

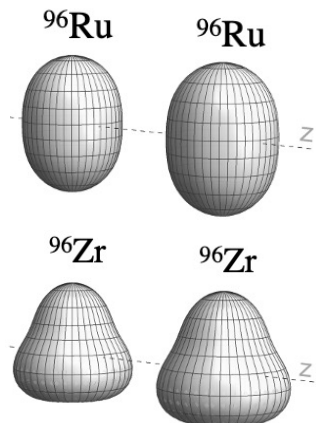
Scaling approach to nuclear structure in high-energy heavy-ion collisions

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Poster Session 2 T14_2, April 6, 2022

[Based on the: arXiv:2111.15559](https://arxiv.org/abs/2111.15559)

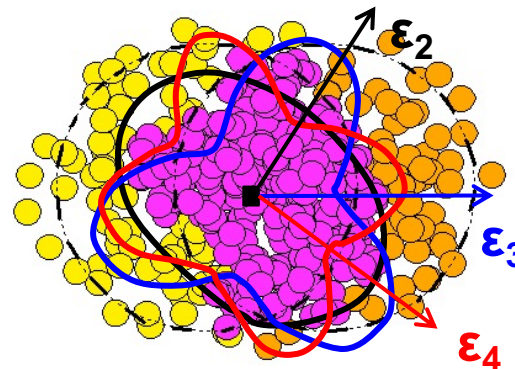
Nuclear Structure



$$\rho(r, \theta, \phi) = \frac{\rho_0}{1 + e^{(r-R_0(1+\sum_n \beta_n Y_n^0(\theta, \phi)))/a_0}}$$

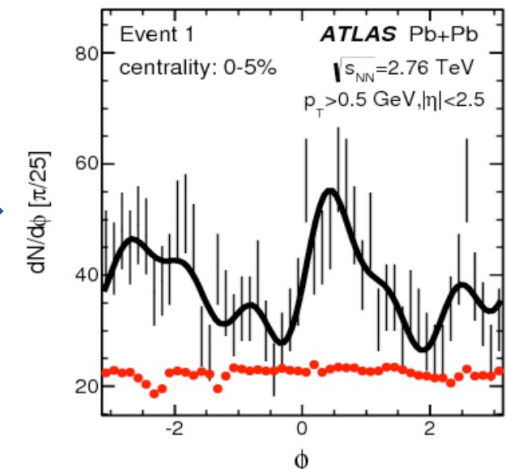
J. Jia, PRC105, 014905(2022)

Initial state

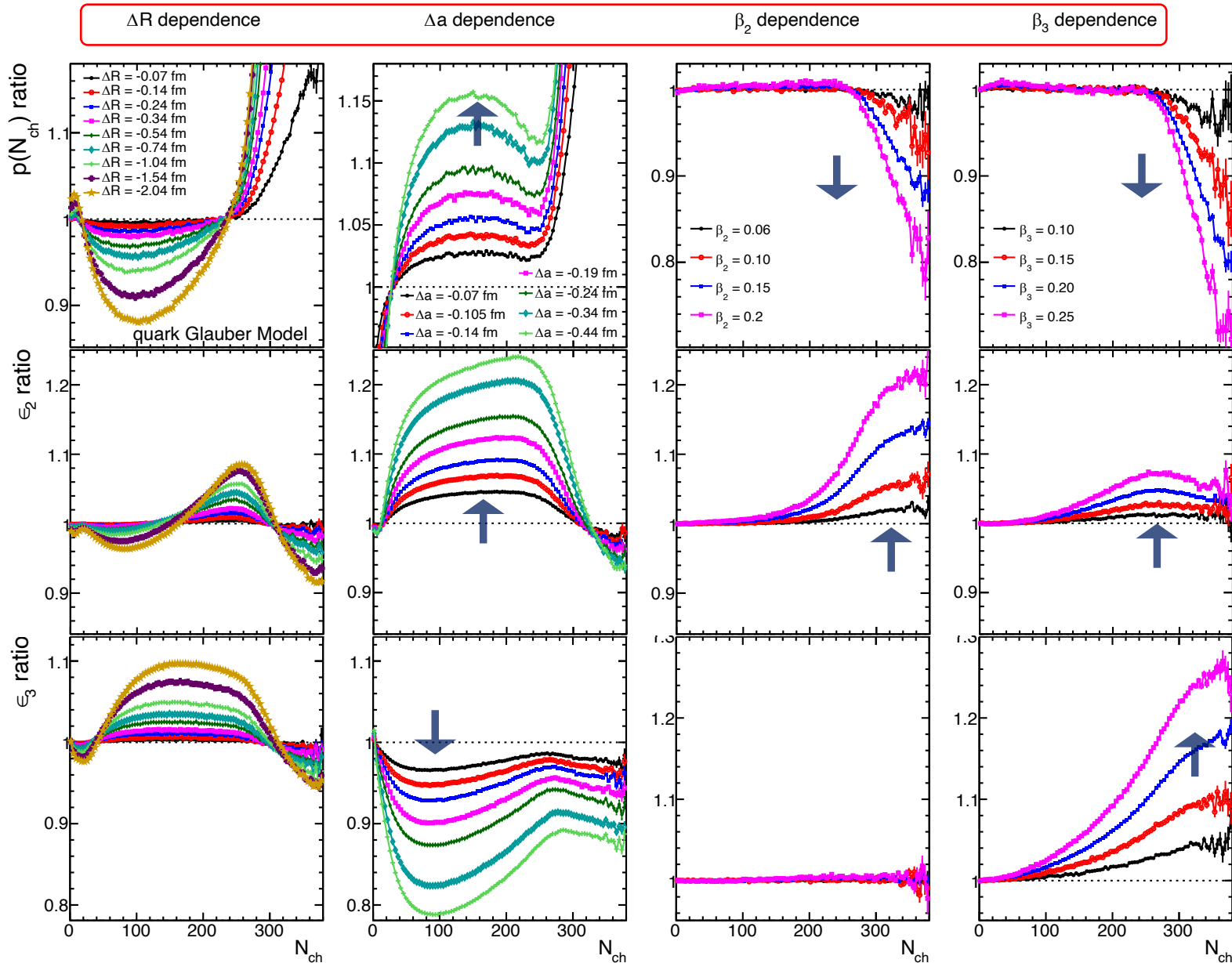


Hydrodynamic

Final state



Nuclear structure effect on the initial state



Species	β_2	β_3	a_0	R_0
Ru	0.162	0	0.46 fm	5.09 fm
Zr	0.06	0.20	0.52 fm	5.02 fm
difference	$\Delta\beta_2^2$	$\Delta\beta_3^2$	Δa_0	ΔR_0
	0.0226	-0.04	-0.06 fm	0.07 fm

R_0 has some effect with unexpected change

$a_0 \uparrow$ in mid-central

$\beta_2 \downarrow$ in central

$\beta_3 \downarrow$ in central

R_0 has some effect with unexpected change

$a_0 \uparrow$ ϵ_2 in mid-central

$\beta_2 \uparrow$ ϵ_2 in from mid-central to central

$\beta_3 \uparrow$ ϵ_2 in mid-central

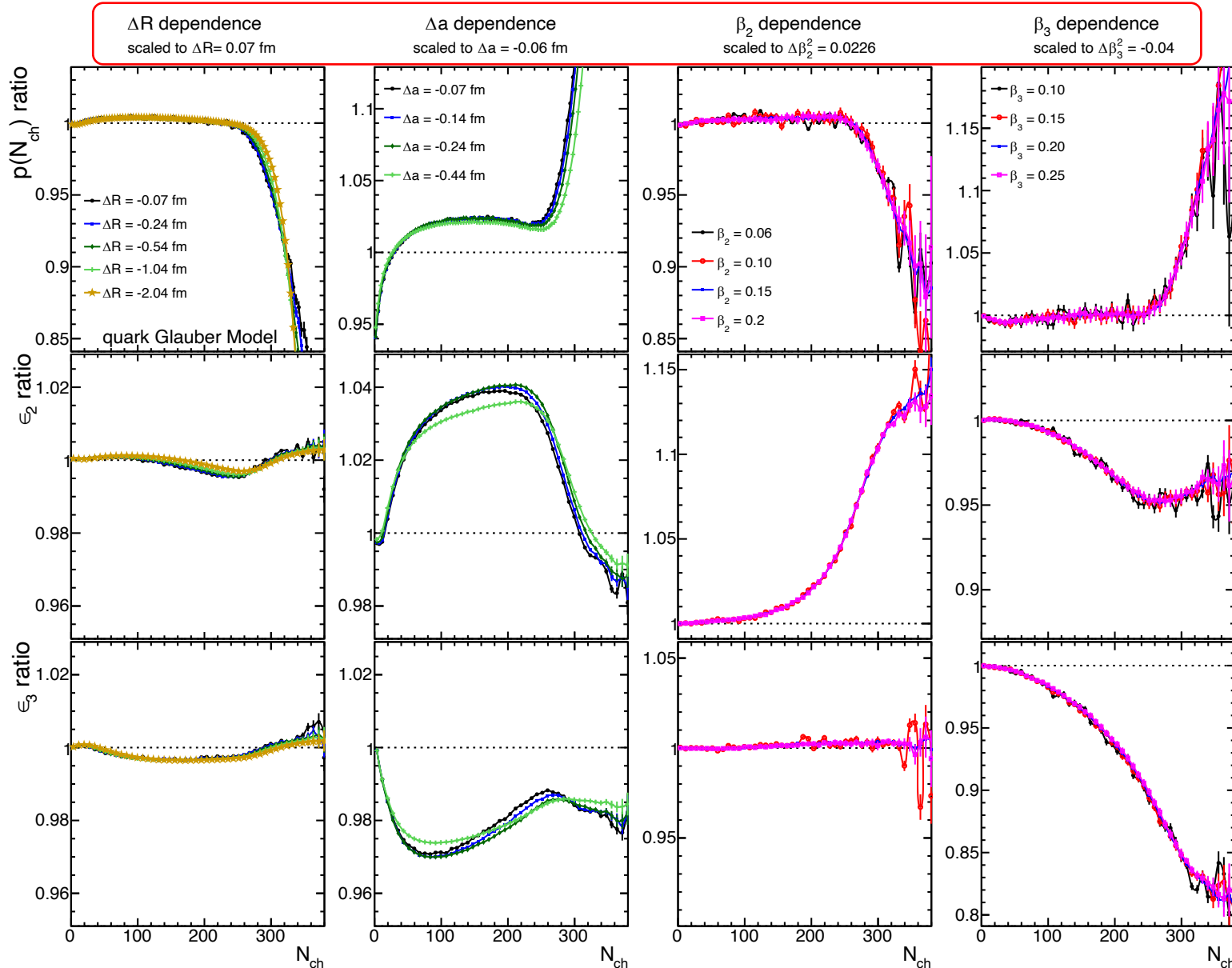
R_0 has some effect with unexpected change

$a_0 \downarrow$ ϵ_3 in mid-central

β_2 has no effect on ϵ_3

$\beta_3 \uparrow$ ϵ_3 in mid-central

Scaling approach to nuclear structure on the initial state



Is it linear?

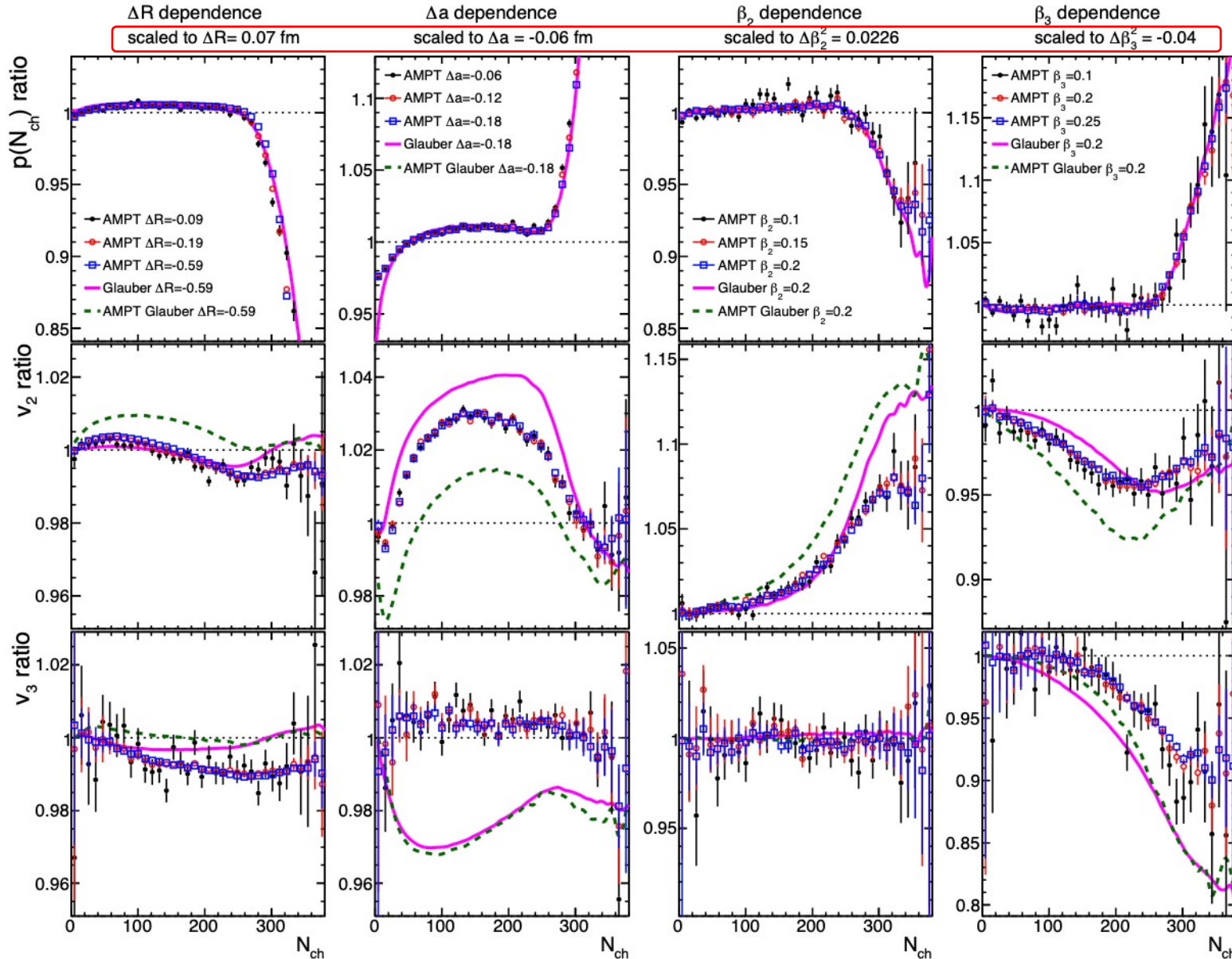


Yes, only scale a constant value

nearly perfect scaling over the wide range of parameter values

$$R_{\mathcal{O}} \equiv \frac{\mathcal{O}_{Ru}}{\mathcal{O}_{Zr}} \approx 1 + c_1 \Delta\beta_2^2 + c_2 \Delta\beta_3^2 + c_3 \Delta R_0 + c_4 \Delta a$$

Scaling approach to nuclear structure on the final state



Species	β_2	β_3	a_0	R_0
Ru	0.162	0	0.46 fm	5.09 fm
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difference	$\Delta \beta_2^2$	$\Delta \beta_3^2$	Δa_0	ΔR_0
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nearly perfect scaling over the wide range of parameter values

c_n can be determined more precisely by using a larger change of these parameters

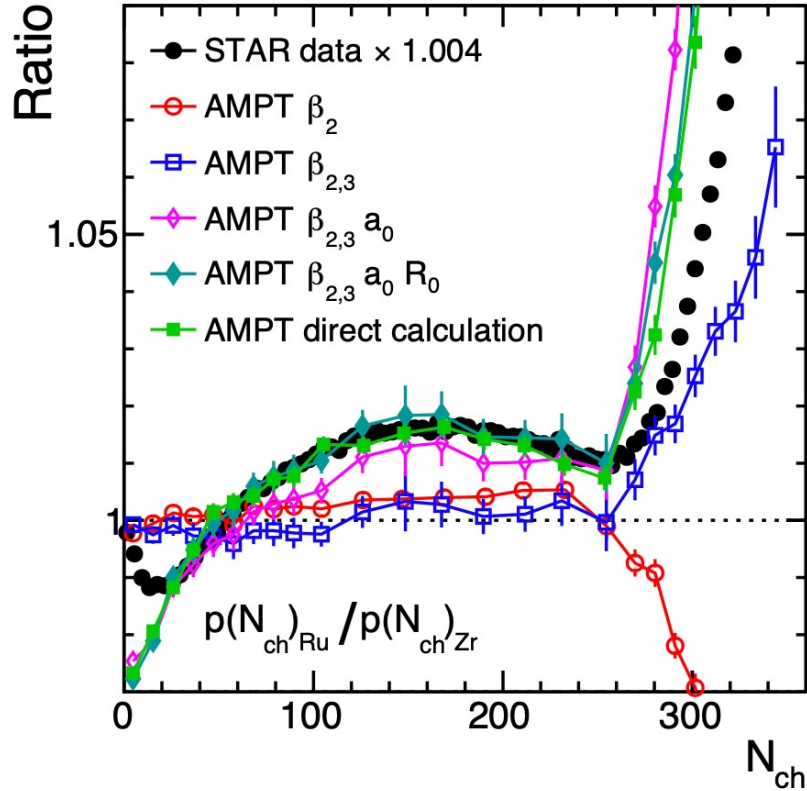
Verifies the relation:

$$\mathcal{O} \approx b_0 + b_1 \beta_2^2 + b_2 \beta_3^2 + b_3 (R_0 - R_{0,\text{ref}}) + b_4 (a - a_{\text{ref}})$$

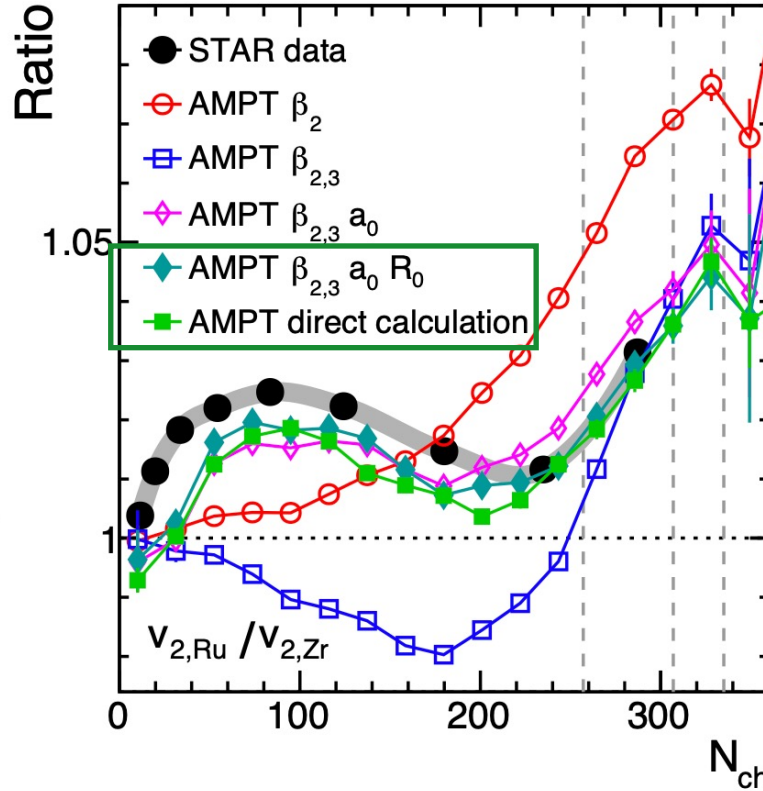
$$R_{\mathcal{O}} \equiv \frac{\mathcal{O}_{\text{Ru}}}{\mathcal{O}_{\text{Zr}}} \approx 1 + c_1 \Delta \beta_2^2 + c_2 \Delta \beta_3^2 + c_3 \Delta R_0 + c_4 \Delta a$$

Nuclear structure via v_n ratio

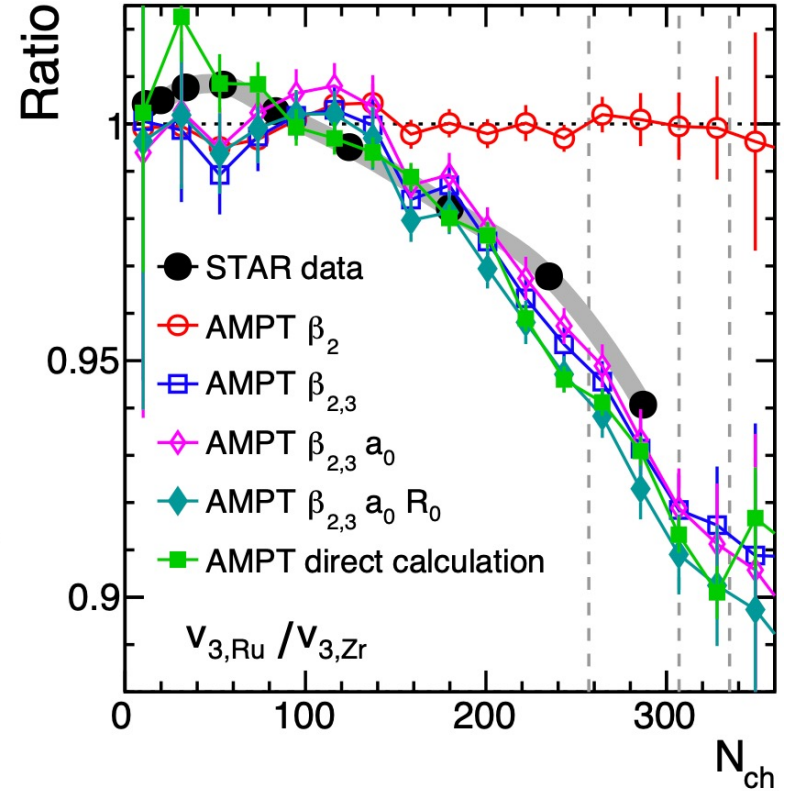
STAR, PRC105, 014901(2022)



5% 1% 0.2%



5% 1% 0.2%



Heavy-ion expectation:

$$v_2^2 = a_2 + b_2 \beta_2^2 + b_{2,3} \beta_3^2, \quad v_3^2 = a_3 + b_3 \beta_3^2$$

$$\frac{v_{2,Ru}^2}{v_{2,Zr}^2} \approx 1 + \frac{b_2}{a_2} (\beta_{2,Ru}^2 - \beta_{2,Zr}^2) - \frac{b_{2,3}}{a_2} \beta_{3,Zr}^2$$

$$\frac{v_{3,Ru}^2}{v_{3,Zr}^2} \approx 1 - \frac{b_3}{a_3} \beta_{3,Zr}^2 < 1$$

Cancelation expected in non-central collisions

- 1) v_2 ratio: large $\beta_{2,Ru}$, negative contribution from $\beta_{3,Zr} \Rightarrow$ Sharper increase in central
- 2) v_3 ratio: strong decrease from $\beta_{3,Zr}$ with negligible $\beta_{2,Ru}$ distortion
- 3) Residual effect due to radial structure, e.g., neutron skin in Zr
- 4) No significant effect due to nuclear size

✓ Direct calculations are same as initial whole $\beta_2, \beta_3, a_2, R_0$ input.

A direct algebra linked to neutron skin

Using relation for WS: $R^2 \equiv \langle r^2 \rangle \approx \left(\frac{3}{5} R_0^2 + \frac{7}{5} \pi^2 a^2 \right) / \left(1 + \frac{5}{4\pi^2} \sum_n \beta_n^2 \right)$

Neutron skin expressed by **R** and **a** parameters for **nucleons** and **protons**:

$$\Delta r_{np} \approx \frac{R^2 - R_p^2}{R(\delta + 1)} \approx \frac{3(R_0^2 - R_{0,p}^2) + 7\pi^2(a^2 - a_p^2)}{\sqrt{15}R_0 \sqrt{1 + \frac{7\pi^2}{3} \frac{a^2}{R_0^2} \left(1 + \delta + \frac{5}{8\pi^2} \sum_n \beta_n^2 \right)}} \quad \delta = (N - Z)/A$$

The difference between two isobar systems can be expressed as:

$$\Delta(\Delta r_{np}) = \Delta r_{np,1} - \Delta r_{np,2} \approx \frac{\Delta Y - \frac{7\pi^2}{3} \frac{\bar{a}^2}{R_0^2} \left(\frac{\Delta Y}{2} + \bar{Y} \left(\frac{\Delta a}{\bar{a}} - \frac{\Delta R_0}{\bar{R}_0} \right) \right)}{\sqrt{15}\bar{R}_0 \left(1 + \bar{\delta} + \frac{5}{8\pi^2} \sum_n \bar{\beta}_n^2 \right)}$$

where $Y \equiv 3(R_0^2 - R_{0,p}^2) + 7\pi^2(a^2 - a_p^2)$ $\Delta x = x_1 - x_2$ $\bar{x} = (x_1 + x_2)/2$

Can obtain skin diff. from ΔR_0 Δa for nucleons and known ΔR_0 Δa for protons

Table from: H.J. Xu et. al., PLB819, 1136453(2021)

	⁹⁶ Ru		⁹⁶ Zr	
	R	a	R	a
p	5.060	0.493	4.915	0.521
n	5.075	0.505	5.015	0.574
p+n	5.067	0.500	4.965	0.556

Direct calc.: $\Delta(\Delta r_{np}) = 0.0296 \text{ fm} - 0.1606 \text{ fm} = -0.1310 \text{ fm}$

Formula: $\Delta(\Delta r_{np}) = -0.1319 \text{ fm}$

<1% difference

Conclusions and Outlooks

1) **Demonstration: the nuclear structure effect on bulk observables.**

2) **AMPT could describe the STAR published data quantitatively.**

3) **A new approach to constrain the collective nuclear structure parameters:**

✓ The final state bulk observables v_2 , v_3 and $p(N_{\text{ch}})$ follow a simple dependences on the variation of parameters:

$$\mathcal{O} \approx b_0 + b_1\beta_2^2 + b_2\beta_3^2 + b_3(R_0 - R_{0,\text{ref}}) + b_4(a - a_{\text{ref}}) \quad R_{\mathcal{O}} \equiv \frac{\mathcal{O}_{\text{Ru}}}{\mathcal{O}_{\text{Zr}}} \approx 1 + c_1\Delta\beta_2^2 + c_2\Delta\beta_3^2 + c_3\Delta R_0 + c_4\Delta a$$

✓ The c_n can be determined precisely in a given model.


✓ The data-model comparison can precisely constrain the dependence of nuclear parameters:

$$\Delta\beta_2^2 = \beta_{2,\text{Ru}}^2 - \beta_{2,\text{Zr}}^2 \quad \Delta\beta_3^2 = \beta_{3,\text{Ru}}^2 - \beta_{3,\text{Zr}}^2 \quad \Delta R_0 = R_{0,\text{Ru}} - R_{0,\text{Zr}} \quad \Delta a = a_{\text{Ru}} - a_{\text{Zr}}$$

✓ Achieve to obtain the difference of neutron skin between two isobar systems:

$$\Delta(\Delta r_{np}) = \Delta r_{np,1} - \Delta r_{np,2} \approx \frac{\Delta Y - \frac{7\pi^2}{3} \frac{\bar{a}^2}{R_0^2} \left(\frac{\Delta Y}{2} + \bar{Y} \left(\frac{\Delta a}{\bar{a}} - \frac{\Delta R_0}{\bar{R}_0} \right) \right)}{\sqrt{15}\bar{R}_0 \left(1 + \bar{\delta} + \frac{5}{8\pi^2} \sum_n \bar{\beta}_n^2 \right)}$$

4) **Unique opportunities by relativistic collisions of isobars as a tool to study nuclear structure.**

Thank you @  QUARK MATTER
KRAKÓW
2022