# Finite Nc and non-eikonal corrections to induced emission

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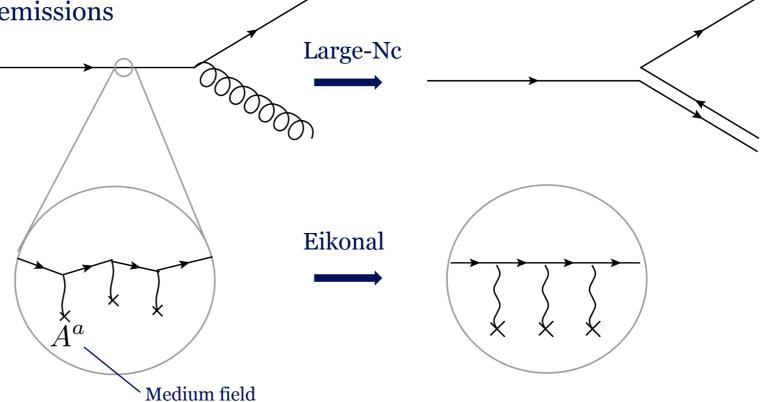
In collaboration with Konrad Tywoniuk

Based on <a href="https://arxiv.org/abs/2107.02542">https://arxiv.org/abs/2107.02542</a>, and ongoing work



## Medium induced emissions

- Interactions with the medium lead to emissions
- Two simplifying approximations:
- 1. Large-Nc approximation
  - Take number of colors (Nc) to infinity
- 2. Eikonal approximation
  - Let partons travel on straight lines
  - Good for high energy
- Important to understand the error of the approximations

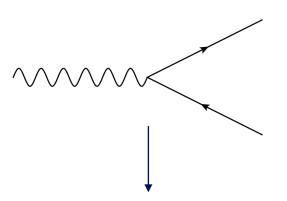


- Constant medium interactions means color continuously rotates
- Effect encapsulated in Wilson line:  $V_R(t, t_0; \mathbf{r}(t)) = \mathcal{P} \exp \left[ ig \int_{t_0}^t \mathrm{d}s \, A^a(s, \mathbf{r}(s)) T_R^a \right]$

## Splitting processes

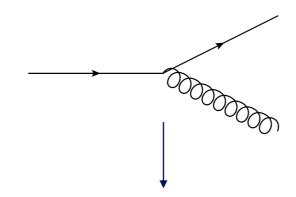
#### Calculate emission spectrum for three processes

Pair production



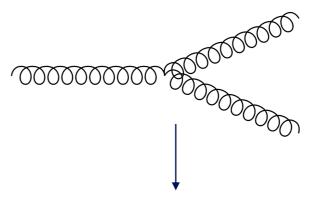
$$rac{1}{N_c}\langle {
m tr}[V_1V_2^\dagger V_{ar{2}}V_{ar{1}}^\dagger]
angle$$

Quark emitting gluon



$$\begin{split} \frac{1}{N_c^2-1} \left\langle \text{tr}[V_{\bar{1}}^\dagger V_1 V_2^\dagger V_{\bar{2}}] \, \text{tr}[V_{\bar{2}}^\dagger V_2] \\ -\frac{1}{N_c} \, \text{tr}[V_{\bar{1}}^\dagger V_1] \right\rangle \end{split}$$

Gluon emitting gluon



$$\frac{1}{N_c^2 - 1} \left\langle \operatorname{tr}[V_{\bar{1}}^{\dagger} V_1 V_2^{\dagger} V_{\bar{2}}] \operatorname{tr}[V_{\bar{2}}^{\dagger} V_2] - \frac{1}{N_c (N_c^2 - 1)} \left\langle \operatorname{tr}[V_1 V_{\bar{1}}^{\dagger}] \operatorname{tr}[V_2 V_{\bar{2}}^{\dagger} V_{\bar{1}} V_1^{\dagger}] \operatorname{tr}[V_{\bar{2}} V_2^{\dagger}] - \operatorname{tr}[V_1 V_{\bar{1}}^{\dagger} V_2 V_{\bar{2}}^{\dagger} V_{\bar{1}} V_1^{\dagger} V_2 V_2^{\dagger}] \right\rangle$$

- Average out medium effects, represented by  $\langle \dots \rangle$
- Leads to increasingly complex medium averaged traces of Wilson lines

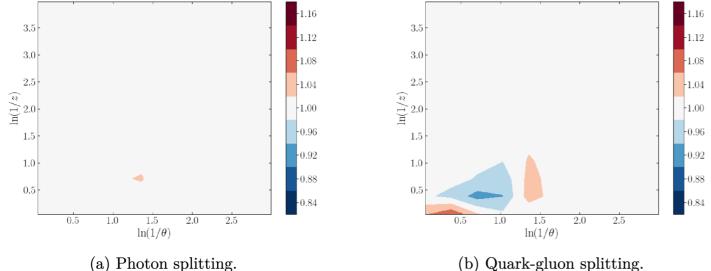
## Finite Nc corrections

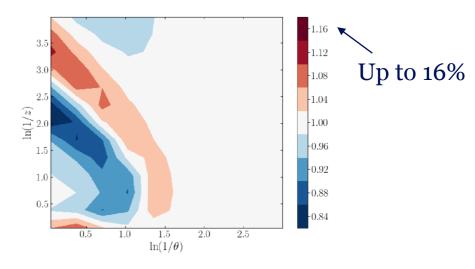
### Wilson line correlators can be calculated through system of differential equations

$$\frac{\mathrm{d}}{\mathrm{d}t} \begin{bmatrix} \langle \operatorname{tr}[VV] \operatorname{tr}[VV] \rangle \\ \langle \operatorname{tr}[VVVV] \rangle \end{bmatrix} = \mathbb{V}(t) \begin{bmatrix} \langle \operatorname{tr}[VV] \operatorname{tr}[VV] \rangle \\ \langle \operatorname{tr}[VVVV] \rangle \end{bmatrix}$$

- The potential matrix V(t) simplifies greatly in the large-Nc limit
  - Example:  $\langle \operatorname{tr}[VV] \operatorname{tr}[VV] \rangle \stackrel{\text{large-}N_c}{\to} \langle \operatorname{tr}[VV] \rangle \langle \operatorname{tr}[VV] \rangle$
- Expect correction around  $\sim 1/N_c^2 \simeq 10\%$

#### Ratio of emission spectrum for large-Nc/finite Nc





(c) Gluon-gluon splitting.

More complex color structure leads to bigger correction!

## Non-eikonal corrections

#### Partons get kicked around by medium

- Path through the medium is not straight
- Important for lower energy
- This leads to path integral over Wilson lines

$$\mathcal{K}_1 \equiv \int \mathcal{D} oldsymbol{u} \int \mathcal{D} oldsymbol{v} \, \mathrm{e}^{i rac{\omega}{2} \int \mathrm{d} s \, (\dot{oldsymbol{u}}^2 - \dot{oldsymbol{v}}^2)} \langle \mathrm{tr}[VV] \, \mathrm{tr}[VV] 
angle 
onumber \ \mathcal{K}_2 \equiv \int \mathcal{D} oldsymbol{u} \int \mathcal{D} oldsymbol{v} \, \mathrm{e}^{i rac{\omega}{2} \int \mathrm{d} s \, (\dot{oldsymbol{u}}^2 - \dot{oldsymbol{v}}^2)} \langle \mathrm{tr}[VVVV] 
angle \, ,$$

• Differential equation from last slide becomes a Schrödinger equation

$$\left(irac{\partial}{\partial t}+rac{\partial_{m{u}}^2-\partial_{m{v}}^2}{2\omega}+i\mathbb{V}(t,m{u},m{v})
ight)igg[m{\mathcal{K}}_1 igg]=i\delta(t-t_2)\delta^2(m{u}-m{u}_2)\delta^2(m{v}-m{v}_2)$$

- Ongoing work: Solving this numerically
- Expect small corrections for high energy, bigger corrections for low energy

t=0.0

