

Probing electromagnetic fields in uRHICs with heavy quarks and leptons from Z^0 decay

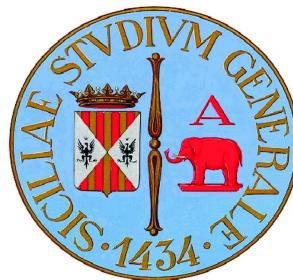
Yifeng Sun

sunyfphy@sjtu.edu.cn

Shanghai Jiao Tong University

Collaborators:

S. Plumari, V. Greco and X.N. Wang

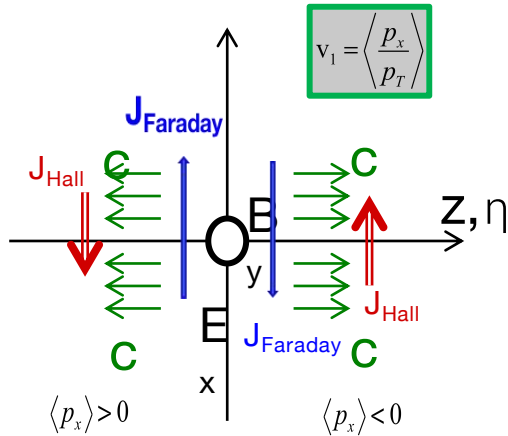


Refs: PLB 816, 136271 (2021); EPJP 136, 726 (2021); PLB 827, 136962 (2022)

Electromagnetic (e.m.) fields in uRHICs

- ❑ **Induces:** CME, CMW, Hyperons polarization splitting
- ❑ **Observables:** v_1 splitting of charged particles
- ❑ **Good Probes:** Heavy Quarks (HQs)

D.E. Kharzeev et al., NPA 803 (2008)
 Y. Burnier et al., PRL 107 (2011), 052303
 STAR, Nature 548 (2017), 62-65
 U. Gürsoy et al., PRC 89 (2014), 054905
 S.K. Das et al., PLB 768 (2017), 260-264



❖ HQs best probe for v_1 induced by e.m. fields:

1. pQCD hard processes
2. negligible thermal production
3. $t_{\text{form}} \approx 0.08 \text{ fm}/c$ when B_y is \approx its maximum and witness of all the QGP evolution
4. $T_{\text{th}}(C) \approx T_{\text{QGP}} \gg T_{\text{e.m}}$ (keep more memory effects)

❖ Delicate balance between E and B

- ✓ E wins \rightarrow negative slope Δv_1 vs y_z between positively and negatively charged particles
- ✓ B wins \rightarrow positive slope Δv_1 vs y_z

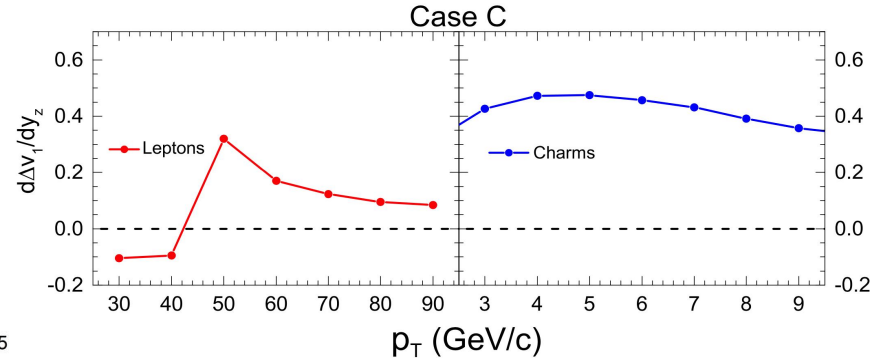
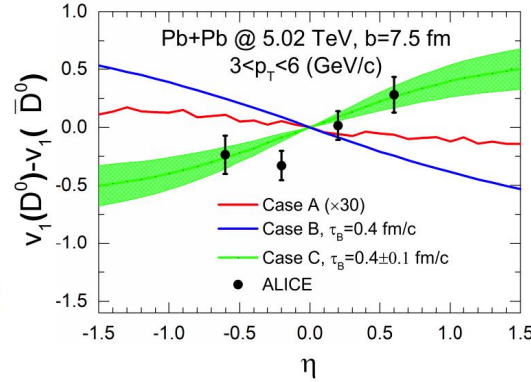
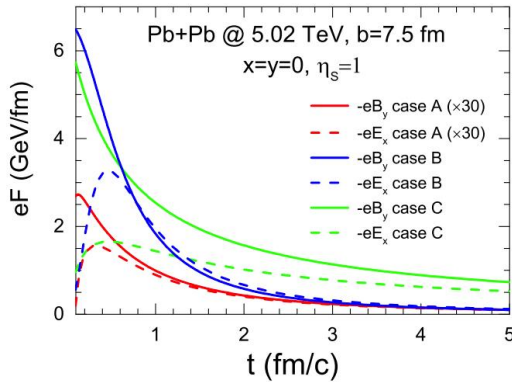
E.M. fields on HQs and leptons from Z^0 decay

❑ A slow decay e.m. fields (Case C) reproduces ALICE data

Sun&Plumari&Greco, PLB 816 (2021), 136271

❑ V_1 splitting of leptons from Z^0 decay has a peculiar pattern

❑ Correlated measurement of c quarks and leptons from Z^0 decay a strong probe of e.m.



Case A

- ✓ E-B fields like U. Gursoy et al., PRC 89 (2014), 054905
- ✓ Medium at $t < 0$ + eq. medium $\sigma_{el} = 0.023 \text{ fm}^{-1}$

Case B and C

$$eB_y(x, y, \tau) = -B(\tau)\rho_B(x, y) \quad B(\tau) = eB_0/(1 + (\tau/\tau_B)^a)$$

$$B(\tau) = eB_0/(1 + \tau^2/\tau_B^2)$$

$$B(\tau) = eB_0/(1 + \tau/\tau_B)$$

- ✓ eB_0 fixed by the value $t=0$ in vacuum; a and τ_B can be tuned
- ✓ E_x is evaluated by the Faraday's Law

$$\nabla \times \mathbf{E} = -\partial \mathbf{B} / \partial t$$

❖ Why leptons from Z^0 decay

1. Clearer observables
2. Separable from other sources
3. $\tau_{\text{decay}}(Z^0) = \tau_{\text{form}}(\text{charm}) = 0.08 \text{ fm/c} \rightarrow$ Strong correlation between $\Delta v_1(D^0, \bar{D}^0)$ and $\Delta v_1(l^+, l^-)$

Signatures of charge dependent flows by e.m. fields

❑ E.M. fields modify charge dependent spectra and flows

Sun&Greco &Plumari, EPJP

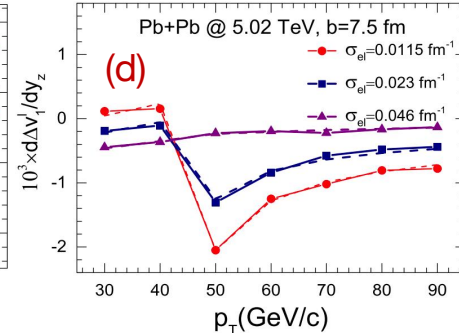
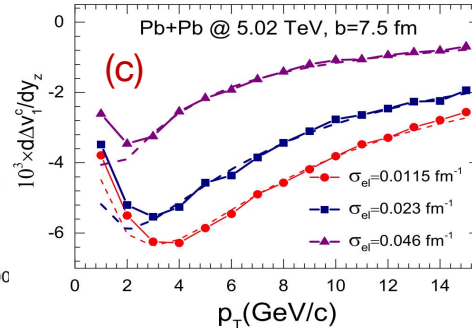
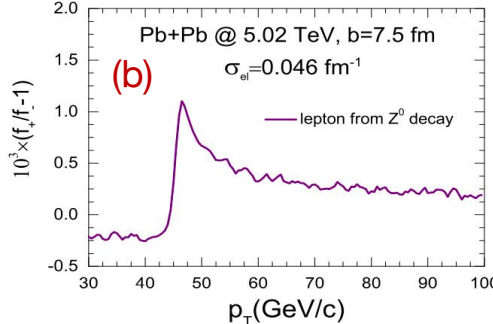
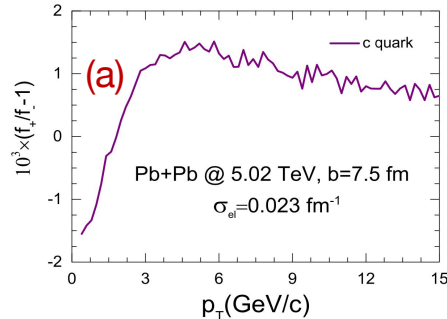
136 (2021), 726

❑ Can be described by simple forms

❖ From the p_T dependence of $d\Delta v_1/dy$ of leptons and charm quarks

- ✓ a ratio for Case A with $\sigma_{el}=0.0115, 0.023 \text{ fm}^{-1}$ are $8.7 \text{ MeV}/6.3 \text{ MeV}=1.4$ and $4.7 \text{ MeV}/3.6 \text{ MeV}=1.3$
- ✓ for Case C is $1.5 \text{ GeV}/0.75 \text{ GeV}=2$
- ✓ close to charge ratio 1.5 for different e.m. fields: 100 times

$$\begin{aligned}
 f' = & f - \left\{ \frac{\partial f(a_1 + b_1)}{\partial p_T} + f \left(-\frac{p_T}{m_T^2} \frac{\partial(a_1 + b_1) \tanh y_z}{\partial y_z} \right. \right. \\
 & + \left. \frac{a_1 + b_1}{p_T} + \frac{2}{m_T} \frac{\partial c_0 / \cosh y_z}{\partial y_z} \right\} \\
 & - \left\{ -f \frac{p_T}{m_T^2} \frac{\partial(a_0 + b_0) \tanh y_z}{\partial y_z} + \frac{\partial(a_0 + b_0) f}{\partial p_T} \right\} \cos \phi \\
 & - \sum_{n=1} \left\{ \frac{\partial f(a_{n+1} + b_{n+1} + a_{n-1} - b_{n-1})}{\partial p_T} \right. \\
 & + \left. f \left[\frac{(n+1)(a_{n+1} + b_{n+1}) - (n-1)(a_{n-1} - b_{n-1})}{p_T} \right. \right. \\
 & - \left. \frac{p_T}{m_T^2} \frac{\partial \tanh y_z (a_{n+1} + b_{n+1} + a_{n-1} - b_{n-1})}{\partial y_z} \right. \\
 & \left. \left. + \frac{2}{m_T} \frac{\partial c_n / \cosh y_z}{\partial y_z} \right] \right\} \cos n\phi. \quad (7)
 \end{aligned}$$



$$f' |_{y_z=0} = f \left[1 - (a_1 + b_1) \frac{\partial \ln f}{\partial p_T} - \frac{2}{p_T} \frac{\partial c_0}{\partial y_z} \right] |_{y_z=0}$$

$$\begin{aligned}
 \frac{d\Delta v_1^c}{dy_z} |_{y_z=0} &= \frac{d\Delta a_0}{dy_z} |_{y_z=0} \left(-\frac{\partial \ln f_c}{\partial p_T} + \frac{2p_T}{m_T^2} \right) - \frac{1}{m_T} \left(\frac{d^2 c_1}{dy_z^2} - c_1 \right) \\
 &= -\alpha \frac{\partial \ln f_c}{\partial p_T} + (2\alpha - \beta) \frac{p_T}{m_T^2}
 \end{aligned}$$

$$\begin{aligned}
 \frac{d\Delta a_0}{dy_z} |_{y_z=0} &= |q|K [\tau_1 g(\tau_1, 0) - \tau_0 g(\tau_0, 0)] \\
 &\simeq -|q|K [\tau_1 B_y(\tau_1, 0) - \tau_0 B_y(\tau_0, 0)].
 \end{aligned}$$

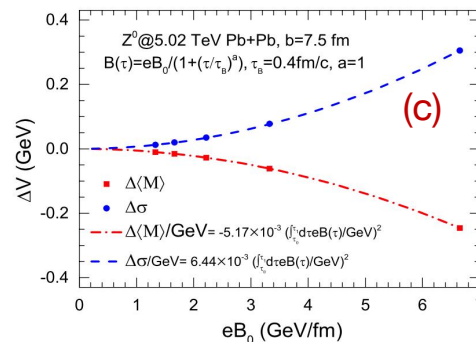
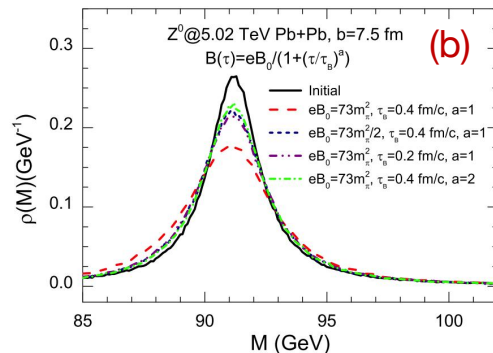
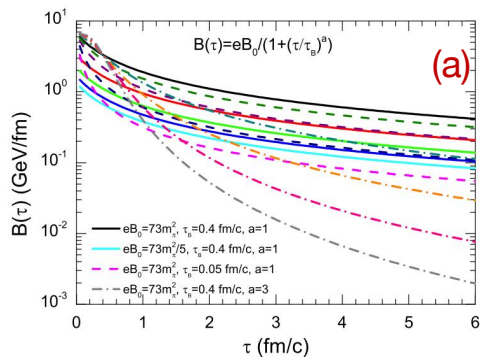
E.m. fields on Z^0 leptonic invariant mass

❑ E.M. fields decreases Z^0 leptonic invariants and increase width

❑ Changes of Z^0 leptonic invariant mass and its width

✓ Depends on the integral of B_y quadratically (approximate)

Sun&Greco&Wang, PLB
827, 136962 (2022)

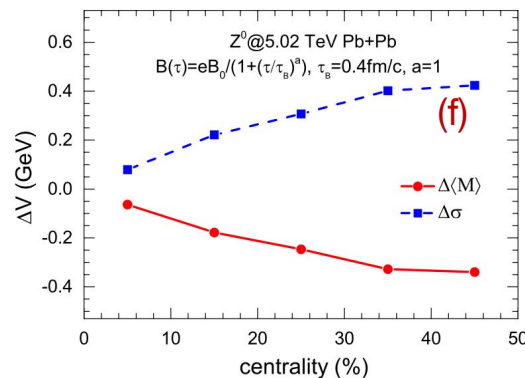
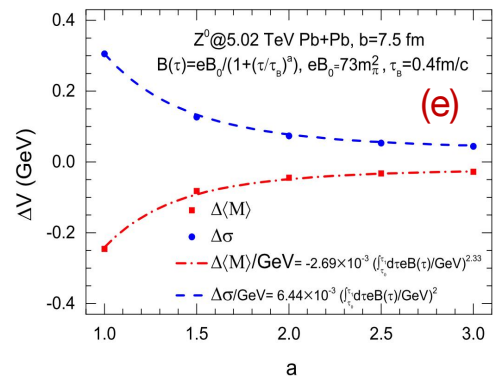
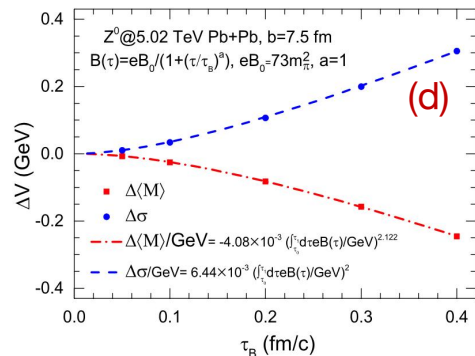


$$eB_y(x, y, \tau) = -B(\tau)\rho_B(x, y)$$

$$\rho_B(x, y) = \exp\left[-\frac{x^2}{2\sigma_x^2} - \frac{y^2}{2\sigma_y^2}\right]$$

$$B(\tau) = eB_0/(1 + (\tau/\tau_B)^a),$$

❖ Small effect of lepton-quarks scattering



$$\Delta\langle M \rangle = -1.9 \text{ MeV}$$

$$\Delta\sigma \leq 0.2 \text{ MeV}$$