

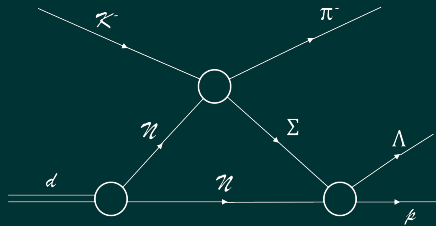
# J-PARC E90: HIGH RESOLUTION SPECTROSCOPY OF THE $\Sigma N$ CUSP BY USING $d(K^-, \pi^-)$ REACTION

YUDAI ICHIKAWA (JAEA)

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14TH INTERNATIONAL CONFERENCE ON HYPERNUCLEAR AND STRANGE PARTICLE PHYSICS

## 1. Reaction mechanism



## 2. Theoretical formula

$$\Gamma(\bar{K}d \rightarrow \Lambda n) \sim \Gamma(\bar{K}N \rightarrow \pi \Sigma) \mathcal{F}_d(q, k_\Sigma) \Gamma(\Sigma N \rightarrow \Lambda n)$$

$$\mathcal{F}_d(q, k_\Sigma) = \int \rho_d(r) \{ \sin(qr) |g(r)\} \{ \exp(ik_\Sigma r) |r\} r^2 dr$$

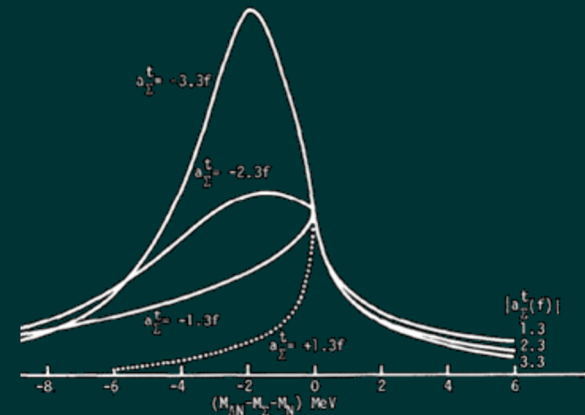
$$\Gamma(\Sigma N \rightarrow \Lambda n) = k_\Sigma |1 - ik_\Sigma A_\Sigma|, \quad A_\Sigma = a - ib$$

$$\text{Above threshold: } R_i = 4\pi b \{ (1 + k_\Sigma b)^2 + (k_\Sigma a)^2 \}$$

$$\text{Below threshold: } R_i = 4\pi b \{ (1 + |k_\Sigma| a)^2 + k_\Sigma^2 b^2 \}$$

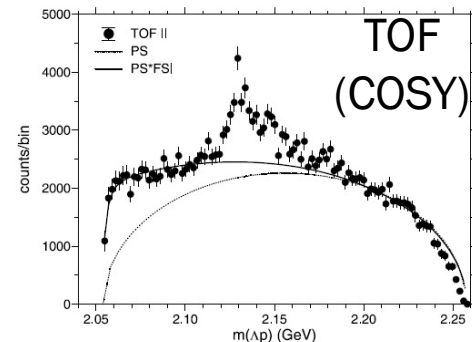
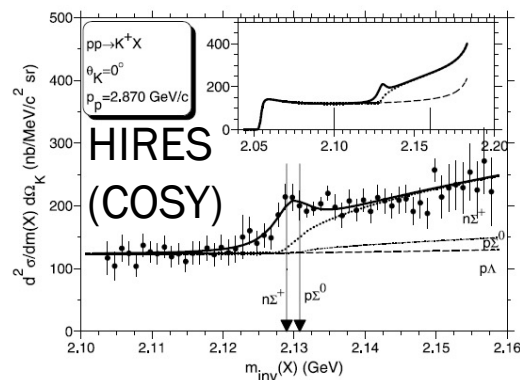
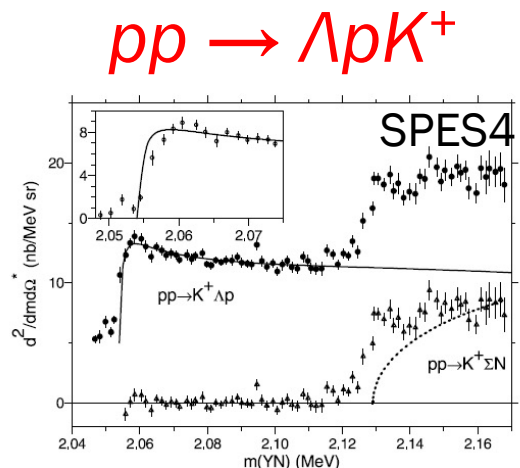
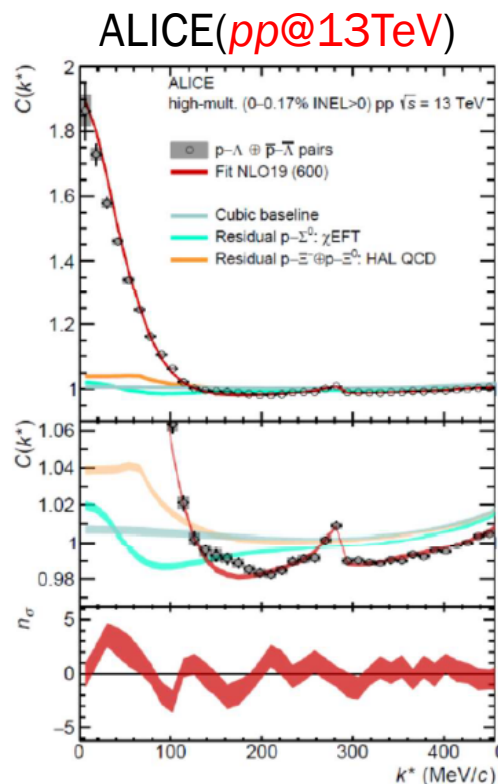
$$(k_\Sigma = i |k_\Sigma|)$$

## 3. Calculated spectrum

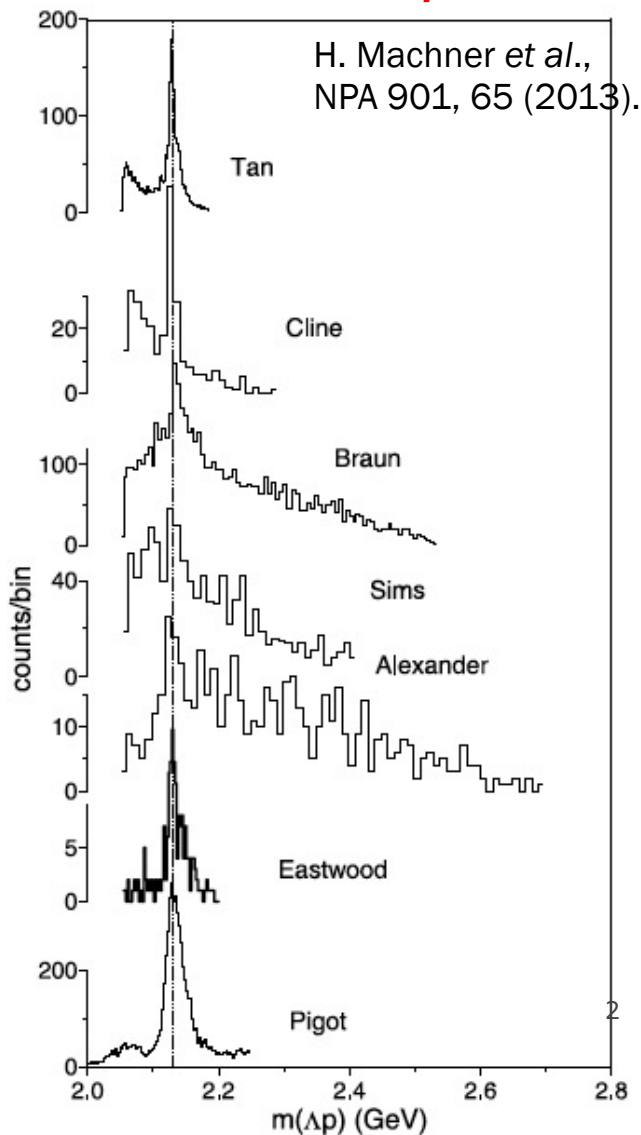


# " $\Sigma N$ CUSP"

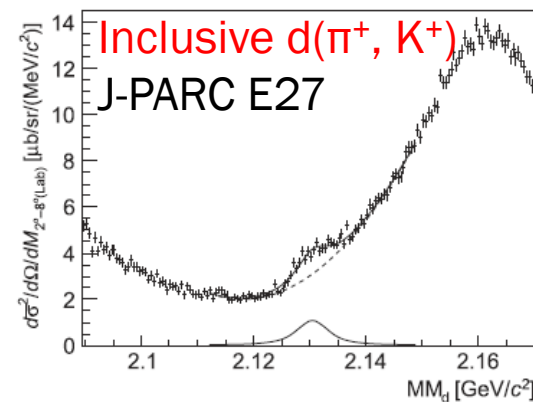
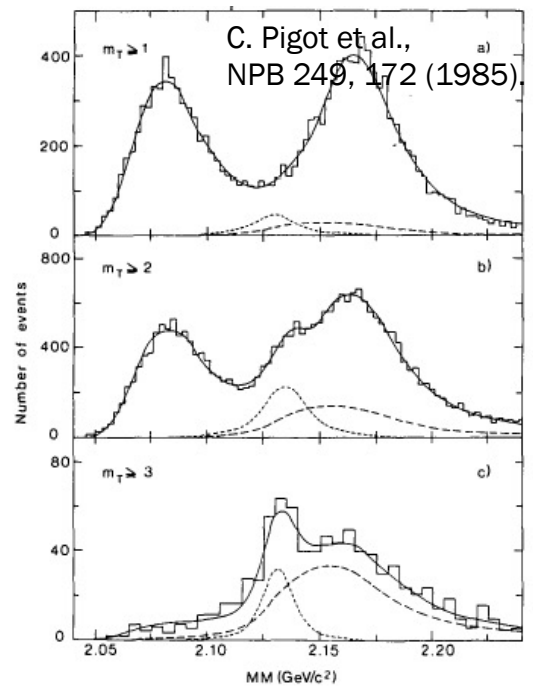
Clear enhancement around  $\Sigma N$  threshold ( $\sim 2.13 \text{ GeV}/c^2$ )



$K^- d \rightarrow \Lambda p \pi^-$



$\pi^+ d \rightarrow \Lambda p K^+$



# PURPOSE OF THE E90 EXPERIMENT

- Deduce the  $\Sigma N$  scattering length of  $(T, S) = (1/2, 1)$  channel by fitting “ $\Sigma N$  cusp” spectrum shape observed in the missing mass of the  $d(K^-, \pi^-)$  reaction.
  - Unstable bound state ( $\Sigma N$  dibaryon)? or Virtual state?

The key of this experiment is the excellent missing-mass resolution thanks to the  $S-2S$  spectrometer (used in E70) and high statistics. We will be able to achieve the best resolution of 0.4 MeV in  $\sigma$ , which is two times better than the past experiment (HIRES at COSY).

# THRESHOLD CUSP

Cusp structure can be expressed by the scattering length (for B'C'),  $A = a + ib$

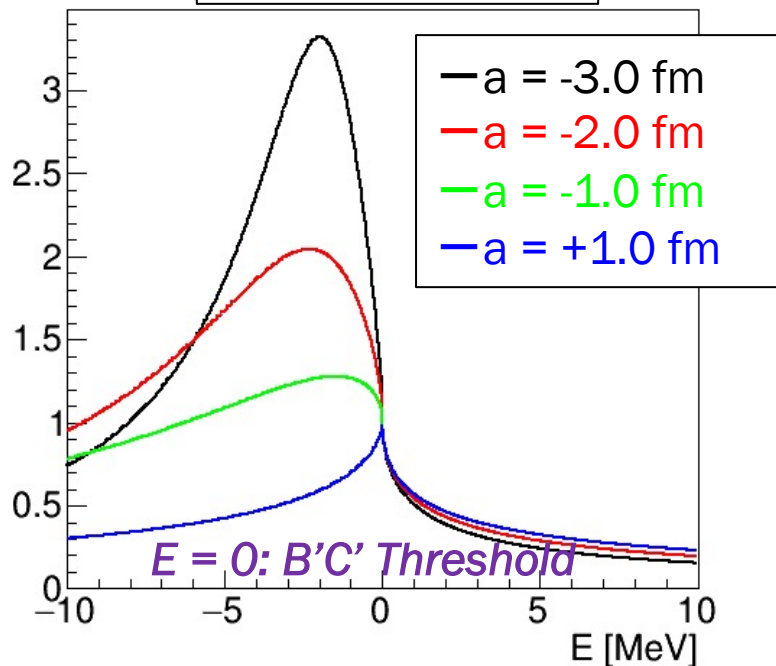
■ Above threshold:  $R = \frac{4\pi b}{\{(1+kb)^2+(ka)^2\}} \sim 1 - 2kb + O(k^2)$

■ Below threshold:  $R = \frac{4\pi b}{\{(1+\kappa a)^2+(\kappa b)^2\}} \sim 1 - 2\kappa a + O(\kappa^2), k = i\kappa$

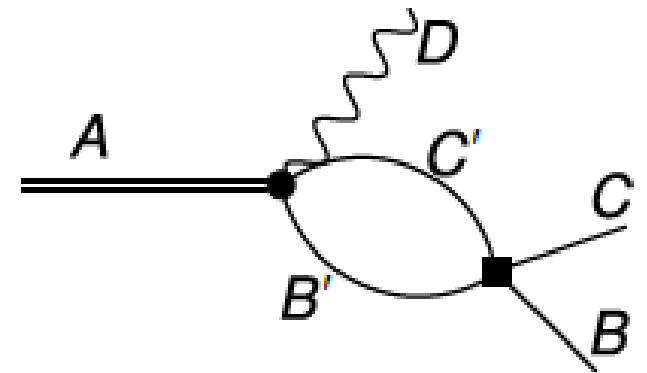
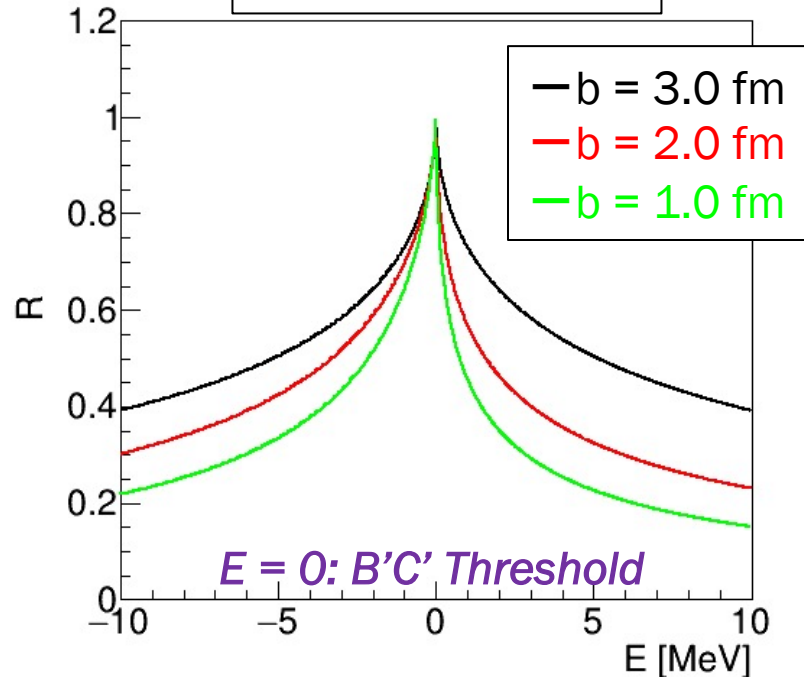
Reduced mass  
 $\mu = m_B m_{C'} / (m_B + m_{C'})$

$k(\text{relative momentum for B'C'}) \sim \sqrt{2\mu E}$

b = 2.0 fm (fixed)



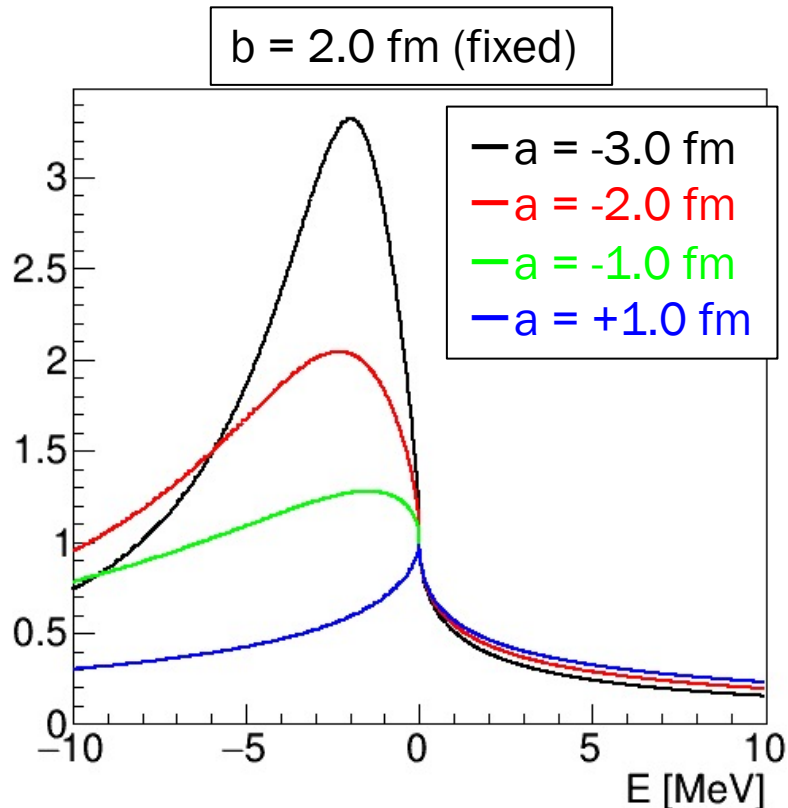
a = +1.0 fm (fixed)



For the “ $\Sigma N$  cusp”,  
 $B' = \Sigma, C' = N, B = \Lambda, C = N$

# “ΣN CUSP”

- “ΣN cusp” is measured by  $K^-d \rightarrow \pi^- \Lambda p$  reaction etc..
  - T:  $T = 1/2$  ( $\Lambda p$  final state)
  - S:  $^3S_1$  is favored, D-target; observed in forward angles



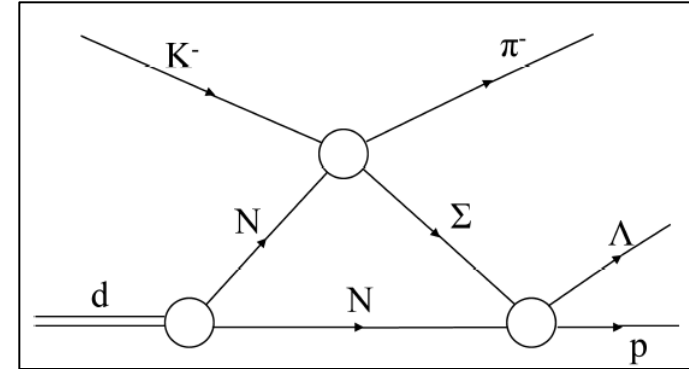
Above threshold:  $R = \frac{4\pi b}{\{(1+kb)^2+(ka)^2\}} \sim 1 - 2kb + O(k^2)$

Below threshold:  $R = \frac{4\pi b}{\{(1+\kappa a)^2+(\kappa b)^2\}} \sim 1 - 2\kappa a + O(\kappa^2)$ ,  $k = i\kappa$   
( $k \sim \sqrt{2\mu E}$ )

**“ΣN Cusp” can be expressed by the ΣN scattering length ( $A_\Sigma = a + ib$ ) of the  $(T, S) = (1/2, ^3S_1)$  channel!!**

# REQUIREMENTS FOR THE “ $\Sigma N$ CUSP” EXPERIMENT

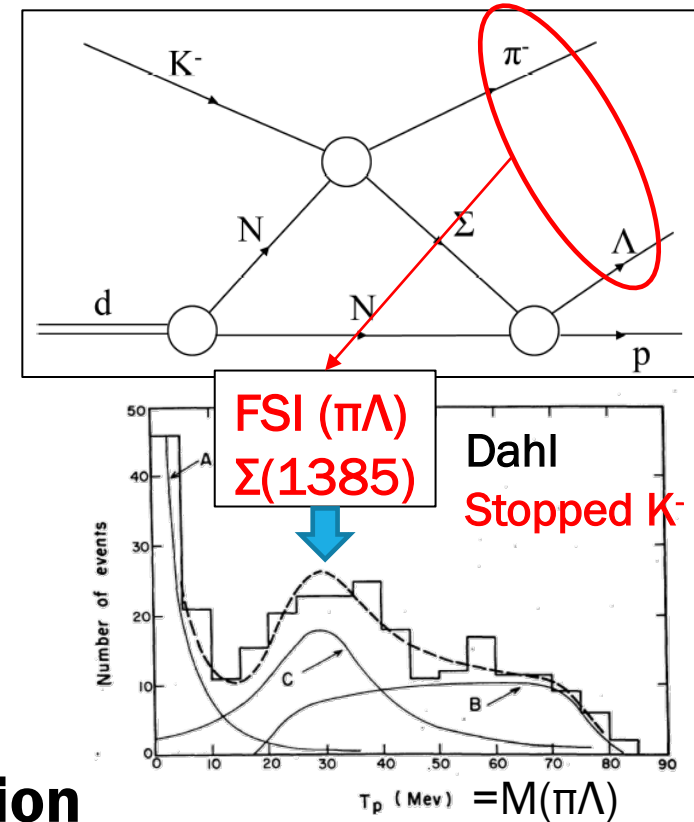
- Good energy resolution ( $\sigma < 1$  MeV)
- High statistics ( $> 10^4$  events)
- Good Signal / Noise (S/N) ratio
- Avoid FSI **except for the  $\Sigma N$** 
  - **×** Stopped  $K^-d \rightarrow \Lambda p \pi^-$  reaction ( $p_\pi \sim p_\Lambda \sim p_p$ )  
FSI:  $\pi\Lambda$ ,  $\pi p$ ,  $YN$  ( $YN$  FSI =  $\Sigma N$  cusp signal)
  - **○** In-flight  $K^-d \rightarrow \Lambda p \pi^-$  reaction ( $p_K \sim p_\pi \gg p_\Lambda \sim p_p$ )  
FSI:  $YN$  ( $YN$  FSI =  $\Sigma N$  cusp signal),  $\bigcirc$  impulse approximation
- Decompose  $^1S_0$  and  $^3S_1$  contribution  
( $K^-d \rightarrow \Lambda p \pi^-$  reaction: extract only  $^3S_1$  contribution by D-target property)



*There was no experiment to satisfy these requirements!!*

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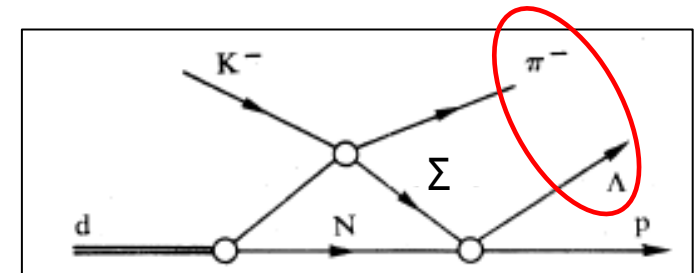
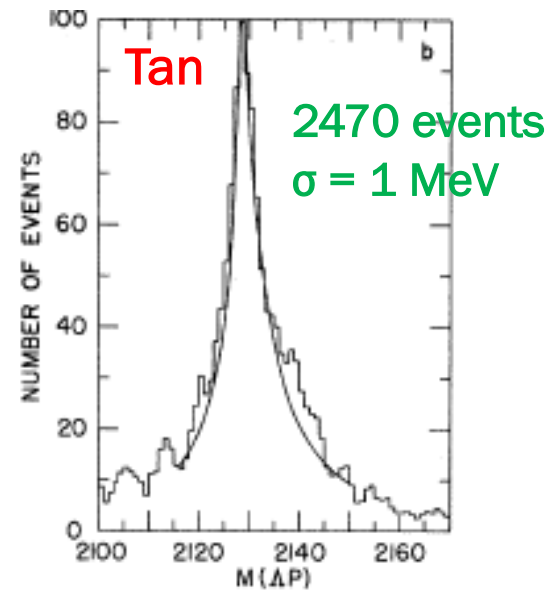
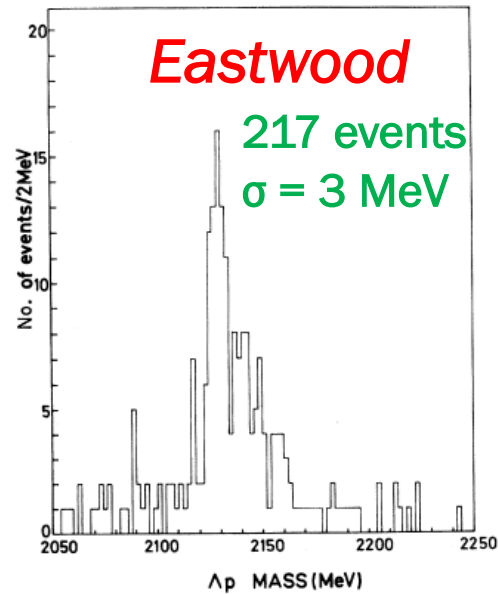
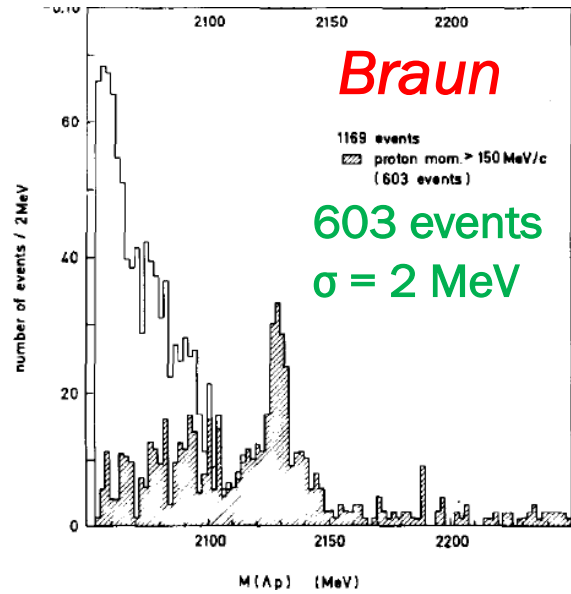
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Bubble chamber

	Reaction	Comments	Statistics	Resolution
Braun	Inflight $d(K^-, \pi^-) \Lambda p$ 680 – 840 MeV/c	Low statistic, worse resolution	603 events ( $\cos \theta > 0.9$ , momcut)	2 MeV
Eastwood	Inflight $d(K^-, \pi^-) \Lambda p$ 1450, 1650 MeV/c	Low statistic worse resolution	217 events ( $\cos \theta > 0.9$ , momcut)	3 MeV
Tan	stopped $d(K^-, \pi^-) \Lambda p$	Large FSI	2470 events	1 MeV
Pigot	Inflight $d(K^-, \pi^-), d(\pi^+, K^+)$	Poor resolution	Uncertain	9.1 MeV ( $d(K^-, \pi^-) 1.4 \text{ GeV/c}$ )
$pp \rightarrow \Lambda p K^+$ (COSY etc)	$pp \rightarrow \Lambda p K^+$	$^1S_0 + ^3S_1$ admixture Worse SN	High	0.8 MeV
ALICE	pp (Femtoscropy)	$^1S_0 + ^3S_1$ admixture	High	No description
J-PARC E27	$d(\pi^+, K^+)$ (Inclusive)	Worse SN (inclusive)	High	1.4 MeV



Stopped  $K^-$  reaction

- Multiple  $K^-$  scattering
- FSI:  $\pi\Lambda, \pi p$  ( $p_\pi \sim p_\Lambda \sim p_p$ )  
[YN FSI = Signal]

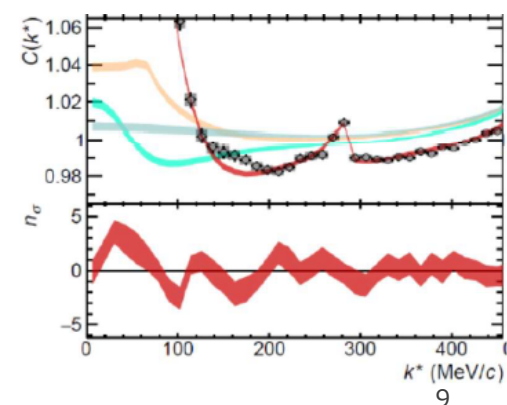
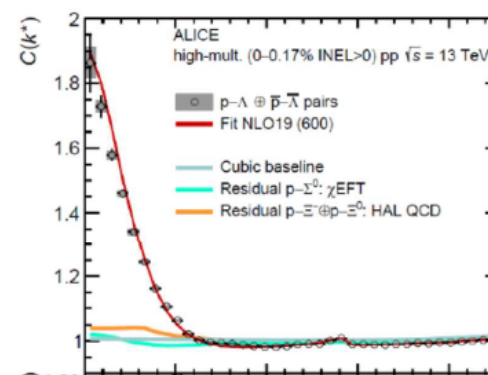
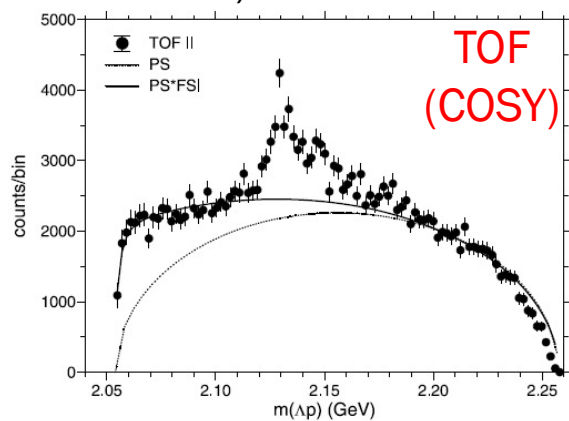
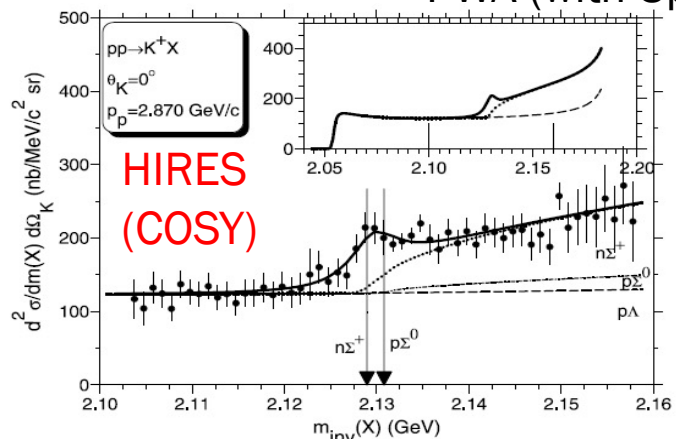


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ALICE	pp (Femtoscscopy)	$^1S_0 + ^3S_1$ admixture	High	No description
J-PARC E27	$d(\pi^+, K^+)$ (Inclusive)	Worse SN (inclusive)	High	1.4 MeV

$pp \rightarrow \Lambda p K^+$ : Good resolution, Worse SN,  $^1S_0 + ^3S_1$  mixed,  
Complicated reaction mechanism (via  $N^*, \Delta^*$ )  
→ PWA (with Spin observable)

ALICE ( $pp @ 13 \text{ TeV}$ , Femtoscopy)  
 $^1S_0 + ^3S_1$  mixed → Spin observable

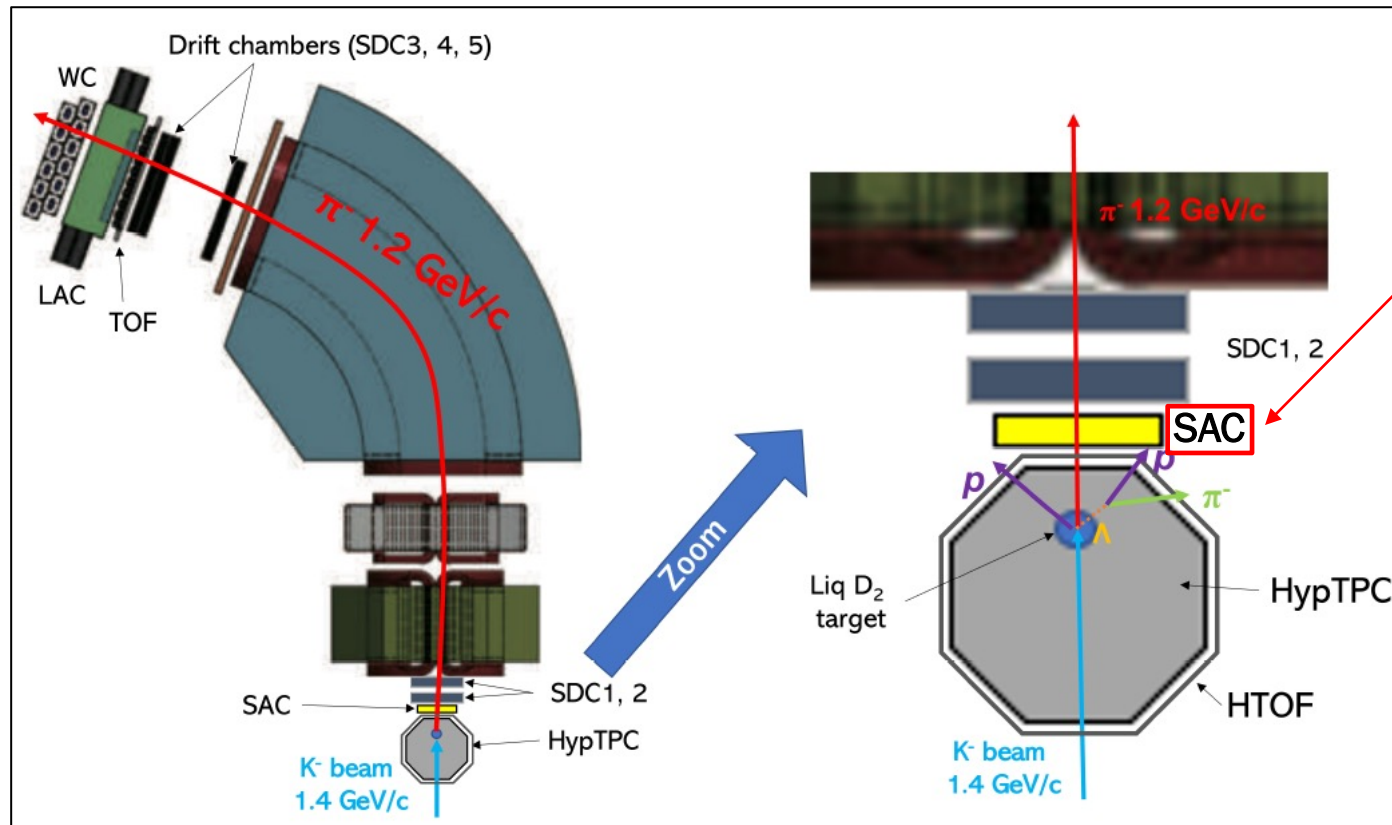


# J-PARC E90 SET UP

- Reaction:  $K^-d \rightarrow \Lambda p \pi^-$  at 1.4 GeV/c
- S-2S(developed for E70):  $\pi^-$  measurements  $\rightarrow$  measurement of missing mass spectrum
  - Good mass resolution:  $\Delta M \sim 0.4 \text{ MeV } (\sigma)$ ,  $(\Delta p/p(\text{K18}))=3.3 \times 10^{-4}(\text{FWHM})$ ,  $\Delta p/p(\text{S-2S})=6.0 \times 10^{-4}(\text{FWHM})$
- HypTPC(developed for E42): Final state ( $\Lambda p$ ) restriction and background suppression

HypTPC:  
Talked by S. Hayakawa  
(6/30)

Momentum transfer  
 $\sim 200 \text{ MeV}/c$



*New detector*

# Key of E90: High resolution $\Delta M = 0.4$ MeV

Model	J04	J04c	J-A	NSC 97f	NSC 89	ND	NF	NB
a [fm]	3.83	3.63	-2.37	-1.03	2.54	2.06	-1.29	-3.0
b [fm]	3.01	3.09	3.74	2.41	0.26	4.64	3.02	1.8
Model	chiral EFT (NLO13)				chiral EFT (NLO19)			
$\Lambda$ [MeV]	500	550	600	650	500	550	600	650
a [fm]	-2.61	-2.44	-2.27	-2.06	-0.95	-0.98	-2.29	-1.95
b [fm]	2.89	3.11	3.29	3.59	4.77	4.59	3.39	3.38

$a > 0$ : Attractive

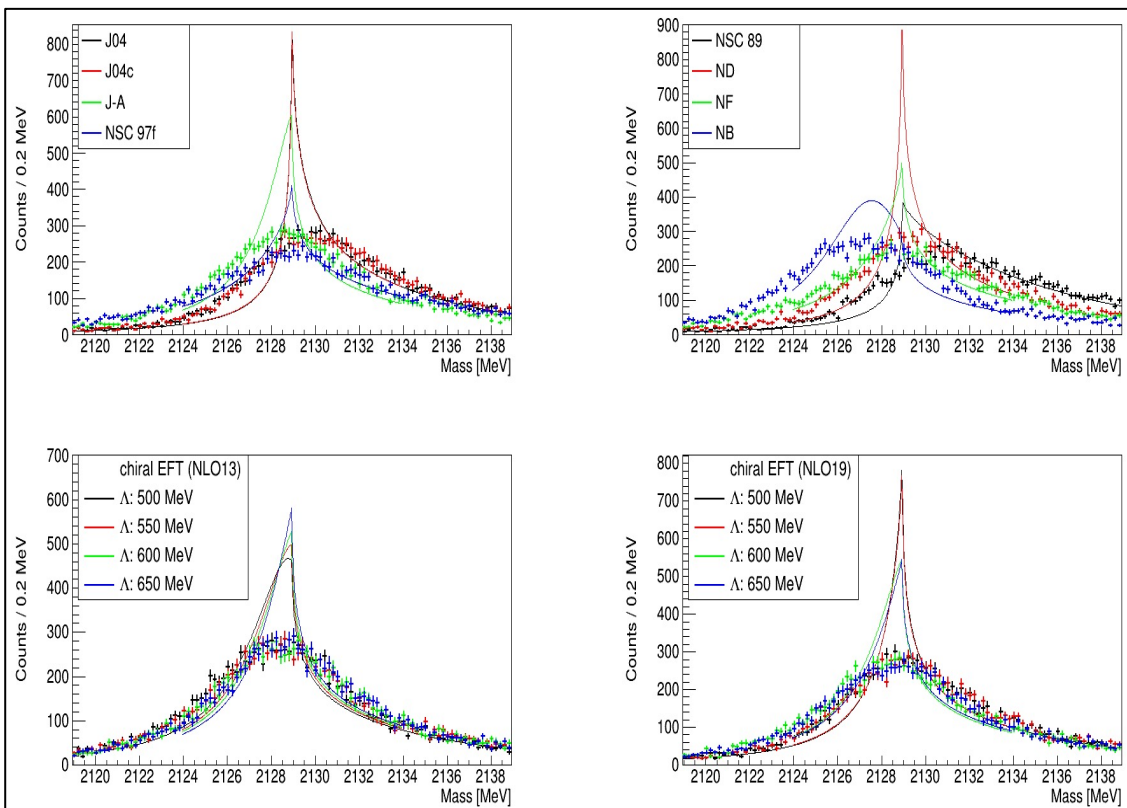
$a < 0$ : Bound state

## $\Sigma N$ -dibaryon !?

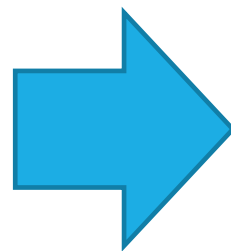
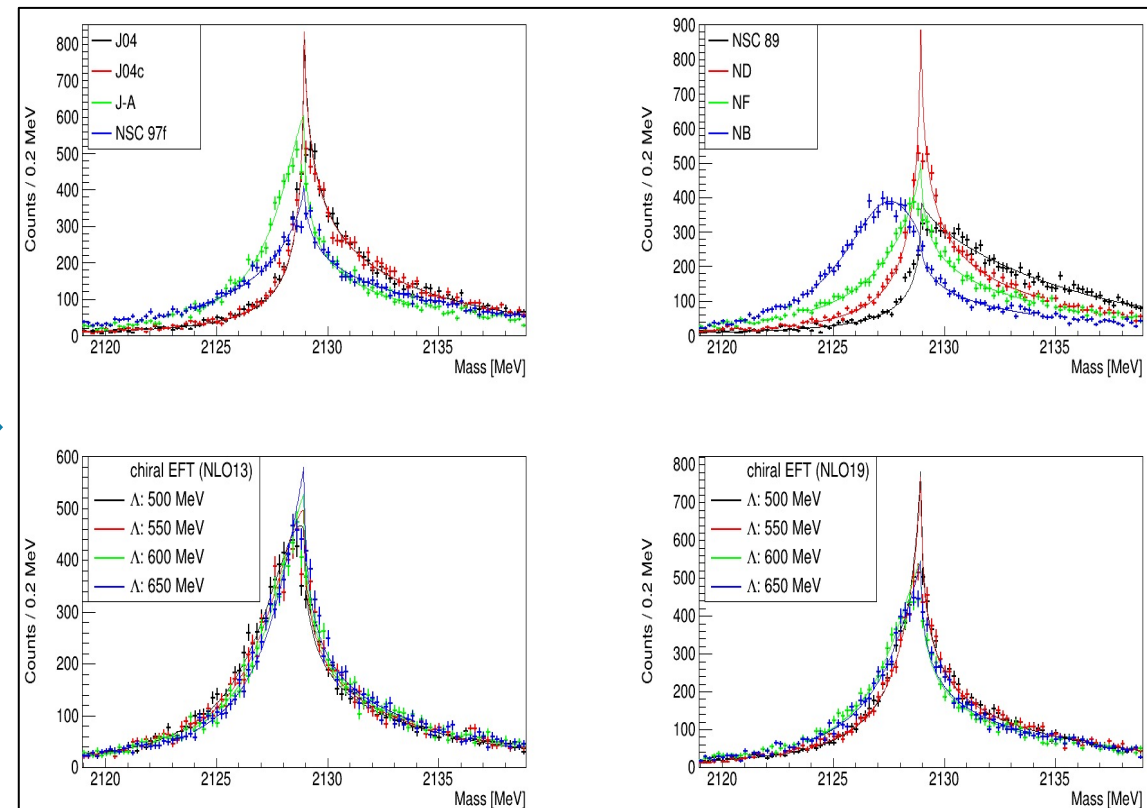
Statistical error  $< 0.3$  fm  
for the  $A_\Sigma = a + ib$  determination

Theoretical Value of  $\Sigma N$  scattering length ( $T=1/2, {}^3S_1$ )

Sensitivity of past experiment ( $\Delta M = 2$  MeV)



Sensitivity of E90 ( $\Delta M = 0.4$  MeV)



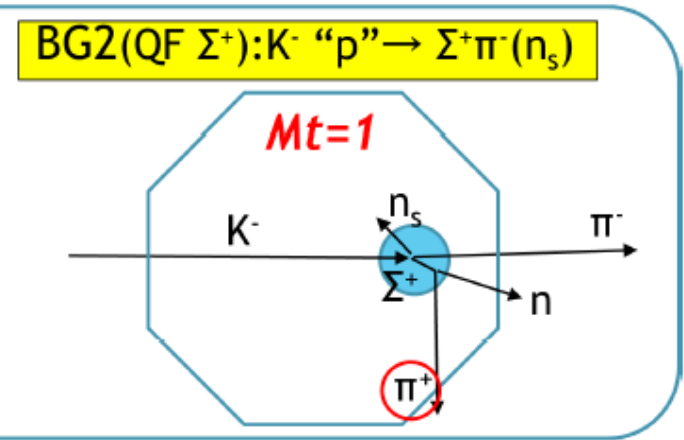
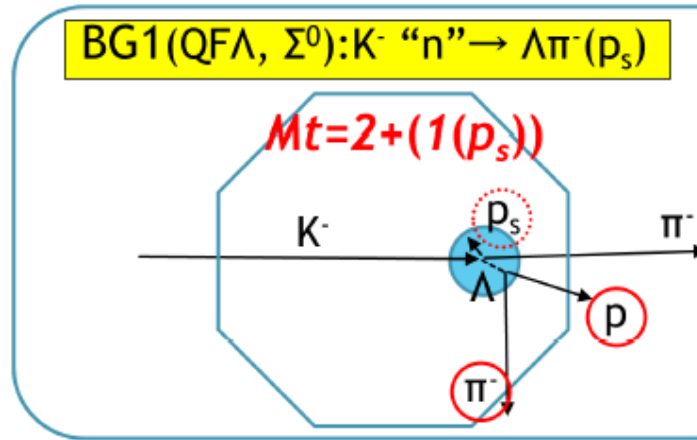
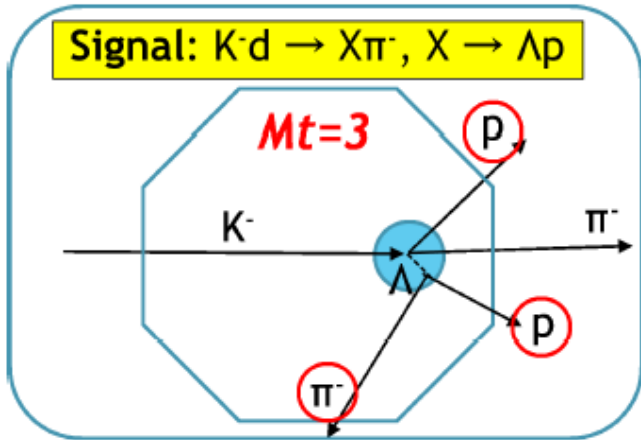
# SUMMARY

- $\Sigma N$  interaction is the important key of the  $B_8 B_8$  interaction and  $(\Lambda, \Sigma)$  hypernuclei.
- “ $\Sigma N$  cusp” can be expressed by the  $\Sigma N$  interaction (scattering length).
  - There are a lot of past experiments to measure the “ $\Sigma N$  cusp” .  
However, the origin of the “ $\Sigma N$  cusp” remains unclear yet.  $\Sigma N$  dibaryon or not?
  - Inflight  $d(K^-, \pi^-)$  reaction has advantage to dedicate ( $T=1/2, {}^3S_1$ ) channel.
- J-PARC E90 will investigate the nature of “ $\Sigma N$  cusp” with the world’s best quality.
  - K1.8 Beam line, S-2S for  $\pi^-$  measurement, and HypTPC for BG suppression.
  - $1.4 \times 10^4$   $\Sigma N$  cusp events are expected in 15 days beam time.
  - 0.4 MeV ( $\sigma$ ) mass resolution will be achieved, 2 times better than past exp.
  - We can deduce scattering length with the statistical error less than 0.3 fm.



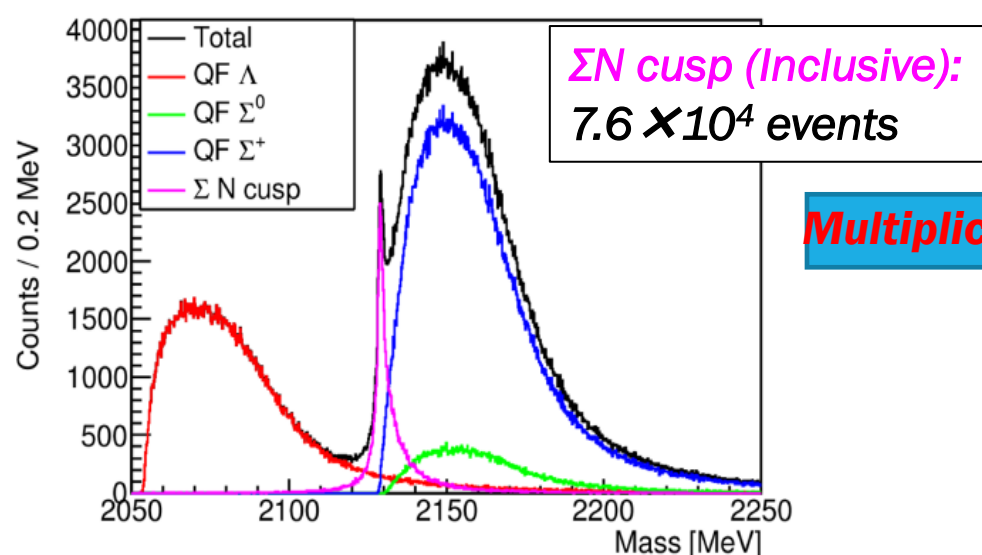
**BACK UP**

# QF BACKGROUND SUPPRESSION BY HYPTPC



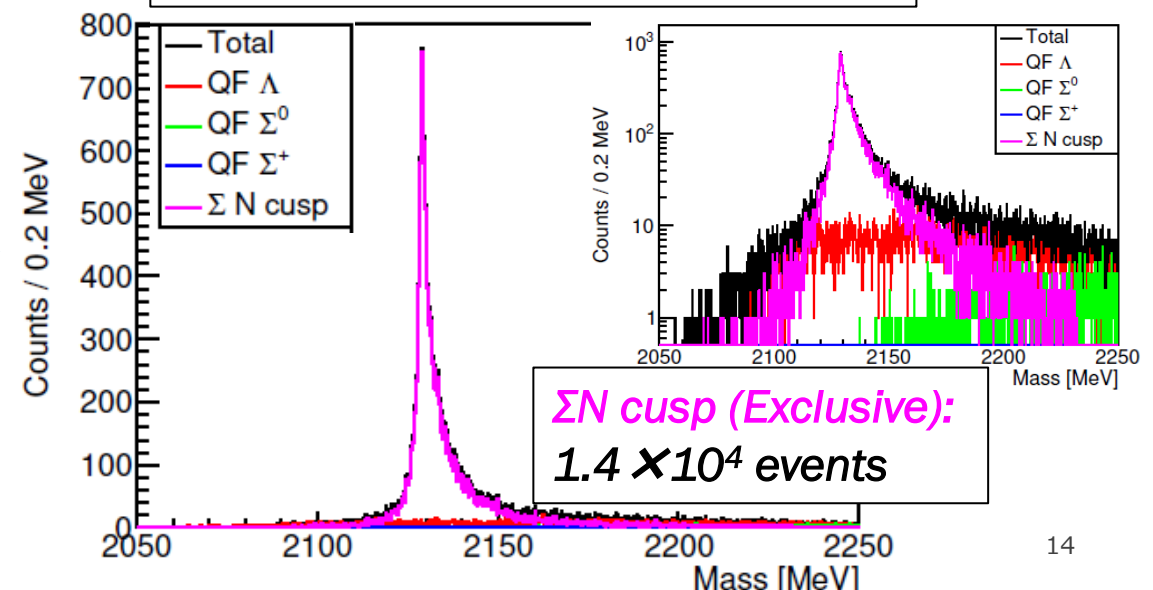
Simulated inclusive spectrum  $d(K^-, \pi^-)$

Expected spectrum for the 15 days beam time



**Multiplicity = 3**

**Multiplicity = 3 without  $(K^-, \pi^-)$**



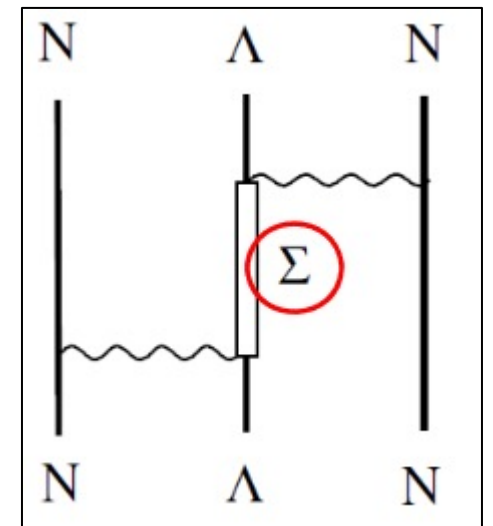
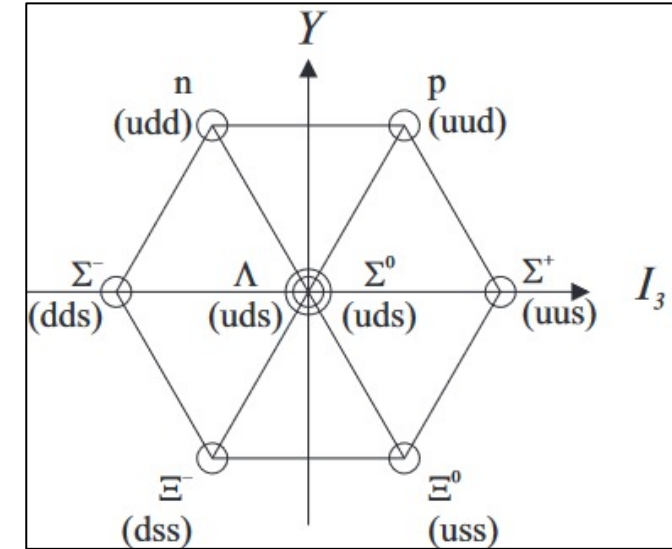


# $\Sigma N$ INTERACTION

- $\Sigma N$  interaction is one of the key to understand  $B_8 B_8$  interaction in  $SU(3)$
- Relation with E40 ( $\Sigma N$  scattering experiment)
  - E40:  $\Sigma N$  scattering ( $p_\Sigma > 470 \text{ MeV}/c$ )  $\rightarrow$  Short range interaction
    - \* $\Sigma N$  scattering experiment in lower momentum is difficult
  - E90: “ $\Sigma N$  cusp”  $\rightarrow$   $\Sigma N$  scattering length (0 energy interaction)
    - $\Sigma N$  scattering length:  $A_\Sigma = a + ib$
    - **a(real part)**  $\rightarrow$  Important for the  $\Sigma$ -hypernuclei
    - **b(imaginary part)**  $\rightarrow$   $\Lambda N$ - $\Sigma N$  coupling strength

$\downarrow$   
 Important for the  $\Lambda$ -hypernuclei

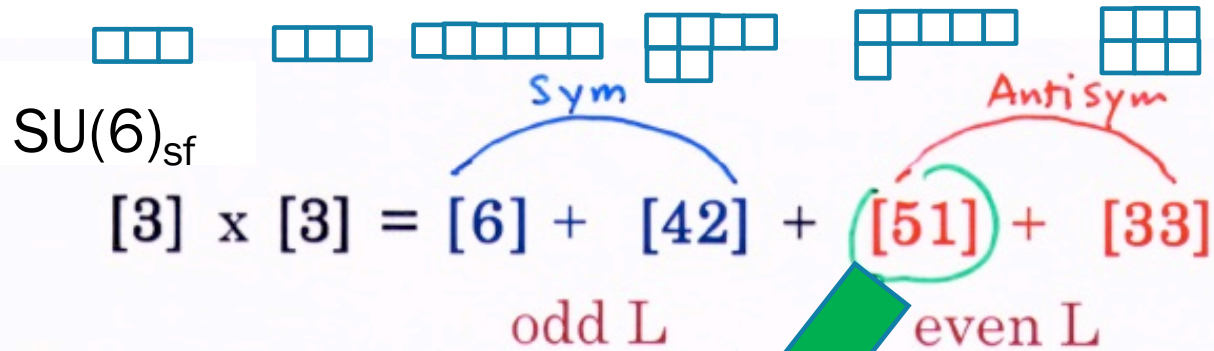
Octet baryon



Complementary

# ΣN INTERACTION

## BB INTERACTION BY QUARK CLUSTER MODEL



$L = 0$

$SU(6)$	orb	color	
$[51]$	$\times [6]$	$\times [222]$	$\neq [1^6]$
	$\times [42]$	$\times [222]$	} = $[1^6]$
$[33]$	$\times [6]$	$\times [222]$	
	$\times [42]$	$\times [222]$	

*Pauli forbidden state ( $L=0$ )*

two-baryon state	[51]	[33]	Pauli
$NN(^1S_0, ^3S_1)$	4/9	5/9	neutral
$\Lambda N(^1S_0, ^3S_1)$	1/2	1/2	neutral
$\Sigma N(T = \frac{1}{2}, ^3S_1)$	1/2	1/2	neutral
$\Sigma N(T = \frac{3}{2}, ^1S_0)$	<b>4/9</b>	<b>5/9</b>	neutral
$\Sigma N(T = \frac{1}{2}, ^1S_0)$	<b>17/18</b>	1/18	unfavored
$\Sigma N(T = \frac{3}{2}, ^3S_1)$	<b>8/9</b>	1/9	unfavored
$H(\Lambda\Lambda - \Xi N - \Sigma\Sigma)$	0	1	avored

*mild repulsion due to the Pauli + CMI*

*Strong Repulsion*

M. Oka Prog.Theor.Phys.Suppl. 137 (2000)



# THRESHOLD CUSP

Cusp structure can be expressed by the **scattering length** (for B'C'),  **$A = a + ib$**

- **B'C' → BC amplitude**

- $f_{B'C',BC} \sim \frac{\sqrt{b}}{1-ikA}$ , **Pole position:**  $k \sim -\frac{i}{A}$

- **(Two body scattering amplitude)**

- $f = \frac{1}{k} \sum_{l=0}^{\infty} (2l+1) e^{i\delta_l} \sin \delta_l P_l(\cos \theta)$

- $\rightarrow (s\text{-wave}) \quad f = \frac{1}{k} e^{i\delta_0} \sin \delta_0 = \frac{1}{k \cot \delta - ik} \rightarrow \frac{a}{1-ika}$

- $k \cot \delta = \frac{1}{a} - \left(\frac{r_{eff}}{2}\right) k^2 + \dots$

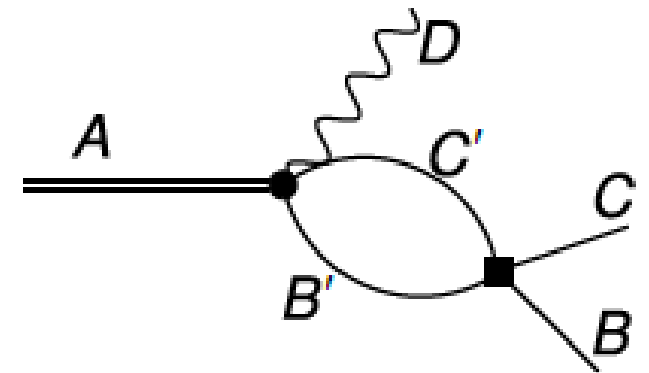
- **Reaction rate (R):**  $\frac{d\sigma^2}{d\Omega dE} \propto |f_{B'C',BC}|^2$

- **Above threshold:**  $R = \frac{4\pi b}{\{(1+kb)^2 + (ka)^2\}}$

- **Below threshold:**  $R = \frac{4\pi b}{\{(1+\kappa a)^2 + (\kappa b)^2\}}$ ,  $k = i\kappa$  (due to analytic continuation)

Reduced mass  
 $\mu = m_{B'} m_{C'} / (m_{B'} + m_{C'})$

$k(\text{relative momentum for B'C'}) \sim \sqrt{2\mu E}$

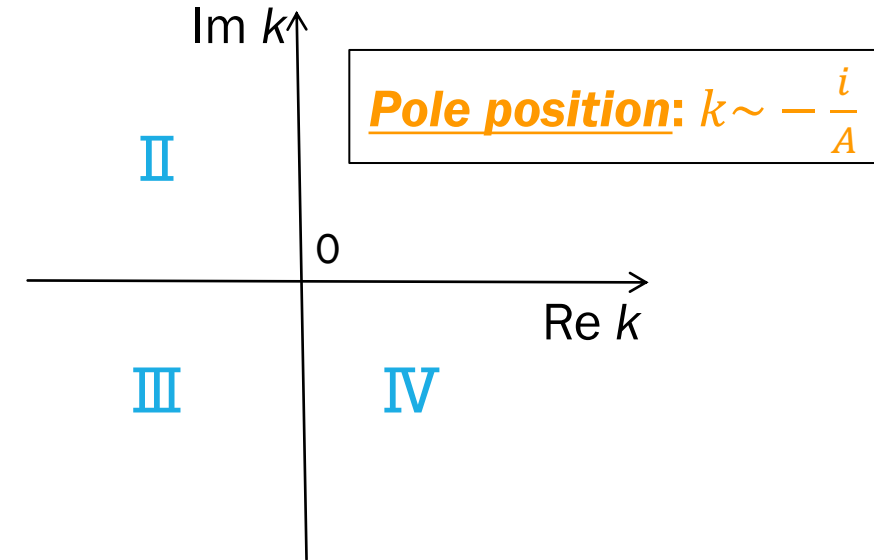
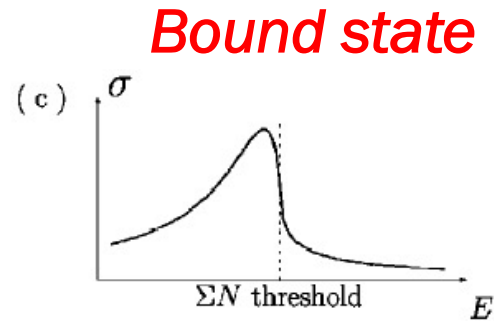
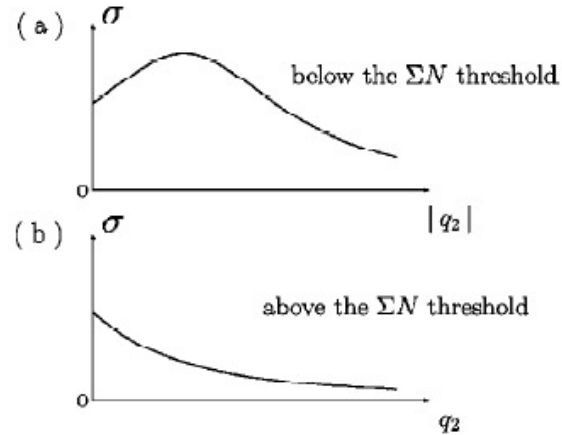


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 $B' = \Sigma, C' = N, B = \Lambda, C = N$

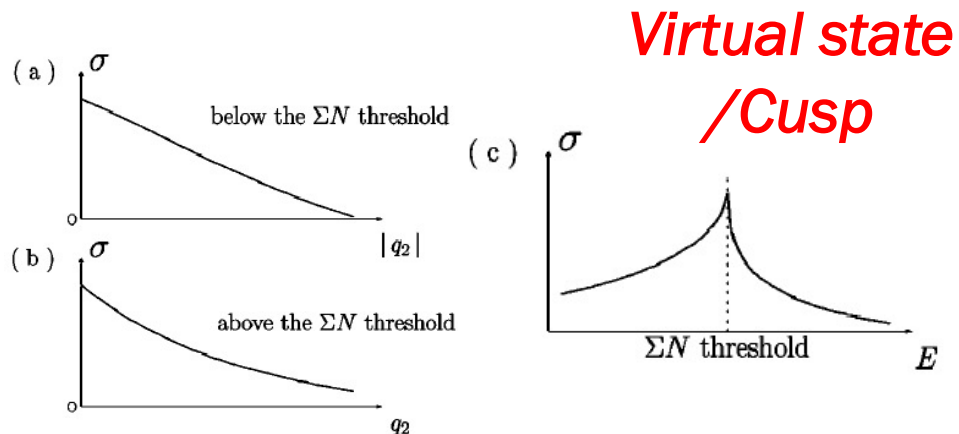
# POLE POSITION vs CROSS SECTION ( $d\sigma/dE$ )

K. Miyagawa and H. Yamamura, PRC 60, 024003 (1999).

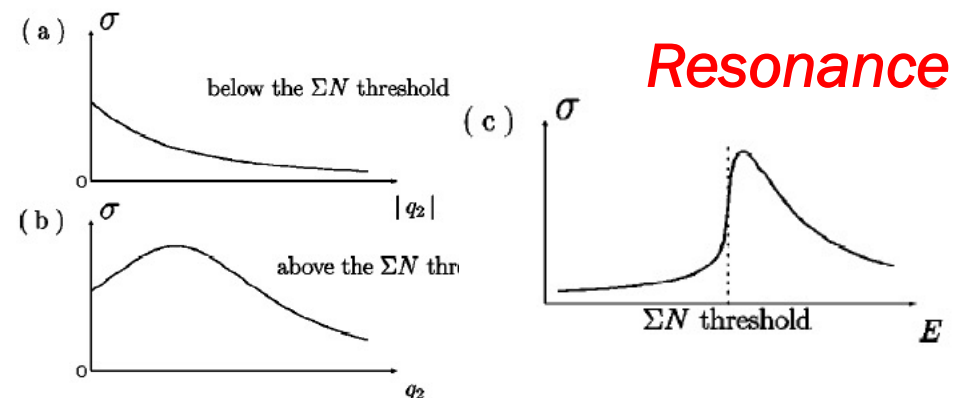
- Pole in (II) quad



- Pole in (III) quad



- Pole in (IV) quad

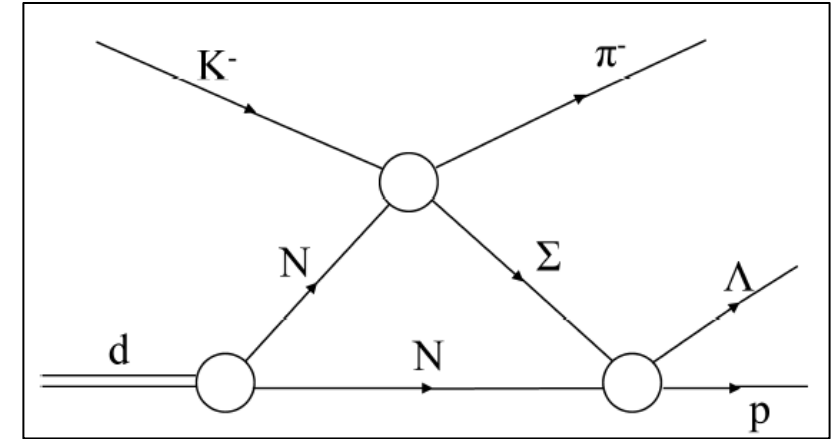


# “ΣN cusp” K<sup>-</sup>d → π<sup>-</sup>Λp reaction

R.H. Dalitz, Nucl. Phys. A354, 101 (1981).

Amplitude of the elementary reaction

$$f(\theta) \propto T(\bar{K}d \rightarrow \Lambda N \pi) \sim \underbrace{T(\bar{K}N \rightarrow \pi \Sigma)}_{\text{Deuteron factor}} \underbrace{F_d(\vec{Q}_\Sigma, k_\Sigma)}_{\text{Amplitude of } \Sigma N \rightarrow \Lambda N \text{ reaction (Important term)}} T(\Sigma N \rightarrow \Lambda N)$$



Deuteron factor

$$F_d(Q_\Sigma, k_\Sigma) = \int \frac{e^{ik_\Sigma r}}{r} e^{i\vec{Q}_\Sigma \cdot \vec{r}} \psi_d(r) d^3r$$

$$\vec{Q}_\Sigma = \vec{q} m_N / (m_N + m_\Sigma)$$

$\psi_d(r)$  : deuteron wave function

Amplitude of  $\Sigma N \rightarrow \Lambda N$  reaction (*Important term*)

$$T(\Sigma N \rightarrow \Lambda N) \propto \frac{\sqrt{b}}{1 - ikA}$$

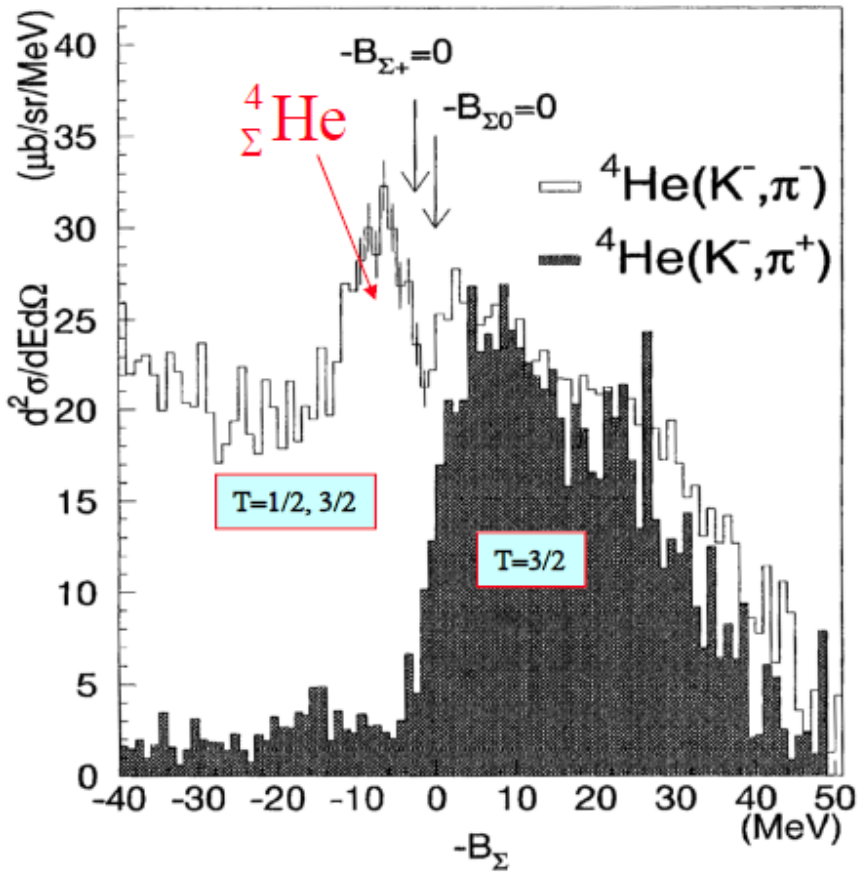
$$\text{Above threshold: } R = \frac{4\pi b}{\{(1+kb)^2 + (ka)^2\}} \sim 1 - 2kb + O(k^2)$$

$$\text{Below threshold: } R = \frac{4\pi b}{\{(1+\kappa a)^2 + (\kappa b)^2\}} \sim 1 - 2\kappa a + O(\kappa^2), \quad k = i\kappa$$

$$(k \sim \sqrt{2\mu E})$$

“ΣN Cusp” can be expressed by the ΣN scattering length ( $A_\Sigma = a + ib$ ) of the  $(T, S) = (1/2, {}^3S_1)$  channel!!

# IMPORTANCE OF $(T, S) = (1/2, 1)$ CHANNEL IN $^4_\Sigma\text{He}$



$\Sigma N$  interaction has strong T and S dependence.  $V_{T=3/2, S=1}$  and  $V_{T=1/2, S=0}$  are expected to be repulsive due to quark Pauli-blocking effect.  $\rightarrow$  No  $\Sigma$ -hypernuclei in large A system.

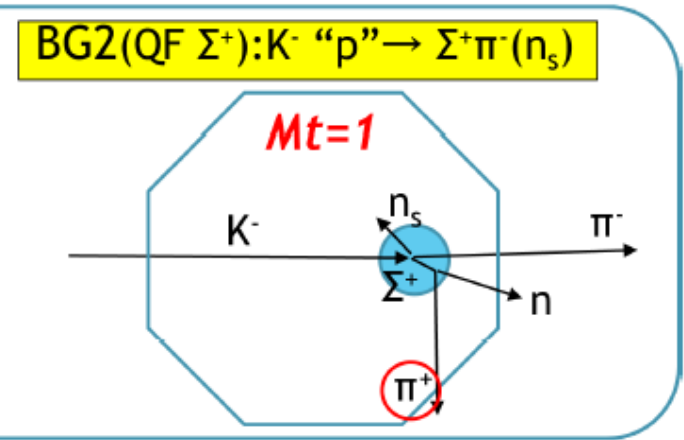
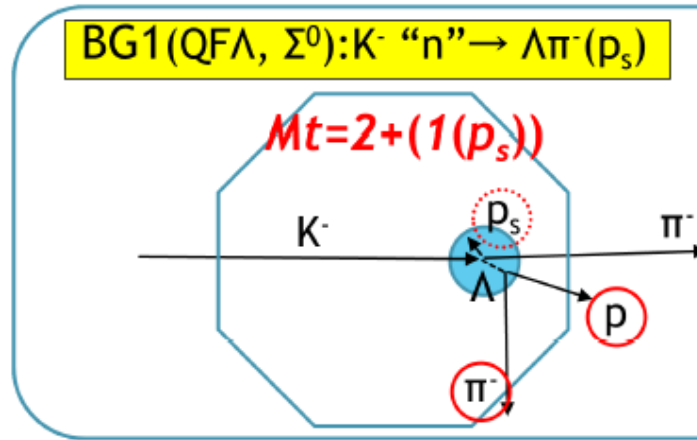
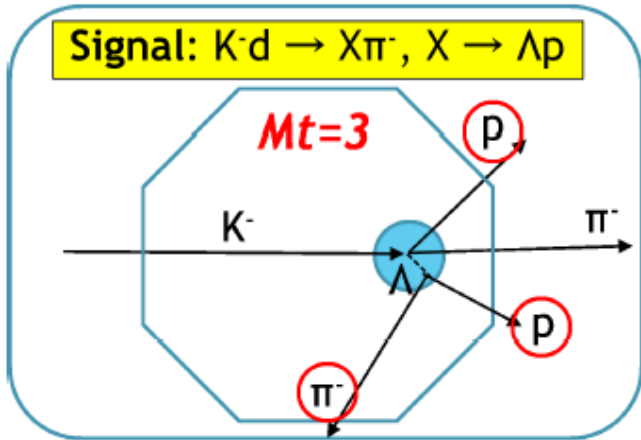
$\Sigma N$  cusp channel:  $V_{T=1/2, S=1}$  (expected to be attractive potential and origin for the  $^4_\Sigma\text{He}$  bound state)

$\Sigma N N N$  (4body) system

$\Sigma N$  (2body) system

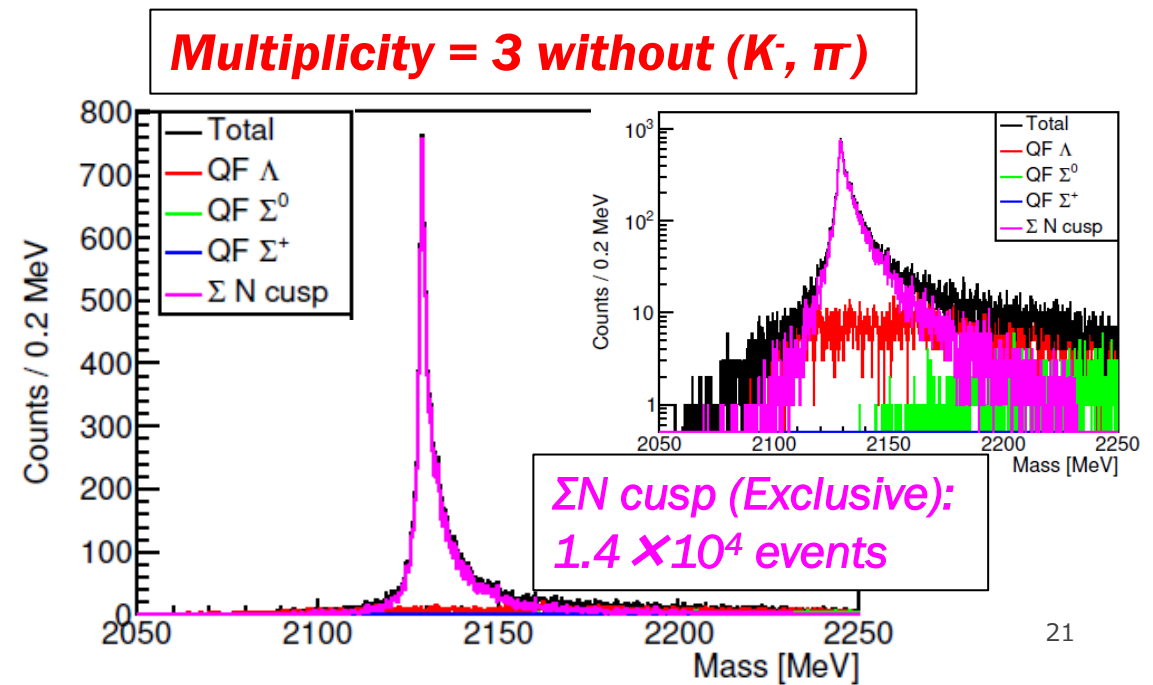
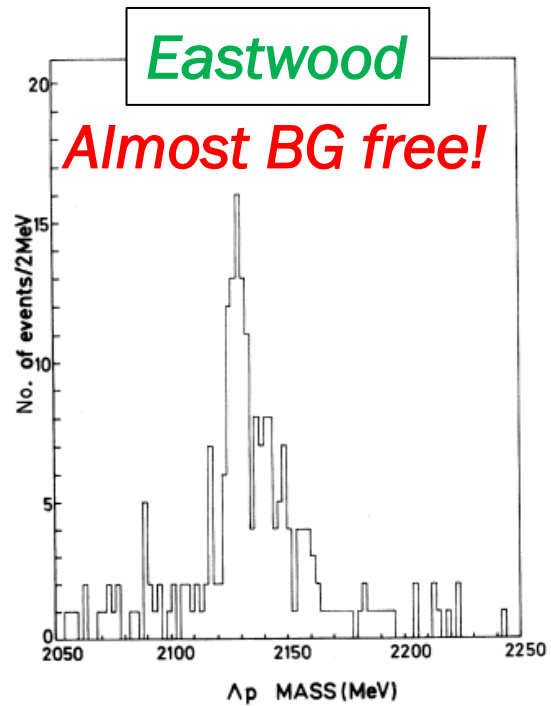
$\Sigma N N N$ (4body) system		$\Sigma N$ (2body) system	
T	S	$\bar{V}_{\Sigma N} (V_{TS})$	
$^4\text{He}(K^-, \pi^+)$ : $\frac{3}{2}$	0	$\frac{5}{18} V_{\frac{3}{2}0}$	$+\frac{1}{2} V_{\frac{3}{2}1}$
$^4\text{He}(K^-, \pi^-)$ : $\frac{1}{2}$	0	$\frac{4}{9} V_{\frac{3}{2}0}$	$+\frac{2}{9} V_{\frac{1}{2}0}$
			$+\frac{1}{18} V_{\frac{1}{2}0}$
			$+\frac{1}{2} V_{\frac{1}{2}1}$

# QF BACKGROUND SUPPRESSION BY HYPTPC



$K^- d \rightarrow \Lambda p \pi^-$   
 @1.45 and 1.65 GeV/c  
 (Bubble chamber)

$\cos\theta_{CM} > 0.9$   
 $p_{\text{proton}} > 150 \text{ MeV}/c$

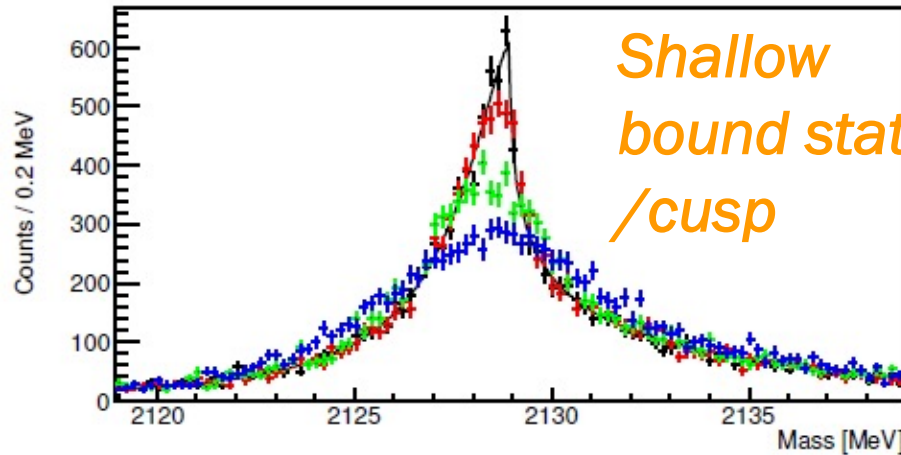


# EXPECTED SPECTRA (**RESOLUTION EFFECT**)

*Good energy resolution is necessary to discuss the cusp shape!!*

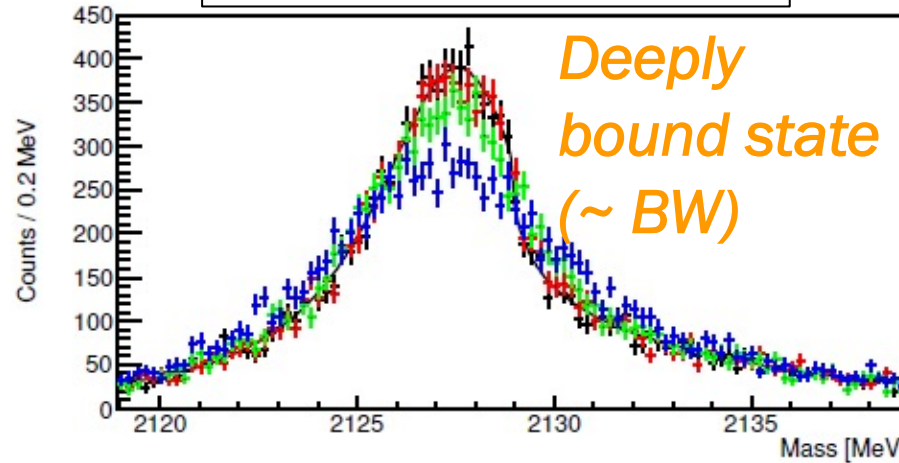
**J-A ( $A_\Sigma = -2.37 + i3.74$  fm)**

*Shallow  
bound state  
/cusp*



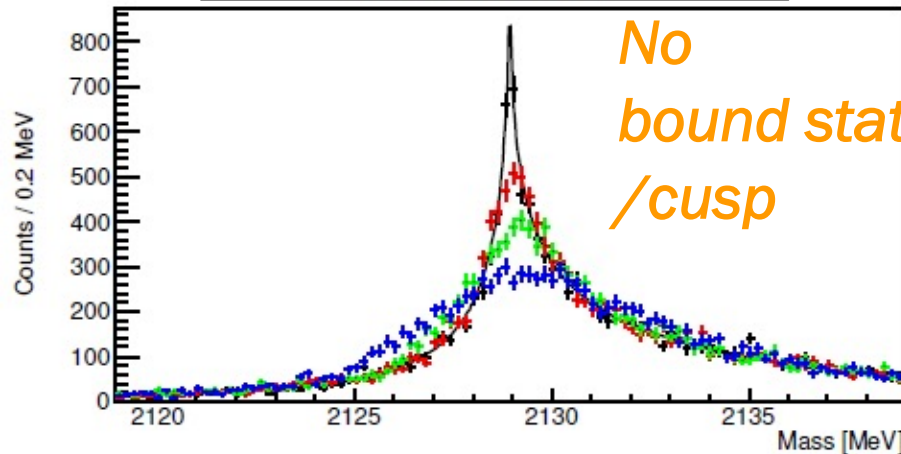
**NB ( $A_\Sigma = -3.00 + i1.8$  fm)**

*Deeply  
bound state  
(~ BW)*



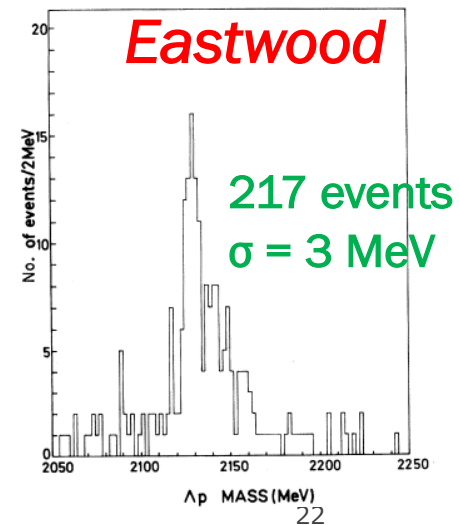
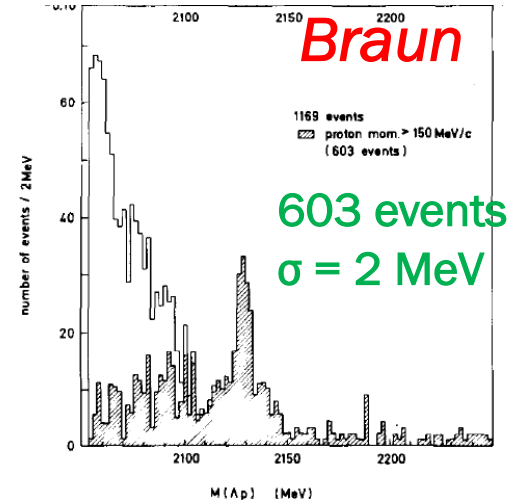
**ND ( $A_\Sigma = 2.06 + i4.64$  fm)**

*No  
bound state  
/cusp*



$1.4 \times 10^4$   
Events  
(P90)

- Ideal
- +  $\Delta M = 0$  MeV
- +  $\Delta M = 0.4$  MeV (P90)
- +  $\Delta M = 1$  MeV
- +  $\Delta M = 2$  MeV (Braun)

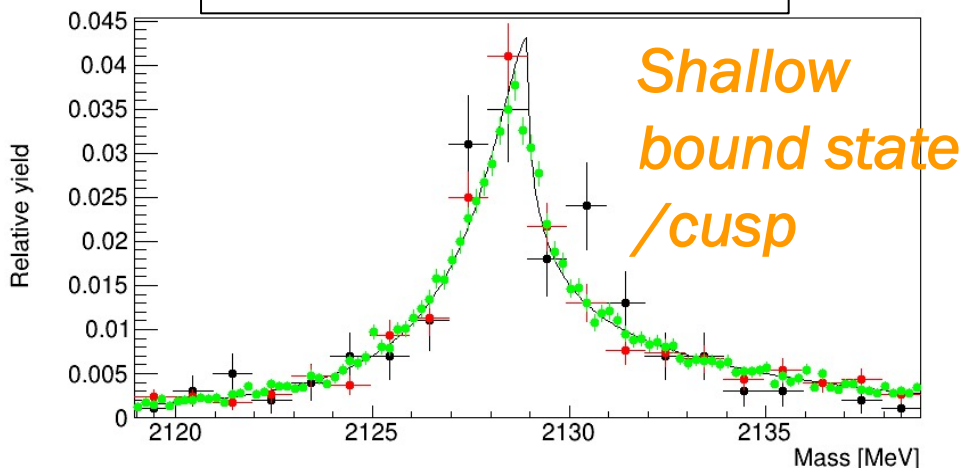




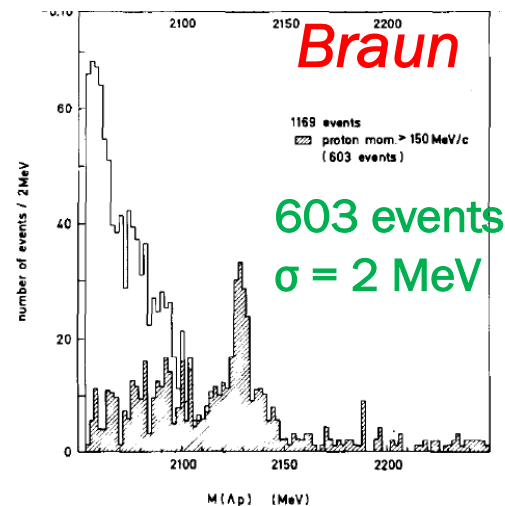
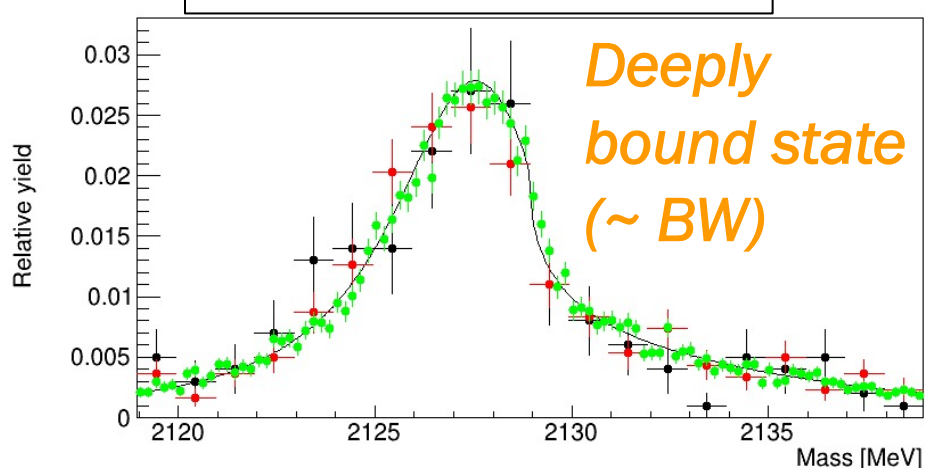
# EXPECTED SPECTRA (STATISTICAL EFFECT)

**>10<sup>4</sup> statistics is necessary!!**

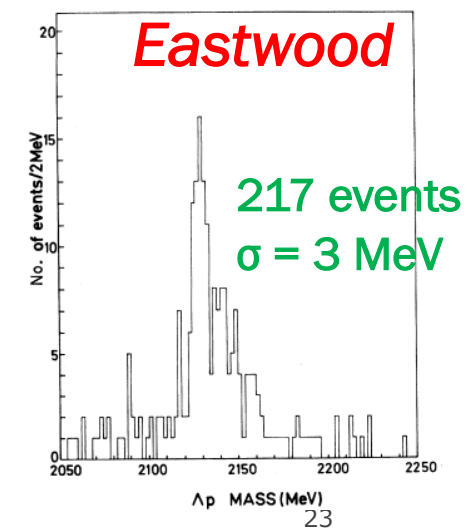
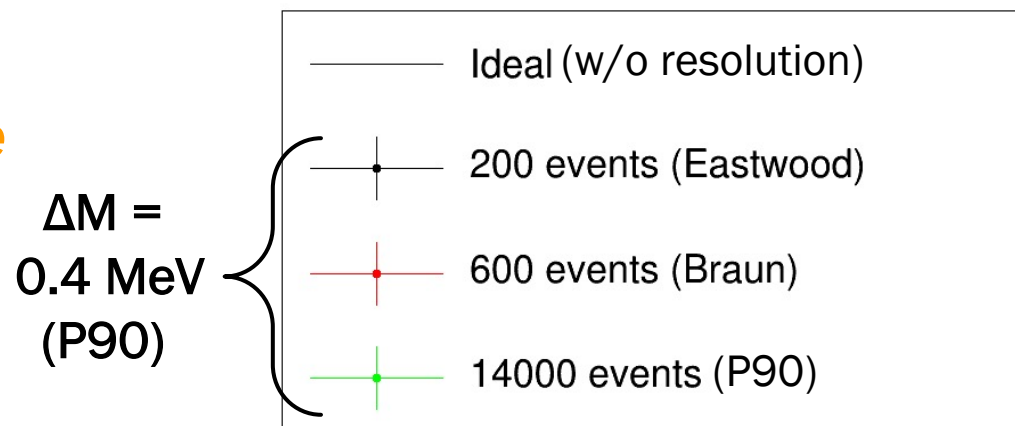
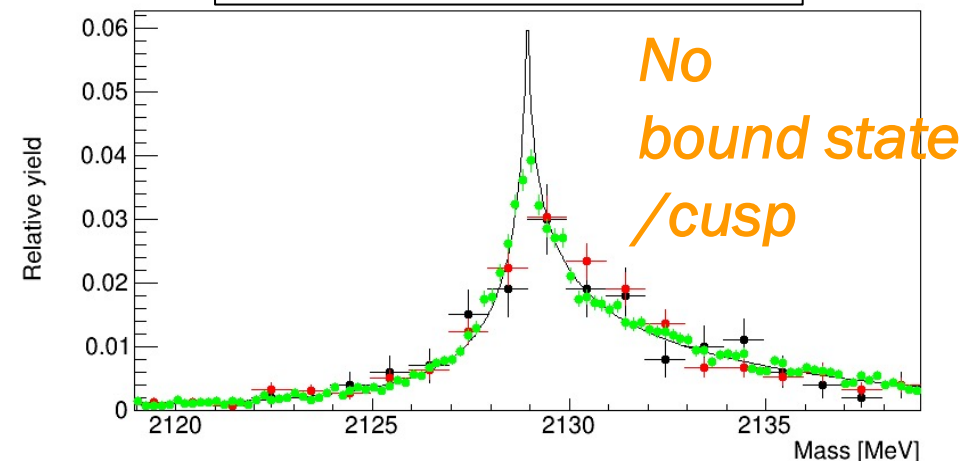
**J-A ( $A_\Sigma = -2.37 + i3.74$  fm)**



**NB ( $A_\Sigma = -3.00 + i1.8$  fm)**

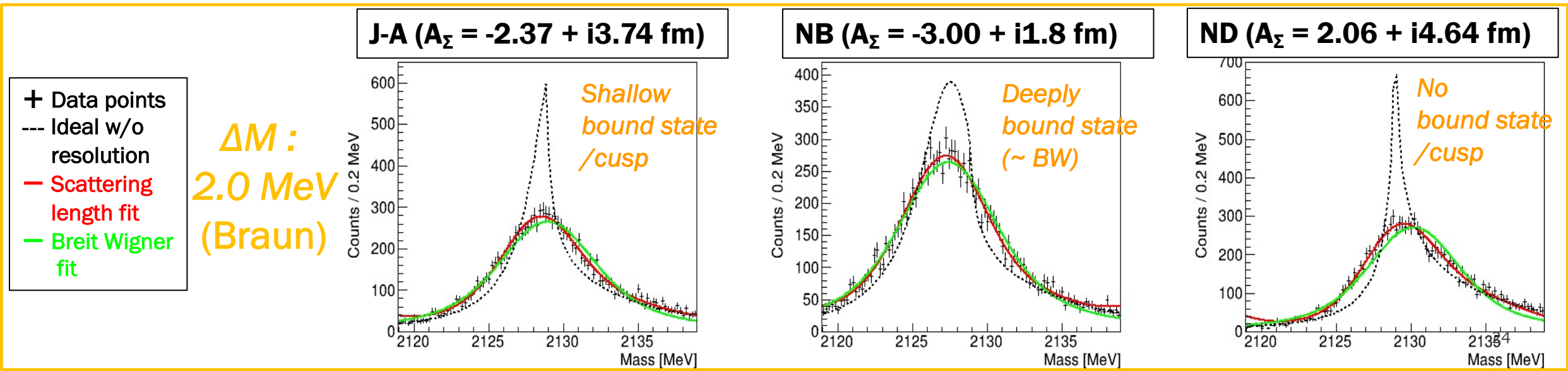
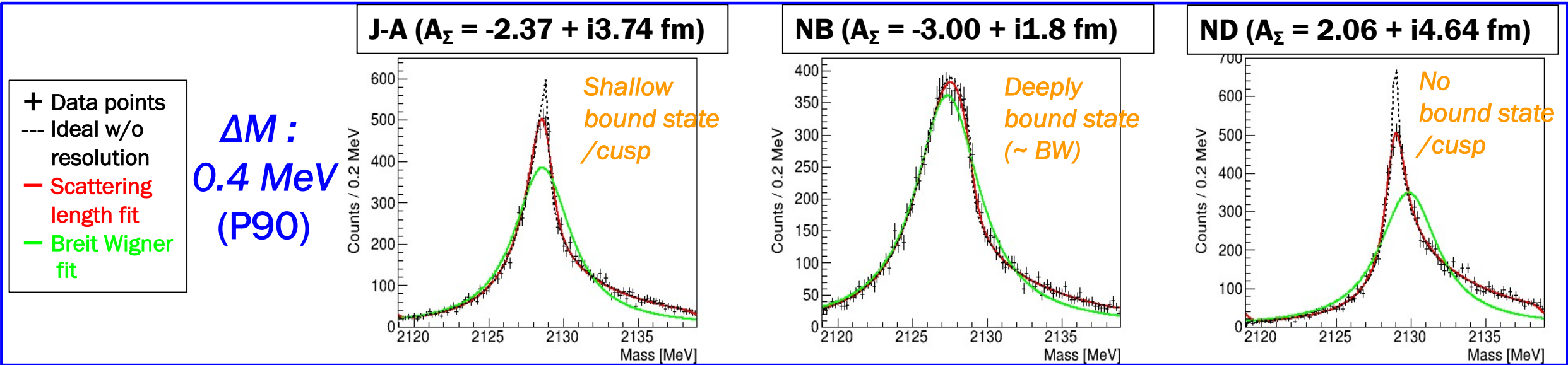


**ND ( $A_\Sigma = 2.06 + i4.64$  fm)**



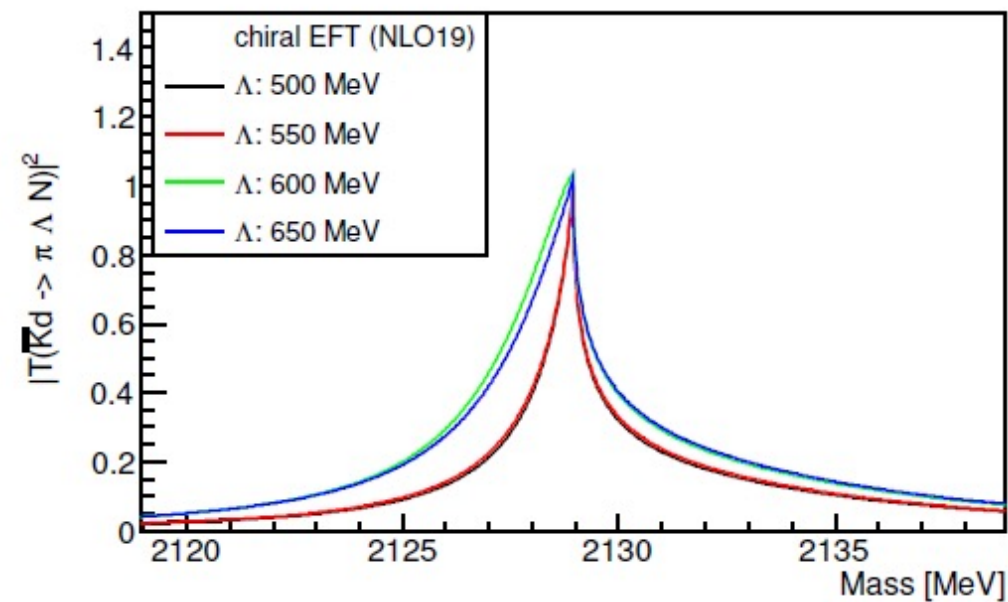
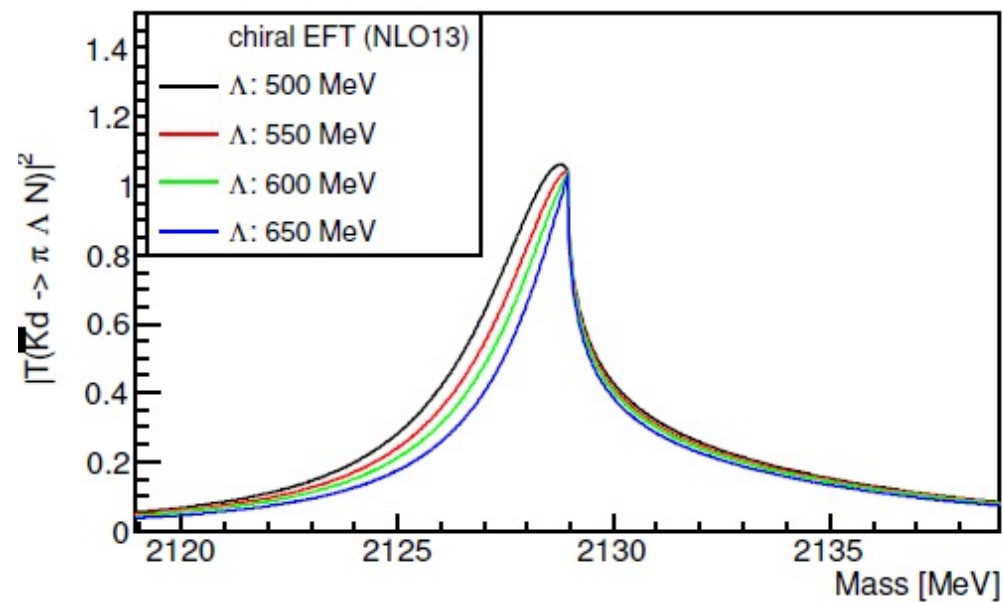
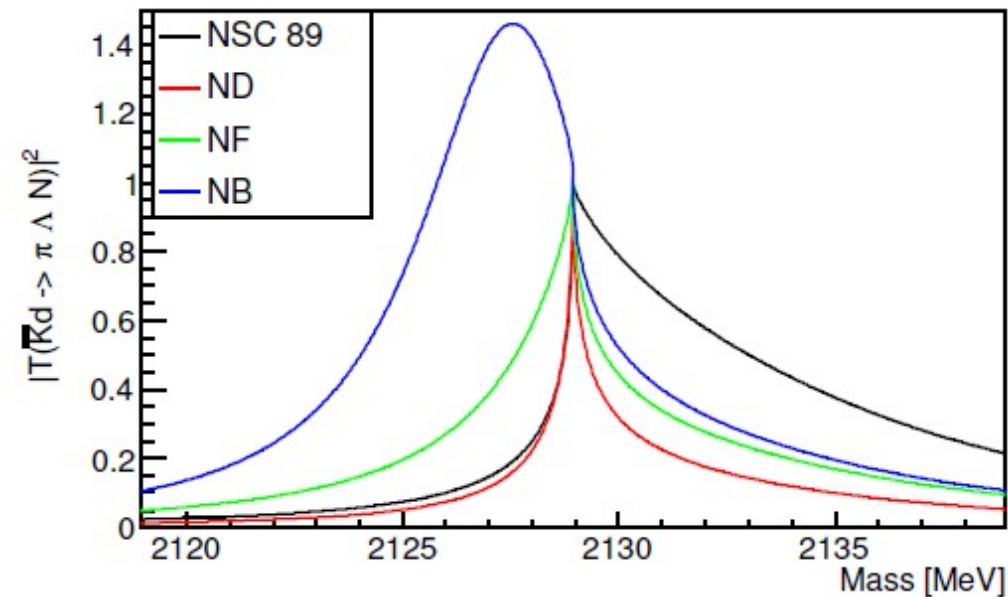
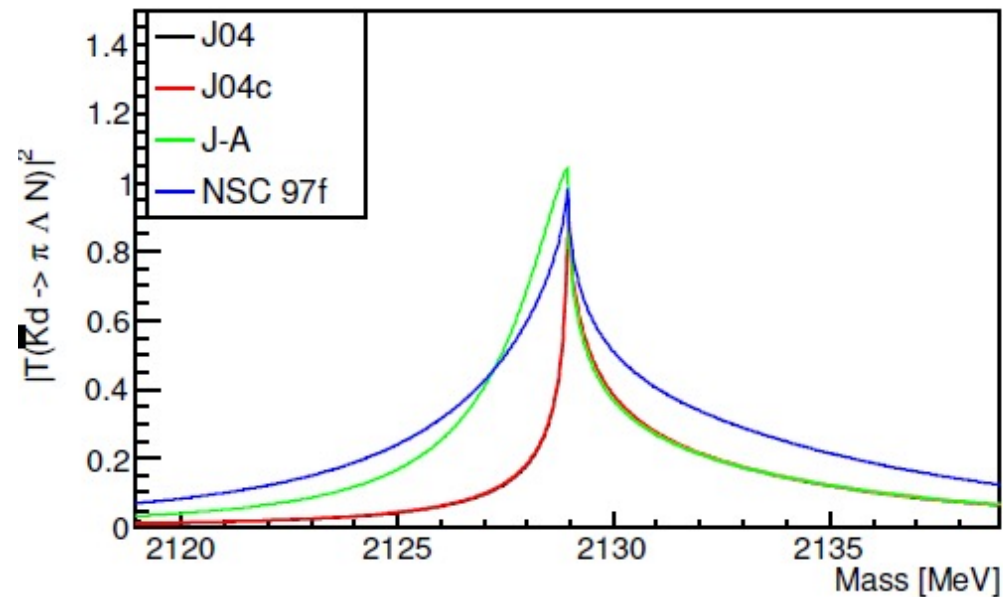
# SCATTERING LENGTH FIT VS BREIT-WIGNER FIT

Significant difference between scattering length fit and Breit-Wigner fit in  $\Delta M = 0.4$  MeV!!  
Statistical error for the scattering length ( $A_\Sigma = a + ib$ ) determination is  $< 0.3$  fm.



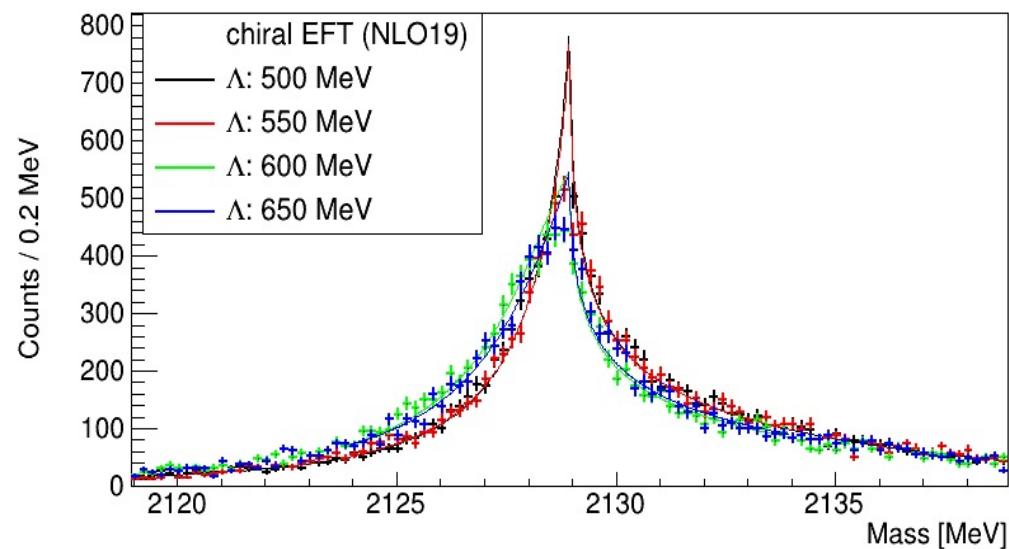
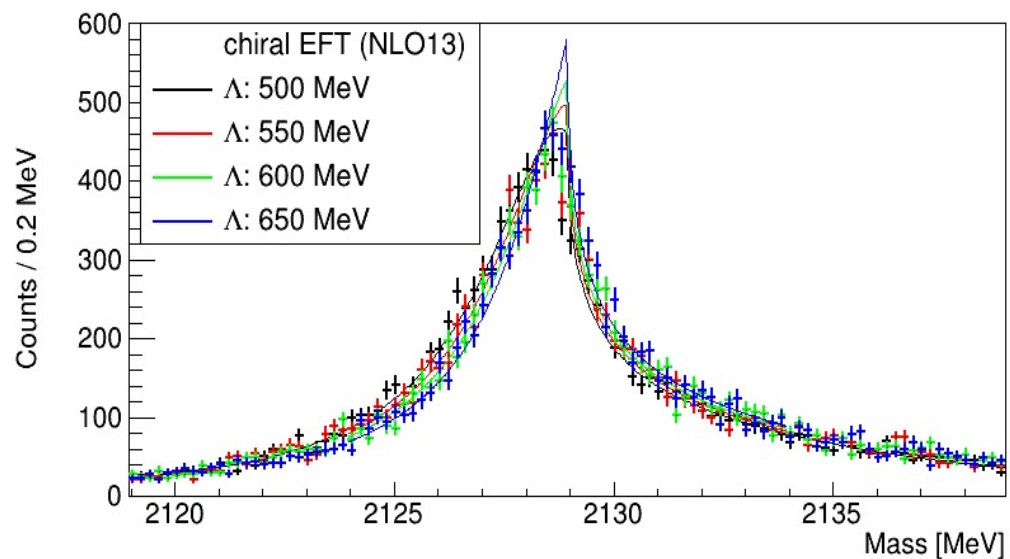
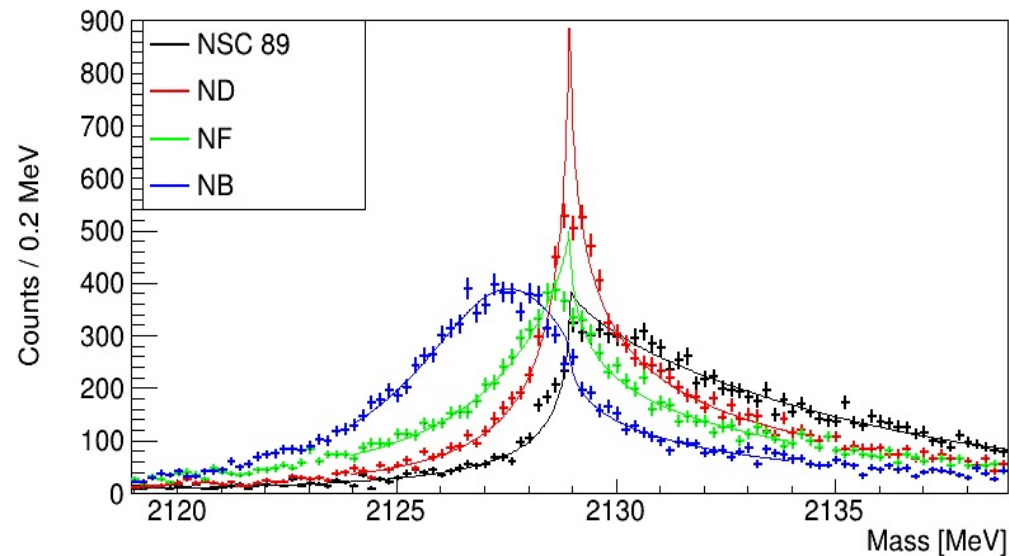
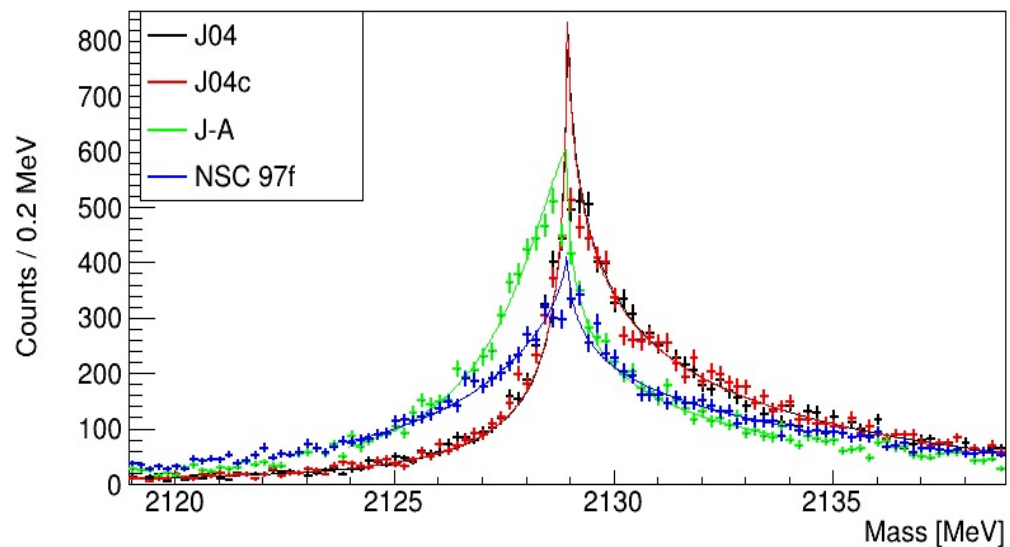


# $\Sigma N(T=1/2, {}^3S_1)$ SCATTERING LENGTH (THEORY)



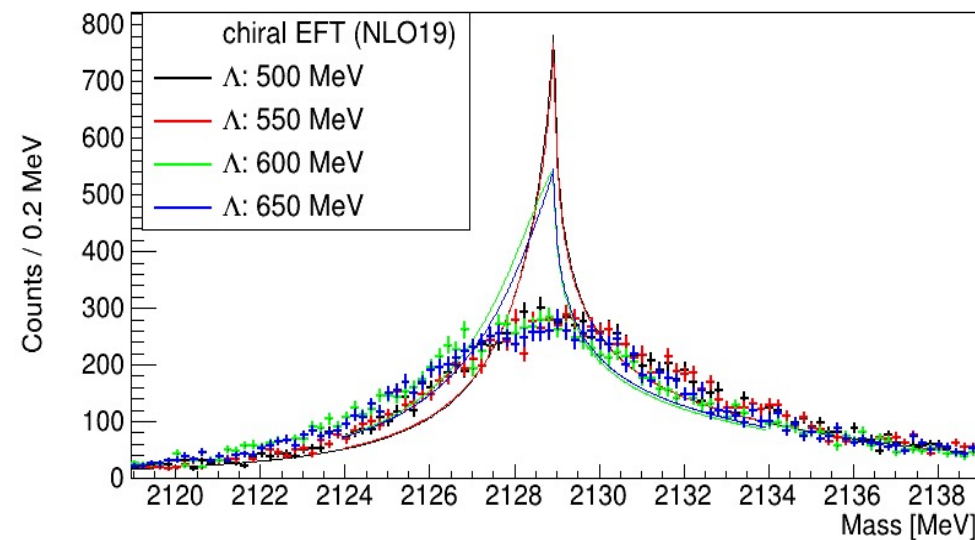
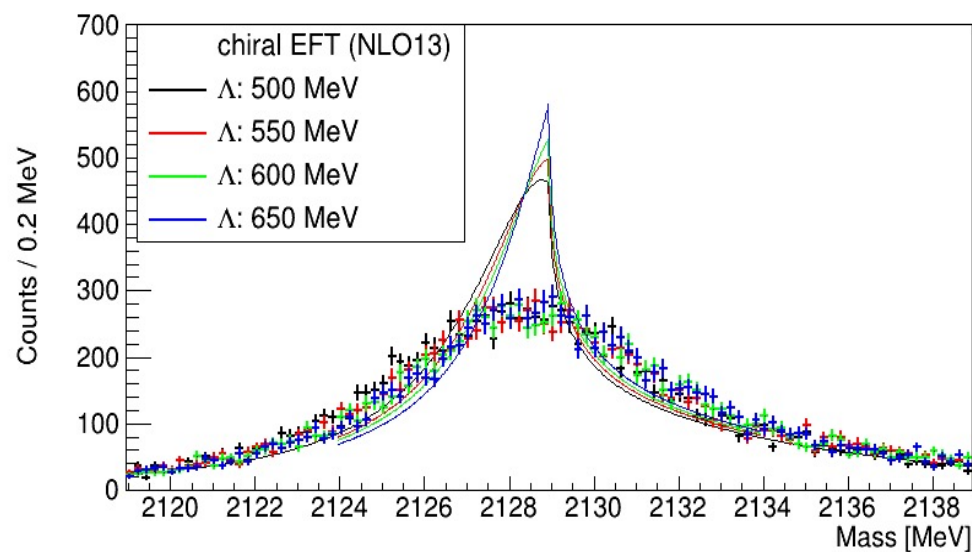
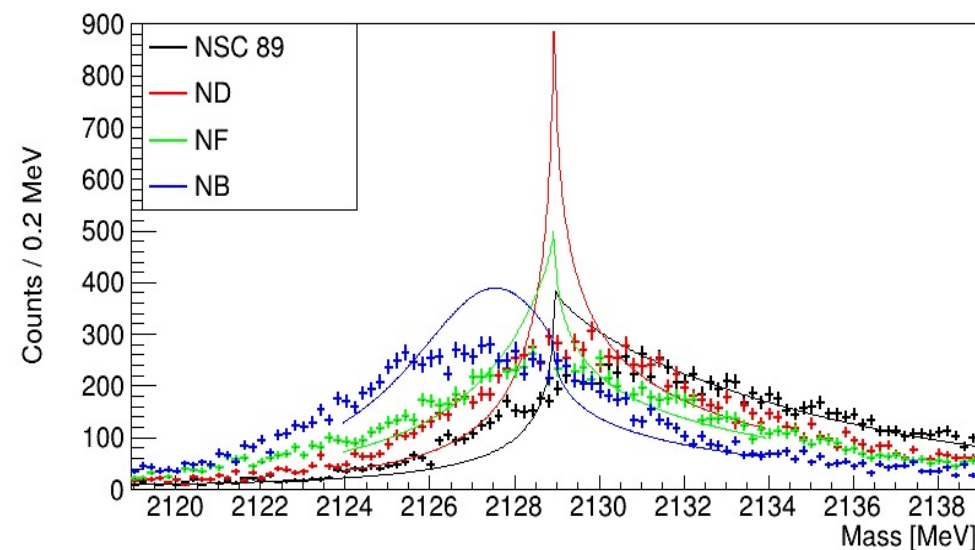
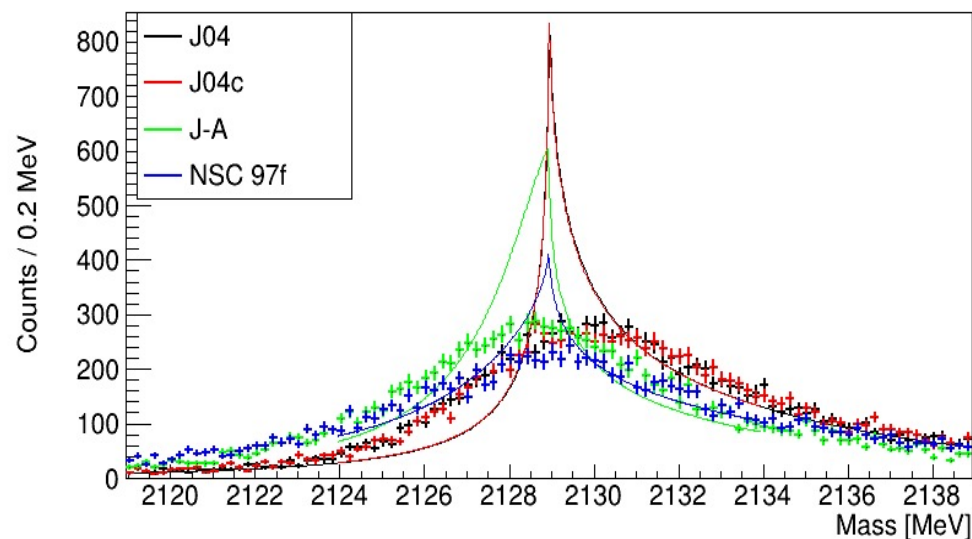
# EXPECTED SPECTRA WITH P90 QUALITY

$\Delta M = 0.4 \text{ MeV}$ ,  $1.4 \times 10^4$  events



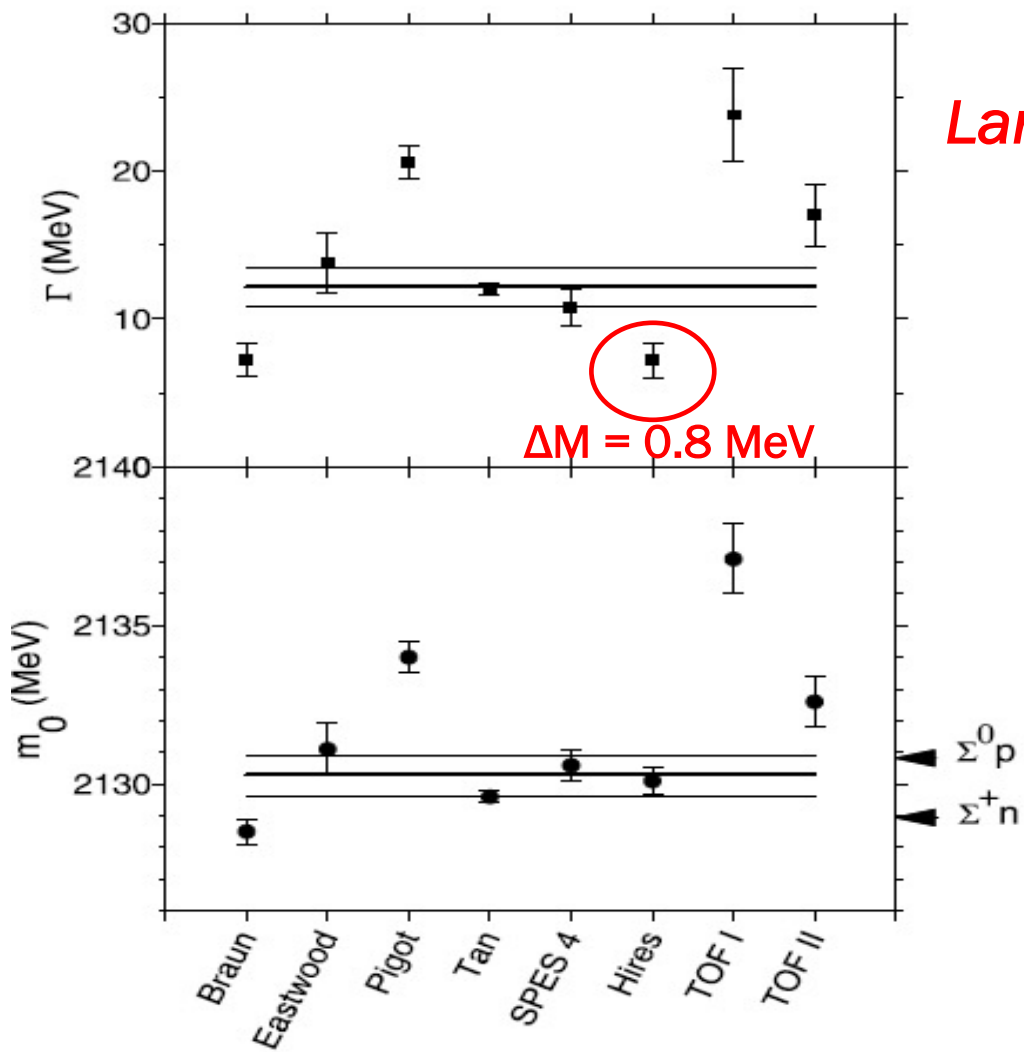
# EXPECTED SPECTRA WITH WORSE RESOLUTION

$\Delta M = 2 \text{ MeV}$ ,  $1.4 \times 10^4$  events



# COMPARISON OF $M_0$ , $\Gamma$ (1 BW FIT)

H. Machner *et al.*, NPA 901, 65 (2013).



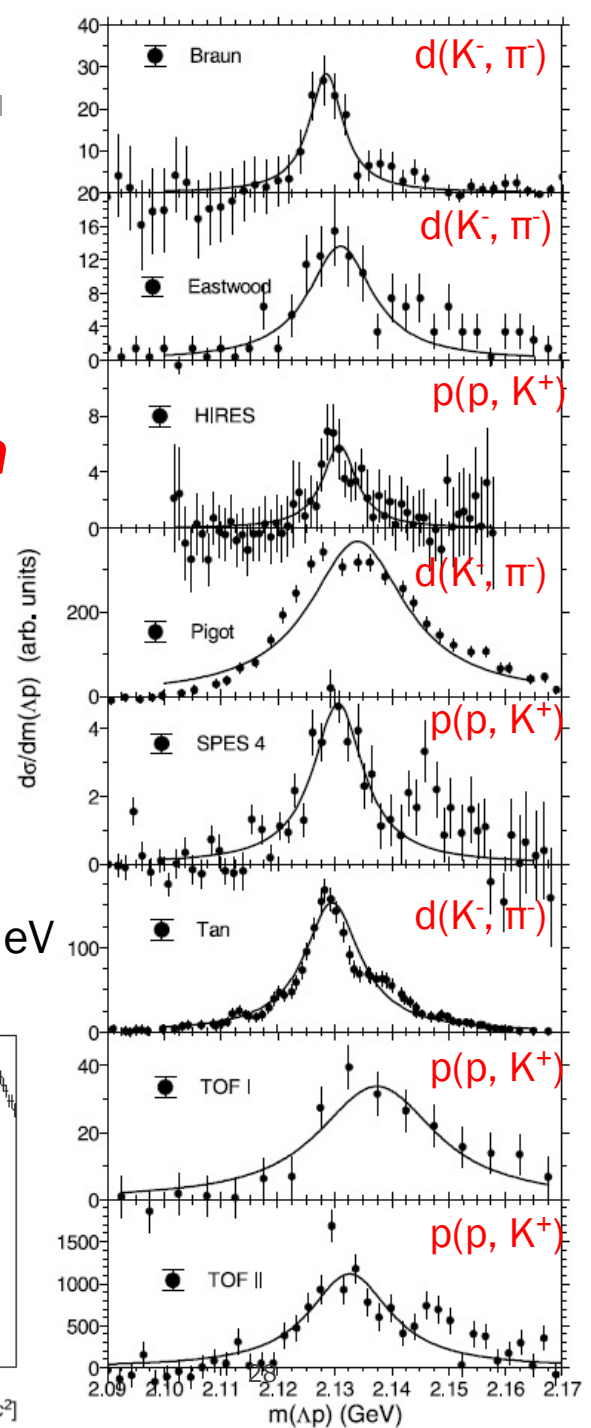
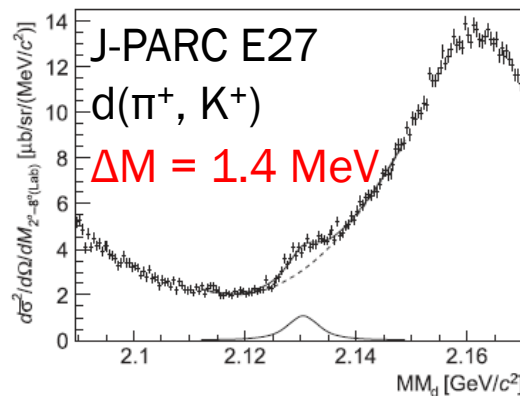
*Large width may be come from the worse resolution.*



*P90 is essential!!*

E27:  $M = 2130.5 \pm 0.4(\text{stat.}) \pm 0.9(\text{syst})$  MeV

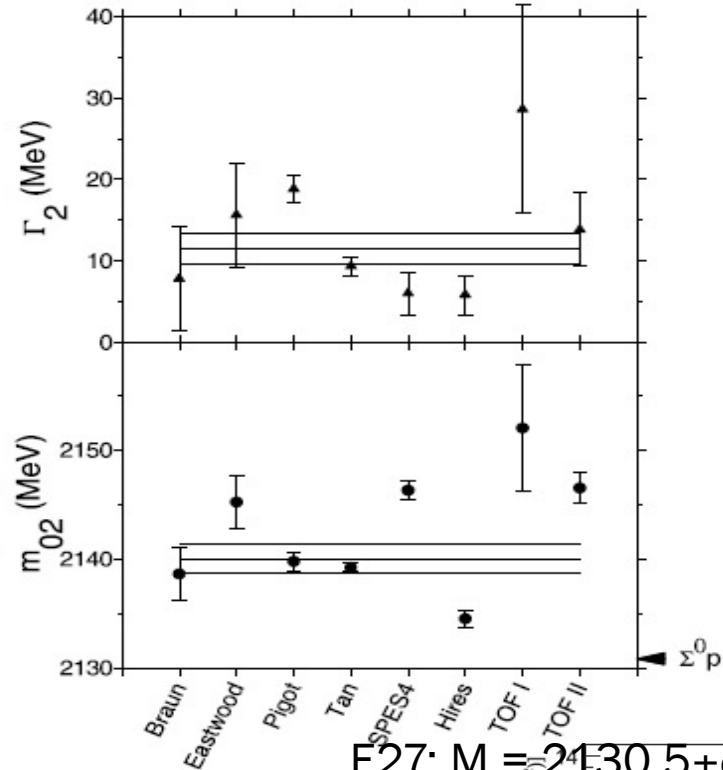
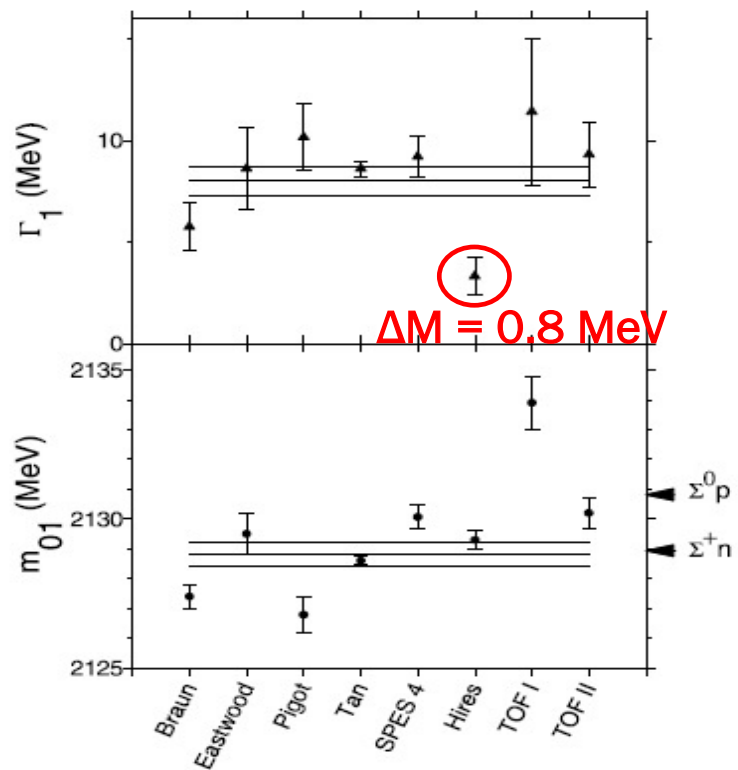
$\Gamma = 5.3^{+1.4}_{-1.2}(\text{stat.})^{+0.6}_{-0.3}(\text{syst.})$  MeV





# COMPARISON OF $M_0$ , $\Gamma$ (2 BW FIT)

H. Machner et al., NPA 901, 65 (2013).



Large width may be come from the worse resolution.



**P90 is essential!!**

E27:  $M = 2130.5 \pm 0.4(\text{stat.}) \pm 0.9(\text{syst.})$  MeV

$\Gamma = 5.3^{+1.4}_{-1.2}(\text{stat.})^{+0.6}_{-0.3}(\text{syst.})$  MeV

J-PARC E27

$\Delta M = 1.4$  MeV

