

# Gluon shadowing and the role of the $c\bar{c}g$ -Fock state in the diffractive photoproduction of $J/\psi$

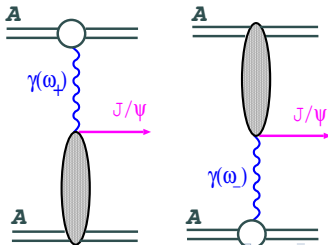
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- The exclusive photoproduction of heavy vector mesons in **ultra-peripheral heavy ion collisions** yields information on the interaction of small color dipoles with nuclei.
- → a probe of the nuclear gluon distribution



# Diffractive photoproduction of $J/\psi$ on protons and nuclei

- in the limit of large photon energy  $\omega$  in the rest frame of the target, the coherence length  $l_c = 2\omega/M_{J/\psi}^2$  becomes much larger than the size  $R$  of the target  $l_c \gg R$
- we can describe photoproduction of the  $J/\psi$  as a splitting of the photon into a  $c\bar{c}$  pair far upstream the target, and an interaction of a color dipole of size  $\mathbf{r}$  formed by quark and antiquark. After the interaction with the target, the scattered  $c\bar{c}$  pair evolves into the outgoing final state vector meson.

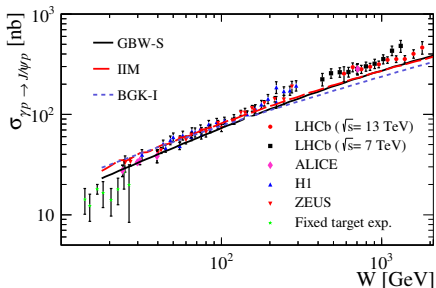
## Forward amplitude

$$A(\gamma A \rightarrow J/\psi A; W, \mathbf{q} = 0) = 2i \int d^2\mathbf{b} \langle J/\psi | \Gamma_A(x, \mathbf{b}, \mathbf{r}) | \gamma \rangle = i \int_0^1 dz \int d^2\mathbf{r} \Psi_{J/\psi}^*(z, \mathbf{r}) \Psi_\gamma(z, \mathbf{r}) \sigma(x, \mathbf{r}).$$

- the main ingredients are the light-front wave functions of photon and vector meson and the dipole cross section
- **diffractive photoproduction of heavy vector mesons gives us information on the interaction of small color dipoles with protons or cold nuclear matter**

# Dipole cross section

- We use dipole cross sections that have been fitted to precise HERA data on the proton structure function  $F_2(x, Q^2)$ .
- a good agreement with  $J/\psi$  diffractive data up to  $\gamma p$ -cm energies of a few hundred GeV is obtained.

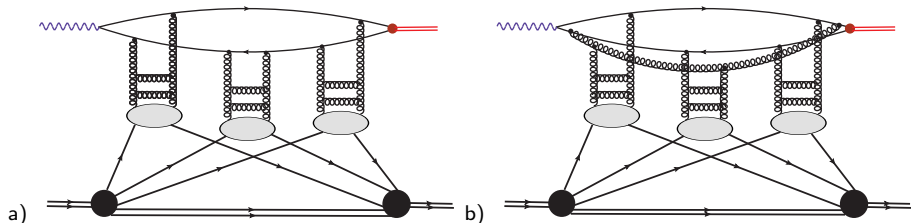


## Dipole-cross section $\leftrightarrow$ gluon distribution

$$\sigma(x, r) = \frac{4\pi\alpha_S}{N_c} \int \frac{d^2\kappa}{\kappa^4} \frac{\partial xg(x, \kappa^2)}{\partial \log(\kappa^2)} \left(1 - e^{i\kappa \cdot r}\right) \approx r^2 \frac{\pi^2\alpha_S}{N_c} xg\left(x, \frac{1}{r^2}\right)$$

- behaviour at small  $r^2$  suggests a proportionality of the diffractive amplitude to the gluon distribution of the target. But: on nuclear targets large higher twists effects from multiple scattering.

## Diffractive photoproduction on nuclei



- rescattering of the  $c\bar{c}$  dipole only can be obtained by a simple Glauber exponential

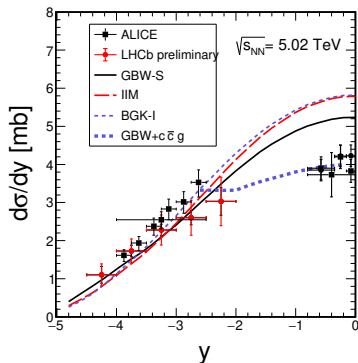
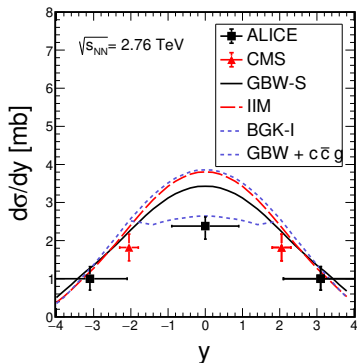
$$\Gamma_A(x, \mathbf{b}, \mathbf{r}) = 1 - S_A(x, \mathbf{b}, \mathbf{r}), \text{ with } S_A(x, \mathbf{b}, \mathbf{r}) = \exp \left[ -\frac{1}{2} \sigma(x, \mathbf{r}) T_A(\mathbf{b}) \right].$$

- the rescattering of the  $c\bar{c}$  pair corresponds to a resummation of higher-twist terms, **not** to a nuclear modification of the nuclear glue! The  $c\bar{c}g$  yields a correction:

$$\Gamma_A(x, \mathbf{r}, \mathbf{b}) = \Gamma_A(x_A, \mathbf{r}, \mathbf{b}) + \log \left( \frac{x_A}{x} \right) \Delta \Gamma_A(x_A, \mathbf{r}, \mathbf{b})$$

- the **nuclear glue** appears only, once the rescattering of  $c\bar{c}g$  states are taken into account. But here the splitting amplitude  $q \rightarrow qg$  in impact parameter space needs an infrared regularization ( **gluon propagation radius**  $R_c$  ).  $x_A \sim 0.01$  also is a parameter of the model.

# Ultraparipheral nuclear collisions



- Glauber-Gribov theory including only rescattering of the  $c\bar{c}$  dipole works well in the forward region (large rapidities).
- In the central rapidity region inclusion of the  $c\bar{c}g$  state introduces additional shadowing which is needed to describe the data. Strong dependence on  $R_c$ , a rather small  $R_c \sim 0.21$  fm is preferred.
- Shadowing due to the  $c\bar{c}g$  state can be (roughly) identified with gluon shadowing of the nuclear pdf. It depends on the infrared regulator. For  $J/\psi$  gluon shadowing is not a prediction of perturbation theory (hard scale  $Q^2 \sim 2.25 \text{ GeV}^2$ ). **Moderate gluon shadowing.**