

Exact equilibrium distributions for massive and massless fermions with rotation and acceleration

Large vorticity in the QGP. Phenomenological consequences: **spin polarization**.

Exact Wigner Function: quantum corrections in relativistic fluids at global equilibrium.

$$W(x, k) = \frac{1}{(2\pi)^3} \int \frac{d^3p}{2\varepsilon} \sum_{n=1}^{\infty} (-1)^{n+1} e^{-n\tilde{\beta}(in\phi)\cdot p} \times$$
$$\left[e^{-in\frac{\phi:\Sigma}{2}} (m + \not{p}) \delta^4 \left(k - \frac{\Lambda^n p + p}{2} \right) + (m - \not{p}) e^{in\frac{\phi:\Sigma}{2}} \delta^4 \left(k + \frac{\Lambda^n p + p}{2} \right) \right]$$

Exact **distribution functions** and **spin vector** at all orders in thermal vorticity.

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Equilibrium density operator

$$\hat{\rho} = \frac{1}{Z} e^{-b_\mu \hat{P}^\mu + \frac{1}{2} \varpi_{\mu\nu} \hat{J}^{\mu\nu}}$$

The **generators of the Poincaré group** appear.

$$\varpi^{\mu\nu} = \epsilon^{\mu\nu\rho\sigma} \frac{\omega_\rho}{T} u_\sigma + \left(\frac{A^\mu}{T} u^\nu - \frac{A^\nu}{T} u^\mu \right)$$

ϖ is the thermal vorticity, a constant antisymmetric tensor.

Analytic continuation, factorization of the density operator.

$$\hat{\rho} = \frac{1}{Z} e^{-b_\mu \hat{P}^\mu - \frac{i}{2} \phi_{\mu\nu} \hat{J}^{\mu\nu}} = \frac{1}{Z} e^{-\tilde{b}_\mu(\phi) \hat{P}^\mu} \underbrace{e^{-\frac{i}{2} \phi_{\mu\nu} \hat{J}^{\mu\nu}}}_{\equiv \hat{\Lambda}}$$

$$\tilde{b}^\mu(\phi) = \sum_{k=0}^{\infty} \frac{1}{(k+1)!} \underbrace{(\phi^{\mu\nu_1} \phi_{\nu_1\nu_2} \dots \phi^{\nu_{k-1}\nu_k})}_{k\text{-times}} b_{\nu_k}$$

Thermal expectation values

The Wigner function is obtained from the **mean number operator**

$$\langle \hat{a}_s^\dagger(p) \hat{a}_r(p') \rangle = 2\varepsilon' \sum_{n=1}^{\infty} (-1)^{2S(n+1)} \delta^3(\Lambda^n p - p') D^{(S)}(W_{rs}(\Lambda^n, p)) e^{-\tilde{b} \cdot \sum_{k=1}^n \Lambda^k p}$$

$W(\Lambda, p)$ is the Wigner rotation. Expectation values are series

$$\langle \hat{T}^{00} \rangle = \frac{3T^4}{8\pi^2} \sum_{n=1}^{\infty} (-1)^{n+1} \phi^4 \frac{\sinh n\phi}{\sinh^5(n\phi/2)} \rightarrow \frac{7\pi^2}{60\beta^4} - \frac{\alpha^2}{24\beta^4} - \frac{17\alpha^4}{960\pi^2\beta^4}$$

Finite result for $\phi \rightarrow i\alpha$ is obtained by **analytic distillation**.



Distribution function

Exact Wigner function can be used to calculate expectation values.

$$W(x, k) = \frac{1}{(2\pi)^3} \int \frac{d^3 p}{2\varepsilon} \sum_{n=1}^{\infty} (-1)^{n+1} e^{-n\tilde{\beta}(in\phi)\cdot p} \times$$
$$\left[e^{-in\frac{\phi:\Sigma}{2}} (m + \not{p}) \delta^4 \left(k - \frac{\Lambda^n p + p}{2} \right) + (m - \not{p}) e^{in\frac{\phi:\Sigma}{2}} \delta^4 \left(k + \frac{\Lambda^n p + p}{2} \right) \right]$$

We infer the distribution function from: $j_+^\mu(x) = \int \frac{d^3 p}{\varepsilon} p^\mu f(x, p) + N^\mu(x, p)$

$$f(x, p) = \frac{1}{2(2\pi)^3} \sum_{n=1}^{\infty} (-1)^{n+1} e^{-n\tilde{\beta}(-n\varpi)\cdot p} \text{tr} \left(e^{n\frac{\varpi}{2}:\Sigma} \right)$$

Differs from commonly used ansatz in kinetic theories:

$$f_A = \text{tr} \left(\left(e^{\beta\cdot p} e^{-\frac{\varpi}{2}:\Sigma} + \mathbb{I} \right)^{-1} \right)$$

Spin-Polarization

Exact spin vector $S^\mu(p) = \frac{1}{2} \frac{\int d\Sigma \cdot p \operatorname{tr}(\gamma^\mu \gamma_5 W_+(x, p))}{\int d\Sigma \cdot p \operatorname{tr}(W_+(x, p))}$

$$S^\mu(p) = \frac{1}{2m} \frac{\sum_{n=1}^{\infty} (-1)^{n+1} e^{-n\tilde{b}(in\phi) \cdot p} \operatorname{tr}\left(\gamma^\mu \gamma_5 e^{-in\frac{\phi \cdot \Sigma}{2}} \not{p}\right) \delta^3(\Lambda^n p - p)}{\sum_{n=1}^{\infty} (-1)^{n+1} e^{-n\tilde{b}(in\phi) \cdot p} \operatorname{tr}\left(e^{-in\frac{\phi \cdot \Sigma}{2}}\right) \delta^3(\Lambda^n p - p)}$$

The leading order in vorticity agrees with the literature.

$$S^\mu(p) \simeq -\frac{1}{8m} \epsilon^{\mu\alpha\beta\rho} \varpi_{\alpha\beta} p_\rho \frac{n_F(1 - n_F)}{n_F}$$

Full summation of the series under investigation...