

# LIFETIME OF THE HYPERTRITON

---

DANIEL GAZDA

Nuclear Physics Institute Řež/Prague

*14th International Conference on Hypernuclear and Strange Particle Physics – HYP2022,  
June 27 – July 1, 2022, Prague, Czech Republic*

# MANY THANKS TO MY COLLABORATORS

- Petr Navrátil  
(TRIUMF, Canada)
- Christian Forssén  
(Chalmers University of Technology,  
Sweden)
- Axel Pérez-Obiol  
(Barcelona Supercomputing Center,  
Spain)
- Avraham Gal, Eli Friedman  
(The Hebrew University of Jerusalem,  
Israel)



## MOTIVATION

---

# MOTIVATION

## Hypertriton

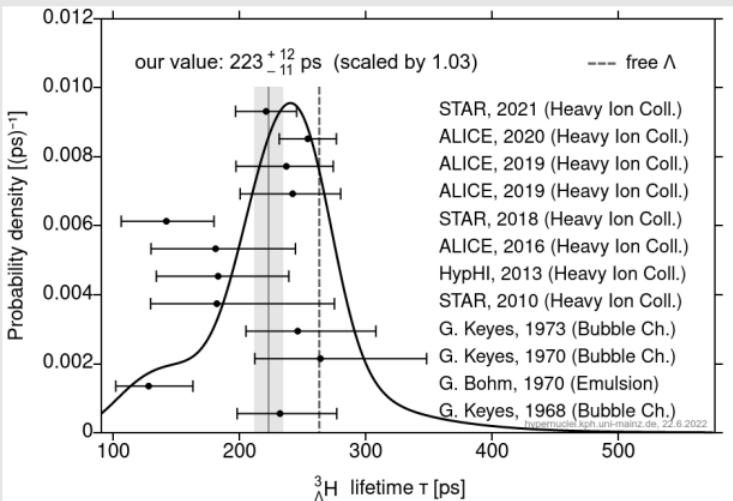
- The lightest bound hypernucleus with spin-parity  $J^\pi = \frac{1}{2}^+$
- A ‘ $\Lambda p n$ ’ bound state with tiny  $\Lambda$  hyperon separation energy  $B_\Lambda = 0.165 \pm 0.044$  MeV, implying a  $\Lambda - {}^2\text{H}$  mean distance  $\approx 10$  fm
- Is expected to have lifetime within few % of the free  $\Lambda$  lifetime  $\tau_\Lambda$ , governed to 99.7% by nonleptonic  $\Lambda \rightarrow N\pi$  weak decay

## Hypertriton lifetime puzzle

- World average of measured  $\tau({}^3\text{H})$  is  $\sim 20\%$  shorter than  $\tau_\Lambda = 263 \pm 2$  ps!
- HypHI  $\tau({}^3\text{H}) = 183^{+42}_{-32} \pm 37$  ps [Rappold et al., NPA 913, 170 (2013)]
- STAR  $\tau({}^3\text{H}) = 142^{+24}_{-21} \pm 29$  ps [Adamczyk et al., PRC 97, 054909 (2018)]
- ALICE  $\tau({}^3\text{H}) = 242^{+34}_{-38} \pm 17$  ps [Acharya et al., PLB 797, 134905 (2019)]
- STAR  $\tau({}^3\text{H}) = 221 \pm 15 \pm 19$  ps [Abdallah et al., PRL 128 202301 (2022)]
- Similar spread with larger uncertainties reported in old emulsion and BC experiments

# MOTIVATION

## $\tau(\Lambda^0 H)$ measurements



[Taken from [hypernuclei.kph.uni-mainz.de](http://hypernuclei.kph.uni-mainz.de)]

## $\tau(\Lambda^0 H)$ calculations

$$1.14 \tau_\Lambda$$

[Rayet, Dalitz (1966)]

$$1.15 \tau_\Lambda$$

[Congleton (1992)]

$$1.06 \tau_\Lambda$$

[Kamada et al. (1998)]

$$1.23 \tau_\Lambda$$

[Gal, Garcilazo (2019)]

$$\approx \tau_\Lambda$$

[Hildenbrand, Hammer (2020)]

# MOTIVATION

## Our aims

- Revisit  $\tau(\Lambda^3\text{H})$  employing  $\Lambda^3\text{H}$  and  $^3\text{He}$  wave functions obtained using **state-of-the-art nuclear and hypernuclear Hamiltonians** (derived from chiral EFT)
- Include **pion final state interactions**, both  $s$ - and  $p$ -wave contributions
- Consider the effect of  **$\Sigma NN$  admixtures** in  $\Lambda^3\text{H}$ , due to  $\Lambda N \leftrightarrow \Sigma N$  coupling
- Study the relation of the **hypertriton lifetime**  $\tau(\Lambda^3\text{H})$  and the  $\Lambda$  hyperon separation energy  $B_\Lambda(^3\text{H})$
- Quantify **theoretical uncertainties**

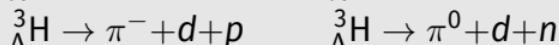
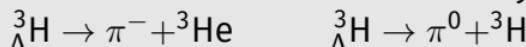
## METHOD

---

# HYPERTRITON LIFETIME

## Hypertriton decay channels

- Mesonic modes due to  $\Lambda \rightarrow \pi N$   
(not Pauli blocked as in heavier hypernuclei)



- Rare non-mesonic modes due to  $\Lambda N \rightarrow NN$



## Hypertriton lifetime $\tau(^3_{\Lambda}\text{H})$

- It is possible to deduce the hypertriton half life  $\tau(^3_{\Lambda}\text{H})$  from two-body  $\pi^-$  decay rate  $\Gamma(^3_{\Lambda}\text{H} \rightarrow ^3\text{He} + \pi^-)$

# HYPERTRITON LIFETIME

From  $\Gamma(\Lambda^3H \rightarrow ^3He + \pi^-)$  to  $\tau(\Lambda^3H)$

- (i) Compute  $\Gamma(\Lambda^3H \rightarrow ^3He + \pi^-)$  two-body  $\pi^-$  decay rate
- (ii) Add contributions from all  $\pi^-$  decay modes using the  $R_3$  branching ratio determined in He BC experiments [Keyes et al., NPB 67, 269 (1973)]

$$\Gamma_{\pi^-}(\Lambda^3H) = \frac{\Gamma(\Lambda^3H \rightarrow ^3He + \pi^-)}{R_3}, \quad R_3 = 0.35 \pm 0.04$$

- (iii) Add contributions from  $\pi^0$  decay modes using  $\Delta I = 1/2$  rule:

$$\Gamma_\pi(\Lambda^3H) = \frac{3}{2} \Gamma_{\pi^-}(\Lambda^3H)$$

- (iv) Add  $\approx 1.5\%$  contribution from  $\Lambda N \rightarrow NN$

[Rayet, Dalitz, NC 46A, 786 (1966); Golak et al., PRC 55, 2196 (1997);

Pérez-Obiol et al., JPCPS 1024, 012033 (2018)]

- (v) Add  $\approx 0.8\%$  contribution from  $\pi NN \rightarrow NN$  pion true absorption estimated from pion optical potential

# HYPERTRITON LIFETIME

**Two-body  $\pi^-$  decay rate**  $\Gamma(\Lambda \rightarrow {}^3\text{He} + \pi^-)$

$$\frac{\Gamma_{\Lambda \rightarrow {}^3\text{He} + \pi^-}}{(G_F m_\pi^2)^2} = 3 \frac{q}{\pi} \frac{M_{{}^3\text{He}}}{M_{{}^3\text{He}} + \omega_\pi} \left[ \mathcal{A}_\Lambda^2 |F^{PV}(\vec{q})|^2 + \mathcal{B}_\Lambda^2 |F^{PC}(\vec{q}, \vec{\sigma})|^2 \left( \frac{k_\pi}{2\bar{M}} \right)^2 \right]$$

with  $\Lambda \rightarrow p\pi^-$  parity-violating  $\mathcal{A}_\Lambda$  and parity-conserving  $\mathcal{B}_\Lambda$  amplitudes accompanied by (hyper)nuclear form factors

$$F^j(\vec{q}, \vec{\sigma}) = \langle \Psi_{{}^3\text{He}} \phi_\pi | \mathcal{O}^j(\vec{q}, \vec{\sigma}) | \Psi_{\Lambda} \rangle$$

$$\mathcal{O}^{PV} = 1,$$

$$\mathcal{O}^{PC} = \vec{\sigma} \cdot \hat{q}$$

- $\phi_\pi$  – pion wave function (plane or distorted wave)
- $\Psi_{{}^3\text{He}}, \Psi_{\Lambda}^3$  –  ${}^3\text{He}$ ,  ${}^3\text{H}$  ground-state wave functions from *ab initio* no-core shell model (NCSM)

## **INPUT WAVE FUNCTIONS FROM NCSM**

---

# AB INITIO NO-CORE SHELL MODEL

Quasi-exact method to solve the nuclear many-body problem:

$$\left[ \sum_{i \leq A} \frac{\hat{\mathbf{p}}_i^2}{2m_i} + \sum_{i < j \leq A-1} \hat{V}_{NN;ij} + \sum_{i < j < k \leq A-1} \hat{V}_{NNN;ijk} + \sum_{i < j = A} \hat{V}_{NY;ij} \right] \Psi = E\Psi$$

- Hamiltonian is diagonalized in a *finite A-particle* harmonic oscillator (HO) basis

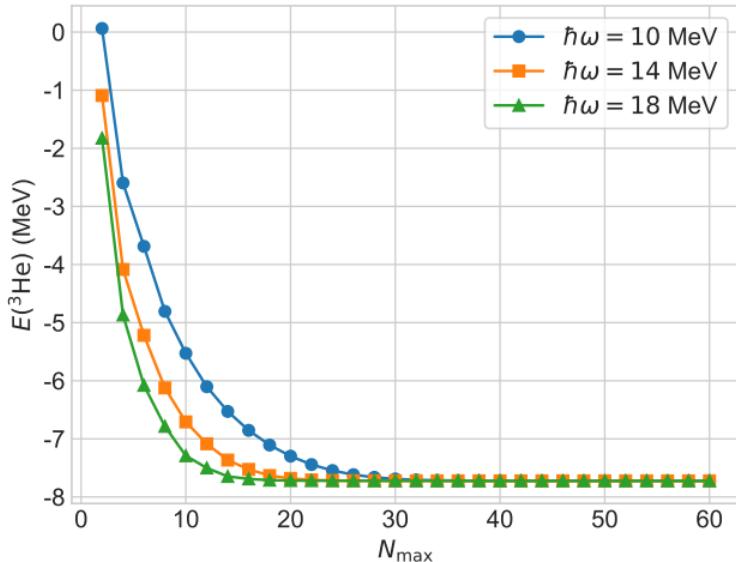
$$\Psi(\mathbf{r}_1, \dots, \mathbf{r}_A) = \sum_{n \leq N_{\max}} \Phi_{n \hbar \omega}^{HO}(\mathbf{r}_1, \dots, \mathbf{r}_A)$$

- Systematically improvable: converges to exact results for  $N_{\max} \rightarrow \infty$

## Input Hamiltonians

- Chiral NNLO<sub>sim</sub>  $NN + NNN$  potential family  
[Carlsson et al., PRX 6, 011019 (2016)]
- Chiral LO YN potential [Polinder et al., NPA 779, 244 (2006)]
  - $\Lambda N - \Sigma N$  mixing explicitly taken into account
  - Coupled-channel  $\Lambda$ -hypernucleus– $\Sigma$ -hypernucleus problem!

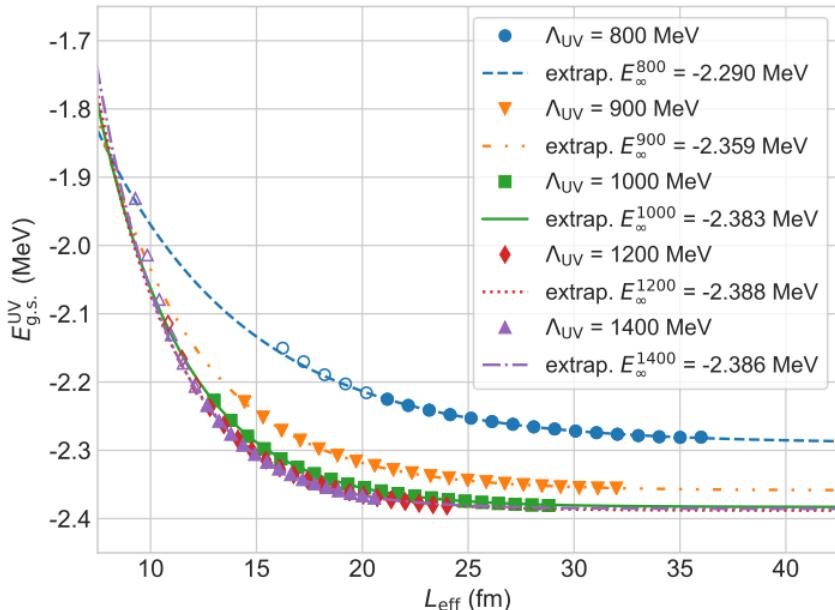
# INPUT ${}^3\text{He}$ WAVE FUNCTIONS FROM NCSM



**Figure 1:**  ${}^3\text{He}$  g.s. energies calculated using NCSM for several HO frequencies  $\omega$  as functions of the model-space truncation  $N_{\max}$ .

- $10^{-3}$  MeV accuracy reached for  $N_{\max} \sim 30$  for a wide range of frequencies  $\omega$
- $E({}^3\text{He}) = -7.723 \text{ MeV}$  for NNLO<sub>sim</sub><sup>(500,290)</sup> (exp.  $-7.718(19) \text{ MeV}$ )

# INPUT ${}^3\Lambda$ WAVE FUNCTIONS FROM NCSM



**Figure 2:**  ${}^3\Lambda$  g.s. energies calculated using NCSM for several  $\Lambda_{UV}$  cutoffs as functions of the IR length scale  $L_{eff}$ .

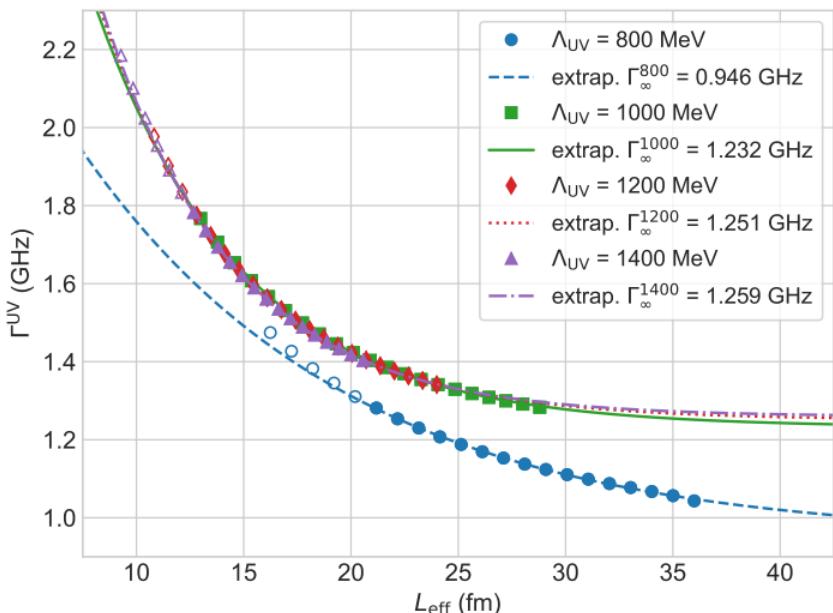
- $L_{eff}$ ,  $\Lambda_{UV}$  functions of  $N_{max}$ ,  $\hbar\omega$ ;  $E^{UV}(L_{eff}) = E_{\infty} + ae^{-bL_{eff}} + \dots$
- UV convergence for  $\Lambda_{UV} \gtrsim 1$  GeV
- $\sim$ keV accuracy reached for  $N_{max} \sim 70$

## **TWO-BODY $\pi^-$ DECAY RATE**

$$\Gamma(\Lambda^0 \rightarrow ^3\text{He} + \pi^-)$$

---

# TWO-BODY $\pi^-$ DECAY RATE $\Gamma(\Lambda \rightarrow {}^3\text{He} + \pi^-)$



**Figure 3:** Calculated two-body decay rates  $\Gamma(\Lambda \rightarrow \pi^- + {}^3\text{He})$  using NCSM wave functions of  ${}^3\text{H}$  and  ${}^3\text{He}$  as functions of the IR length scale  $L_{\text{eff}}$  for several values of the  $\Lambda_{\text{UV}}$  cutoff.

- UV convergence reached for  $\Lambda_{\text{UV}} = 1 \text{ GeV}$
- Convergence with  $L_{\text{eff}}$  ( $N_{\text{max}}, \hbar\omega$ ) is slower than for the g.s. energies  $\rightarrow \Gamma^{UV}(L_{\text{eff}}) = \Gamma_{\infty}^{UV} + a e^{-b L_{\text{eff}}}$  extrapolation

## **PION FINAL STATE INTERACTIONS**

---

# PION FINAL STATE INTERACTIONS

## $\pi^-$ -nucleus interaction

- Influences the emitted  $\pi^-$  in  ${}^3\text{H} \rightarrow \pi^- + {}^3\text{He}$
- Understood in terms of  $\pi^-$ -nucleus optical potentials constrained by fits to  $\pi^-$ -atom level shifts and widths from Ne to U
  - Reproduces 1S level shift and width of  $\pi^-$  atoms of  ${}^3\text{He}$
- Supplemented by  $\pi N$  and  $\pi A$  scattering to extrapolate from near-threshold to  $q = 114.4$  MeV in the  $\pi^- - {}^3\text{He}$  c.m. system

## $\pi^-$ distorted waves in ${}^3\text{H} \rightarrow {}^3\text{He} + \pi^-$

- $\phi_\pi(\vec{r}; q)$  plane wave replaced by distorted wave
- Interplay of s- and p-wave parts of the optical potential produces robust attractive  $\pi^-$  FSI
- Increases  $\Gamma({}^3\text{H} \rightarrow {}^3\text{He} + \pi^-)$  by 15 %!

[A. Pérez-Obiol, DG, E. Friedman, A. Gal, Phys. Lett. B 811, 135916 (2020)]

# $\Sigma NN$ ADMIXTURES IN ${}^3\Lambda H$

---

# $\Sigma NN$ ADMIXTURES IN ${}^3\Lambda H$

## ${}^3\Lambda H$ structure

- Strong interaction  $\Lambda N \leftrightarrow \Sigma N$  transitions couple  $\Lambda NN$  and  $\Sigma NN$  hypernuclear sectors

$$|{}^3\Lambda H\rangle = \alpha |\Lambda pn\rangle + \beta |\Sigma^0 pn\rangle + \gamma |\Sigma^- pp\rangle + \delta |\Sigma^+ nn\rangle$$

- $\Sigma NN$  contributes  $\lesssim 0.5\%$  to the norm

## ${}^3\Lambda H$ decay

- New  $\Sigma$  hyperon two-body decay channels

$$\Sigma^- \rightarrow n + \pi^- \text{ and } \Sigma^0 \rightarrow p + \pi^-$$

become available in  ${}^3\Lambda H \rightarrow {}^3He + \pi^-$

- Amplitudes

$$\mathcal{A}_\Lambda F^{PV} \rightarrow \mathcal{A}_\Lambda F_{l=0}^{PV} + \frac{1}{3}(\sqrt{2}\mathcal{A}_{\Sigma^-} + \mathcal{A}_{\Sigma^0})F_{l=1}^{PV}$$

interfere in  $\Gamma_{{}^3\Lambda H \rightarrow {}^3He + \pi^-} \propto (\mathcal{A}_\Lambda |F^{PV}|)^2$

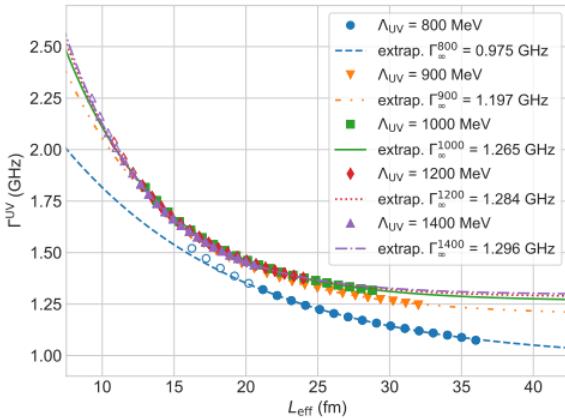
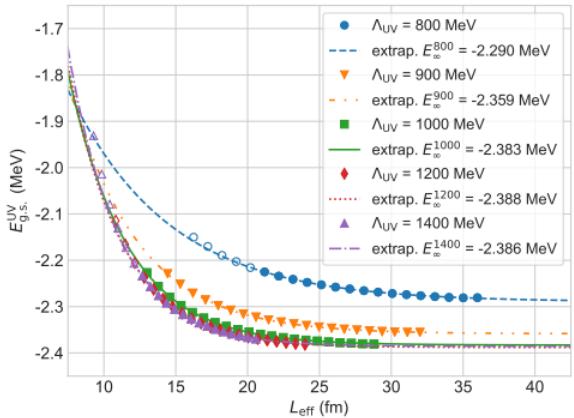
- Two-body  $\pi^-$  decay rate found to be reduced  $\gtrsim 10\%$

**RELATIONSHIP OF  $\Gamma(^3\Lambda \rightarrow ^3\text{He} + \pi^-)$**   
**TO  $B_\Lambda$**

---

# RELATIONSHIP OF $\Gamma(\Lambda^3\text{H} \rightarrow {}^3\text{He} + \pi^-)$ TO $B_\Lambda$

- $B_\Lambda({}^3\text{H}) = 165 \pm 44 \text{ keV}$ , not known precisely
- Use the  $\Lambda_{\text{UV}}$  dependence of  $B_\Lambda$  and  $\Gamma(\Lambda^3\text{H} \rightarrow {}^3\text{He} + \pi^-)$



- Correlation between  $B_\Lambda$  and  $\Gamma(\Lambda^3\text{H} \rightarrow {}^3\text{He} + \pi^-)$  at different  $\Lambda_{\text{UV}}$  seems robust (despite of missing UV corrections in the extrapolation scheme)

## RELATIONSHIP OF $\Gamma(\Lambda \rightarrow {}^3\text{He} + \pi^-)$ TO $B_\Lambda$ (cont.)

$\Lambda_{UV}$ (MeV)	$B_\Lambda$ (keV)	$\Gamma(\Lambda \rightarrow {}^3\text{He} + \pi^-)$ (GHz)	$\tau(\Lambda)$ (ps)	
800	69	0.975	$234 \pm 27$	(a)
900	135	1.197	$190 \pm 22$	(b)
1000	159	1.265	$180 \pm 21$	(b)
—	410	1.403	$163 \pm 18$	(c)

- (a) Agrees with recent ALICE lifetime measurement and also with [Kamada et al., PRC 57, 1595 (1998)]
- (b) Agrees with HypHI lifetime measurement
- (c) Has substantial overlap with STAR(18) lifetime value when extrapolated to  $B_\Lambda^{\text{STAR}} = 0.41 \pm 0.12 \pm 0.11$  MeV (almost coincides when  $R_3^{\text{STAR}}$  is used)

[A. Pérez-Obiol, DG, E. Friedman, A. Gal, Phys. Lett. B 811, 135916 (2020)]

## **QUANTIFYING UNCERTAINTIES**

### **(Ongoing)**

---

# QUANTIFYING UNCERTAINTIES

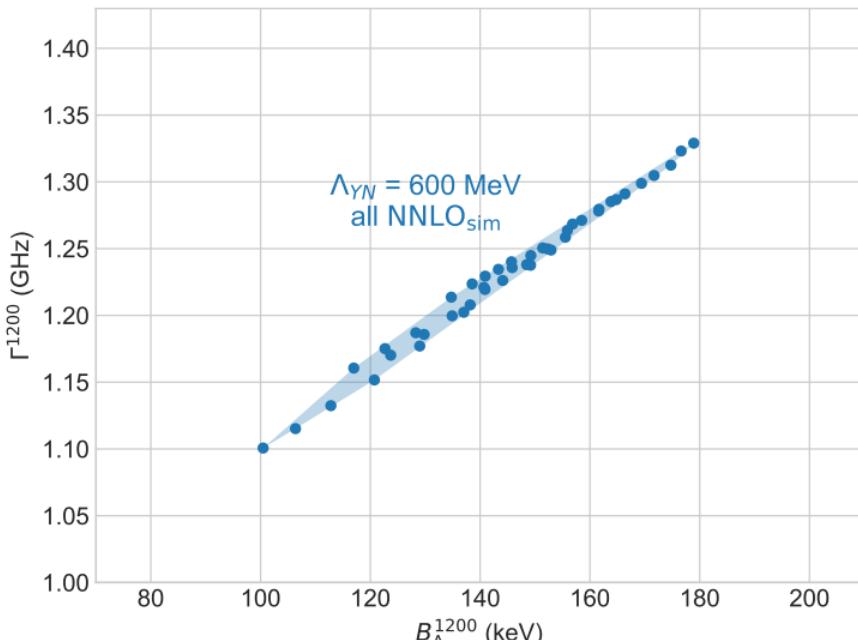
## Theoretical (hyper)nuclear-structure uncertainties in $\Gamma(^3\Lambda \rightarrow ^3\text{He} + \pi^-)$

- Method
  - Schrödinger equation solver, extrapolation
  - Under control for  $A = 3$  (at least for energies)
  - Methods are more precise than input Hamiltonians
- Model
  - **YN interaction** – poor data base of scattering data suffering from large uncertainties; cutoff dependence as a diagnostic tool(?)
  - **NN+NNN interaction** – rich data base of low-energy observables

## The NNLO<sub>sim</sub> family of NN+NNN potentials

- Parameters fitted to reproduce simultaneously  $\pi N$ ,  $NN$ , and  $NNN$  low-energy observables
- Family of 42 Hamiltonians where the experimental uncertainties propagate into the LECs of the  $\chi$ EFT Lagrangian
- All Hamiltonians give equally good description of the fit data
- Note that  $\Delta E(^3\text{He}/^3\text{H}) \approx 0$  (fitted) while  $\Delta E_{g.s.}^{(^4\text{He})} \approx 1.5 \text{ MeV}$

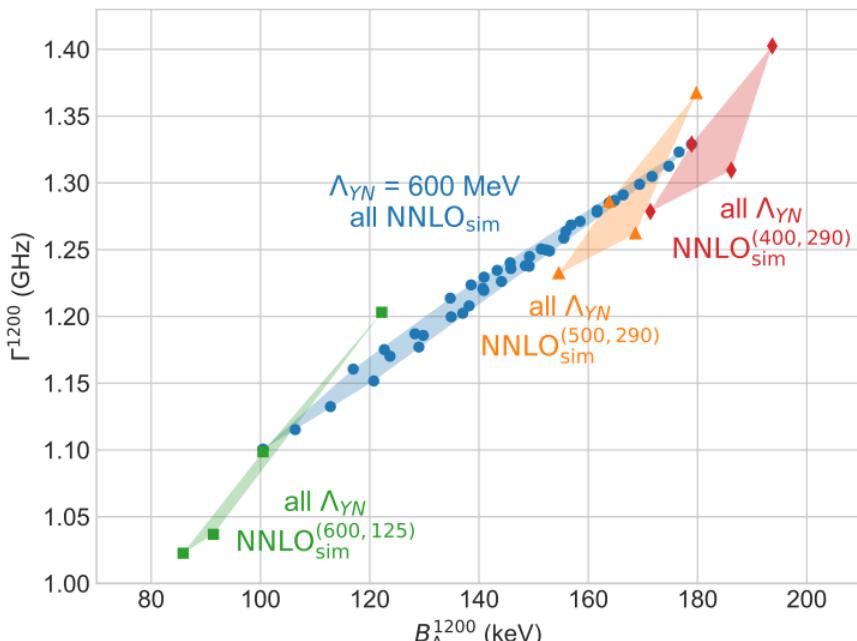
# (HYPER)NUCLEAR-STRUCTURE UNCERTAINTIES in $\Gamma(\Lambda^{\text{3H}} \rightarrow \text{He}^3 + \pi^-)$



**Figure 4:** Calculated two-body decay rates  $\Gamma(\Lambda^{\text{3H}} \rightarrow \pi^- + \text{He}^3)$  and  $\Lambda$  separation energies  $B_{\Lambda}$  for all 42 NNLO<sub>sim</sub> Hamiltonians.

- $\Delta B_{\Lambda}(\text{NNLO}_{\text{sim}}) \approx 80$  keV
- $\Delta \Gamma(\Lambda^{\text{3H}} \rightarrow \text{He}^3 + \pi^-)(\text{NNLO}_{\text{sim}}) \approx 0.35$  GHz

# (HYPER)NUCLEAR-STRUCTURE UNCERTAINTIES in $\Gamma(\Lambda \rightarrow ^3\text{He} + \pi^-)$



**Figure 4:** Calculated two-body decay rates  $\Gamma(\Lambda \rightarrow \pi^- + ^3\text{He})$  and  $\Lambda$  separation energies  $B_{\Lambda}$  for all 42 NNLO<sub>sim</sub> Hamiltonians.

- $\Delta B_{\Lambda}(\text{NNLO}_{\text{sim}}) \approx 80 \text{ keV}$
- $\Delta \Gamma(\Lambda \rightarrow ^3\text{He} + \pi^-)(\text{NNLO}_{\text{sim}}) \approx 0.35 \text{ GHz}$

## **SUMMARY**

---

## Hypertriton lifetime

- Performed new microscopic three-body calculation of two-body decay rate  $\Gamma(\Lambda^3\text{H} \rightarrow {}^3\text{He} + \pi^-)$
- Using the  $\Delta I = 1/2$  rule and a branching ratio  $R_3$  from experiment we deduced the value of hypertriton lifetime  $\tau(\Lambda^3\text{H})$
- Pion FSI increase the  $\Lambda^3\text{H}$  decay rate  $\Gamma(\Lambda^3\text{H})$  by  $\sim 15\%$
- $\Sigma NN$  admixtures in  $\Lambda^3\text{H}$  decrease the  $\Gamma(\Lambda^3\text{H})$  by  $\sim 10\%$
- $\tau(\Lambda^3\text{H})$  varies strongly with the poorly known  $\Lambda$  separation energy  $B_\Lambda$  – it is possible to correlate each of the reported lifetime values from ALICE, HypHI, and STAR(18) to its own underlying  $B_\Lambda$  value
- New experiments proposed at MAMI, Jlab, J-PARC, and ELPH will hopefully pin down  $B_\Lambda$  to better than 50 keV and lead to a resolution of the ‘hypertriton lifetime puzzle’

**Thank you!**