



# Nuclear modification of hard scattering processes in small systems at PHENIX

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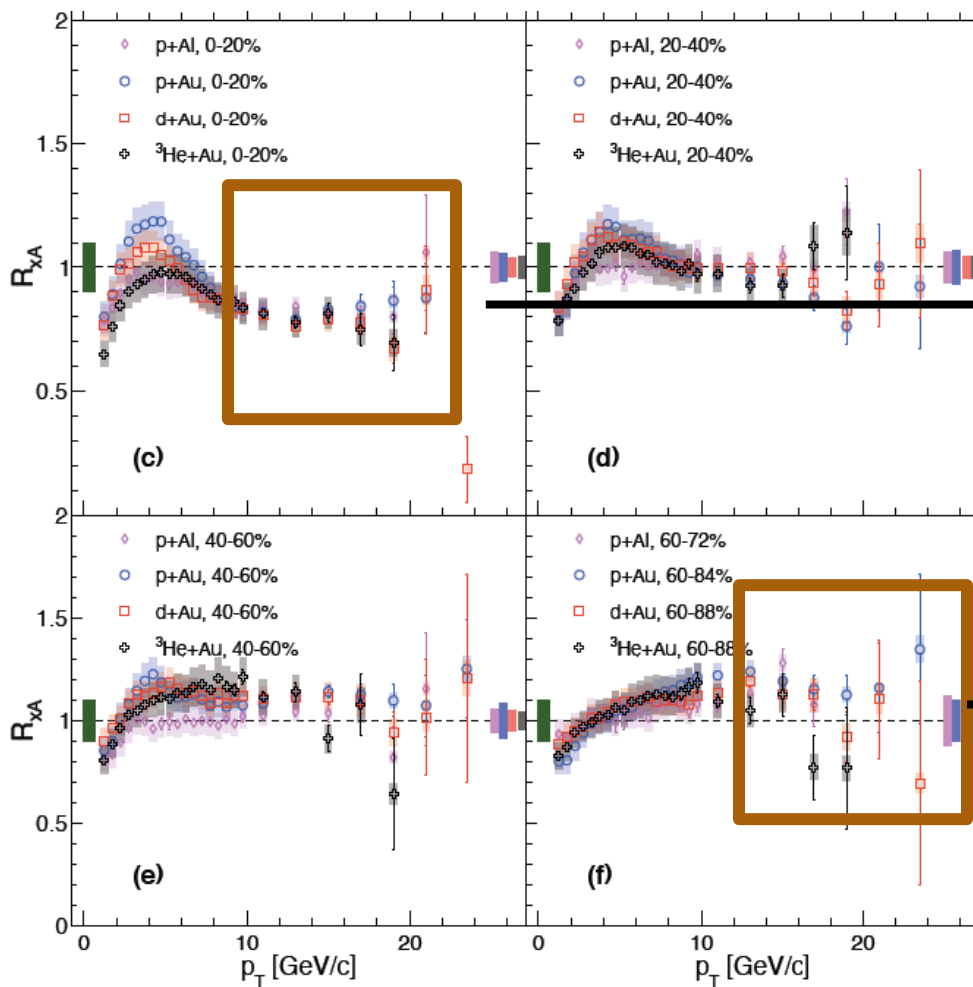
NIVEDITHA RAM

07/04/2022

QUARK MATTER 2022

# What to expect in this talk

Using direct photons at high  $p_T$  to measure the number of binary collisions ( $N_{\text{coll}}$ ) in a system



Observed suppression in central collisions remains even after correcting for bias in  $N_{\text{coll}}$  determination

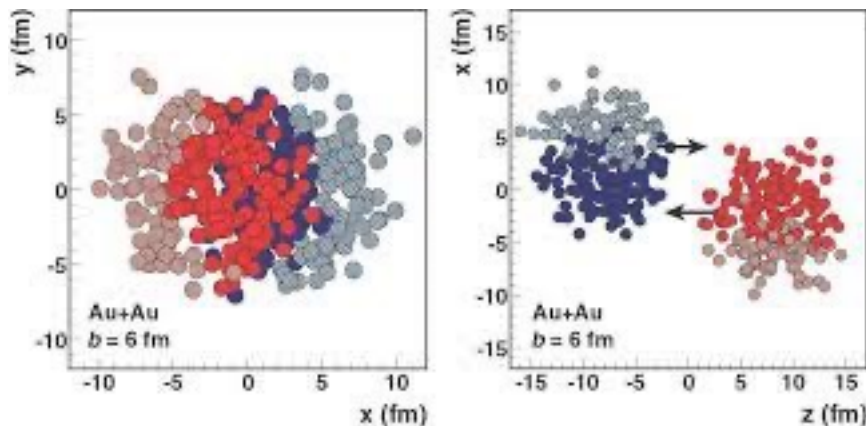
Observed enhancement in peripheral collisions is an artifact of bias in the  $N_{\text{coll}}$  determination using Glauber model for small systems

Definition

# Nuclear Modification Factor

$$R_{AB}(p_T) = \frac{\left(\frac{d^2 N}{dp_T d\eta}\right)_{AB}}{\langle N_{coll} \rangle_{AB} * \left(\frac{d^2 N}{dp_T d\eta}\right)_{pp}} = \frac{Y(AB) \text{ d+Au}}{\langle N_{coll} \rangle_{AB} * Y(pp) \text{ p+p}}$$

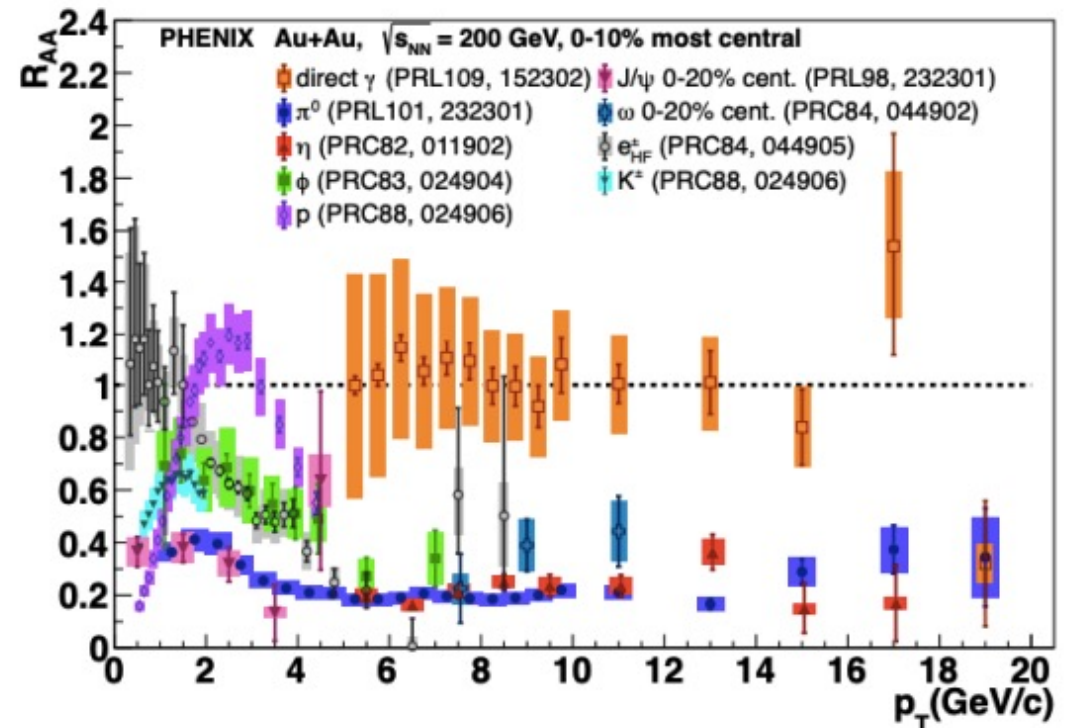
How is heavy ion collision different from a scaled p+p collision



$R_{AB}(p_T) < 1$  is a signature of QGP

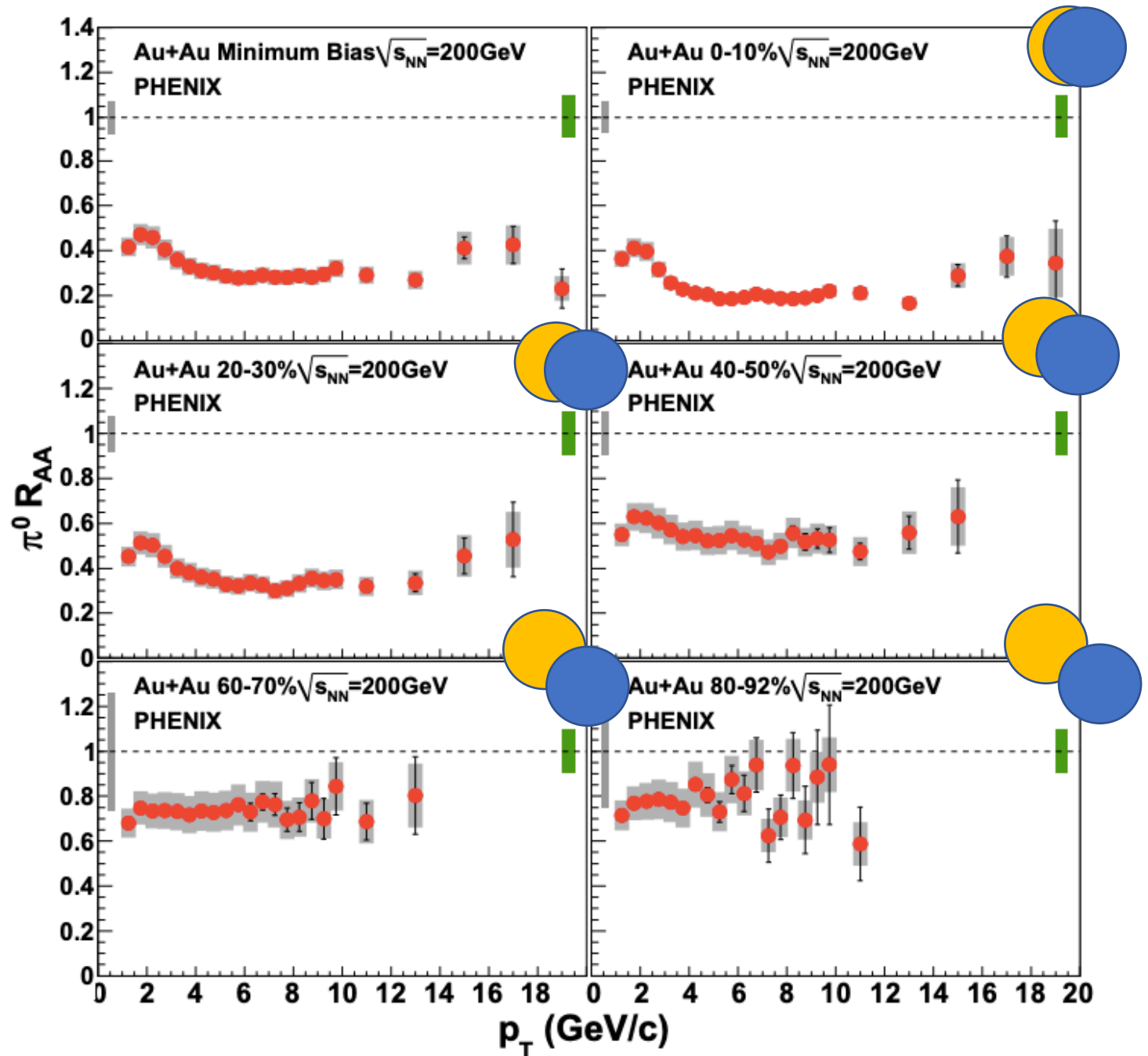
1)  $\pi^0, \eta, \phi, J/\psi, \omega$  interact with the QGP  $\rightarrow R_{AB}(p_T)$  is suppressed

2) Direct photon is transparent to the QGP  $\rightarrow R_{AB}(p_T)$  is unity



# Centrality binned $R_{AA}$ of $\pi^0$ in Au+Au collisions

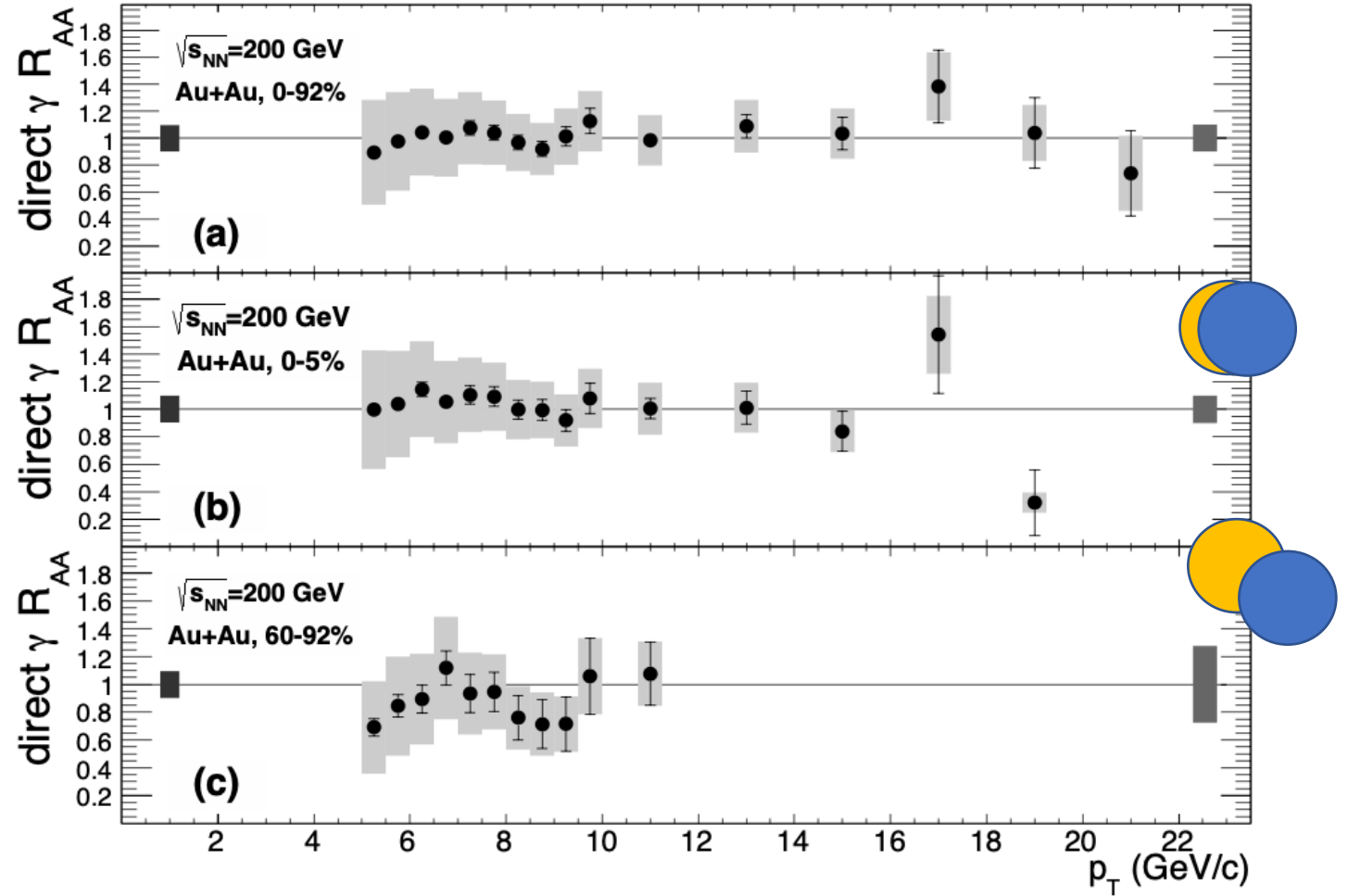
- Most central collisions show the most suppression.
- Vanishing suppression at peripheral collisions
- Intuitive trend in centrality for a system with QGP creation



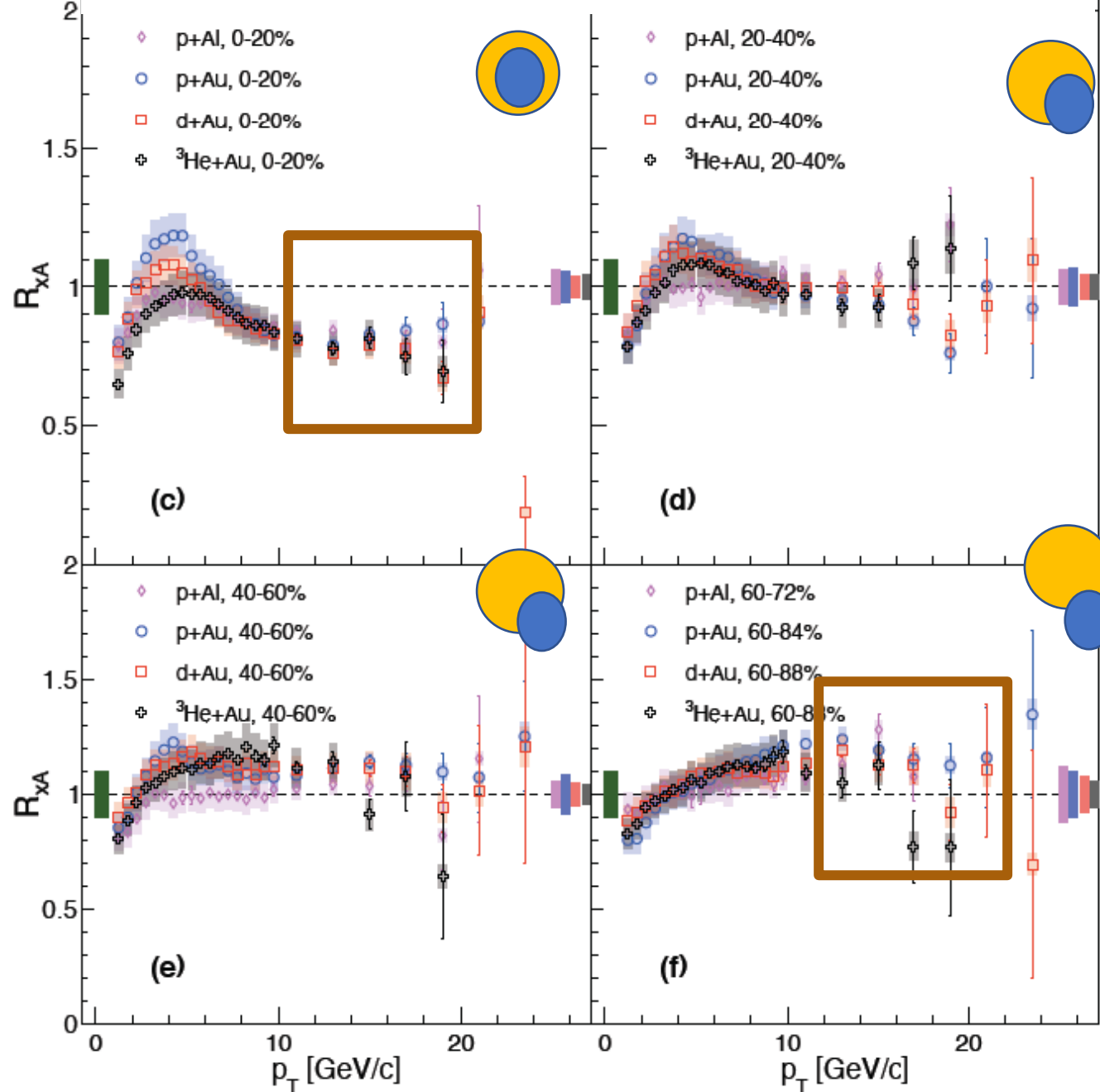
# Centrality binned $R_{AA}$ of direct $\gamma$ in Au+Au collisions

- Unity at all centralities.
- Direct photons are transparent to QGP
- Using high  $p_T$  direct photons, we obtain :

$$N_{coll}^{exp} = \frac{Yield_{AA}^{\gamma}}{Yield_{pp}^{\gamma}}$$



# Puzzling behavior of $R_{AA}$ in small systems



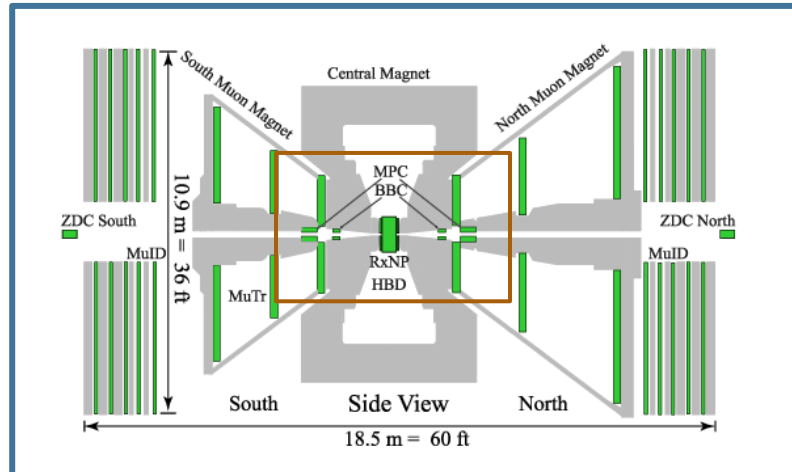
- Suppression in central collisions  $\rightarrow$  Formation of QGP droplets?
- Enhancement in peripheral collisions?

Is the determination of  $N_{\text{coll}}$  in different event classes biased

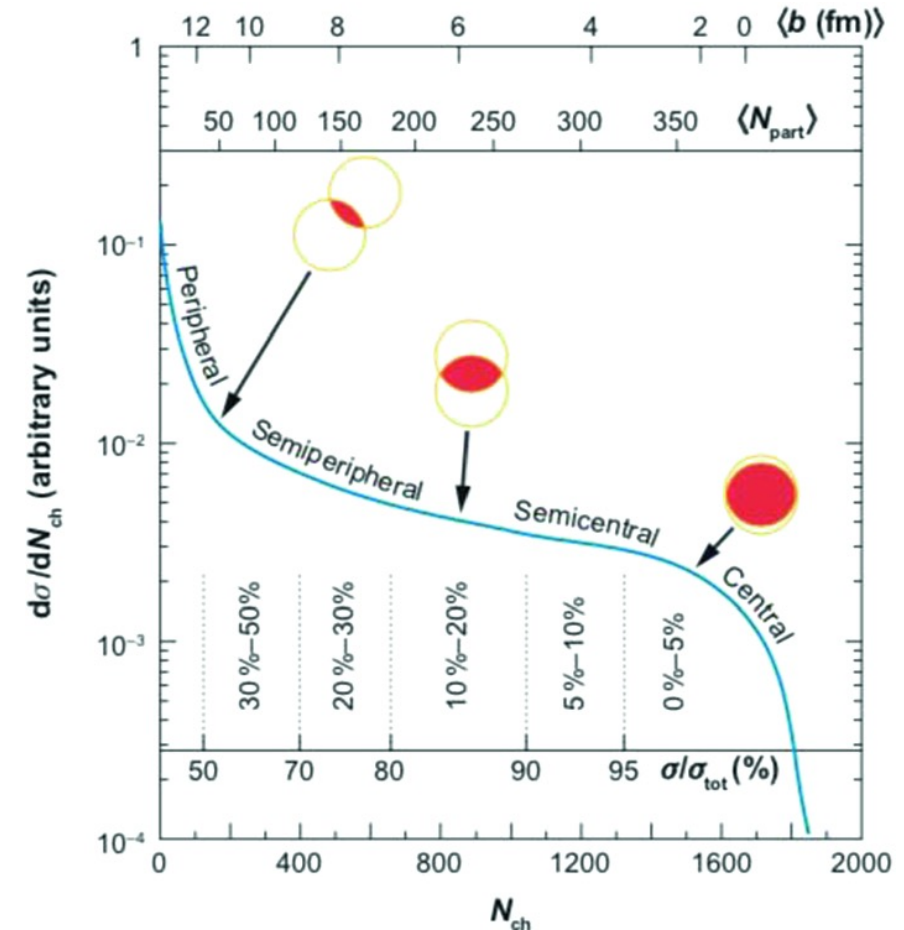
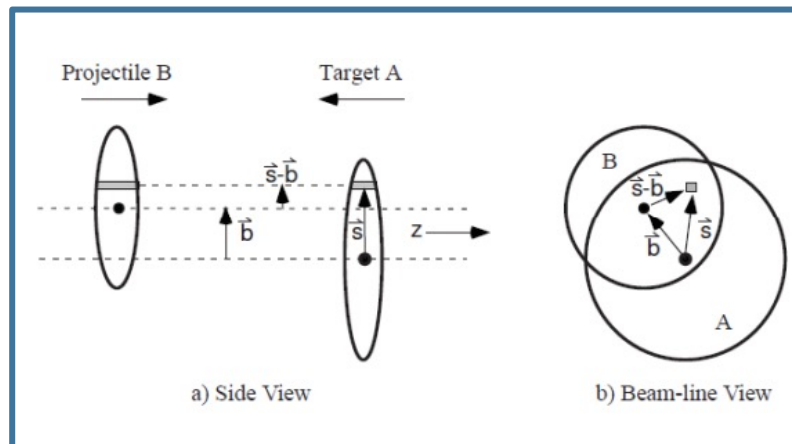
# Measurement of average number of binary collisions from bulk observables

$$\langle N_{coll} \rangle_{AB} * Y(pp)$$

Number of charged particle from BBC  
 $3.1 < |\eta| < 4$

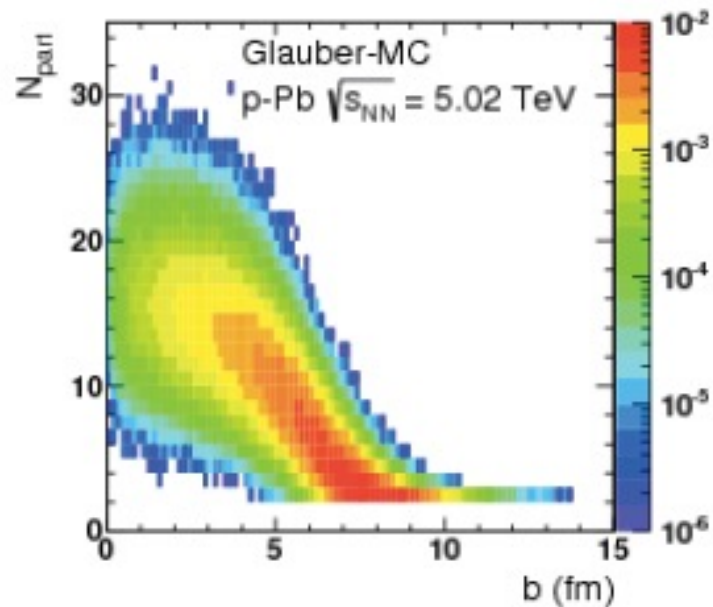
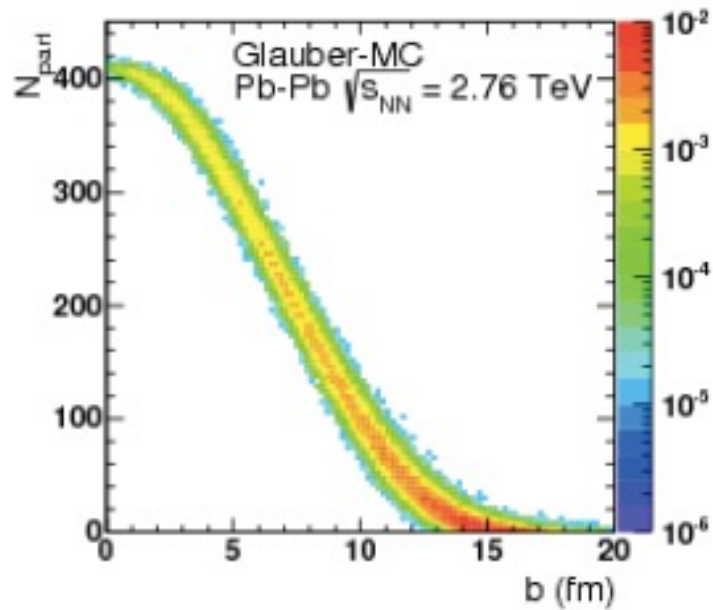


Standard Glauber model gives mapping of charged particle in forward region to number of binary collisions of the event.





# Is Glauber model valid for small systems?

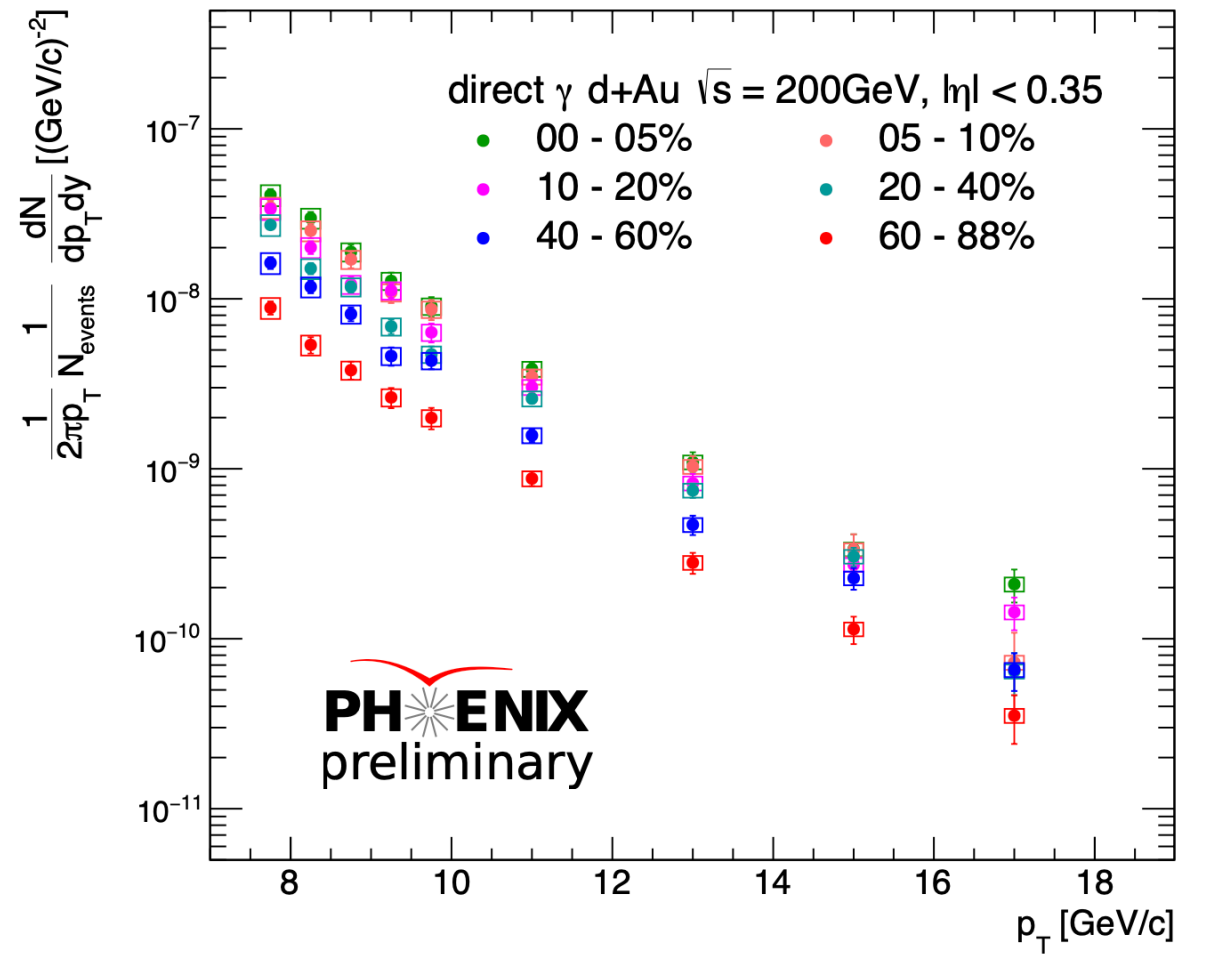
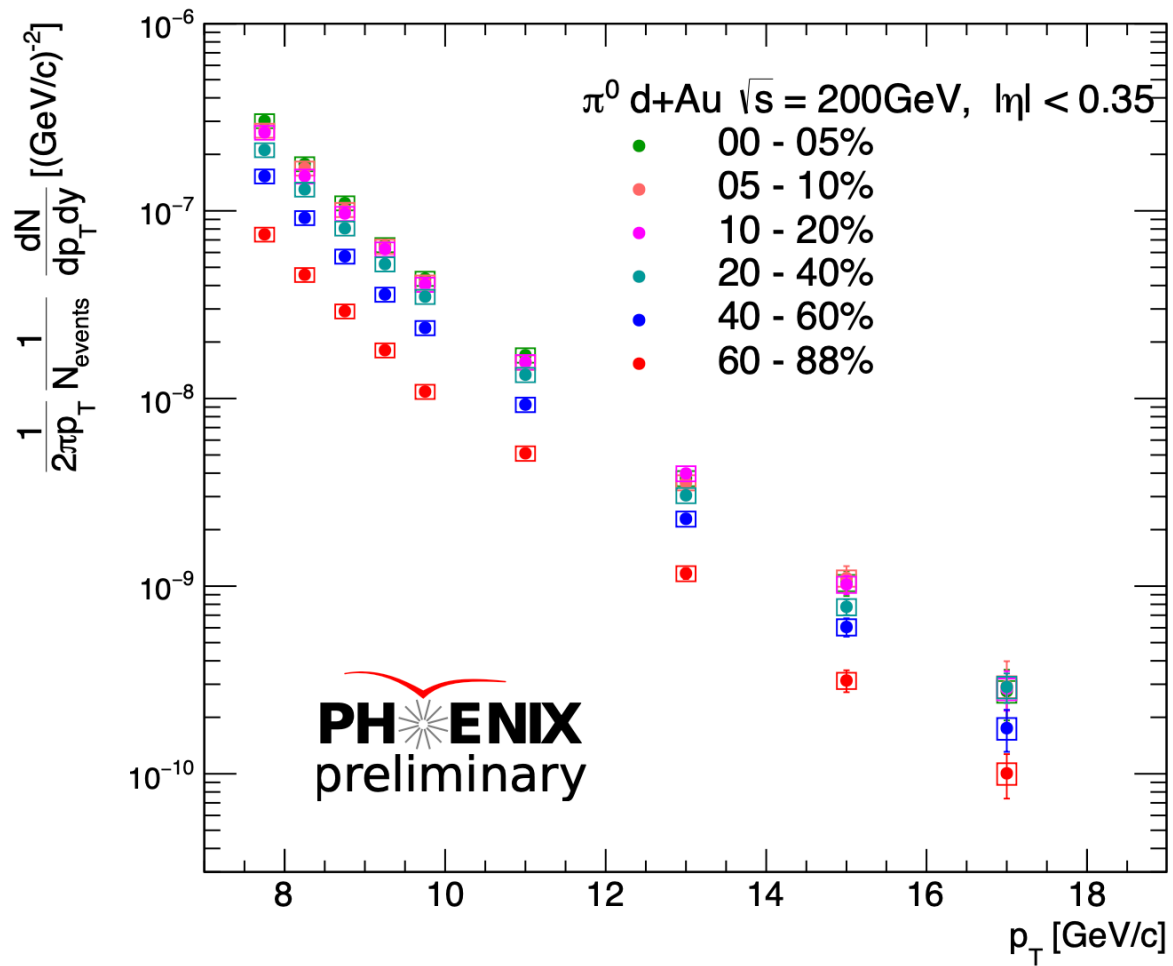


- ❖ 0-20% in Pb+Pb : average impact parameter of 3 fm with a very small variance.
- ❖ 0-20% in p+Pb : average impact parameter of 3 fm but with a large variance.
- ❖ **Cannot draw equivalent physics conclusions about central p+Pb and Pb+Pb events.**

In addition to this, there are additional biases and differences which will be discussed next.

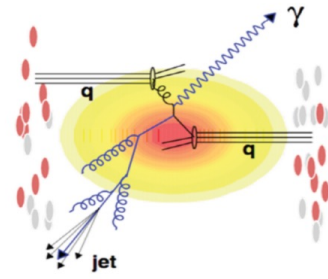
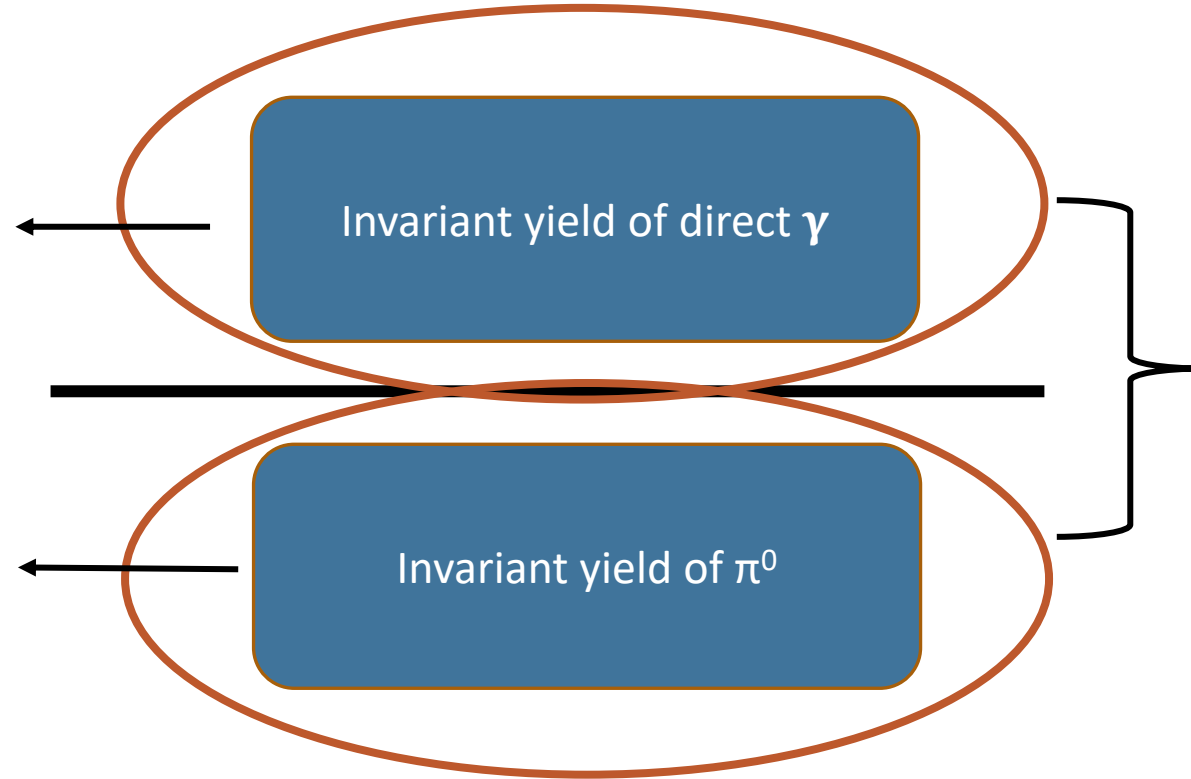
How do we study the centrality bias in experiments?

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Invariant yield of  $\pi^0$  and direct  $\gamma$

No dependence on centrality  
Dependence on centrality



Dependence on centrality  
NO dependence on centrality

Hypothesis #1 : The centrality dependence of  $\pi^0$  is from final state effect

Hypothesis #2 : The centrality dependence of  $\pi^0$  is from bias in determination of  $N_{\text{coll}}$  in different centralities

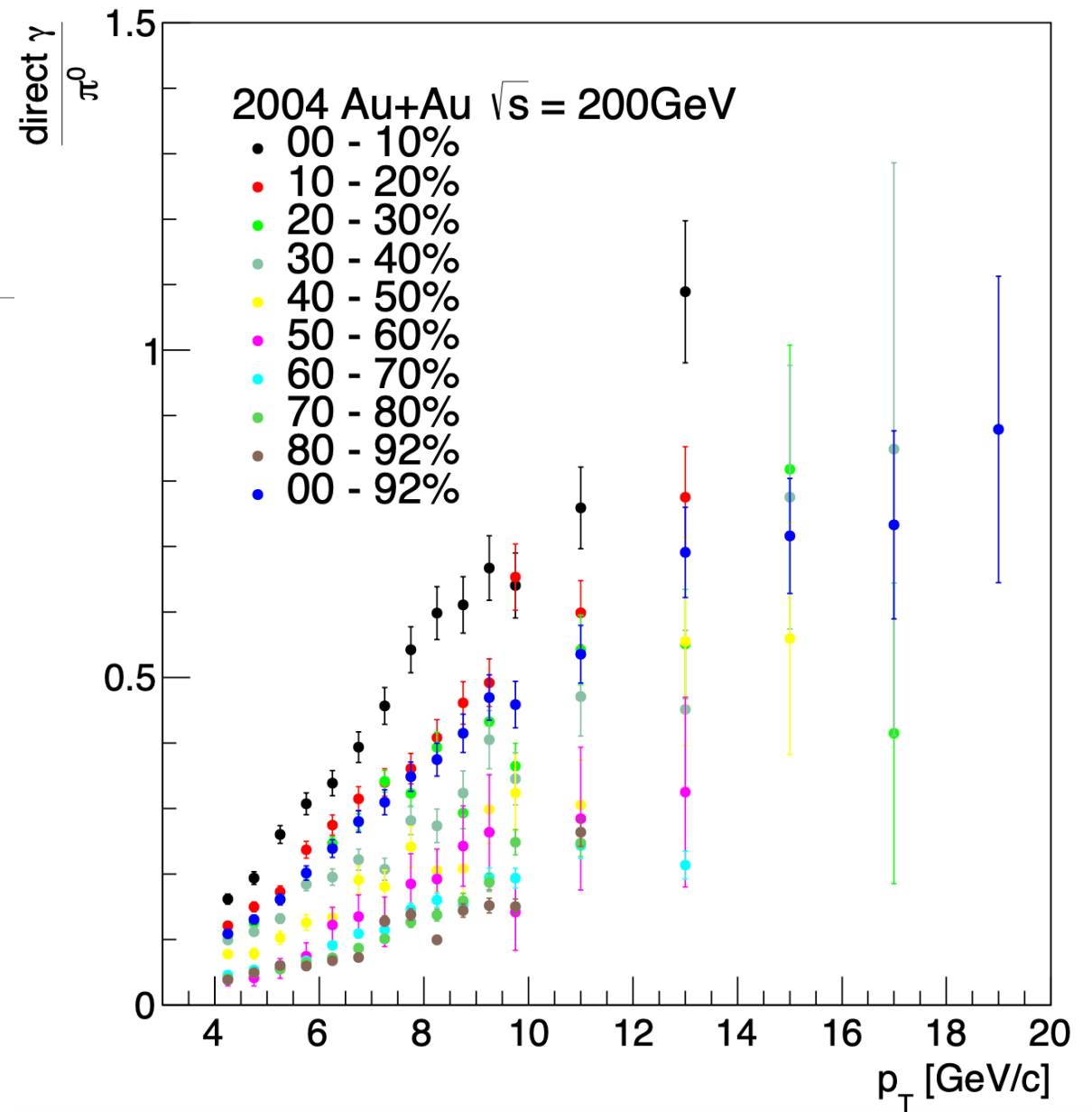
# Ratio of direct $\gamma$ over $\pi^0$

Clear centrality dependence



Hypothesis #1 :

The centrality dependence of  $\pi^0$  is from final state effect



d+Au

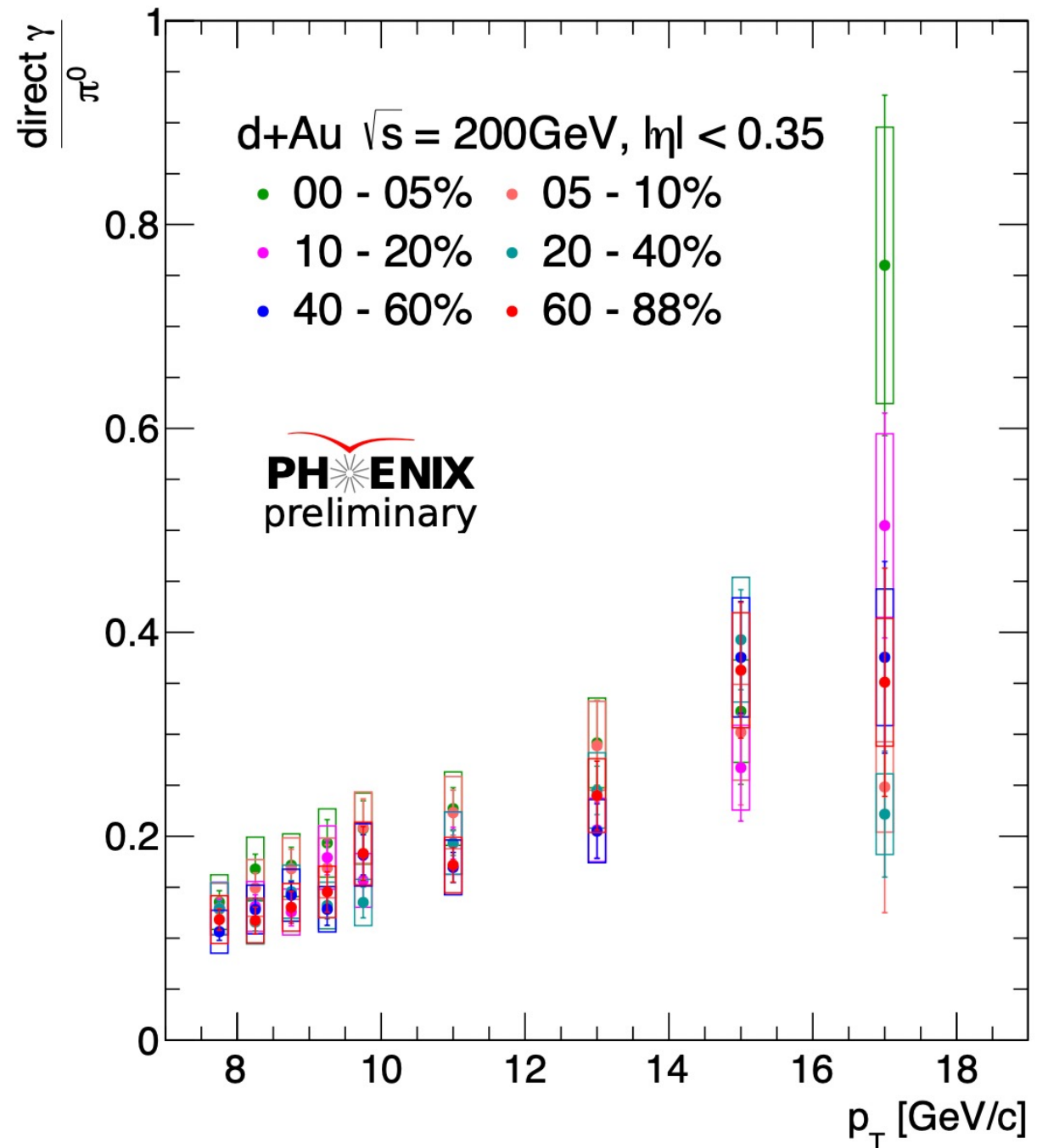
# Ratio of direct $\gamma$ over $\pi^0$

to first order, NO clear centrality  
dependence



Hypothesis #2 :

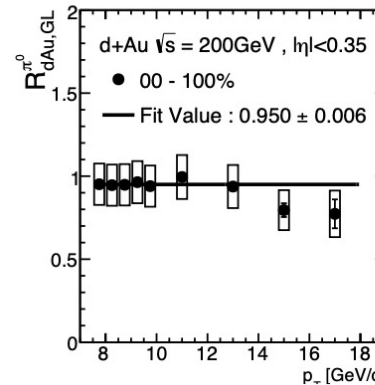
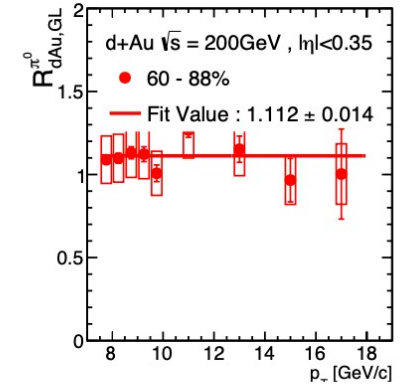
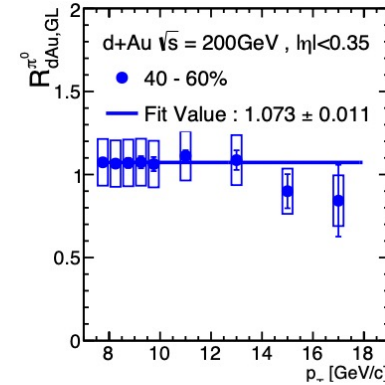
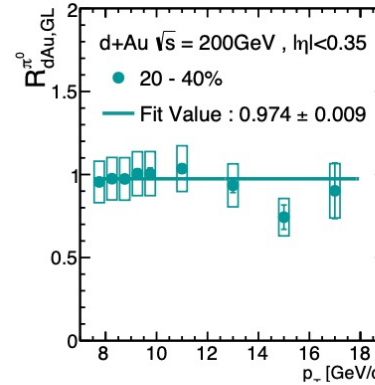
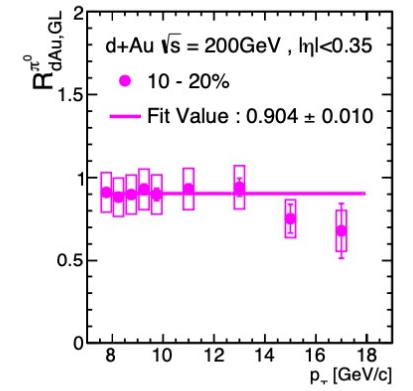
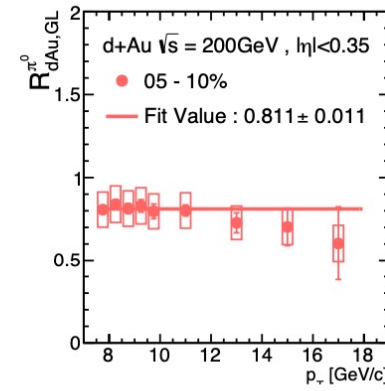
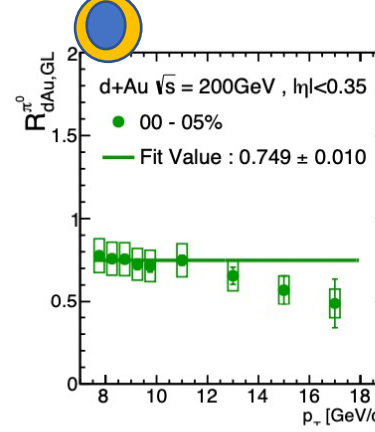
The centrality dependence of  $\pi^0$  is  
from bias in determination of  $N_{\text{coll}}$  in  
different centralities



# Nuclear Modification Factor of $\pi^0$ s

$$R_{AB,GL}(p_T) = \frac{\left(\frac{d^2N}{dp_T d\eta}\right)_{AB}}{\langle N_{coll}^{GL} \rangle_{AB} * \left(\frac{d^2N}{dp_T d\eta}\right)_{pp}} = \frac{Y(AB)}{\langle N_{coll}^{GL} \rangle_{AB} * Y(pp)}$$

- ❖ There is a centrality dependence of  $\pi^0$
- ❖ The most central events are suppressed ( $<1$ ) and peripheral events are enhanced ( $>1$ )
- ❖ In the given  $p_T$  range, to first order  $R_{dAu}$  appears to be flat.

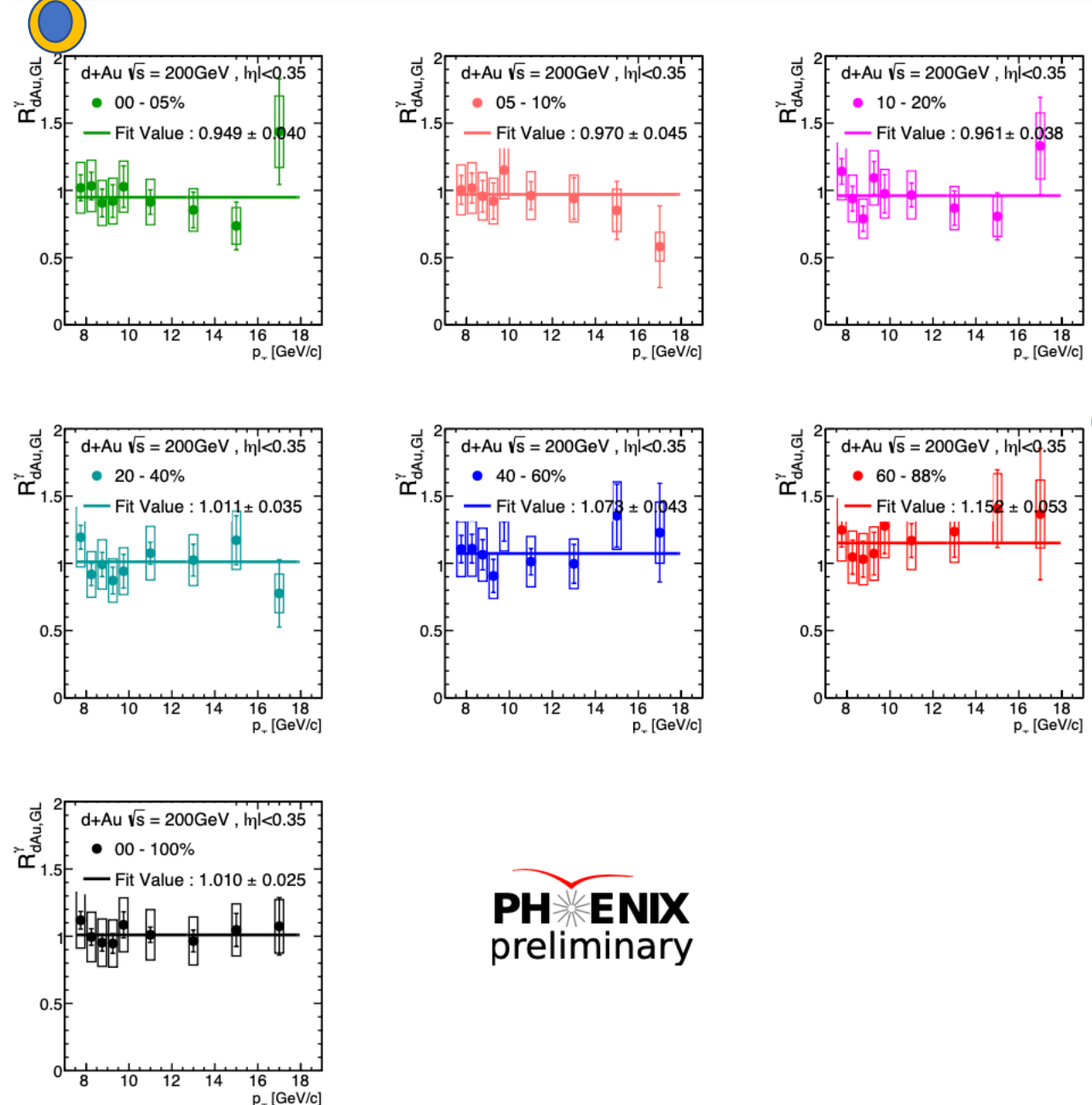


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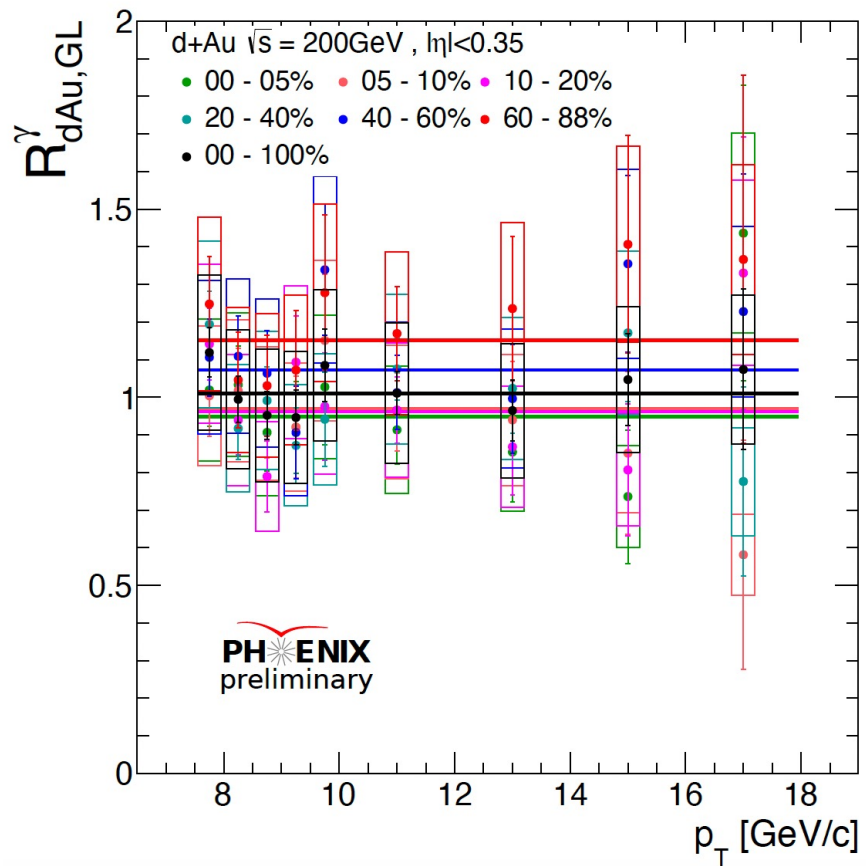
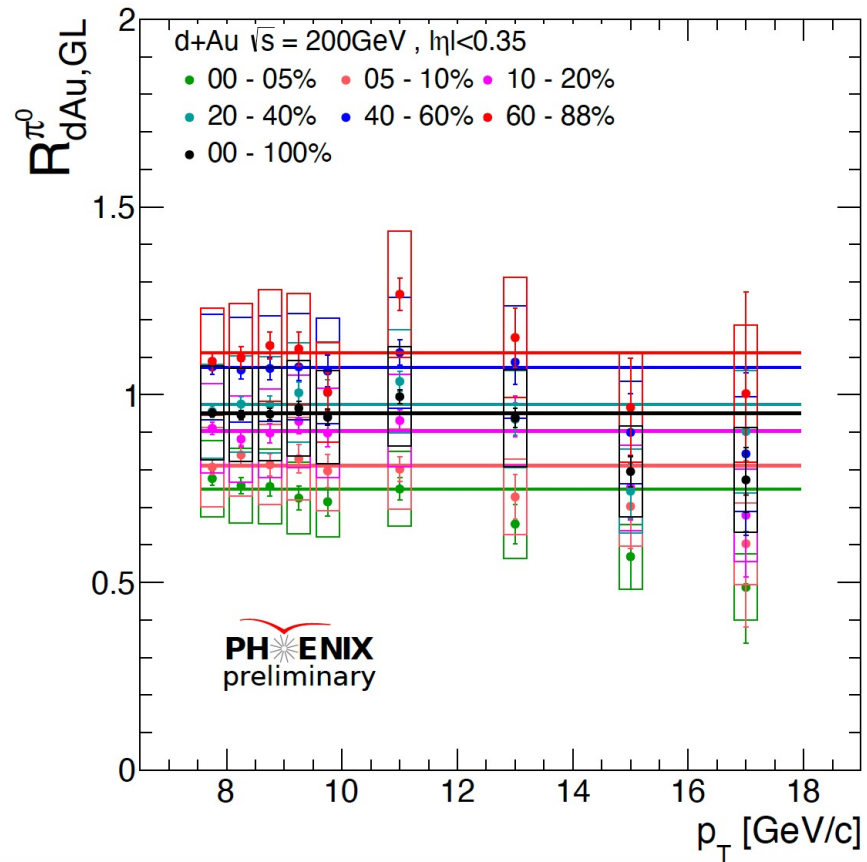
# Nuclear Modification Factor of direct $\gamma$ s

$$R_{AB,GL}(p_T) = \frac{\left(\frac{d^2N}{dp_T d\eta}\right)_{AB}}{\langle N_{coll}^{GL} \rangle_{AB} * \left(\frac{d^2N}{dp_T d\eta}\right)_{pp}} = \frac{Y(AB)}{\langle N_{coll}^{GL} \rangle_{AB} * Y(pp)}$$

- ❖ There is a centrality dependence of direct  $\gamma$ s
- ❖ The most central events are suppressed ( $<1$ ) and peripheral events are enhanced ( $>1$ )
- ❖ In the given  $p_T$  range, to first order  $R_{dAu}$  appears to be flat.







The straight lines are fit to the datapoints at different centralities

- ❖ In central events the  $\pi^0$ s are suppressed and no suppression observed in direct  $\gamma$ .
- ❖ In most peripheral events, the degree of enhancement of  $\pi^0$  matches that of direct  $\gamma$ .

## Nuclear Modification Factor of $\pi^0$ s and direct $\gamma$ s

# Experimentally determined $N_{coll}$

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$$\langle N_{coll}^{exp} \rangle = \frac{\left( \frac{d^2 N_\gamma}{dp_T d\eta} \right)_{dAu}}{\left( \frac{d^2 N_\gamma}{dp_T d\eta} \right)_{pp}}$$

High  $p_T$  direct  $\gamma$ s are transparent to QGP and scales exactly with the number of binary collisions in an event sample.

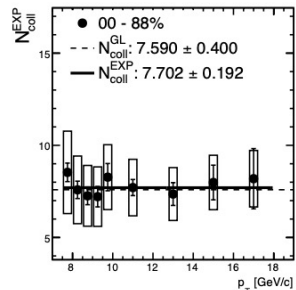
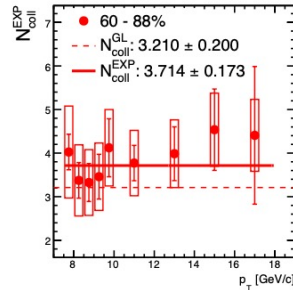
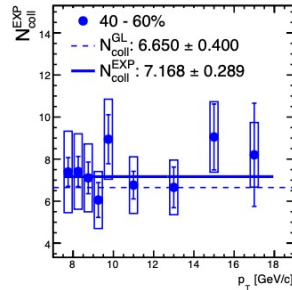
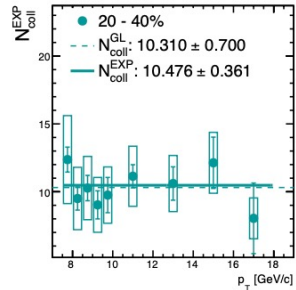
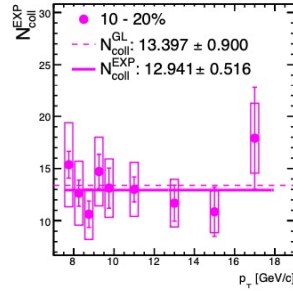
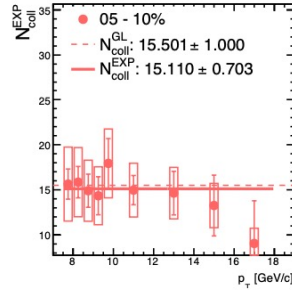
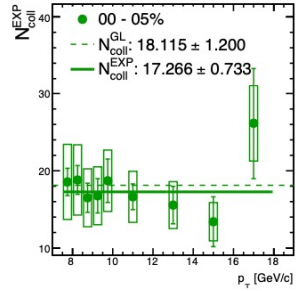
Better way to measure  $N_{coll}$  than standard Glauber model.

Can also be used to test the validity of any other modified Glauber models.

Note : on the right hand side, the numerator and the denominator are  $p_T$  dependent. The value on the left is obtained by dividing each  $p_T$  point and then doing a linear fit to get an average value.

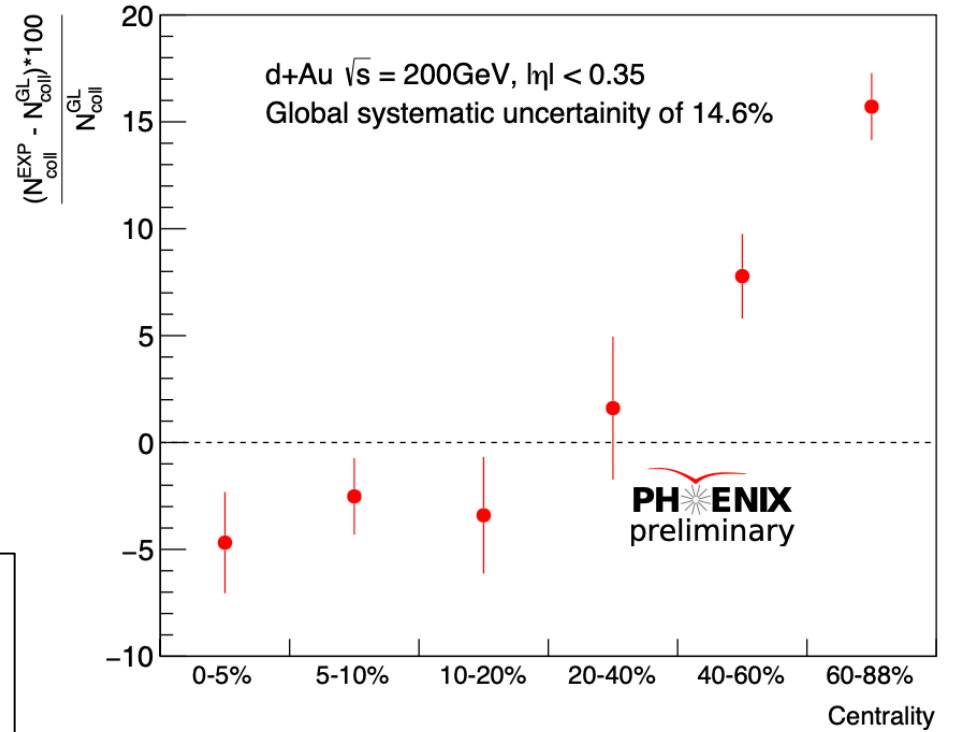
# Deriving $N_{coll}$ from direct $\gamma$ s

Bias in  $N_{coll}$  is observed in peripheral d+Au collisions



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$$\langle N_{coll}^{exp} \rangle = \frac{\left( \frac{d^2 N_\gamma}{dp_T d\eta} \right) dAu}{\left( \frac{d^2 N_\gamma}{dp_T d\eta} \right) pp}$$



# Deriving unbiased $R_{dAu}^{\pi^0}$

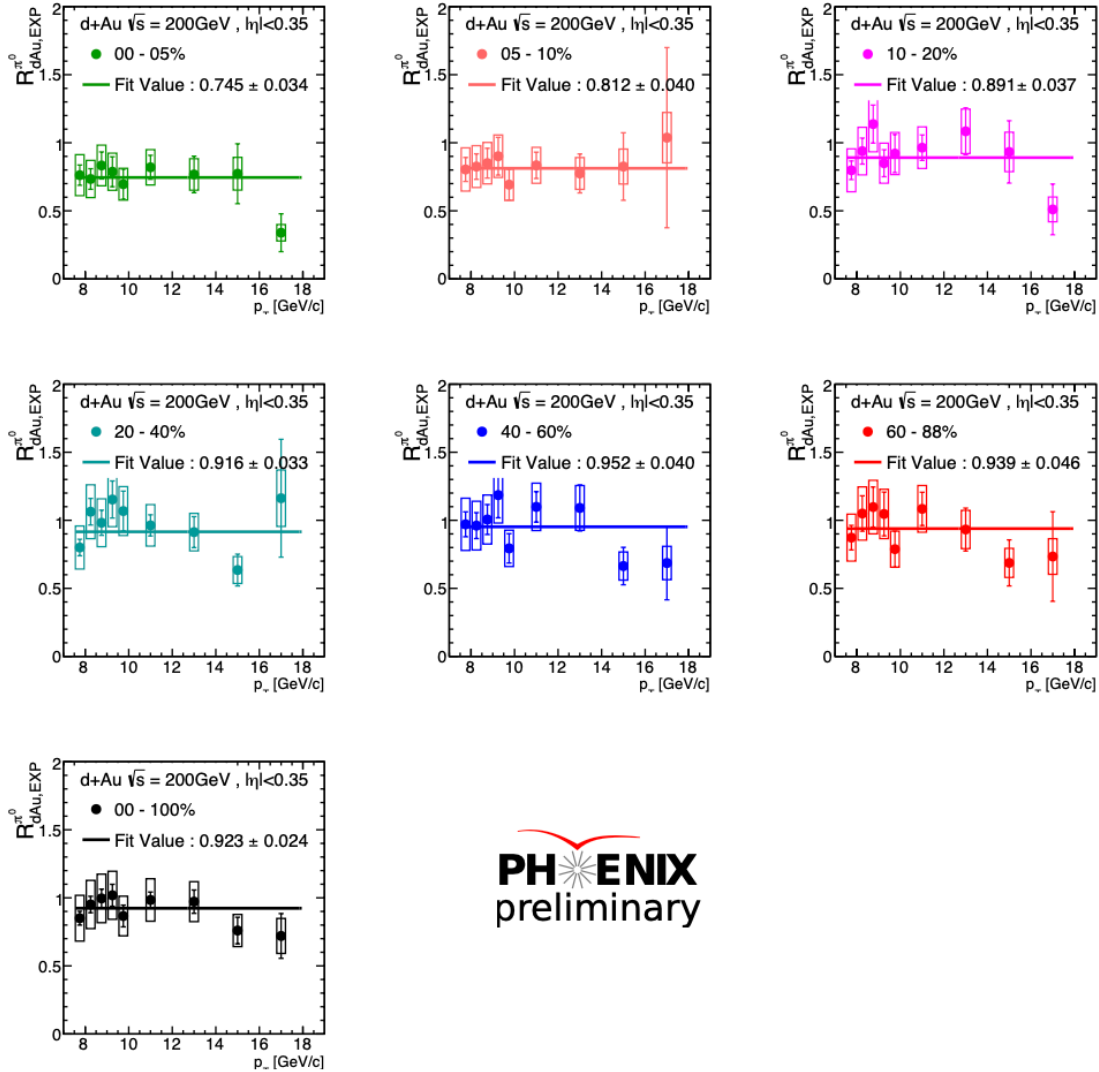
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$$R_{dAu,EXP}^{\pi^0} = \frac{R_{dAu,GL}^{\pi^0}}{R_{dAu,GL}^{\gamma}} = \frac{(Y_{dAu}^{\pi^0}/Y_{pp}^{\pi^0})}{(Y_{dAu}^{\gamma}/Y_{pp}^{\gamma})} = \frac{(Y_{dAu}^{\pi^0}/Y_{dAu}^{\gamma})}{(Y_{pp}^{\pi^0}/Y_{pp}^{\gamma})}$$

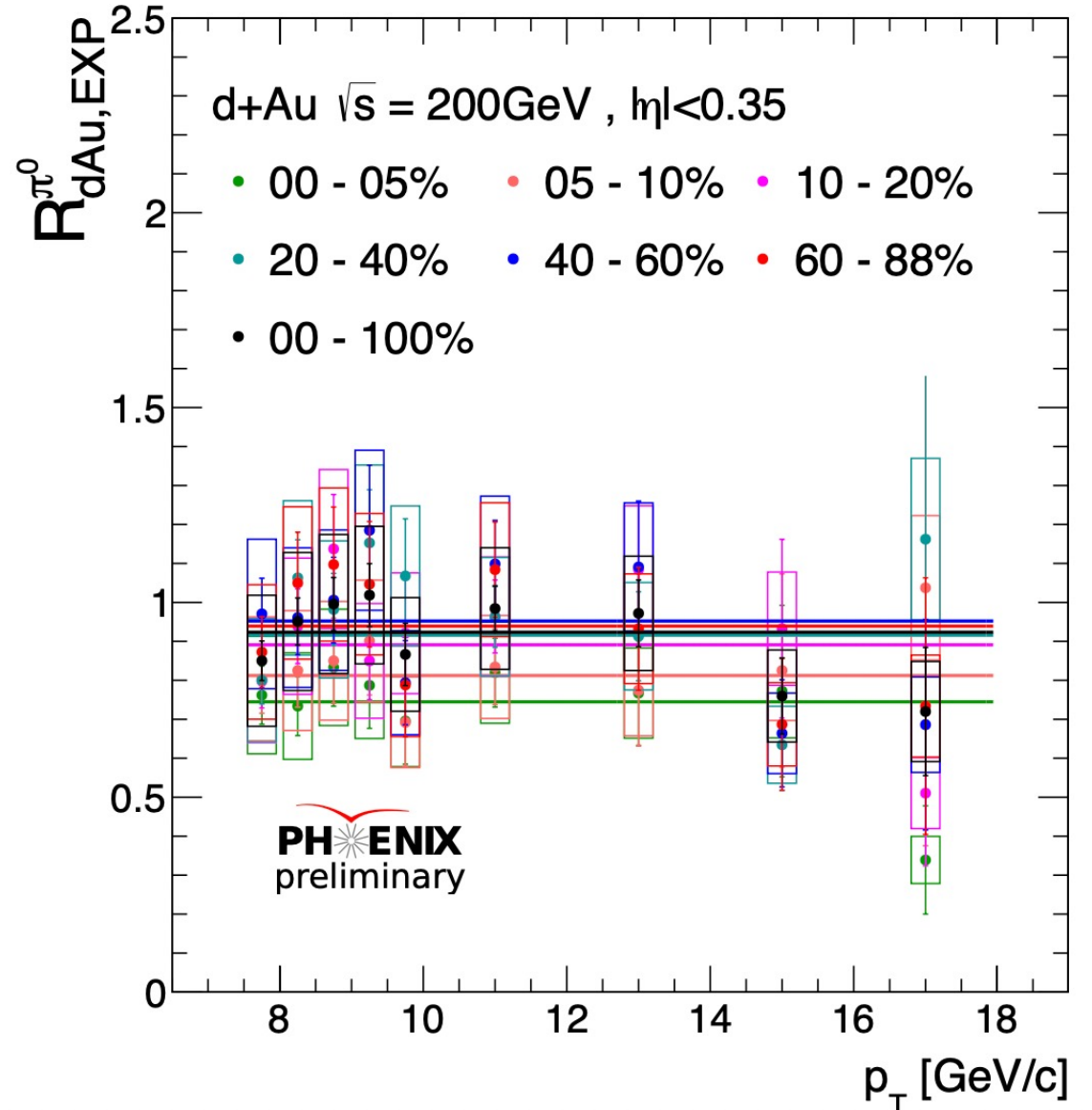
Using experimentally derived Ncoll = Renormalising  $R_{dAu,GL}^{\pi^0}$  with  $R_{dAu,GL}^{\gamma}$

# Unbiased $R_{dAu}^{\pi^0}$

$$R_{dAu,EXP}^{\pi^0} = \frac{R_{dAu,GL}^{\pi^0}}{R_{dAu,GL}^{\gamma}} = \frac{(Y_{dAu}^{\pi^0}/Y_{pp}^{\pi^0})}{(Y_{dAu}^{\gamma}/Y_{pp}^{\gamma})} = \frac{(Y_{dAu}^{\pi^0}/Y_{dAu}^{\gamma})}{(Y_{pp}^{\pi^0}/Y_{pp}^{\gamma})}$$

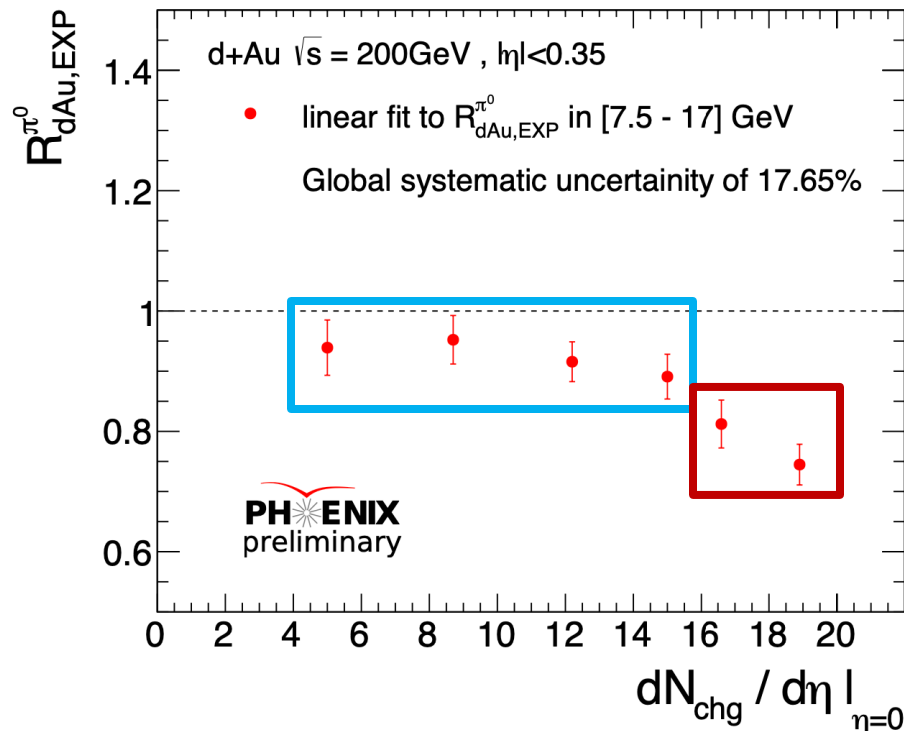


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$$R_{dAu,EXP}^{\pi^0} = \frac{R_{dAu,GL}^{\pi^0}}{R_{dAu,GL}^{\gamma}} = \frac{(Y_{dAu}^{\pi^0}/Y_{pp}^{\pi^0})}{(Y_{dAu}^{\gamma}/Y_{pp}^{\gamma})} = \frac{(Y_{dAu}^{\pi^0}/Y_{dAu}^{\gamma})}{(Y_{pp}^{\pi^0}/Y_{pp}^{\gamma})}$$

# First evidence of suppression in $R_{dAu,EXP}^{\pi^0}$



- ❖ After correcting for the bias, the  $R_{dAu,EXP}^{\pi^0}$ 
  - ❖ Does not show any enhancement in peripheral collisions
  - ❖ Does show suppression in central collisions
- ❖ There is no identified source of systematic error which has centrality dependence.
- ❖ Detailed study including,
  - ❖ ongoing analysis of p+Au,  $^3\text{He}+\text{Au}$
  - ❖ initial state effects on the production mechanism of  $\pi^0$ s and direct  $\gamma$ s

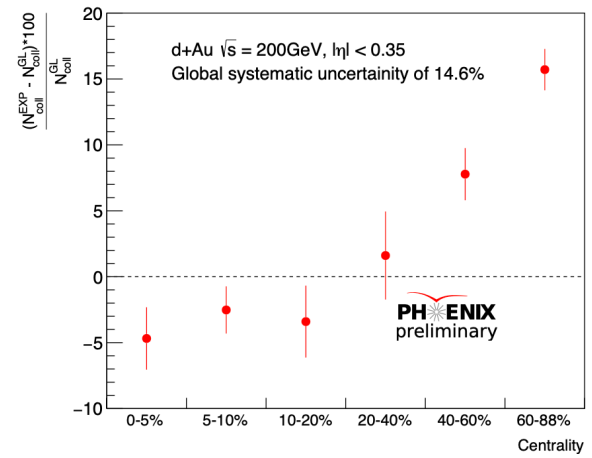
is necessary to understand whether this observed suppression is an initial or final state effect.



# SUMMARY

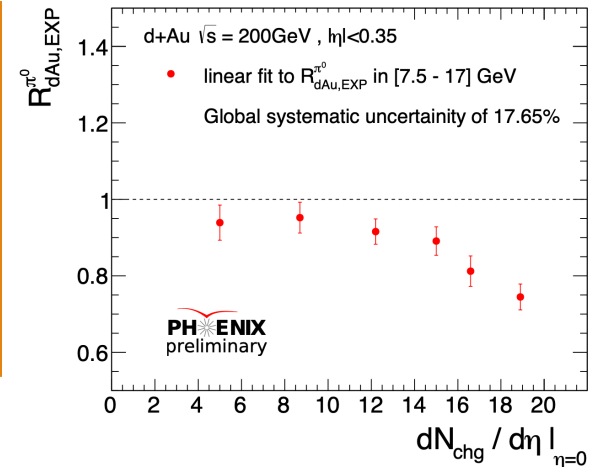
Bias in centrality determination in small system collisions

- ❖ Determination of  $N_{\text{coll}}$  by using Glauber Model is biased. Specifically for peripheral events
- ❖ Experimentally derived  $N_{\text{coll}}$  using high  $p_T$  direct  $\gamma$ s
- ❖ Future analysis in p+Au and  $^3\text{He}+\text{Au}$  system will provide more clarification on the origin of centrality bias



Evidence of  $R_{dAu}^{\pi^0}$  suppression in most central

- ❖ Normalized nuclear modification factor of  $\pi^0$ s using direct  $\gamma$ s
- ❖ Most central events show suppression of  $\sim 15\%$
- ❖ Models studying the initial state effect in the production of  $\pi^0$ s and direct  $\gamma$ s will shed light on the origin of the observed suppression





THANK YOU FOR YOUR ATTENTION



# Backup

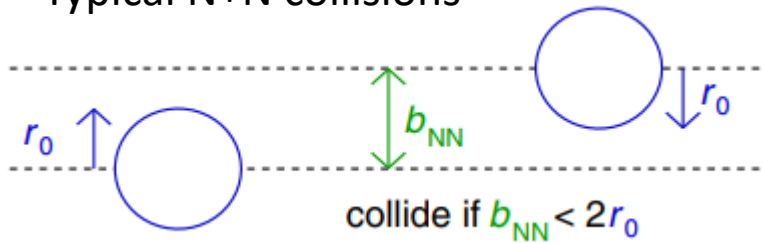
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# Origin of the bias...?

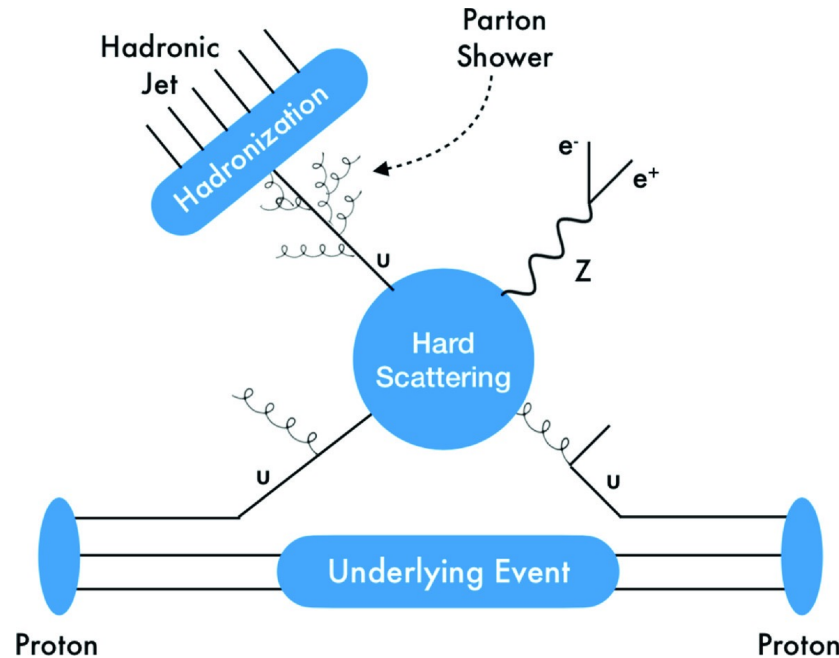
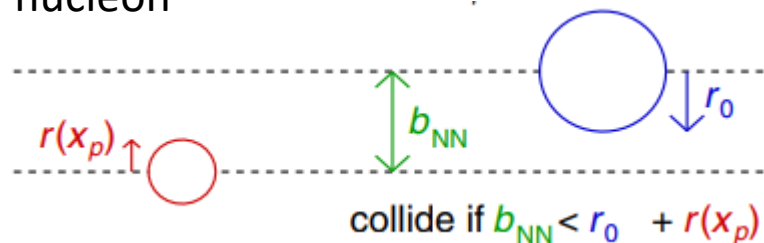
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# high-x (effective) size fluctuations

Typical N+N collisions



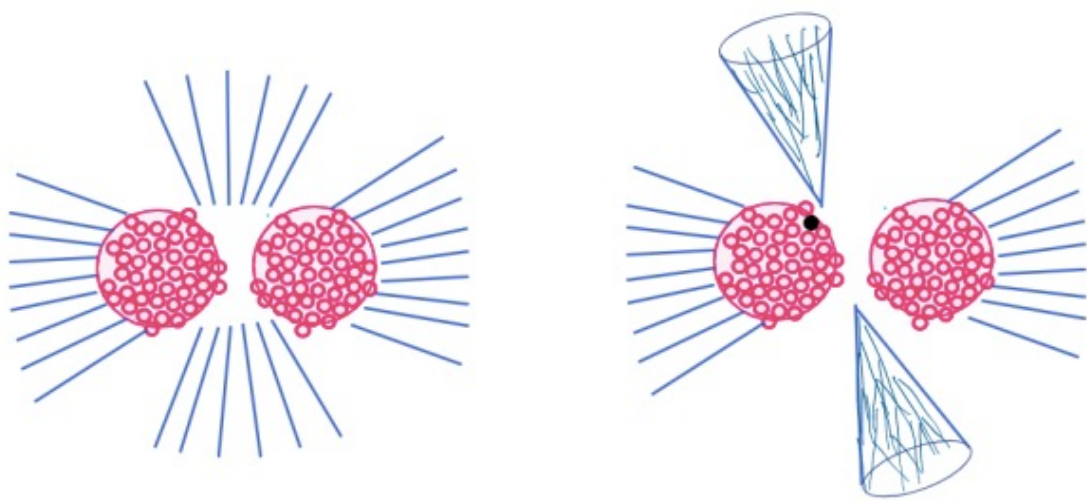
N+N collisions with large- $x_p$  projectile nucleon



The high-X parton creates the hard scattering event. But the underlying event is severely depleted.

This can be thought of as  
 a) energy conservation or  
 b) change in the cross-section of the nuclei due to the presence of high-X parton.

Phys. Rev. C 94 (2016), 024915

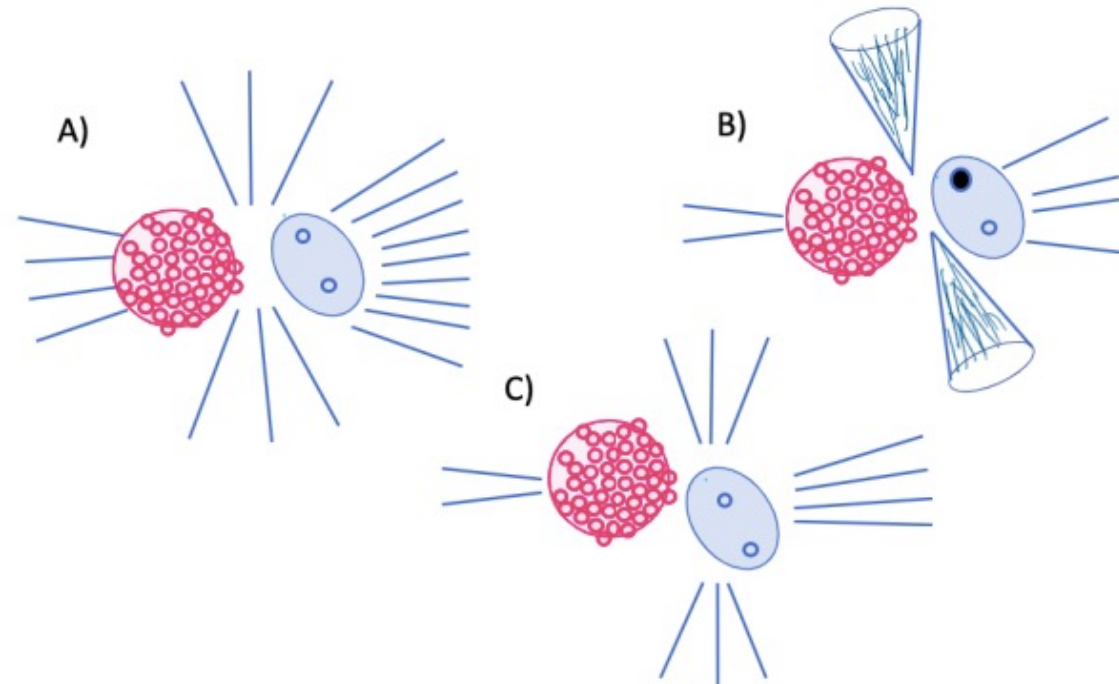


In a heavy-ion collision, the presence of one high- $X$  parton nuclei, creates the jets, but the average underlying event isn't affected as there are several other partons for interactions.

In a d+Au collision, the presence of one high- $X$  parton depletes the underlying event and there are not enough other interactions to compensate for this.

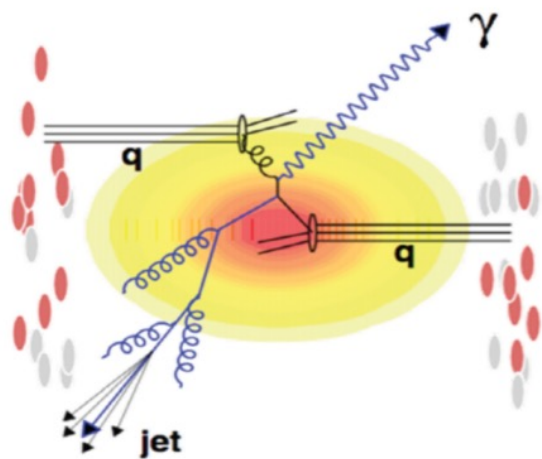
Thus a central d+Au event will look like a peripheral d+Au event.

This is a  $p_T$  (or  $x$ ) dependent change. The bin-shift is larger at higher momentum.



**This shrinking nucleon model has a prediction for  $R_{dAu}(x)$  and thus we can compare it to our data.**

$R_{dAu}^{\pi^0}$  and  $R_{dAu}^{\gamma}$  as a function of parton momentum  $x$



$$x_p = 2p_T^{\text{jet}} / \sqrt{s_{NN}} \approx 2p_T^{\pi^0} / (0.75 * \sqrt{s_{NN}})$$

$$x_p = 2p_T^{\text{direct}\gamma} / \sqrt{s_{NN}}$$

Does our data fit the expectation from the “shrinking nucleon” picture?

