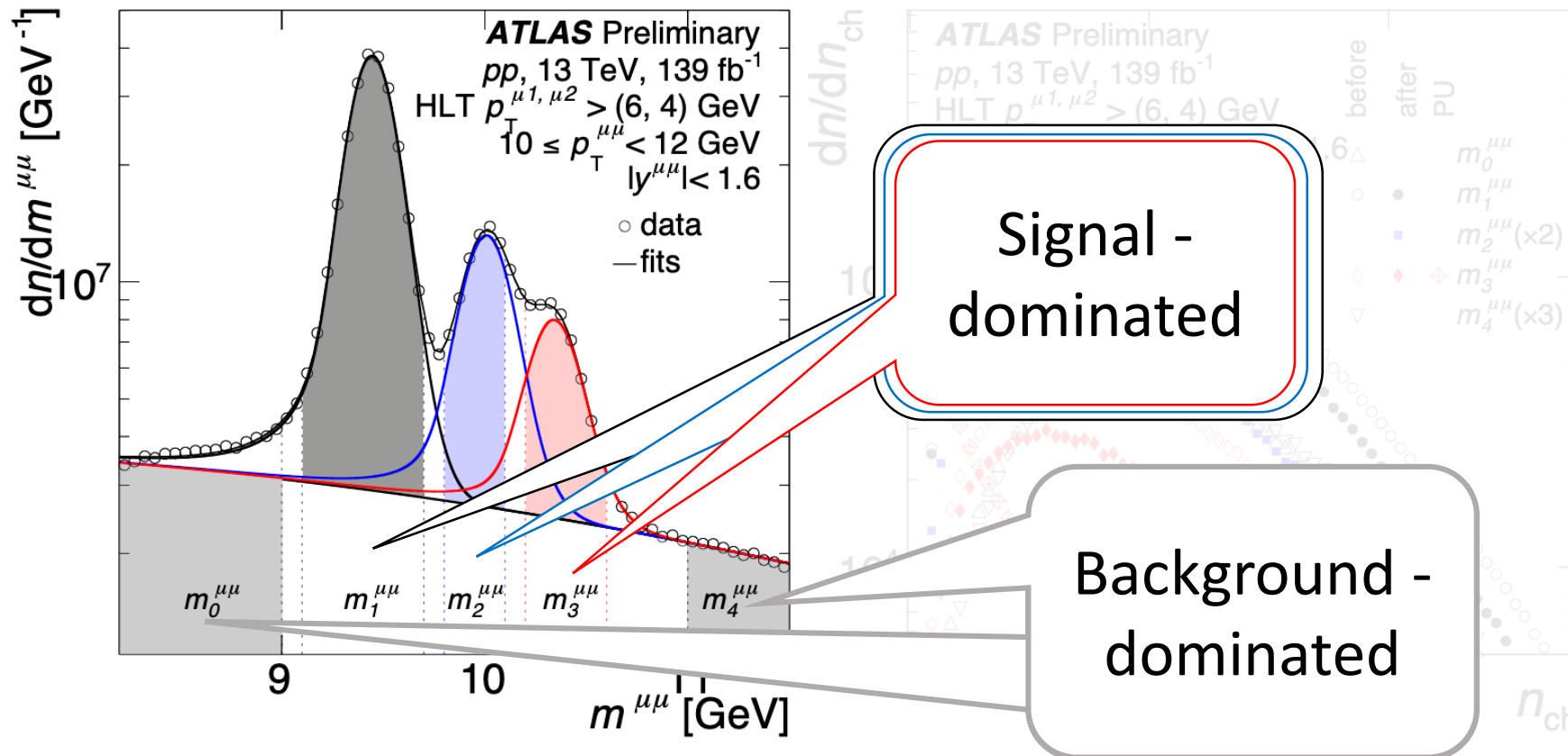
- Entire Run-2  $pp$  13 TeV collisions,  $139 \text{ fb}^{-1}$
- Events with  $Y(nS) \rightarrow \mu\mu$ ,  $|y| < 1.6$ , all triggers, all  $p_T$
- Charged hadron  $n_{\text{ch}}$ ,  $dn_{\text{ch}}/dp_T$  and  $dn_{\text{ch}}/d\Delta\phi$ .  $\Delta\phi = \phi^Y - \phi^h$
- Fully corrected, including pileup



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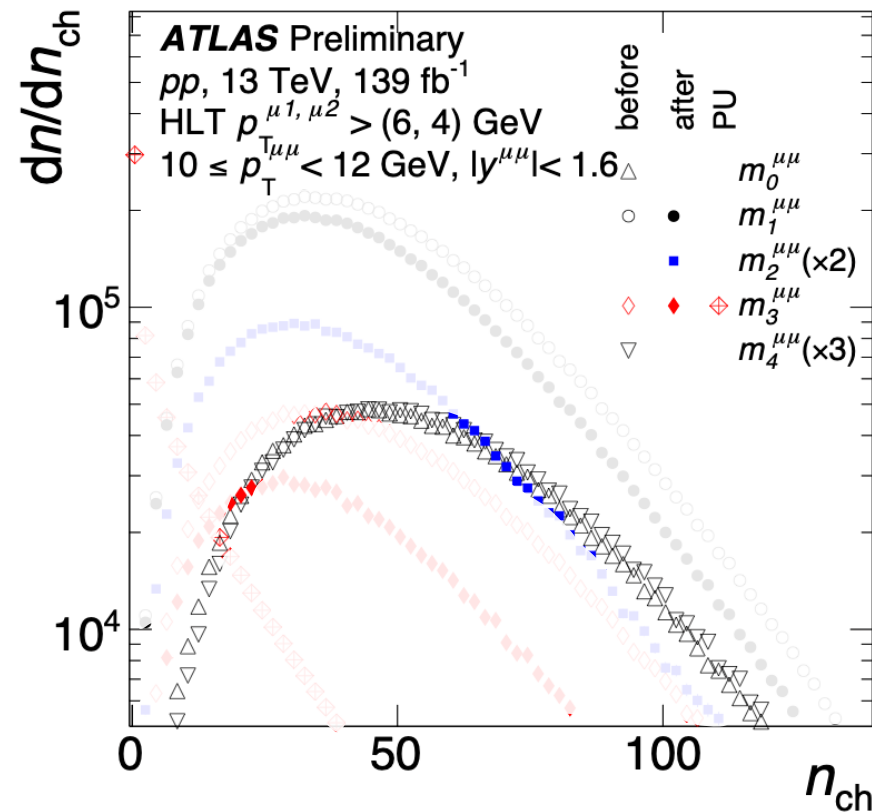
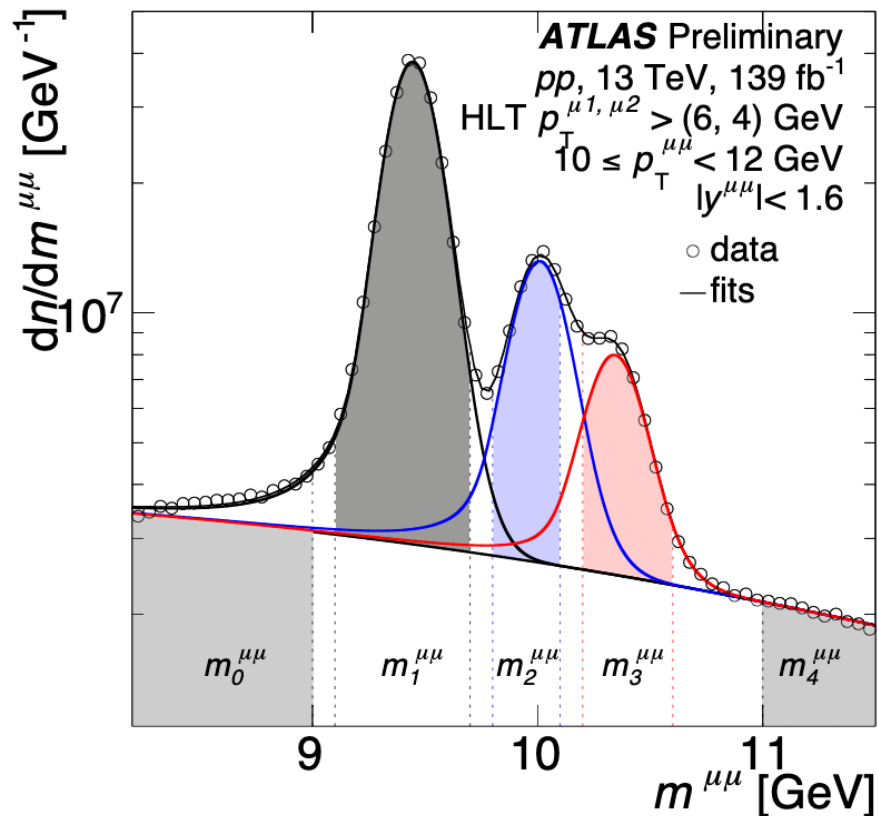
- Select 5 mass regions



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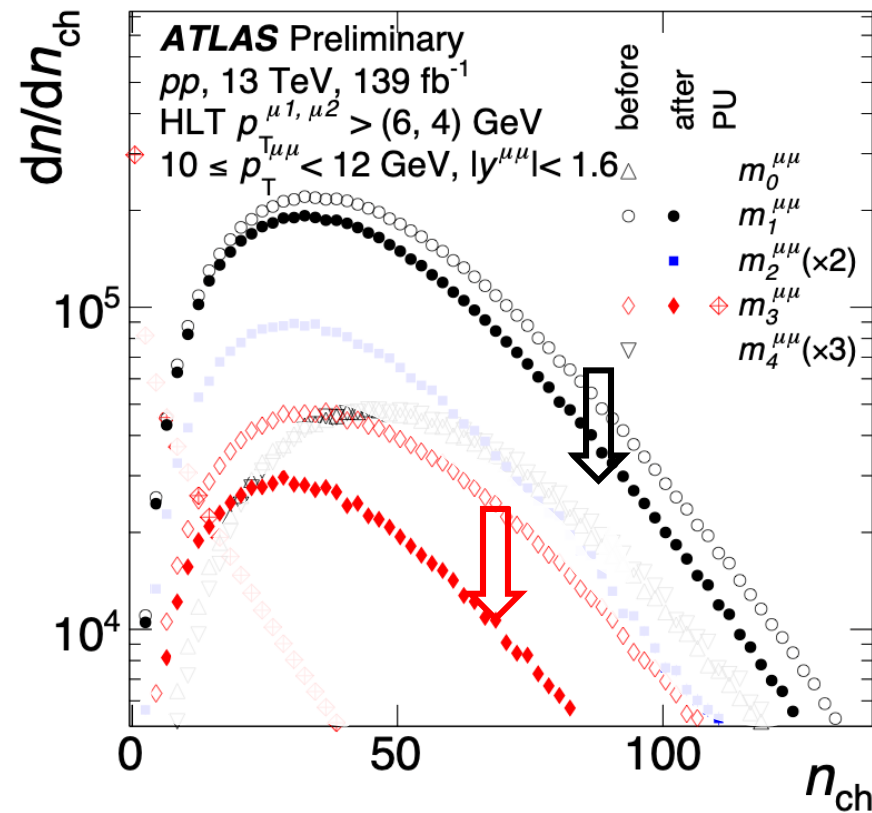
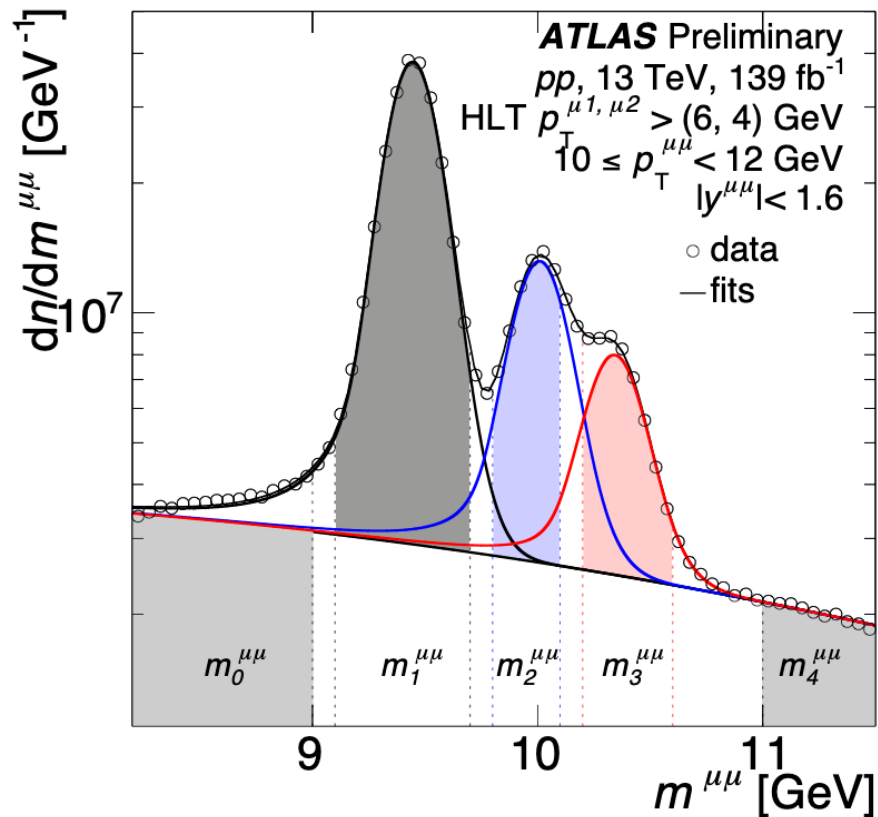
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- Bkg shapes are similar – interpolate



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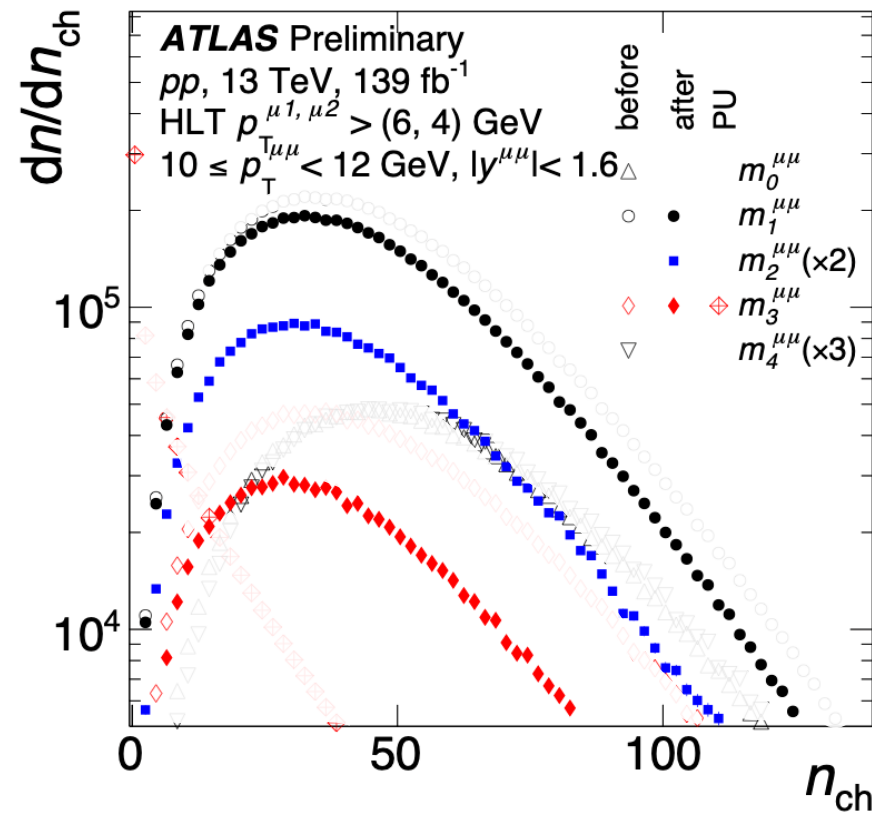
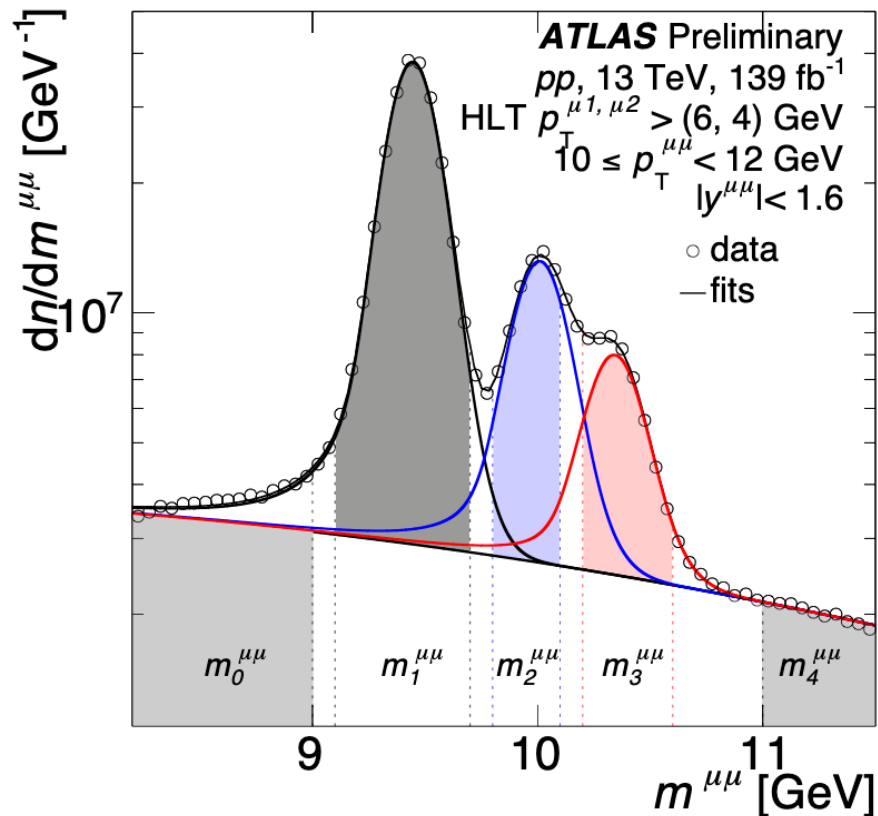
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- Define 3+2 regions
- Bkg shapes are similar – interpolate
- Bkg subtraction for  $Y(1S)$  and  $Y(3S)$

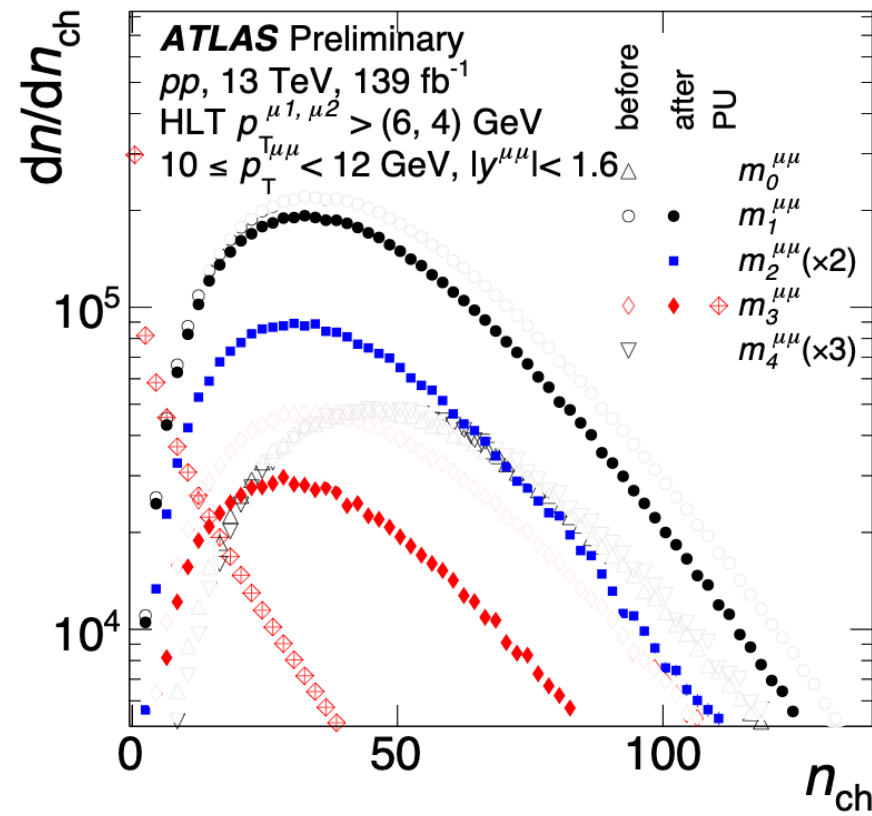
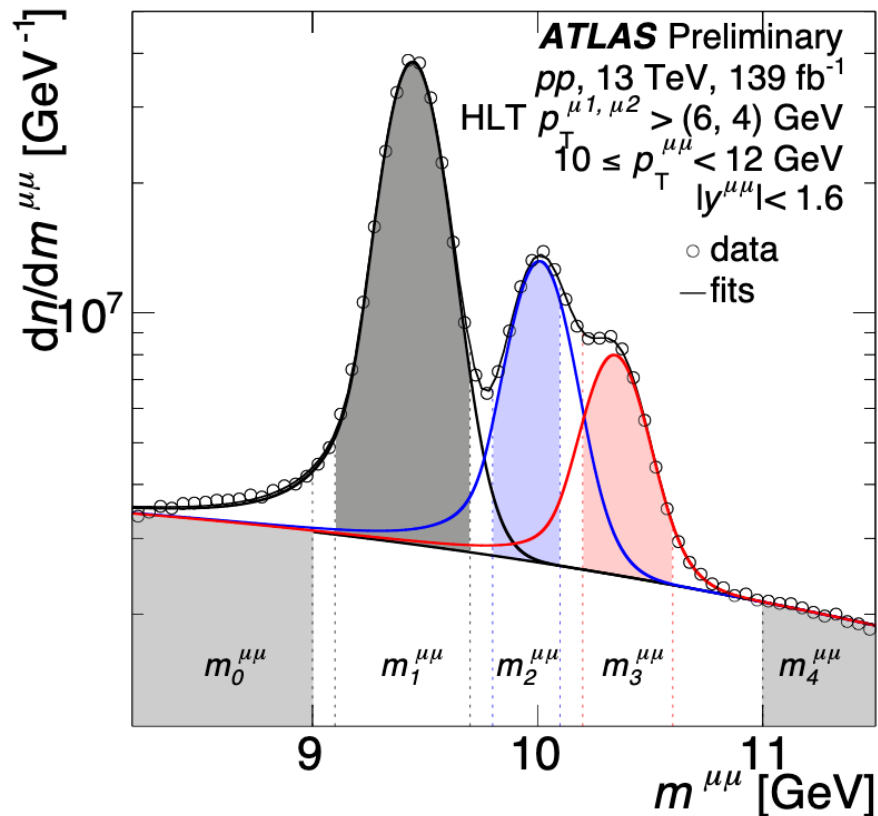


- After subtraction  $n_{\text{ch}}$  look different

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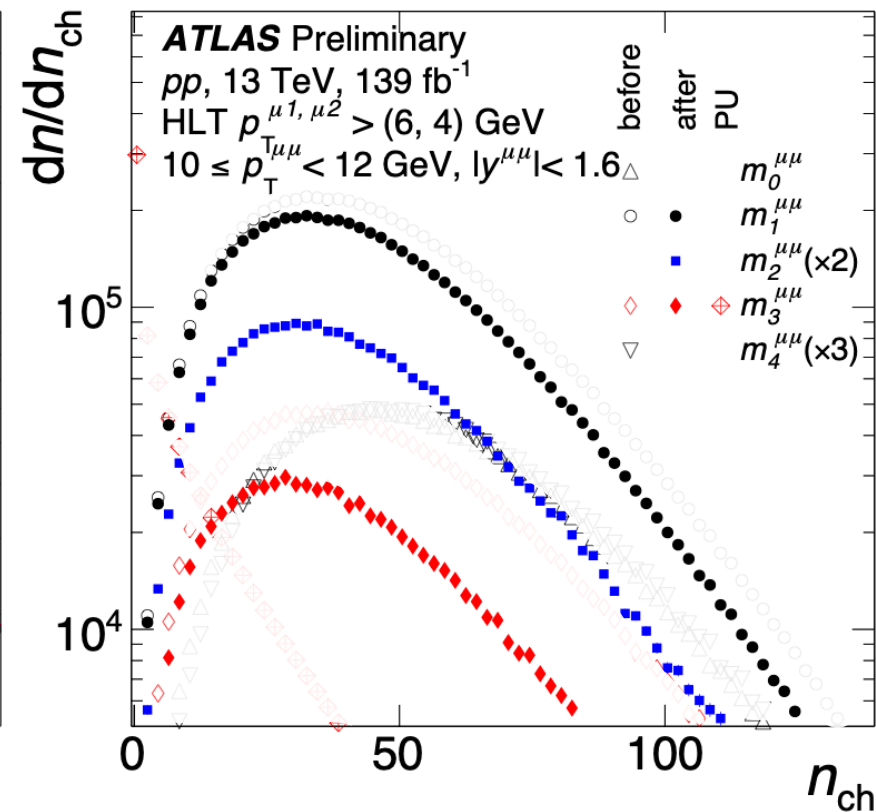
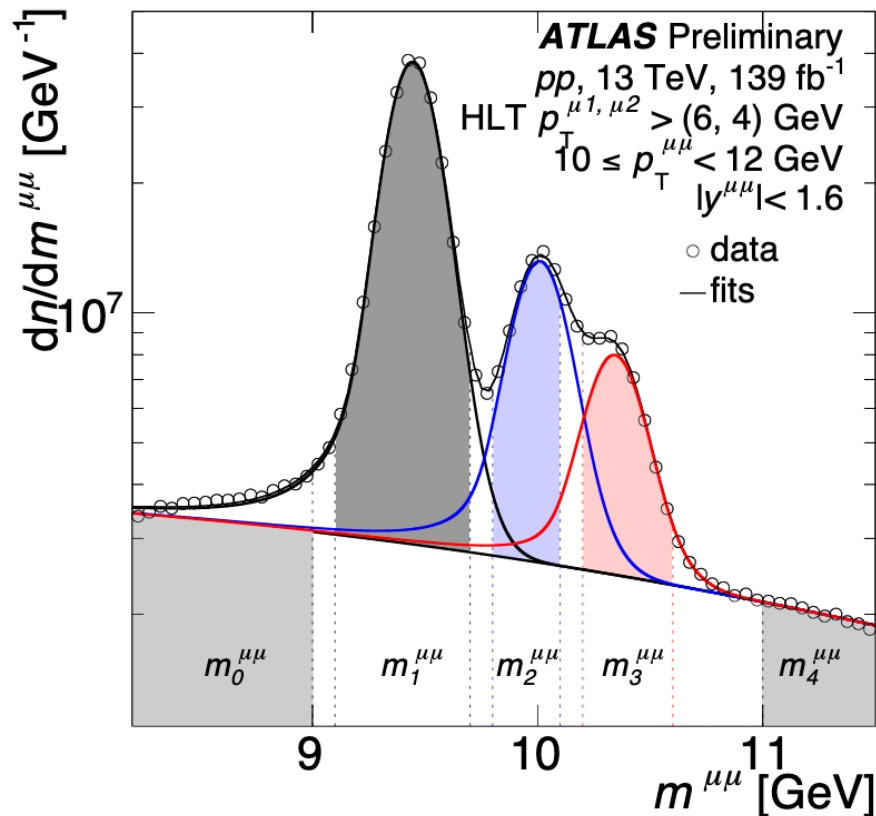
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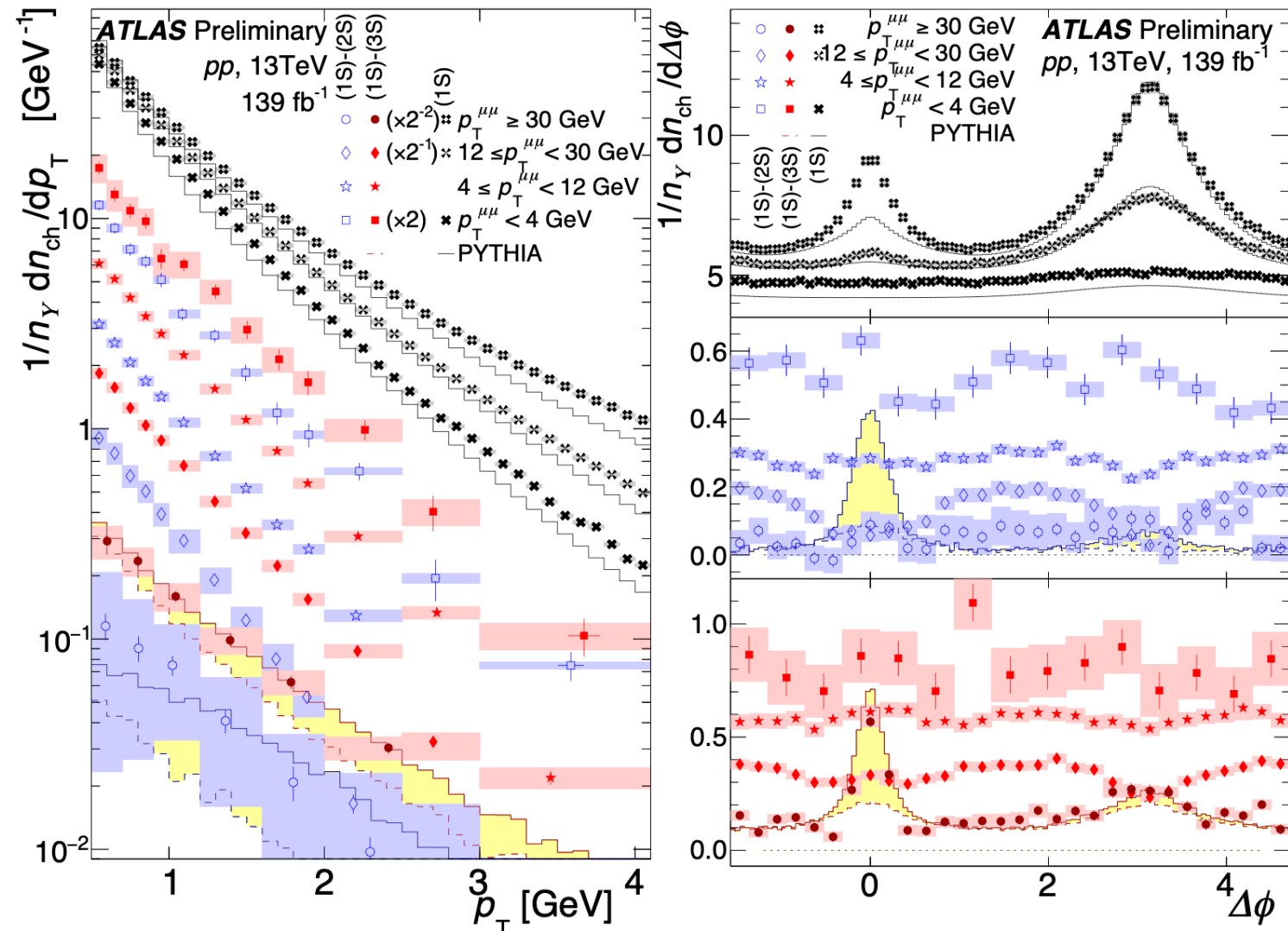
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Direct measurement of  $n_{\text{ch}}$   $dn_{\text{ch}}/dp_T$   $dn_{\text{ch}}/d\Delta\phi$

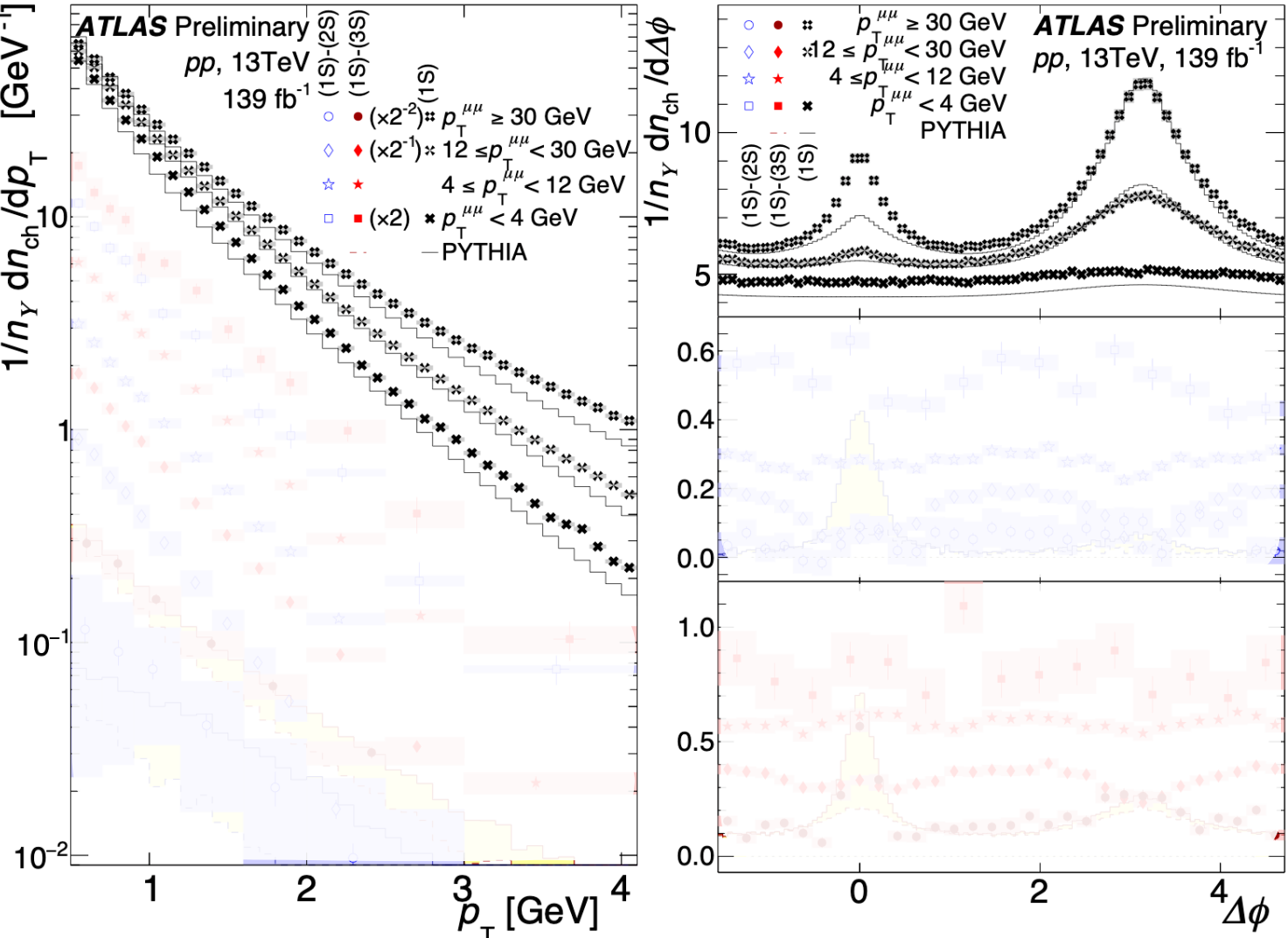


# Kinematic distributions



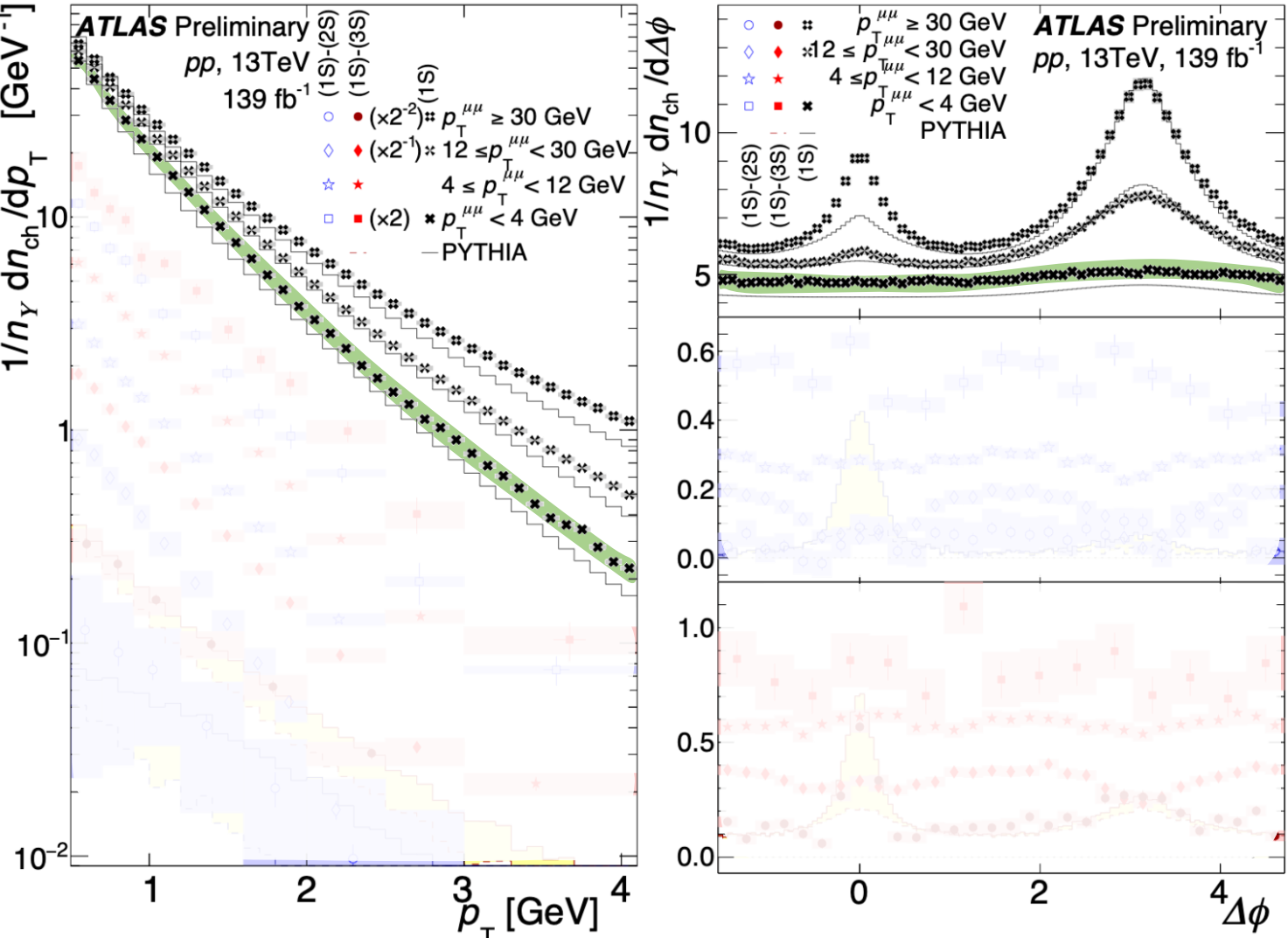
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- Distributions for Y(1S)
- Pythia does not describe well



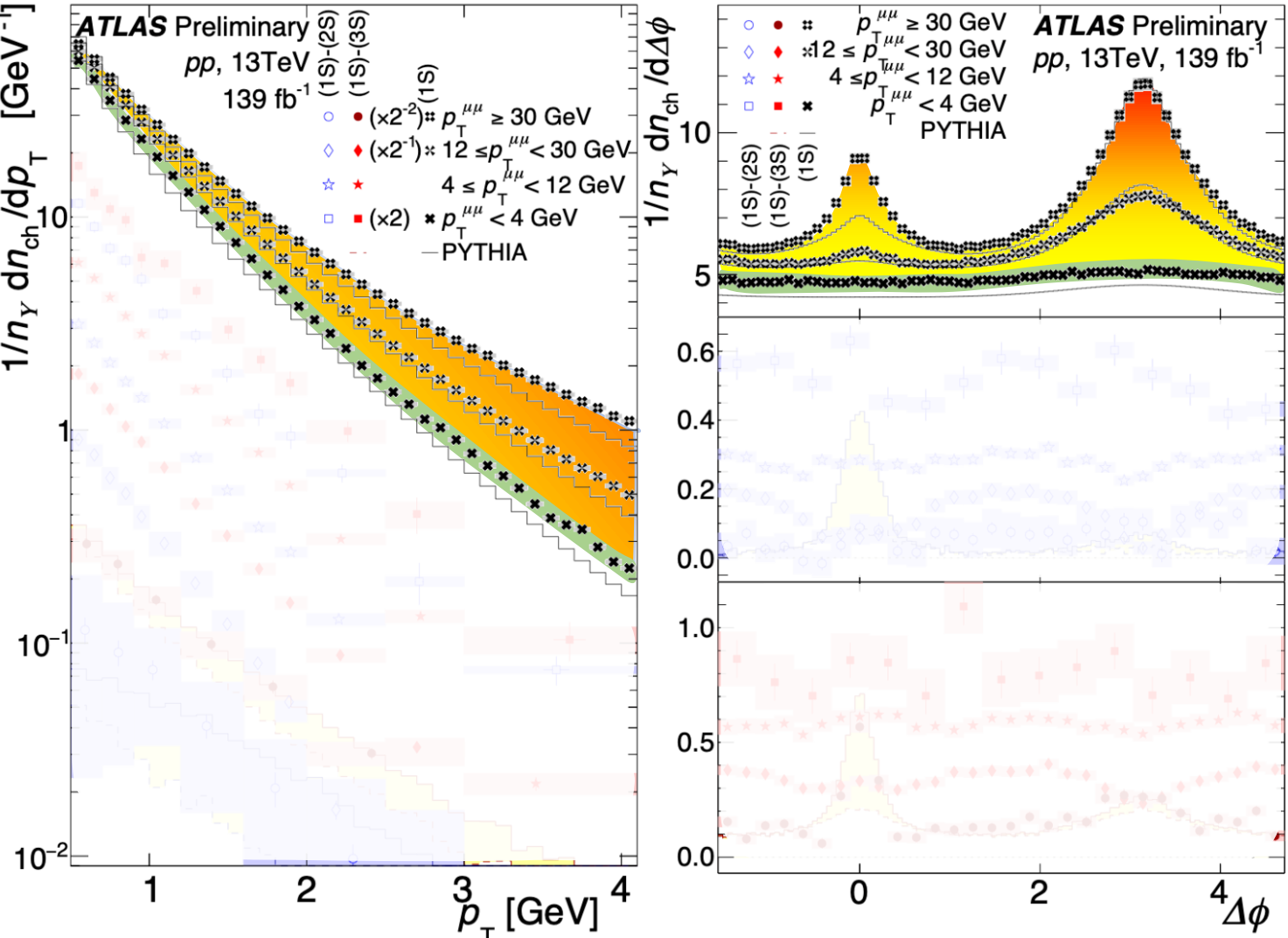
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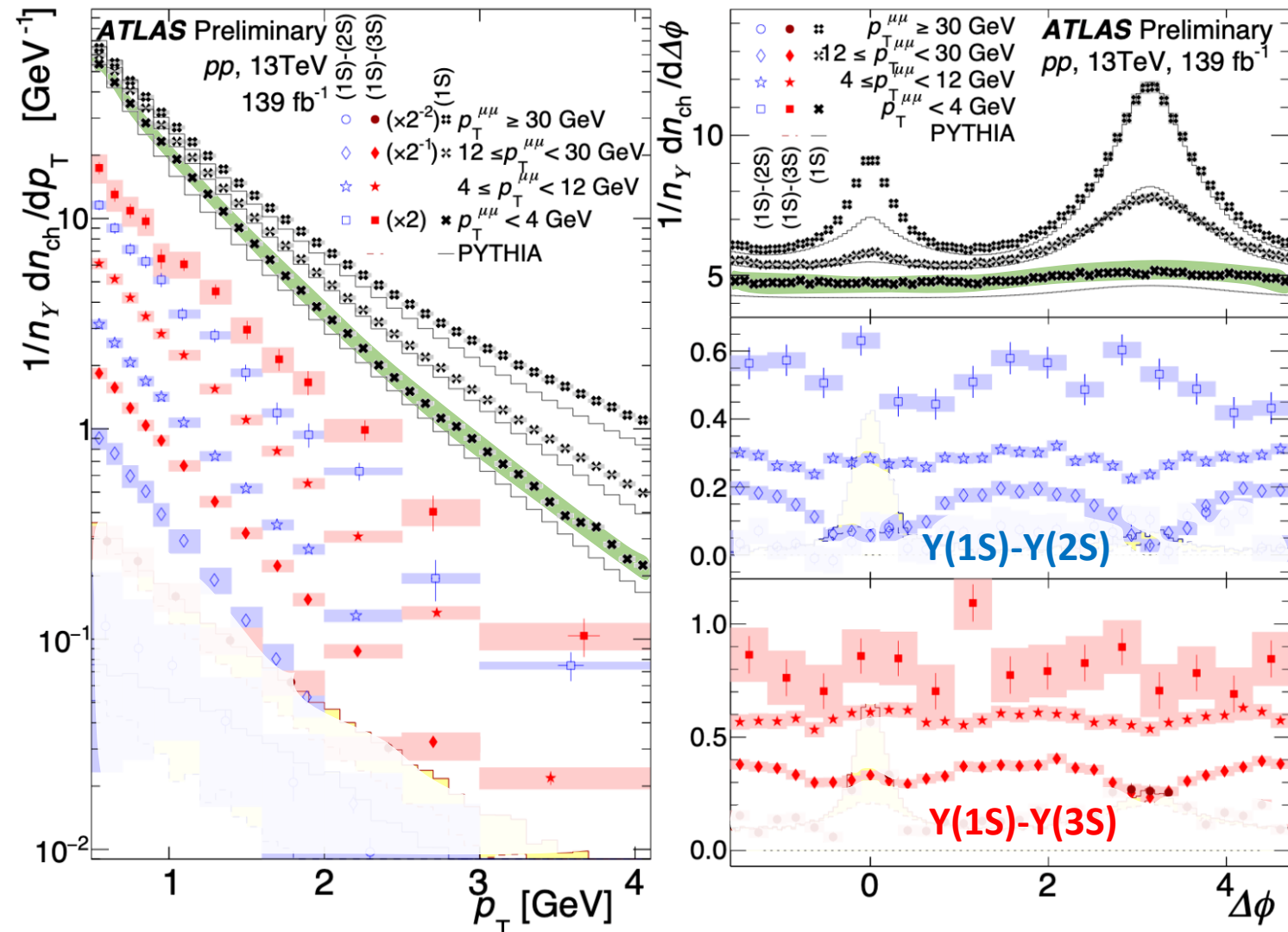


# Kinematic distributions

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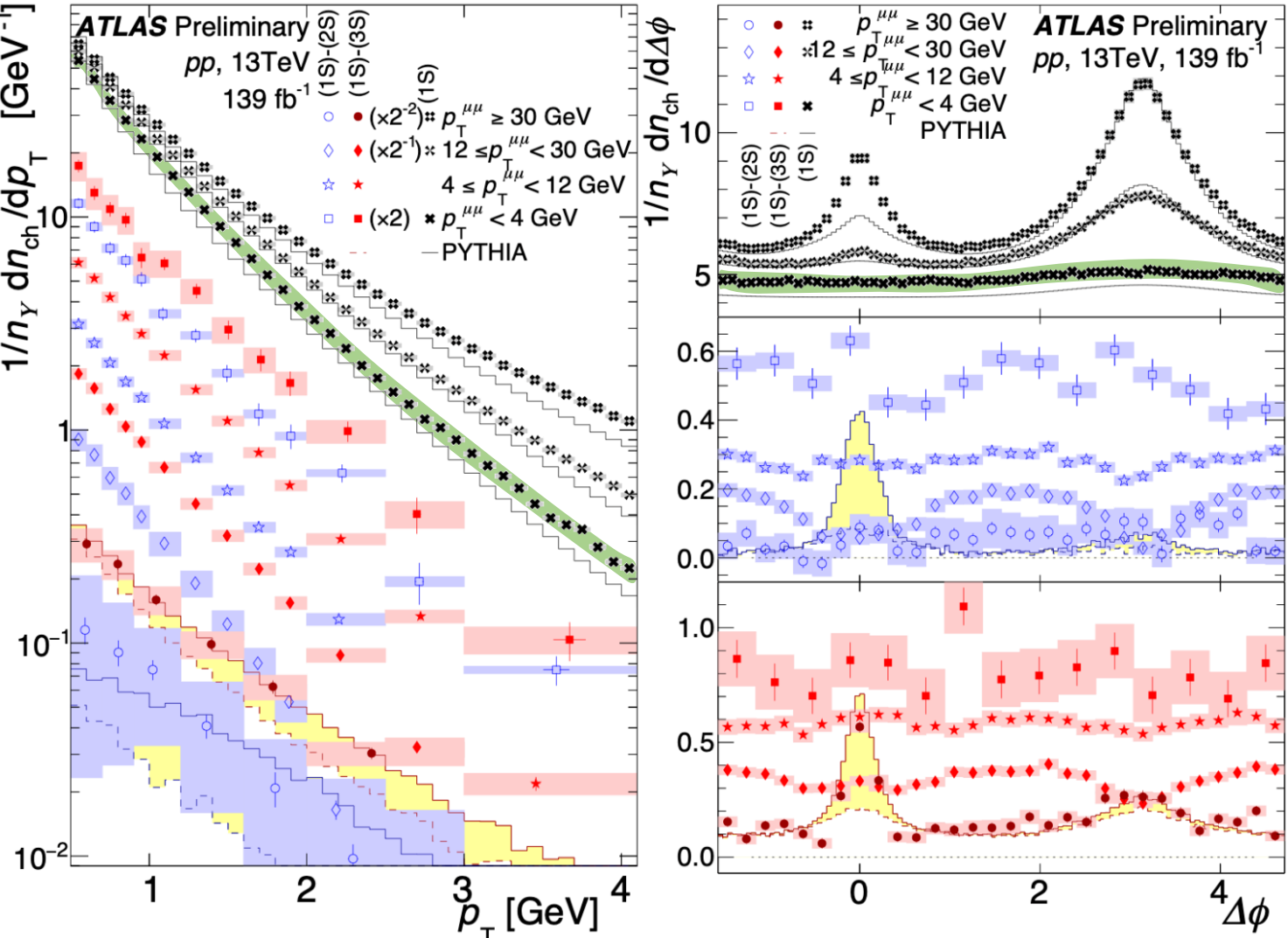
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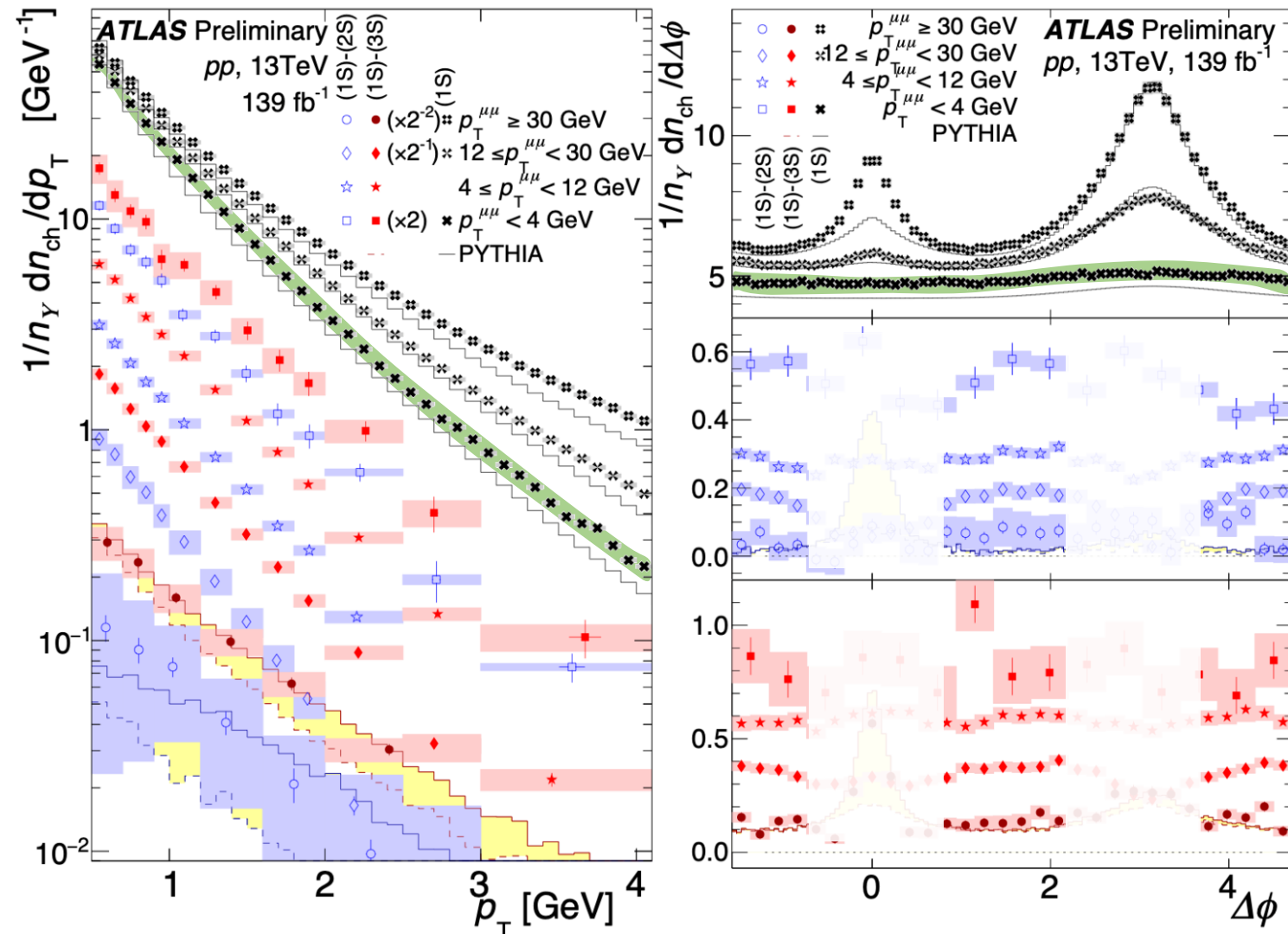
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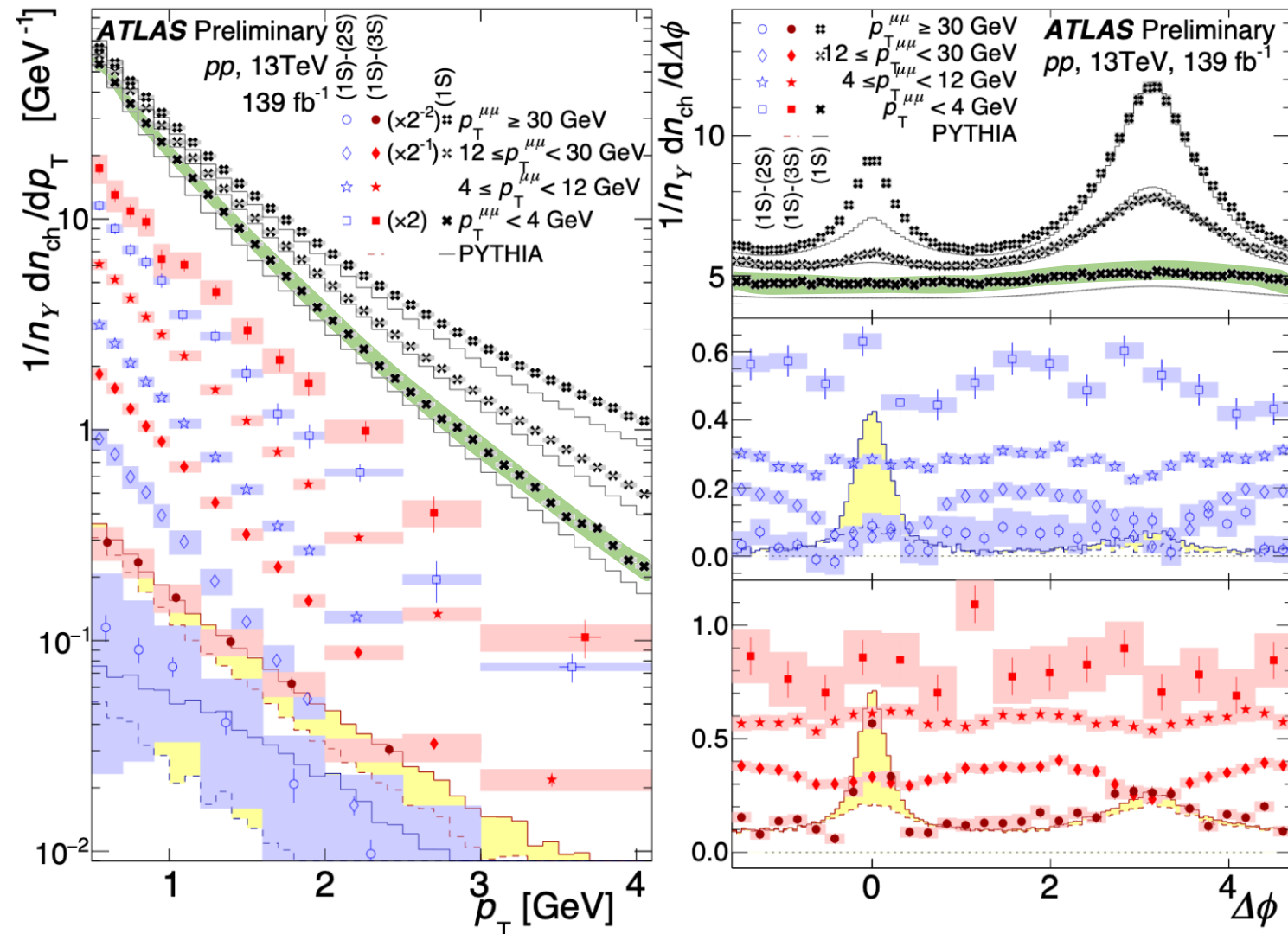


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- Away from jets there are regions with charged particles.

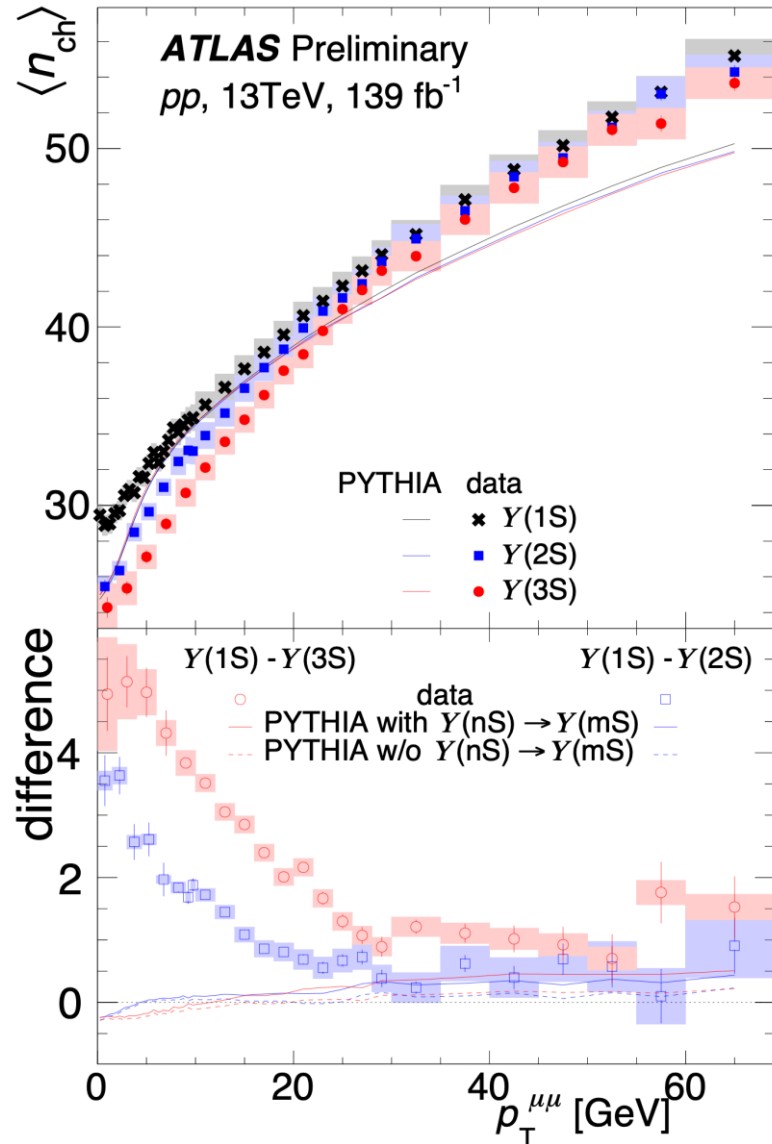
# Kinematic distributions



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- Subtracted distributions look like UE at rather high  $Y(nS) p_T$ . At the highest  $p_T$  there are feed-downs
- Away from jets there are regions with charged particles
- This suggests that the effect is related to the UE.



# Mean $n_{ch}$ values associated with different $Y$ states



- Strong difference in the multiplicity of the UE for different  $Y(nS)$  states is observed.
- It does not exist in pythia and pythia needs to be improved to better model  $Y$  production
- The effect is strongest at  $p_T = 0$  and diminishes with increasing  $p_T$ , but still visible at 20-30 GeV.
- At the lowest  $p_T$ ,  
 for  $Y(1S)-Y(2S)$  the difference is  $3.6 \pm 0.4$   
 for  $Y(1S)-Y(3S)$  the difference is  $4.9 \pm 1.1$

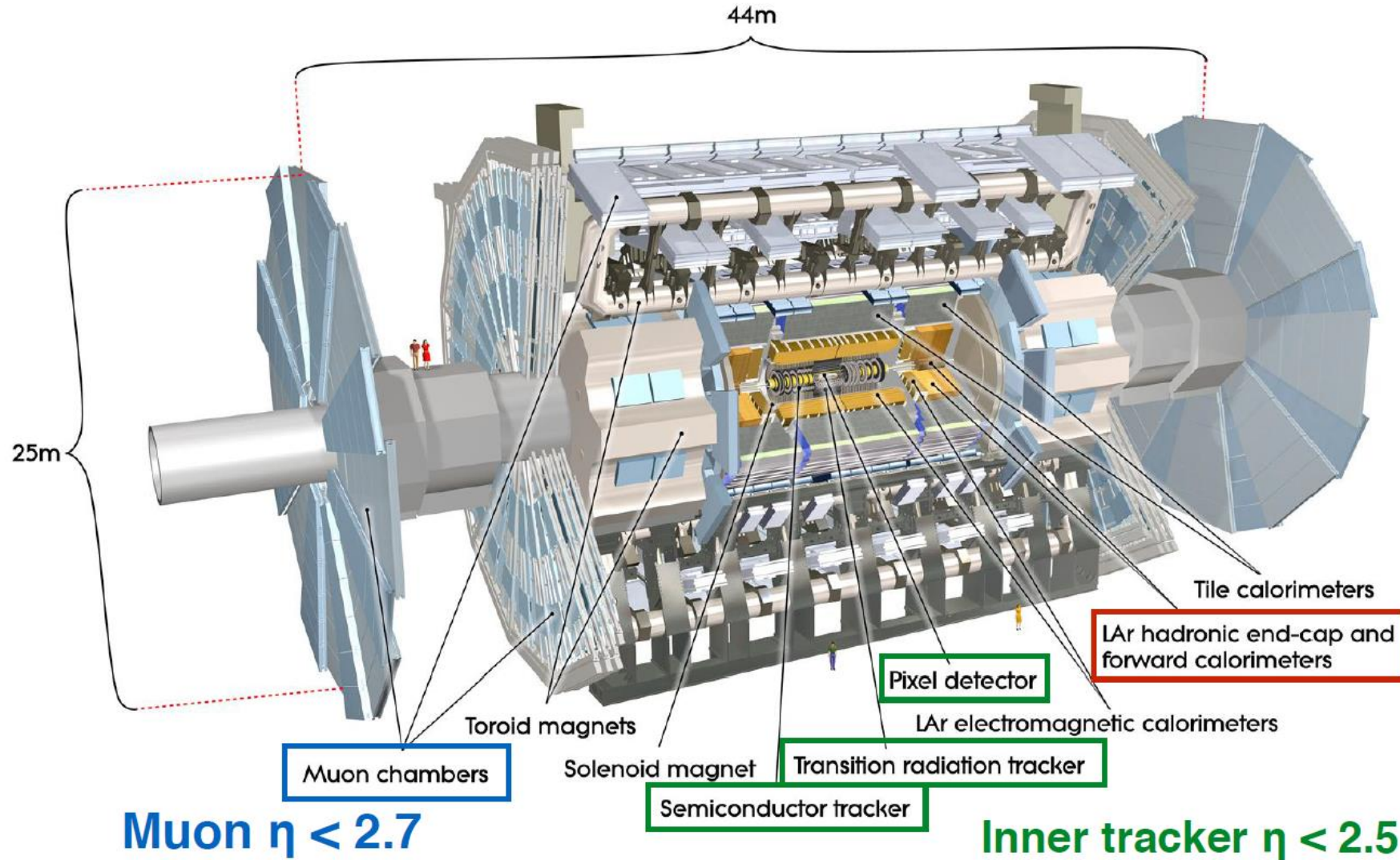
# Conclusions

- Two measurements with similar effects in very different systems.
- In Pb+Pb:  $R_{AA}$  and  $\rho_{AA}$  decrease with increasing centrality, and show no strong dependence on  $p_T$ .
- More suppression for excited states supporting a sequential melting scenario.
  - Models describe the data well and incorporate deconfinement as a key ingredient in the suppression of the Y yields.
- In pp: Difference in  $\langle n_{ch} \rangle$  distributions in events tagged with Y is observed.
- Events with excited Upsilon states have lower  $\langle n_{ch} \rangle$  than with the ground state
  - The difference is largest at low  $p_T$  and almost disappears at  $p_T > 30$  GeV.
- One possible explanation is higher suppression of excited Y states in high multiplicity events, consistent with sequential melting scenario.

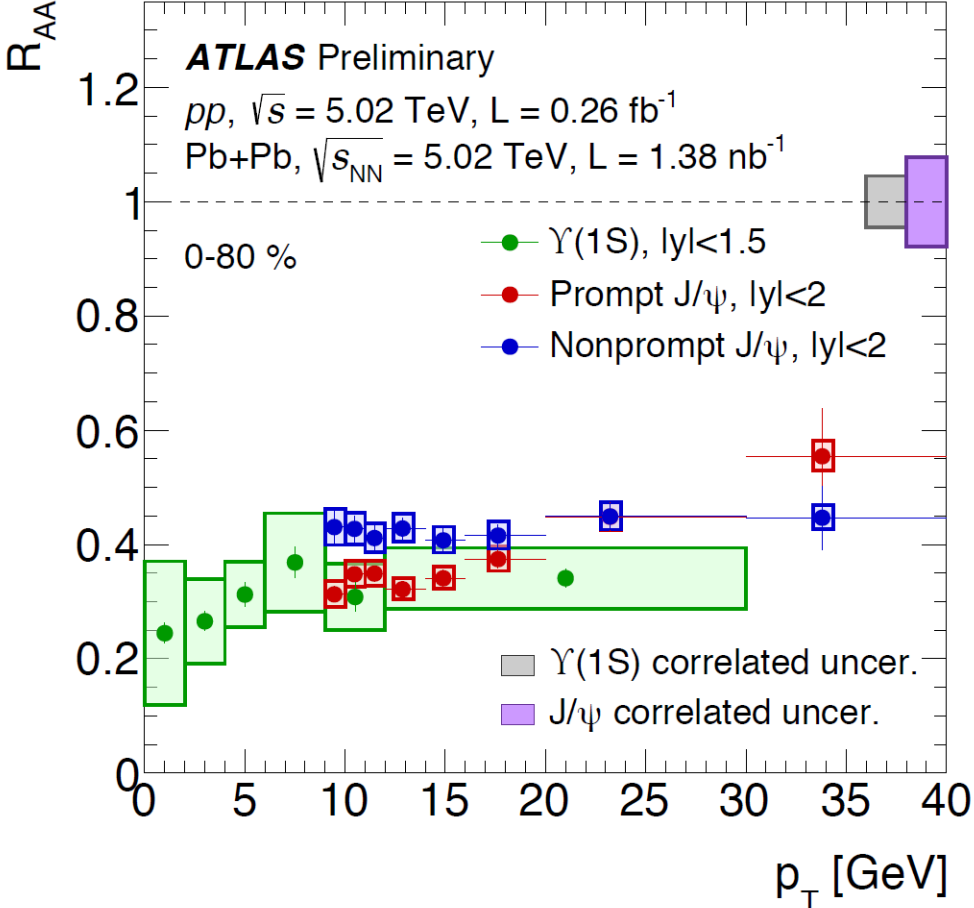
# Backup Slides

# ATLAS Detector

$Y(nS)$  are measured via dimuon decay channel at midrapidity.



# Bottomonium vs. Charmonium



$J/\psi$  (EPJC 78 (2018) 762):  $9 < p_T < 40 \text{ GeV}, |y| < 2$

Similar amount of suppression for  $Y(1S)$  and prompt  $J/\psi$  although  $Y(1S)$  is more tightly bound

Hard to draw a firm conclusion due to different regenerations, feed-down effects, etc.

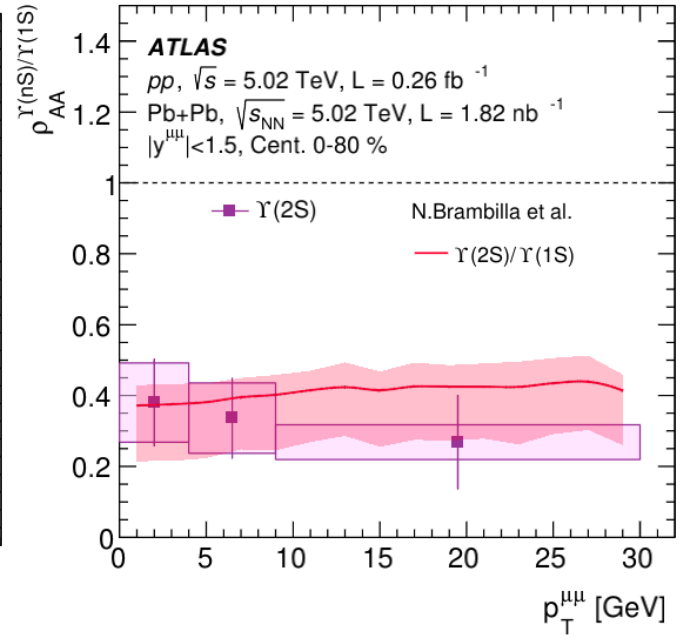
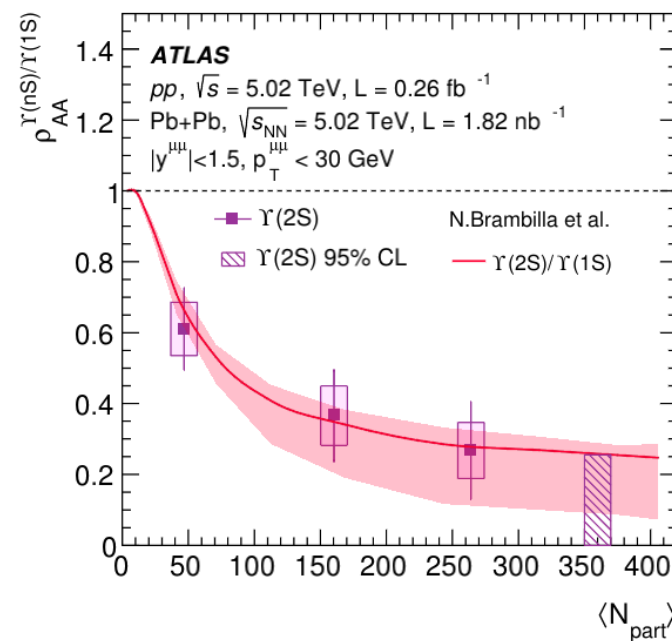
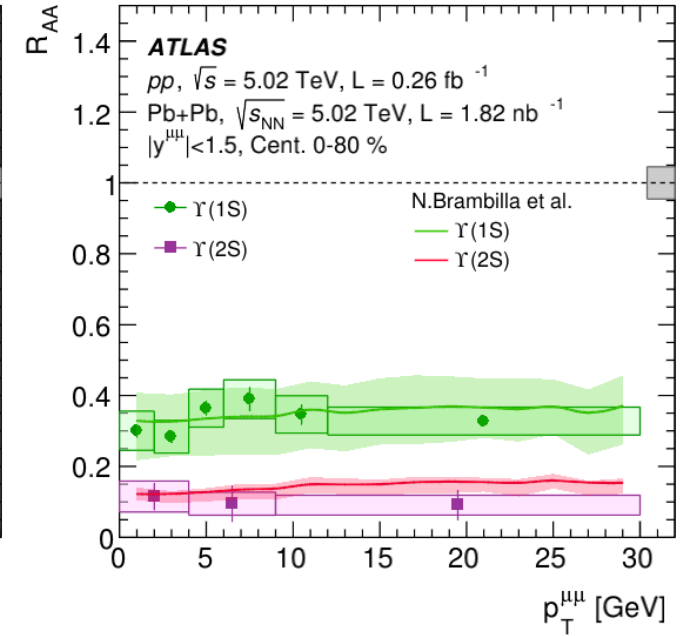
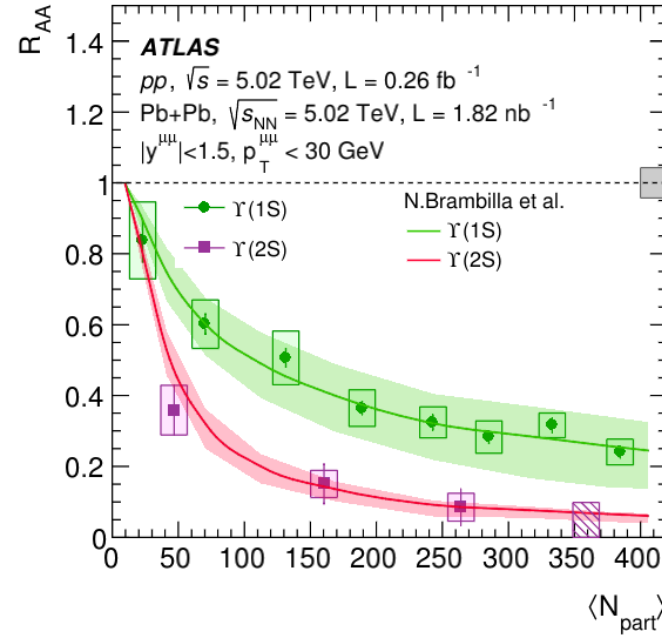
Non-prompt  $J/\psi$   $R_{AA}$  reflects b-quark energy loss

# Theory 1

N. Brambilla et al.,  
*“Bottomonium production in heavy-ion collisions using quantum trajectories: Differential observables and momentum anisotropy.”*  
**Phys. Rev. D 104 (2021) 094049,**  
**arXiv: 2107.06222**

Potential NRQCD and the formalism of open quantum systems to numerically solve the Lindblad equation.

Heavy-quark interactions with the strongly coupled medium are encoded in the two nonperturbative transport coefficients: the heavy-quark momentum diffusion coefficient and its dispersive counterpart.



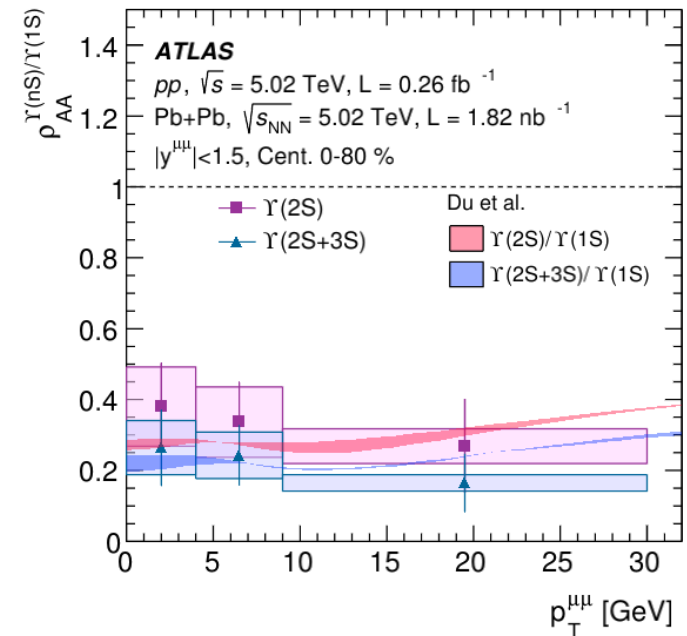
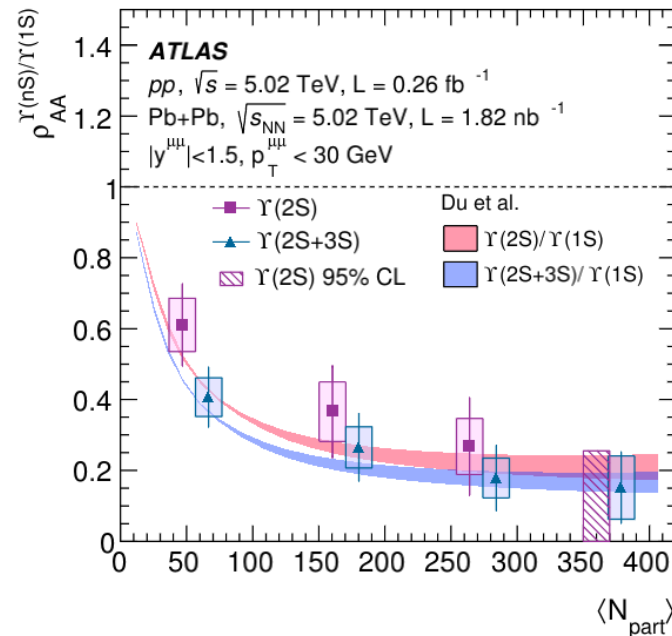
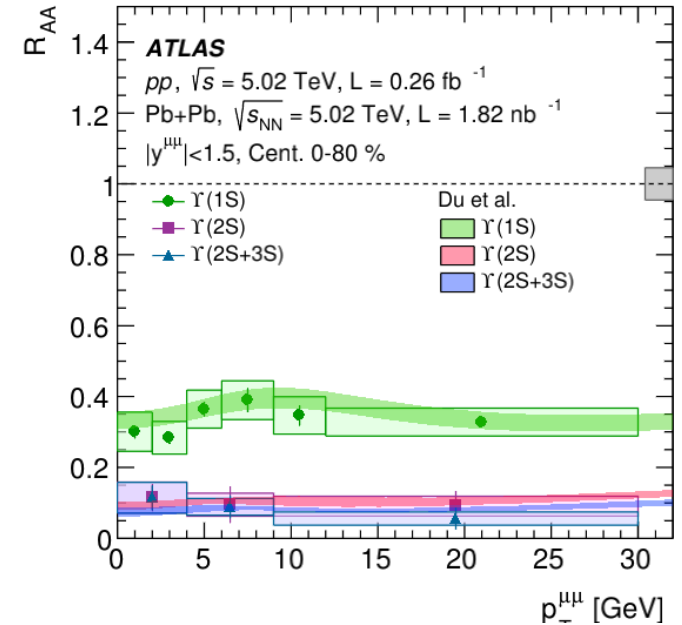
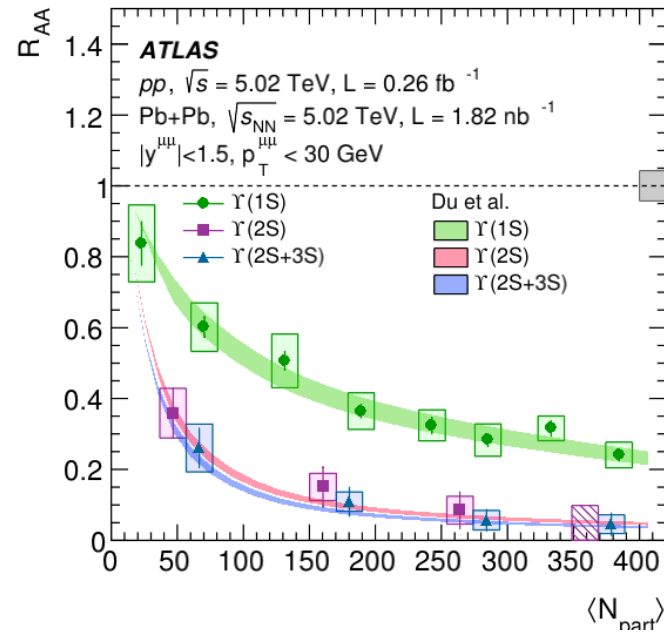
# Theory 2

M. H. X. Du and R. Rapp,

*“Color Screening and Regeneration of Bottomonic in High-Energy Heavy-Ion Collisions.”*

**Phys. Rev. C 96 (2017) 054901,**  
**arXiv: 1706.08670**

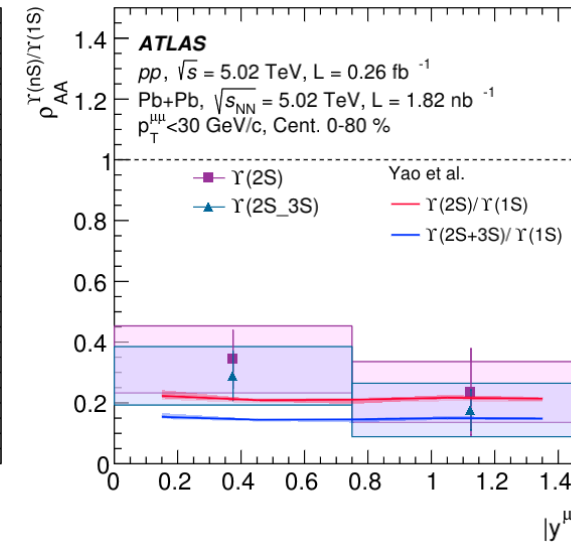
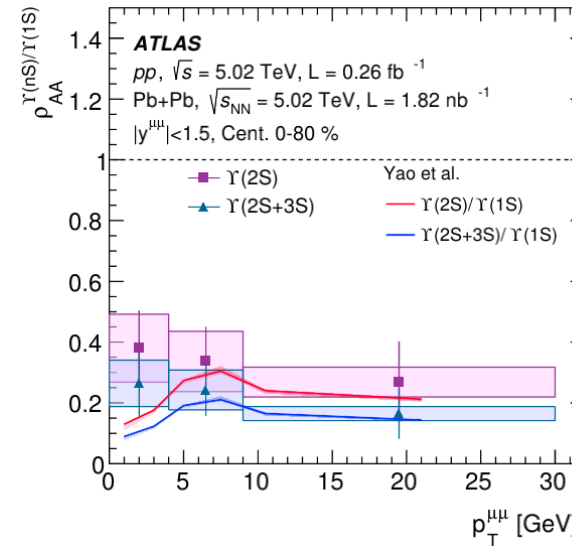
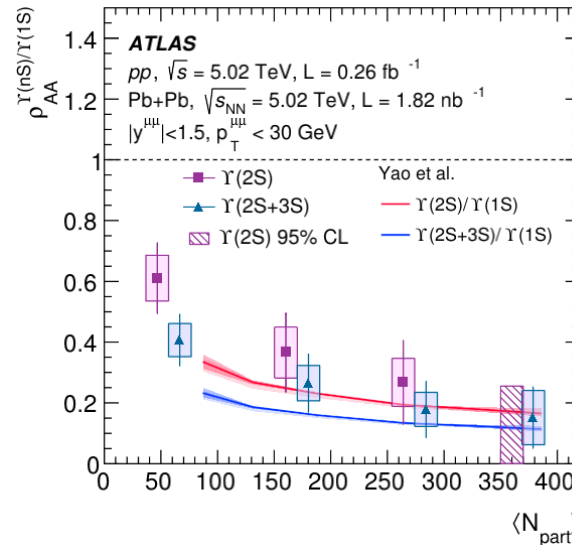
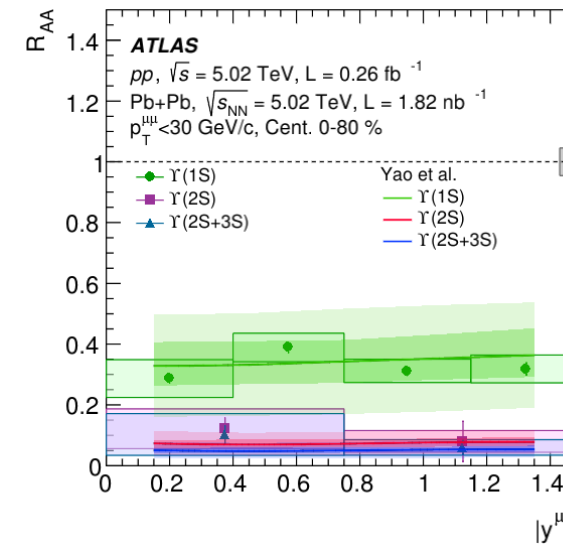
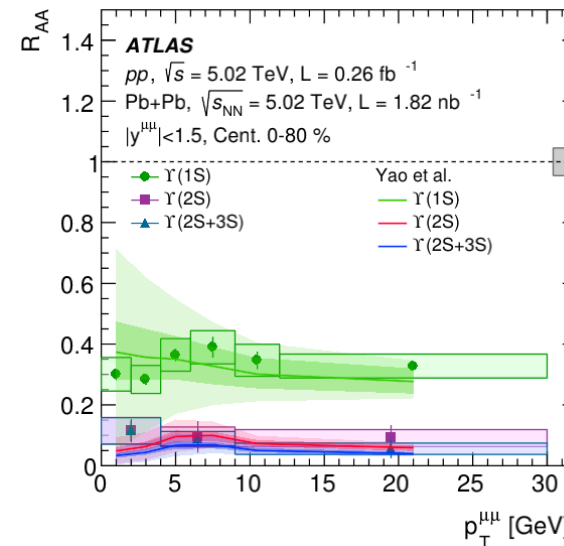
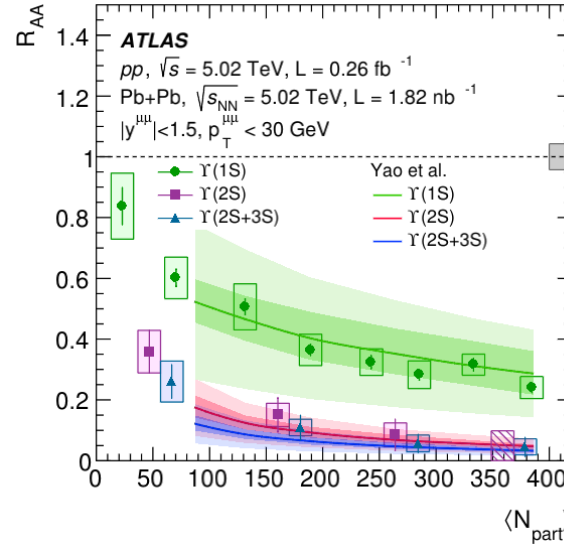
Kinetic-rate equation approach including regeneration and has four dimensionless parameters which characterize the temperature dependence of the pertinent screening masses. These parameters are extracted through the fits to the data already available at RHIC and LHC.



# Theory 3

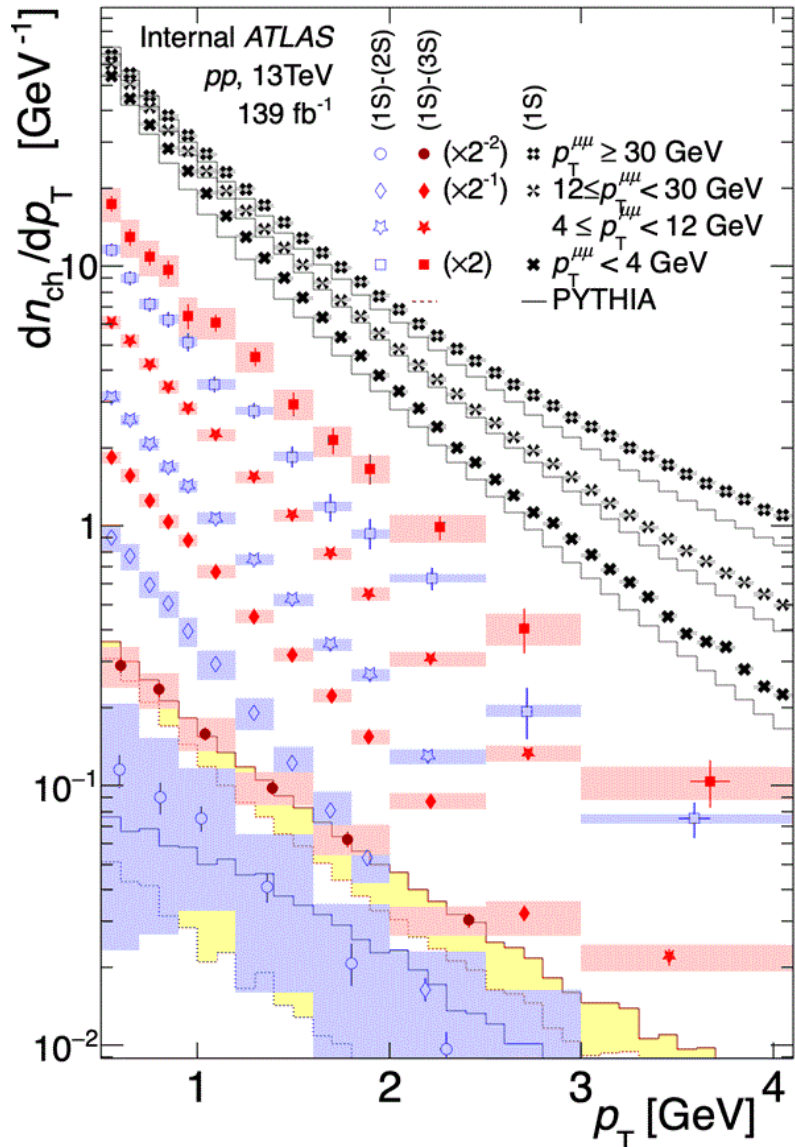
X. Yao et al.,  
*“Coupled Boltzmann Transport Equations of Heavy Quarks and Quarkonia in Quark-Gluon Plasma.”*  
**JHEP 2021 (2021) 46,**  
**arXiv: 2004.06746**

Use of coupled transport equations for open heavy-flavor and quarkonium in order to describe their transport inside the QGP, including regeneration.  
 The model depends on the choice of nPDF and two coupling constant parameters.





# $p_T$ dependence of $\langle n_{ch} \rangle$ in $\Upsilon$ events



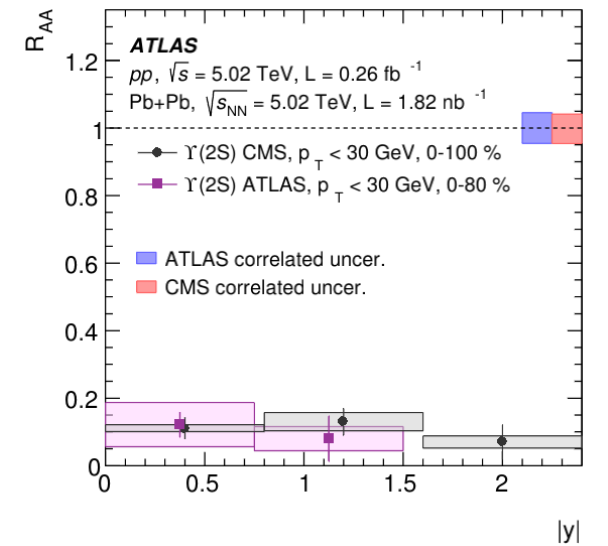
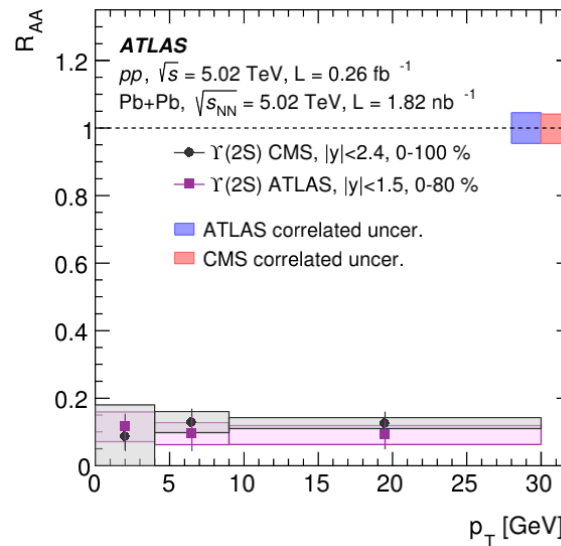
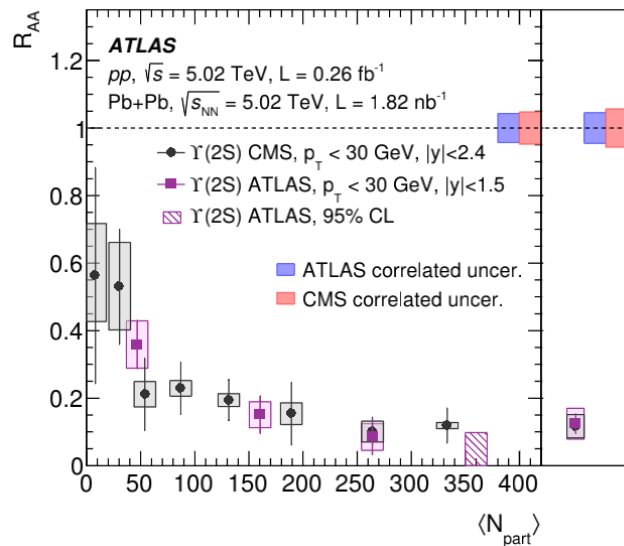
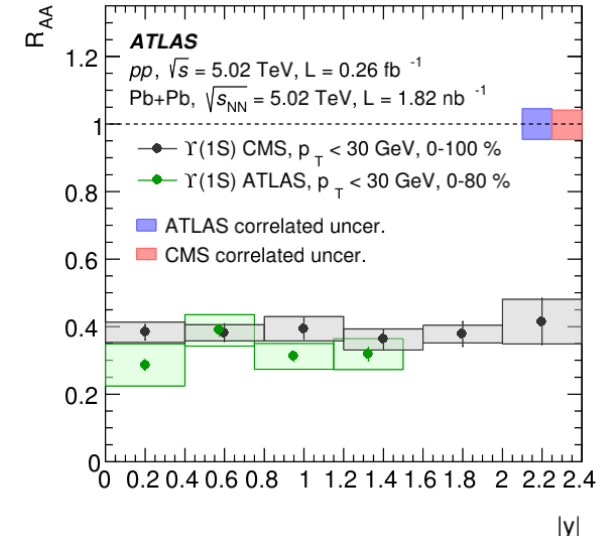
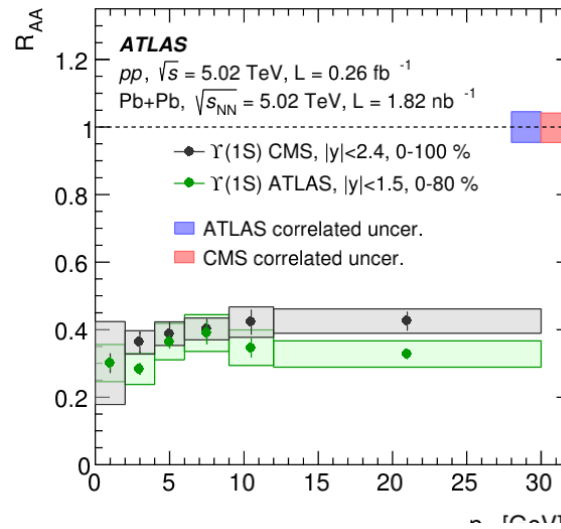
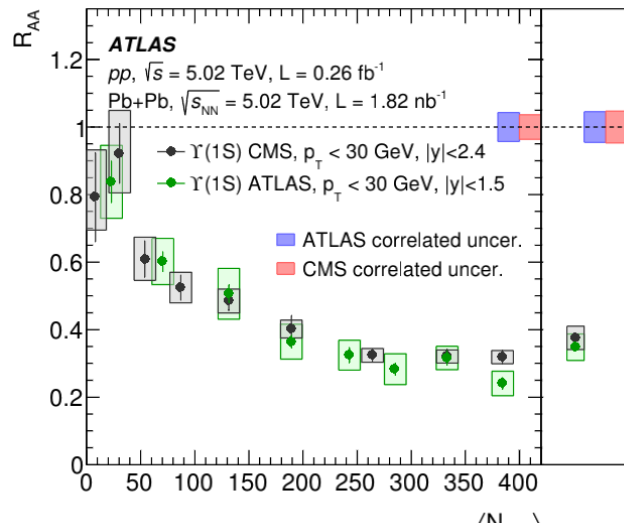
For  $p_T < 4$  GeV, the distribution resembles particles coming from the UE.

For  $p_T < 30$  GeV, subtracted distributions are consistent in shape with  $\Upsilon(1S)$  distribution measured in the lowest  $p_T$ .

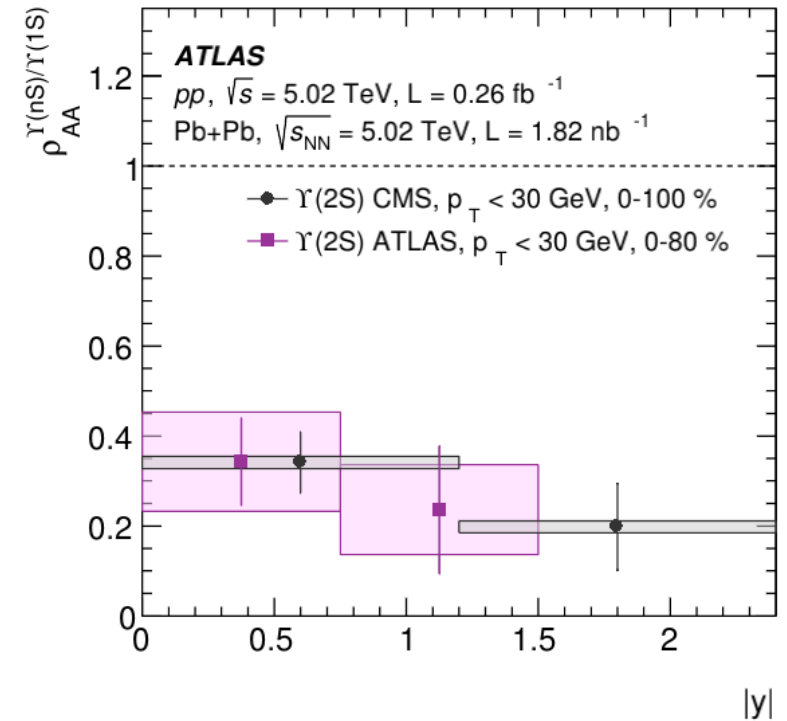
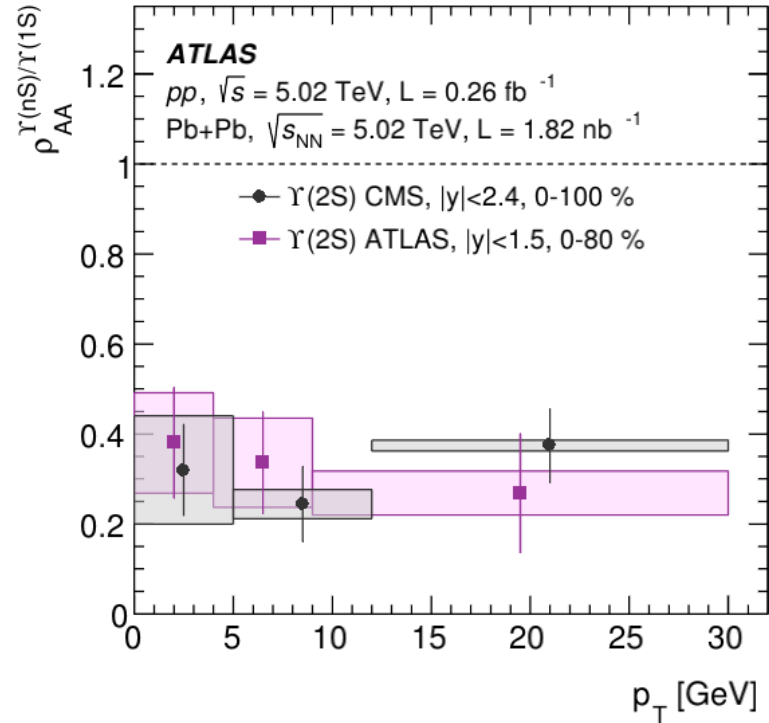
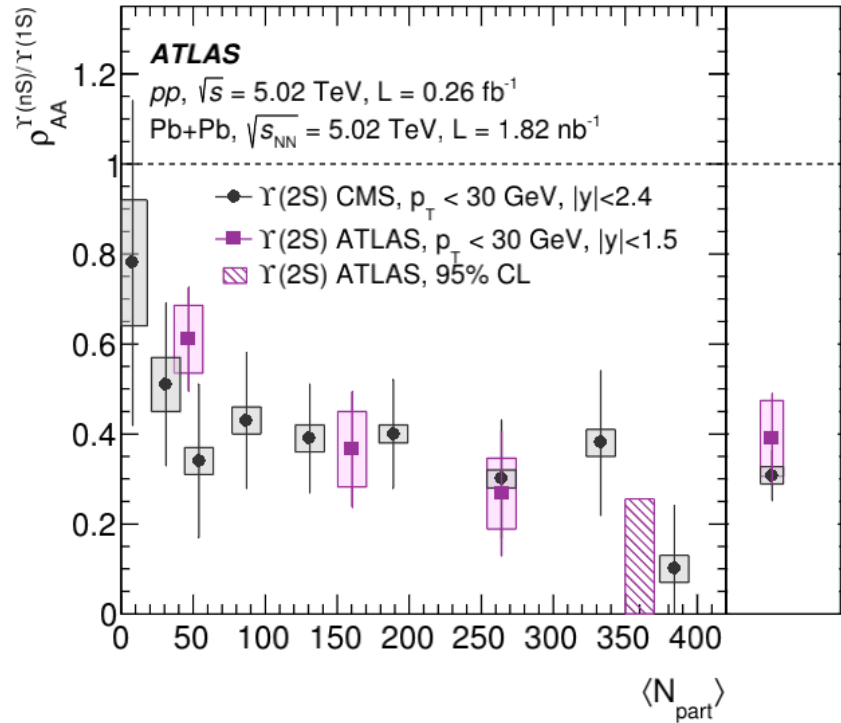
Above 30 GeV, subtracted distributions gets harder, which is explained by feed-down decay processes.

# Comparison with CMS ( $R_{AA}$ )

CMS: “Measurement of nuclear modification factors of  $Y(1S)$ ,  $Y(2S)$ , and  $Y(3S)$  mesons in PbPb collisions at  $\sqrt{s_{NN}} = 5.02$  TeV”, Phys. Lett. B 790 (2019) 270



# Comparison with CMS (double ratio)



# Miscellaneous plots

