



29TH INTERNATIONAL  
CONFERENCE ON ULTRARELATIVISTIC  
NUCLEUS - NUCLEUS COLLISIONS

APRIL 4-10, 2022  
KRAKÓW, POLAND

## Combined constraints from jet and hadron quenching to $\hat{q}$

---

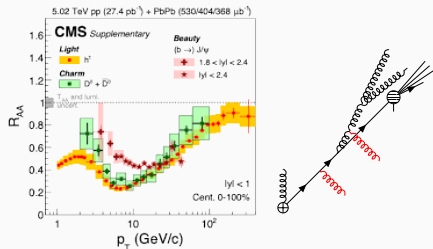
Weiyao Ke (LANL), in collaboration with Xin-Nian Wang (LBNL)

*Based on W Ke, X-N Wang JHEP 05, 041 (2021)*

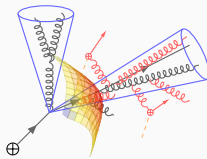
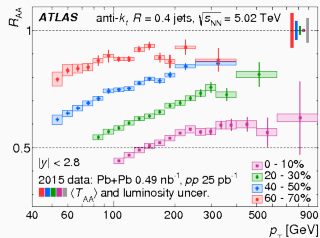


# Hadron & jet probe different aspects of parton dynamics in the QGP

## Single hadron suppression at high- $p_T$



## Single jet suppression



## Probe modification/loss of large- $z$ partons

- Induced radiations that modify  $D(z)$  at  $zE \gg T$ .
- Energy loss from soft rad. & collisions,  $\omega \sim T$ .

## Sensitive to redistribution of "lost energy" by

- Collisions, induced radiations.
- Collective excitations.

**This work:** - study hadron & jet within the LIDO parton transport model.  
 - a consistent transport parameter for jet and hadron from Bayesian analysis.  
 - basis for predicting other jet modifications:  $R$ -dependence, fragmentation, shape.

& many other predictions for, [dijet T. Rinn ATLAS](#) , [b-jet,  \$\gamma\$ -jet S. Araya, ATLAS](#) , [Y. Go, ATLAS](#)

# Method: LIDO transport model approach for hadron and jet

Hard parton transport  $f_H = f(t, x, p)\Theta(p \cdot u > 4T)$ ,  $f_s = e^{-p \cdot u/T}$

$$\frac{df_H}{dt} = \Theta(p \cdot u > 4T) \{ \mathcal{D}f_H + \mathcal{D}_{12}f_H \rightarrow \text{small-}q \text{ diffusion \& diff.-induced rad.} \\ \mathcal{C}_{22}f_H + \mathcal{C}_{23}f_H \} \rightarrow \text{large-}q \text{ collision \& coll.-induced rad.}$$

1. Soft diffusion<sup>1</sup>:  $\mathcal{D} = -\eta \nabla_p - \frac{\hat{q}_s}{2} \nabla_p^2$



2. Large- $q$  collision:  $\frac{d\sigma}{d^2\mathbf{q}_\perp} \propto \frac{\alpha_s^2(\mathbf{q}_\perp)}{q_\perp^4} \Theta(\mathbf{q}_\perp^2 - Q_c^2)$

$$\mathcal{C}_{22}f_H = \int_{\mathbf{k}, \mathbf{q}} \left[ \frac{d\sigma}{d^2\mathbf{q}} f_s(\mathbf{k}) f_H(\mathbf{p} - \mathbf{q}) + \dots \right]$$



▷ Combine to the jet transport parameter

$$\hat{q}_F = \underbrace{\alpha_s C_F T m_D^2 \ln \frac{Q_c^2}{m_D^2}}_{\hat{q}_s} + \int_{Q_c^2} \mathbf{q}_\perp^2 \frac{d\sigma}{d^2\mathbf{q}_\perp} d^2\mathbf{q}_\perp$$

<sup>1</sup>In J. Ghiglieri, G. D. Moore, D. Teaney JHEP 03, 095(2016), separation requires  $m_D \ll Q_c \ll T$ . we take  $Q_c = 2m_D$

# Method: LIDO transport model approach for hadron and jet

Hard parton transport  $f_H = f(t, x, p)\Theta(p \cdot u > 4T)$

$$\frac{df_H}{dt} = \Theta(p \cdot u > 4T) \{ \mathcal{D}f_H + \mathcal{D}_{12}f_H \rightarrow \text{small-}q \text{ diffusion \& diff.-induced rad.} \\ \mathcal{C}_{22}f_H + \mathcal{C}_{23}f_H \} \rightarrow \text{large-}q \text{ collision \& coll.-induced rad.}$$

3. Diffusion-induced 1 to 2 radiation:

$$\mathcal{D}_{12}f(x) = \int \frac{dz}{1-z} z d^2\mathbf{k}_\perp \frac{\alpha_s P_{ij}(z)}{2\pi^2 \mathbf{k}_\perp^2} \frac{\hat{q}_s}{k_\perp^2} f\left(\frac{x}{1-z}\right) + \dots$$

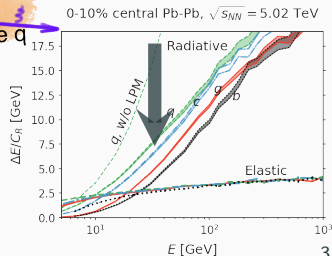
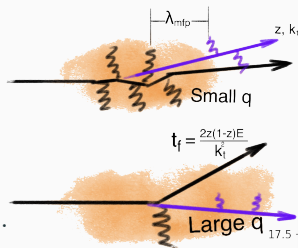
4. Large- $q$  collision-induced 2 to 3 radiation.

$$\mathcal{C}_{23}(x) = \int_{\mathbf{k}, \mathbf{q}} \frac{dz}{1-z} \frac{d\sigma_{23}}{dz d^2\mathbf{k}_\perp d^2\mathbf{q}} f_s(\mathbf{k}) f_H\left(\frac{x}{1-z}, \mathbf{p} - \mathbf{q}\right) + \dots$$

▷ Landau-Pomeranchuk-Midgal effect implemented.

Radiation suppressed by  $\#\lambda_{\text{mfp}}/\tau_f$ .

[WK, Y Xu, S Bass, PRC100 064911 (2019)].



# Method: A model for collective excitation induced by energy loss

- Energy-momentum deposition to soft sector:

$$\frac{d\delta p^\mu}{dt}(t, x) = \int_{\mathbf{p}} \Theta(p \cdot u < 4T) p^\mu \frac{d}{dt} f_H(t, x, p)$$

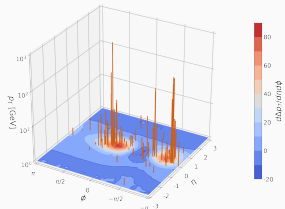
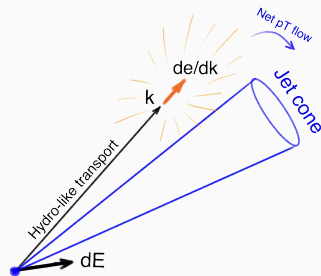
- An ideal-hydro response:

$$\frac{de}{d\Omega_{k'}} = \frac{\delta p^0 + \hat{k}' \cdot \delta \vec{p} / c_s}{4\pi}, \quad \frac{d\vec{p}}{d\Omega_{k'}} = \frac{3(c_s \delta p^0 + \hat{k}' \cdot \delta \vec{p}) \hat{k}'}{4\pi}$$

- Freeze-out to massless particles w/ radial flow  $v_\perp$   
 $\Rightarrow$  corrections to momentum density in the cone:

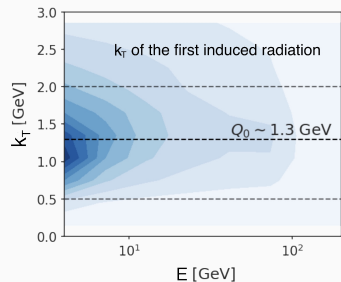
$$\frac{d\Delta p_T}{d\phi d\eta} = \int \frac{3}{4\pi} \frac{\frac{4}{3} \sigma u_\mu - \hat{p}_\mu}{\sigma^4} \delta p^\mu(\hat{k}) \frac{d\Omega_{\hat{k}}}{4\pi}$$

$$\sigma = \gamma_\perp [\cosh(\eta - \eta_s - \eta_{\hat{k}}) - v_\perp \cos(\phi - \phi_{\hat{k}})]$$



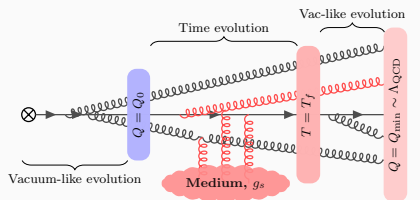
# Method: Merging vacuum-like evolution and transport equation

- Medium effects take place in a more restricted  $\mathbf{k}_\perp$  region:
  - Collisions  $|\mathbf{k}_\perp| \sim m_D = 0.4 \dots 1.2$  GeV.
  - Induced radiation  $|\mathbf{k}_\perp| \sim 1$  GeV.
- A “sudden” transition from DGLAP to transport at  $Q_0$ .
  - A reasonable  $Q_0^2 \approx \langle \mathbf{k}_\perp^2 \rangle = \int_{t_0}^{\tau_f} \hat{q}(t) dt \propto t_0 T_0^3$  in fast-expanding medium.
  - $Q_0$ : weak energy dependence; change in different medium.



Systems	Pb-Pb 5 TeV		Au-Au 0.2 TeV	Xe-Xe 5.44 TeV
	0-5%	40-50%	0-5%	0-5%
$5t_0 T_0^3$ [GeV <sup>2</sup> ]	1.1	0.55	0.46	0.96

# Uncertainties: two separation scales affect the extraction of $g_s, \hat{q}$

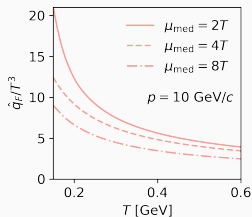
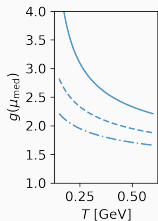


Objective: determine “jet-medium coupling  $g_s$ ”  
or “jet transport parameter  $\hat{q}$ ”.

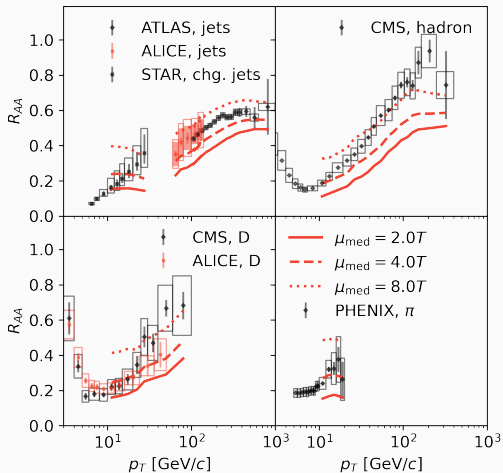
Uncertainties:

- $0.5 < Q_0 < 2.0$  GeV: separates vacuum-like and transport evolution.
- $0.15 < T_f < 0.17$  GeV: color source = 0 for  $T < T_f$ .
- $0.7\pi T < \mu_{\text{med}} < 4\pi T$ : controls in-medium  $g_s$ :

$$\frac{g_s^2(\mathbf{k}_\perp)}{4\pi} = \frac{4\pi}{9} \ln^{-1} \left[ \frac{\max\{\mathbf{k}_\perp^2, \mu_{\text{med}}^2\}}{\Lambda^2} \right]$$



# Uncertainties: $\mu_{\text{med}}$ or $g_s(\max\{k_T, \mu_{\text{med}}\})$



## Experimental data:

[STAR charged jet: PRC 102, 054913(2020)]

[ALICE jet: PRC 101 034911(2020)]

[ATLAS jet: PLB 790 108-128(2019)]

[CMS  $D$ : PLB 287 474-496(2018)]

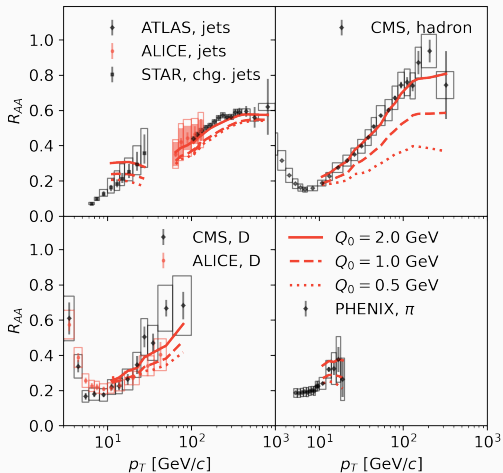
[CMS  $h$ : JHEP 04, 039(2017)]

[PHENIX  $\pi$ : PRC 87, 034911(2013)]

Changing the coupling strength by  
varying  $\mu_{\text{med}} = 2T, 4T, 8T$  GeV



# Uncertainties of $Q_0$ : advantage of using jet $R_{AA}$ to calibrate $\hat{q}$



## Experimental data:

[STAR charged jet: PRC 102, 054913(2020)]

[ALICE jet: PRC 101 034911(2020)]

[ATLAS jet: PLB 790 108-128(2019)]

[CMS  $D$ : PLB 287 474-496(2018)]

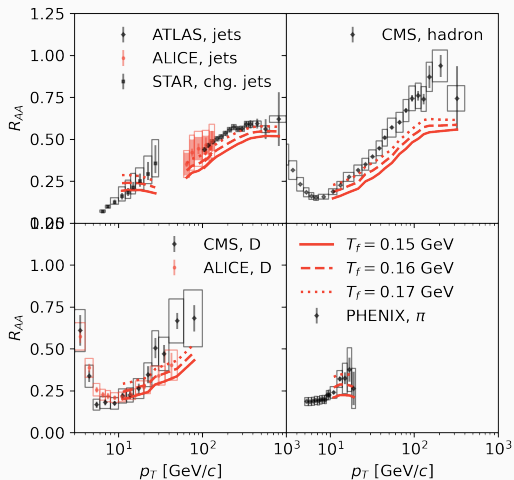
[CMS  $h$ : JHEP 04, 039(2017)]

[PHENIX  $\pi$ : PRC 87, 034911(2013)]

Test the variation of  $Q_0 = 0.5, 1.0, 2.0$  GeV.

- Light hadron  $R_{AA}$  are very sensitive to  $Q_0$ .
- Jet and heavy-flavor  $R_{AA}$  at the LHC energy are the least sensitive.

# Uncertainties: the QGP termination temperature $T_f$



## Experimental data:

[STAR charged jet: PRC 102, 054913(2020)]

[ALICE jet: PRC 101 034911(2020)]

[ATLAS jet: PLB 790 108-128(2019)]

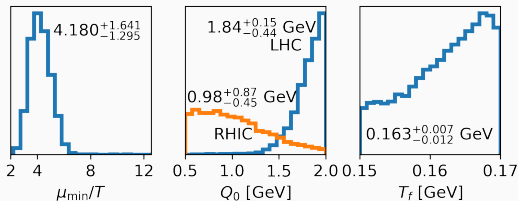
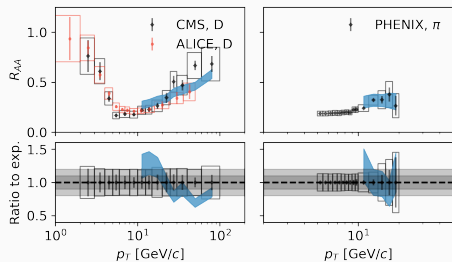
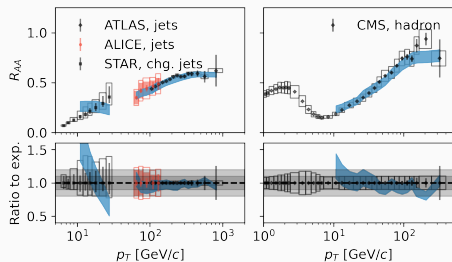
[CMS  $D$ : PLB 287 474-496(2018)]

[CMS  $h$ : JHEP 04, 039(2017)]

[PHENIX  $\pi$ : PRC 87, 034911(2013)]

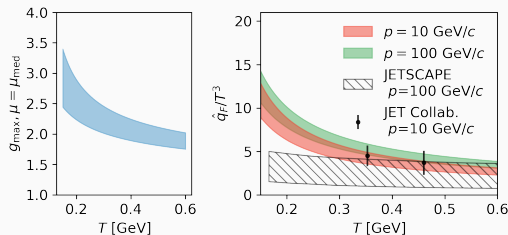
Change  $T_f = 0.15, 0.16, 0.17$  GeV.  $\Leftrightarrow$  effectively change color density near  $T_c$ .

# Results: Bayesian analysis of $\mu_{\text{med}}$ , $Q_0^{\text{LHC}}$ , $Q_0^{\text{RHIC}}$ , $T_f$



- $Q_0^{\text{LHC}}$  varies independently from  $Q_0^{\text{RHIC}}$ .  
 $Q_0^{\text{LHC}} > Q_0^{\text{RHIC}}$  is consistent with the expectation from  $T_0^{\text{LHC}} > T_0^{\text{RHIC}}$ .
- Favors higher  $T_f$  than the pseudo-critical  $T_c$ .
- Running of  $g_s$  in medium saturates around  $k_{\perp} > \mu_{\text{med}} \approx 4.2T$  (or  $1.3\pi T$ ).

# Results: jet-medium coupling $g$ and jet transport parameter $\hat{q}$



Left: maximum coupling at different temperature  $g_s(\mu_{\text{med}})$

Right:  $\hat{q}$  at  $p = 10$  and  $100$  GeV for a quark.

Compared to [JET Collab: PhysRevC.90.014909 (2014), JETSCAPE : PRC 104, 024905 (2021)] using inclusive hadrons.

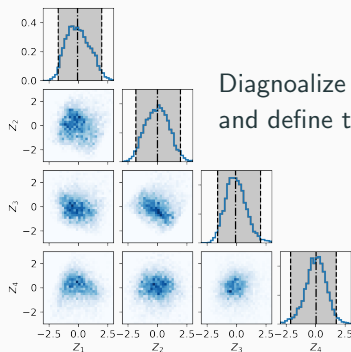
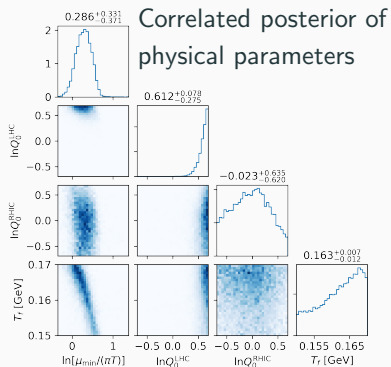
- Results consistent with JET collaboration at high  $p$ .

- Higher than the recent JETSCAPE Collaboration analysis.

Possible reason: JETSCAPE include medium corrections to the DGLAP stage ( $\mathbf{k}_{\perp}^2 > Q_0^2$ )

# ”Representative parameter sets” for the study of other observables

To make “quick” predictions: we defined central + error parameter sets

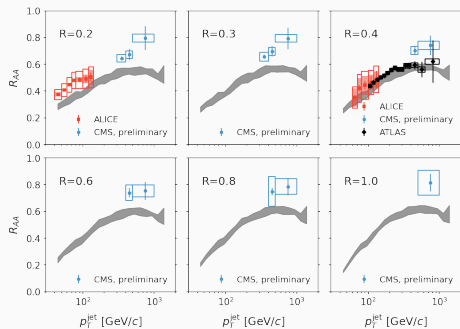


Diagonalize the covariance matrix and define the  $50^{+47.5\%}_{-47.5\%}$  bands.

Transform the median and  $+47.5\%$  edges back to physical parameters  $\nabla$

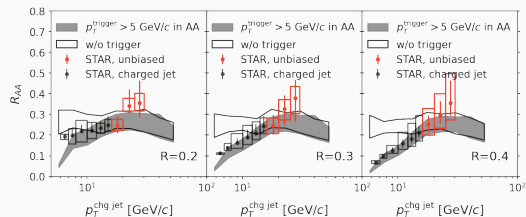
#	$\mu_{\text{med}}$	$Q_0^{\text{LHC}}$ [GeV]	$Q_0^{\text{RHIC}}$ [GeV]	$T_f$ [GeV]
0	1.32	1.81	0.97	0.16
1	1.90	1.64	0.83	0.153
...	...	...	...	...

# Test the transport of energy: cone-size dependence of jet $R_{AA}$



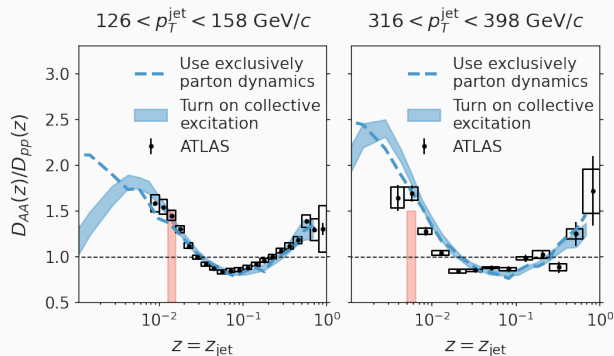
[CMS: JHEP 05, 284(2021)]

- LHC: LIDO predicts  $R_{AA}$  increased by 10% from  $R = 0.2$  to  $R = 1.0$  at  $p_T^{\text{jet}} = 500$  GeV.
- RHIC: Weak  $R$ -dependence in the **unbiased** region.  
Triggering bias well understood from simulation.



$\Delta$  Unbiased region (red) are in sensitive to the high- $p_T$  hadron trigger. The triggering bias is also understood from the simulation. [STAR: PRC 102, 054913(2020)]

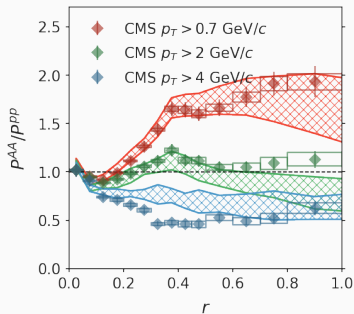
# Test the transport of energy: fragmentation function



[ATLAS: PRC 98, 024908(2018)]

- Calculations that treats everything with partonic dynamics well describes the fragmentation at  $zp_T^{\text{jet}} > 2$  GeV (red bands).
- Use collective excitations to redistribution soft particles improves at  $p_T \lesssim 2$  GeV.

# Test the transport of energy: a detailed look at low- $p_T$ particles



[CMS: JHEP05, 006(2018)]

◁ Jet shape with different minimum hadron  $p_T$

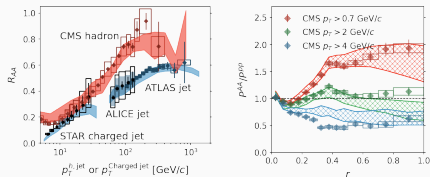
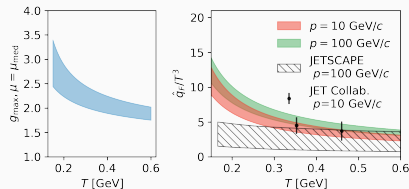
- Energy is shifted to particles at lower  $p_T$  and larger  $r$ .
- Discrepancy appears within the cone for  $p_{T,cut} = 4$  GeV
  - Can this be fixed by fine-tuning of parameters?
  - Suggest missing physics? Such as coalescence shifting intermediate- $p_T$  hadrons to higher  $p_T$ .



# Summary and outlook

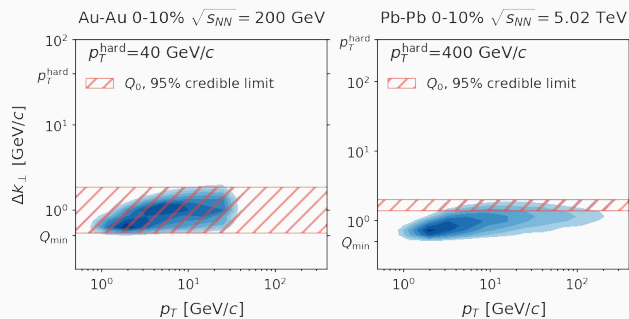
- Jet and hadron quenching are sensitive to different aspects of parton dynamics.
  - Hadron: total energy loss.
  - Jet: redistribution of the lost energy. Require a modeling of collective excitation.
  - Jet  $R_{AA}$ : less dependent on the separation scale between vacuum-like evolution & transport equation.

- Extract  $g_s/\hat{q}$  from jet ( $R = 0.4$ ) and hadron ( $h/D$ )  $R_{AA}$  at RHIC and LHC central AA collisions.  $\Rightarrow$
- The resulting jet cone-size dependence is weak. Undershoot CMS data; consistent with STAR data.
- The redistribution of low momentum particles around the jet tested with fragmentation function and jet shape.



Questions?

# A consistency check: compare $Q_0^{\text{RHIC}}$ and $Q_0^{\text{LHC}}$ to the medium $k_{\perp}$



- Compare the radiative  $k_{\perp}$  distribution with the separation scale  $Q_0$ .
- At LHC: most in-medium activity happens below  $Q_0^{\text{LHC}}$ .
- At the RHIC,  $Q_0^{\text{RHIC}}$  is comparable to typical  $k_{\perp}$ .