

Updated experimental insight into the strong interaction between antikaons and nucleons

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HYP-2022, July 1st, 2022



This project has received funding from the European Union's Horizon 2020 research and innovation programme under the Marie Skłodowska-Curie grant agreement No 754496

KbarN interaction

KbarN interaction: building block of non-perturbative regime of QCD

The presence of the strange quark has dramatic consequences:

KN and KbarN interactions are very different and with very different cross sections

- excitation of resonances below threshold
- strong coupled channel dynamics
- absorption from one or more nucleons
- attractiveness in KbarN gives rise to bound states

$\Lambda(1405)$ is an “old object” not fitting in the standard 3-quark picture

Theoretical framework

Theoretical approaches:

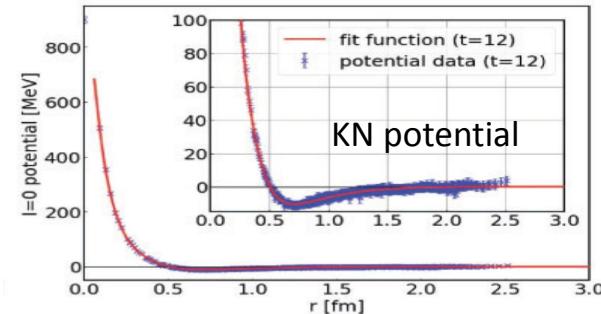
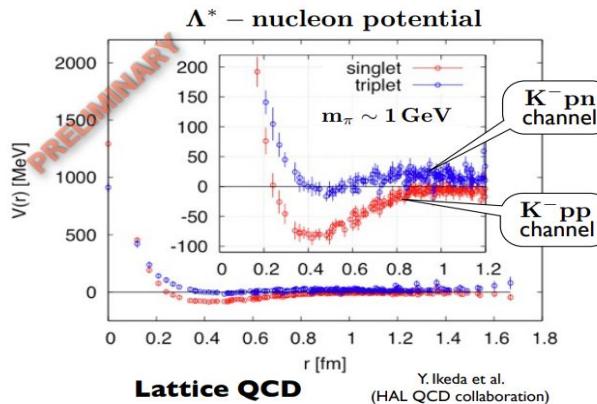
- meson exchange
- phenomenological
- chiral SU(3) dynamical
- Lattice QCD (?)

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⇒ “Lattice QCD Evidence that the $\Lambda(1405)$ Resonance is an Antikaon-Nucleon Molecule”, Hall et al., PRL 114, 132002 (2015)



Murakami, Akahoshi, Aoki
PTEP 2020, 093B03 (2020)

⇒ MB systems with annihilation available K. Murakami arXiv:2111.15563v1 [hep-lat]

Theoretical framework

Theoretical approaches:

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Data is crucial to test (+feed) this approaches.

Theoretical frame

Theoretical approaches:

- meson exchange
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Data is crucial to test (+feed)

Data fitting by Chiral SU(3).

- Going to NLO ($N^2\text{LO?}$), s+p waves \Rightarrow more parameters to be fixed (by data)
- Adding **new data** helps to improve the model
- Adding **more precise data** helps to improve the model
- Adding **data at different energies** helps to improve the model

- Next to leading order (NLO), just considering the **contact term**

$$\mathcal{L}_{\phi B}^{(2)} = b_D \langle \bar{B} \{\chi_+, B\} \rangle + b_F \langle \bar{B} [\chi_+, B] \rangle + b_0 \langle \bar{B} B \rangle \langle \chi_+ \rangle + d_1 \langle \bar{B} \{u_\mu, [u^\mu, B]\} \rangle \\ + d_2 \langle \bar{B} [u_\mu, [u^\mu, B]] \rangle + d_3 \langle \bar{B} u_\mu \rangle \langle u^\mu B \rangle + d_4 \langle \bar{B} B \rangle \langle u^\mu u_\mu \rangle$$

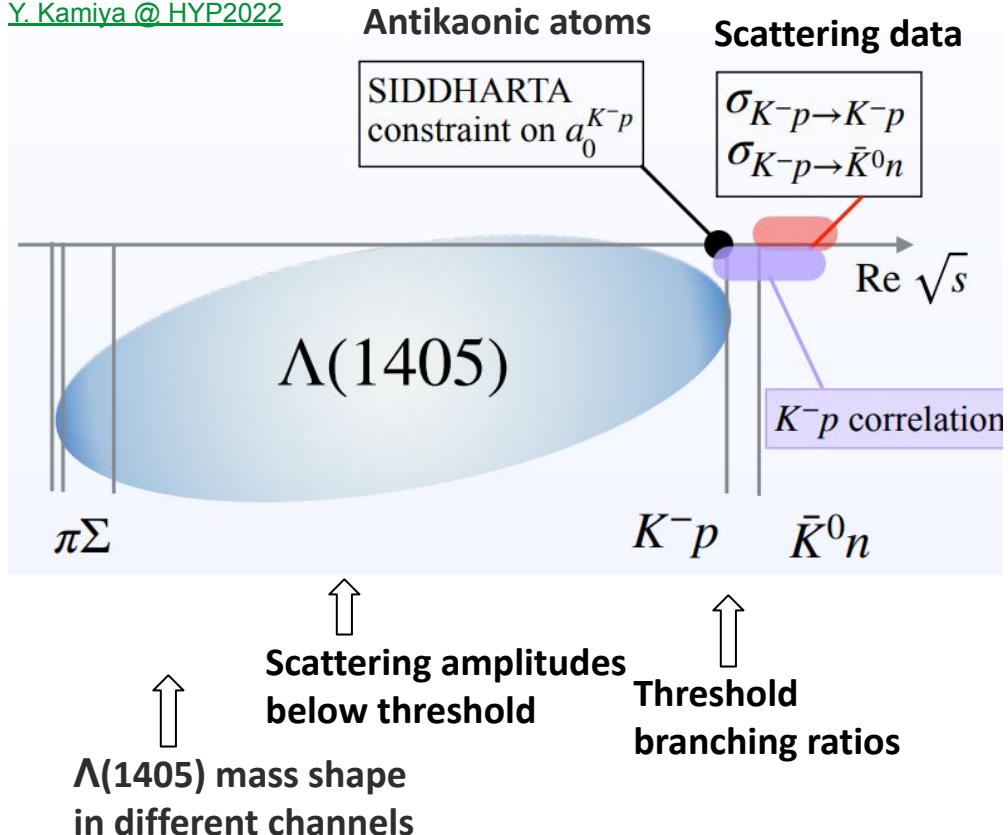
New terms taken into account

$$\left. \begin{aligned} & -\frac{g_1}{8M_N^2} \langle \bar{B} \{u_\mu, [u_\nu, \{D^\mu, D^\nu\} B]\} \rangle - \frac{g_2}{8M_N^2} \langle \bar{B} [u_\mu, [u_\nu, \{D^\mu, D^\nu\} B]] \rangle \\ & -\frac{g_3}{8M_N^2} \langle \bar{B} u_\mu \rangle \langle [u_\nu, \{D^\mu, D^\nu\} B] \rangle - \frac{g_4}{8M_N^2} \langle \bar{B} \{D^\mu, D^\nu\} B \rangle \langle u_\mu u_\nu \rangle \\ & -\frac{h_1}{4} \langle \bar{B} [\gamma^\mu, \gamma^\nu] B u_\mu u_\nu \rangle - \frac{h_2}{4} \langle \bar{B} [\gamma^\mu, \gamma^\nu] u_\mu [u_\nu, B] \rangle - \frac{h_3}{4} \langle \bar{B} [\gamma^\mu, \gamma^\nu] u_\mu \{u_\nu, B\} \rangle \\ & -\frac{h_4}{4} \langle \bar{B} [\gamma^\mu, \gamma^\nu] u_\mu \rangle \langle u_\nu, B \rangle + h.c. \end{aligned} \right\}$$

- $b_0, b_D, b_F, d_1, d_2, d_3, d_4, g_1, g_2, g_4, h_1, h_2, h_3, h_4$ are not well established, so they should be treated as parameters of the model!

Available experimental data

[Y. Kamiya @ HYP2022](#)

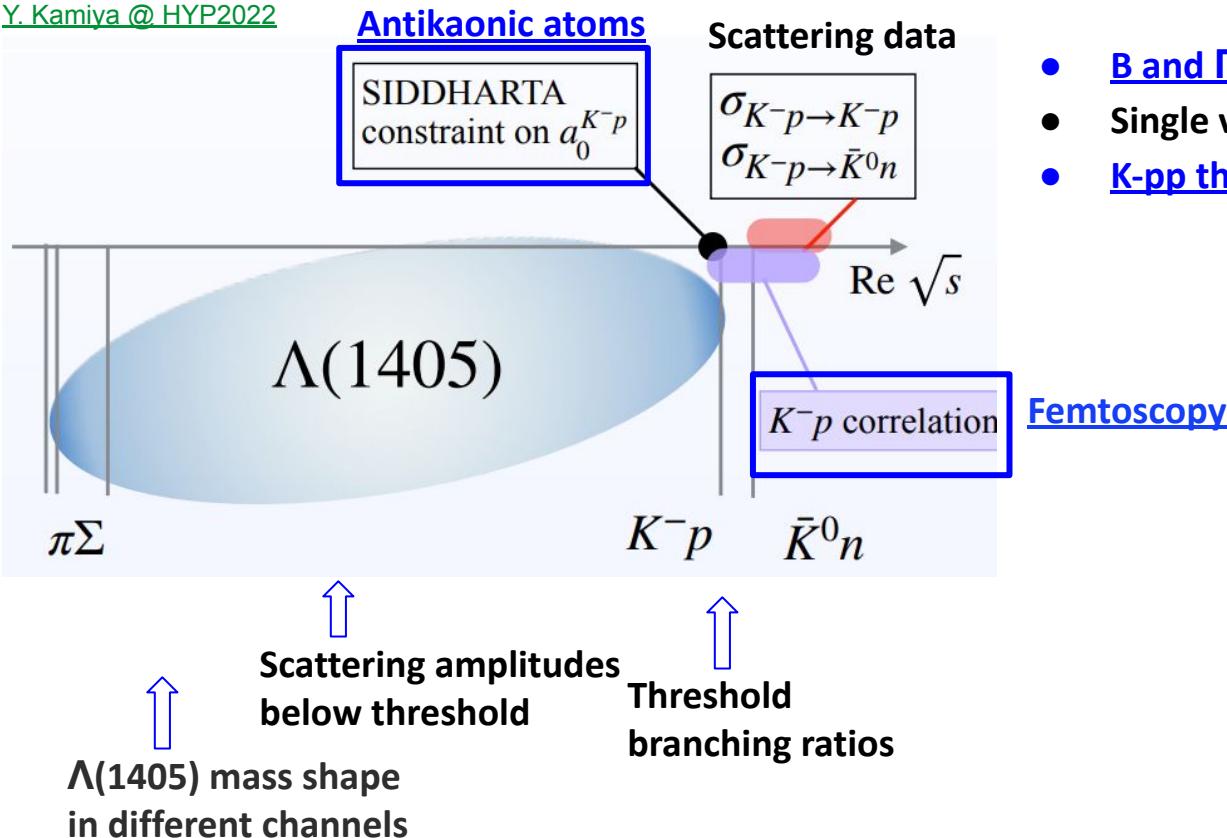


- B and Γ of kaonic nuclear states
- Single vs Multi-nucleonic absorption rates
- K-pp three body femtoscopy

Femtoscopy

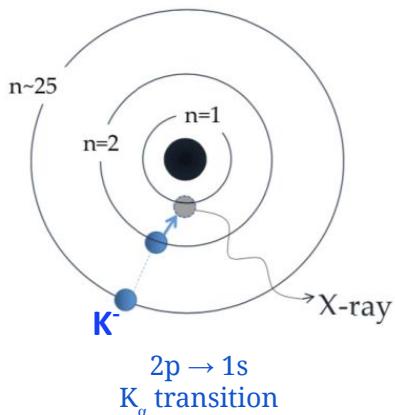
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- B and Γ of kaonic nuclear states
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antikaonic hydrogen: SIDDHARTA



shift(ϵ), width(Γ) with respect to e.m. value caused by attractive/repulsive strong interaction and the presence of inelastic channels

Measurement of the **shift(ϵ) and width(Γ) induced by the strong interaction** in the lowest level atomic transition.

$$\epsilon_{1s} = -283 \pm 36(\text{stat}) \pm 6(\text{syst}) \text{ eV}$$

$$\Gamma_{1s} = 541 \pm 89(\text{stat}) \pm 22(\text{syst}) \text{ eV},$$

SIDDHARTA Coll., PLB 704 (2011) 113

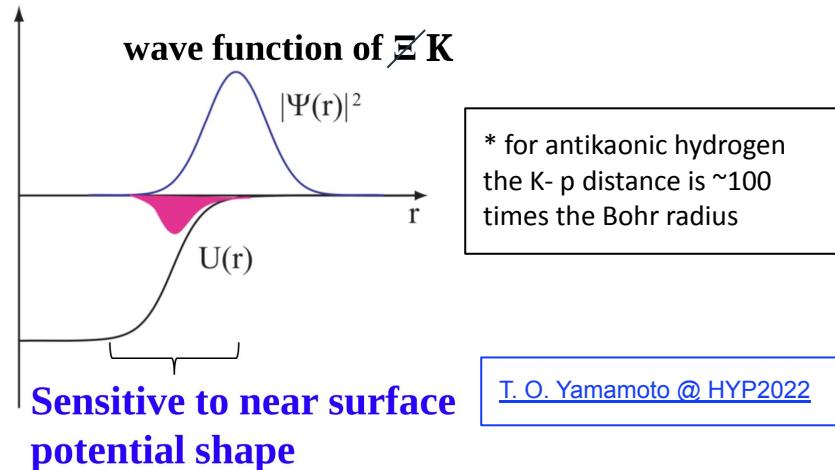
Translated via Desser-type Formula into a **$K^- p$ scattering length** that is an average of the $K\bar{p}N$ scattering lengths for $I=0$ and $I=1$

$$\epsilon_{1s} - \frac{i}{2} \Gamma_{1s} = -2\alpha^3 \mu_c^2 a_p (1 - 2\alpha \mu_c (\ln \alpha - 1) a_p)$$

$$a_{K^- p} = \frac{a_0(I=0) + a_1(I=1)}{2}$$

KbarN at threshold and low momentum

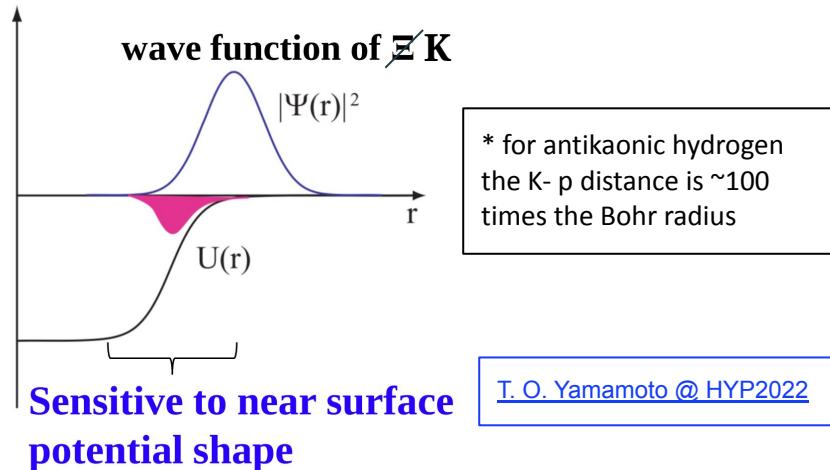
SIDDHARTA: antiKaonic Hydrogen



The overlap of the kaon wavefunction with the nucleon delivers insight into the effects of the strong interaction, competing with Coulomb effects

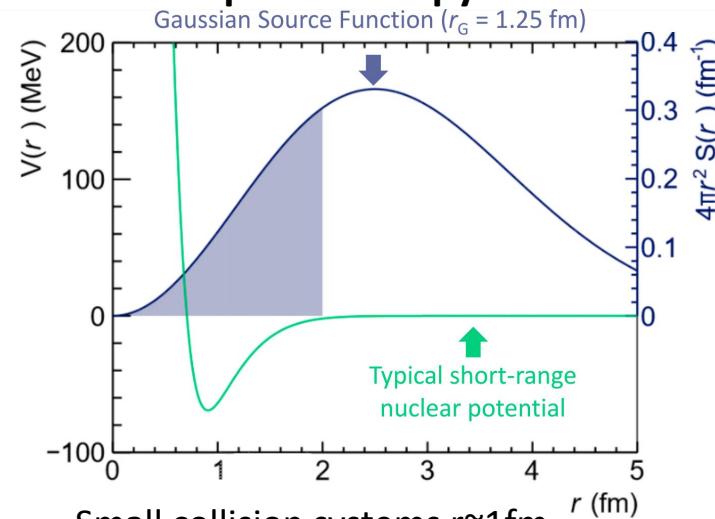
KbarN at threshold and low momentum

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ALICE: $K^- p$ femtoscopy



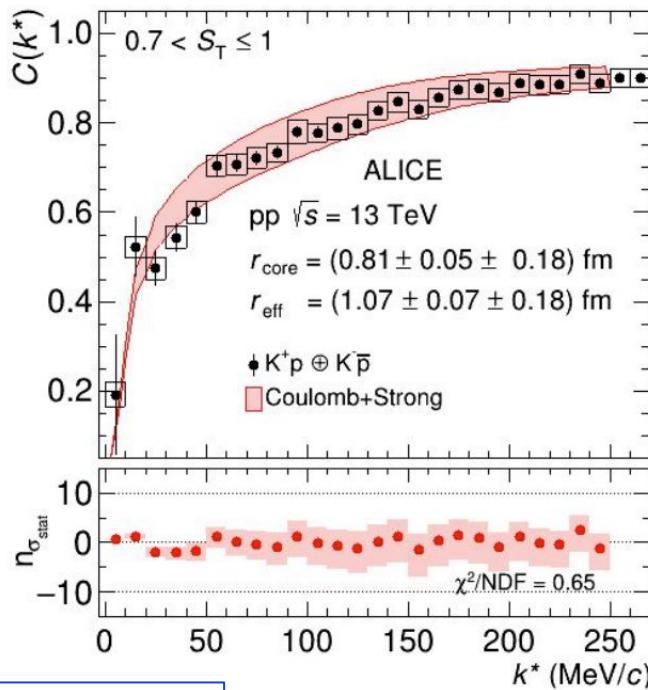
\Rightarrow effect of the interaction is enhanced

$$C(k^*) = \int S(r^*) |\Psi(k^*, \vec{r}^*)|^2 d^3r^*$$

Deliver different observables \Leftrightarrow scattering lengths can be obtained from both
(via Deser-type and Lednický–Lyuboshitz formulae)

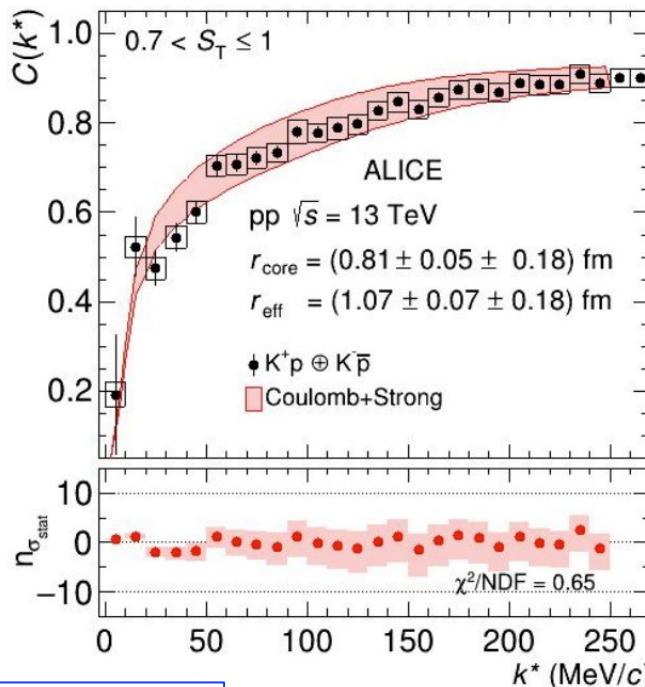
KbarN Femtoscopy with ALICE

An crucial PLUS: K^+p interaction for crosscheck
and experimental determination of the radius



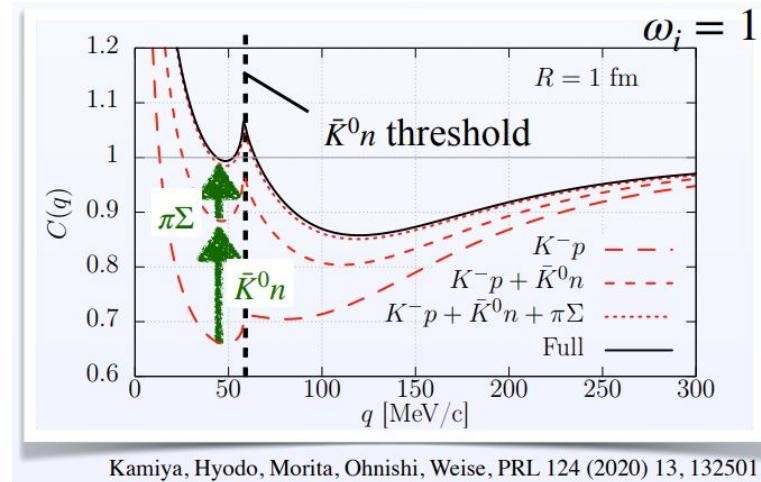
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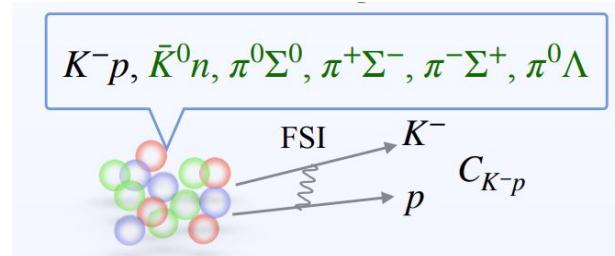
R. Lea @ HYP2022

Effects of coupled channels enhanced by small source
- less important for large sources of HIC



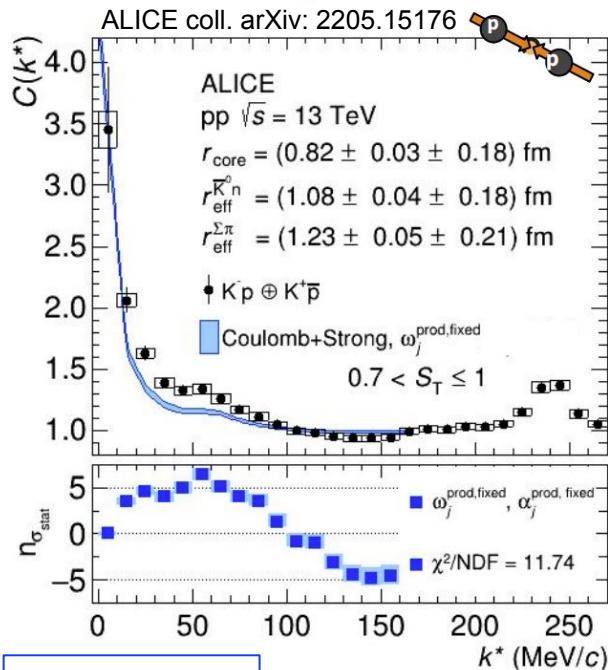
Kamiya, Hyodo, Morita, Ohnishi, Weise, PRL 124 (2020) 13, 132501

$K^-p, \bar{K}^0n, \pi^0\Sigma^0, \pi^+\Sigma^-, \pi^-\Sigma^+, \pi^0\Lambda$



Y. Kamiya @ HYP2022

KbarN Femtoscopy with ALICE in pp collisions

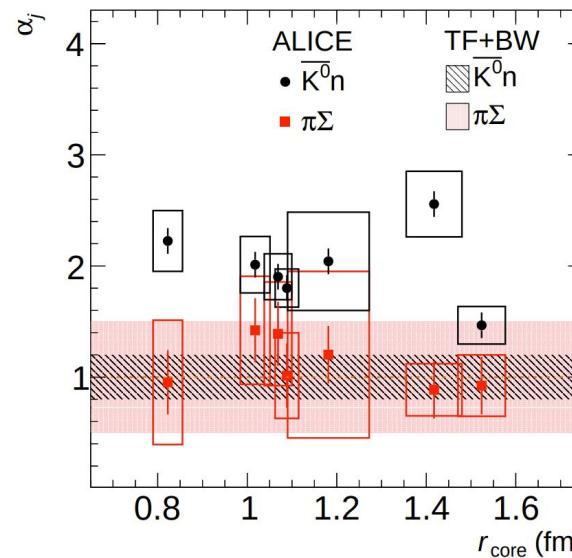


R. Lea @ HYP2022

Small systems: pp collisions r~1fm

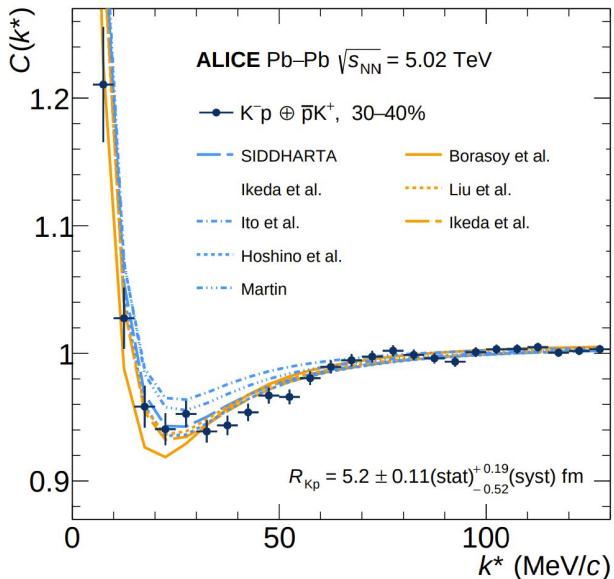
Coupled channels in the correlation function: production yields (thermal model) + production p_T spectrum (blast-wave) + pair kinematics

⇒ Provides a quantitative test of coupled channels in the theory



KbarN Femtoscopy with ALICE in PbPb collisions

ALICE Coll., PLB 822 (2021) 136708



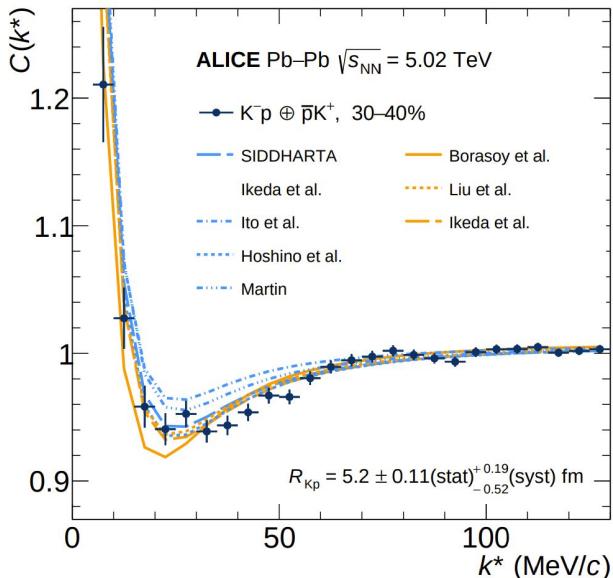
Large systems (HIC): PbPb collisions, up to $r \sim 9$ fm

Strength of coupled channels significantly reduced

- Cyan line: Kyoto model with SIDDHARTA constraint
- Orange line: Fit to the scattering parameters (better χ^2)

KbarN Femtoscopy with ALICE in PbPb collisions

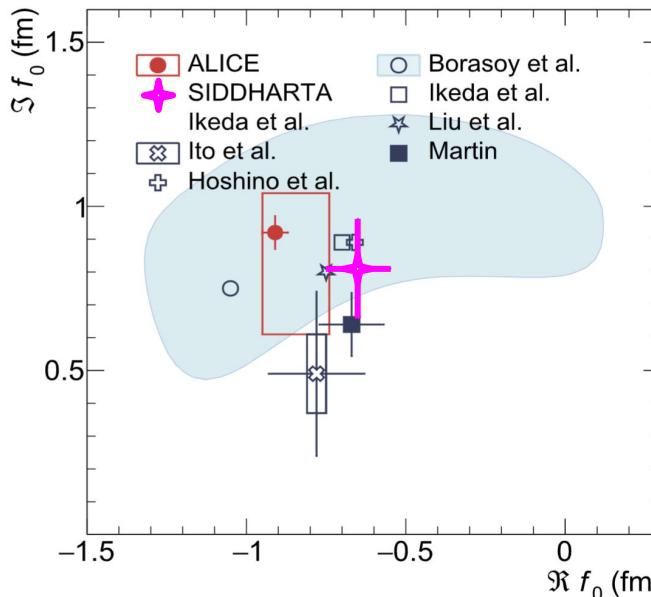
ALICE Coll., PLB 822 (2021) 136708



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Accessing KbarN |=1 interaction

Full isospin dependence needs K-d interaction measurements:

$$a_{K^-d} = \frac{1}{2} \frac{m_N + m_K}{m_N + \frac{m_K}{2}} (3a_1 + a_0) + C$$

A experimental constraint exists from pp→d K⁰bar K⁺ (COSY)

K-d scatt. length |Re|≤1.3 fm, Im≤1.3 fm
 V. Kleber et al., PRL 91 (2003) 172304
 A. Sibirtsev et al PLB 601 (2004) 132

SIDDHARTA-2 with new experimental setup

[J. Zmeskal @ HYP2022](#)

→measurement of **antikaonic deuterium**

→very challenging! low yield of signal.

The interpretation of the SIDDHARTA-2 results will require some effort too:

- Deser type formulae with model dependence
- Better test potentials by **full three body calculations**

[T. Hyodo@ HYP2022](#)

Potential	$\Delta E - i\Gamma/2$ [eV]
$V_{KN-\pi\Sigma}^{1,\text{SIDD}}$	767 – 464 <i>i</i> [1]
$V_{\bar{K}N-\pi\Sigma}^{2,\text{SIDD}}$	782 – 469 <i>i</i> [1]
$V_{\bar{K}N-\pi\Sigma-\pi\Lambda}^{\text{chiral}}$	835 – 502 <i>i</i> [1]
Kyoto $\bar{K}N$	670 – 508 <i>i</i> [2]

Accessing KbarN |I=1 interaction with ALICE

- $K_S^0 p$ correlation

$$|K_S^0 p\rangle = [|\bar{K}^0 p\rangle - |K^0 p\rangle]/\sqrt{2}$$

$\bar{K}N, I = 1$

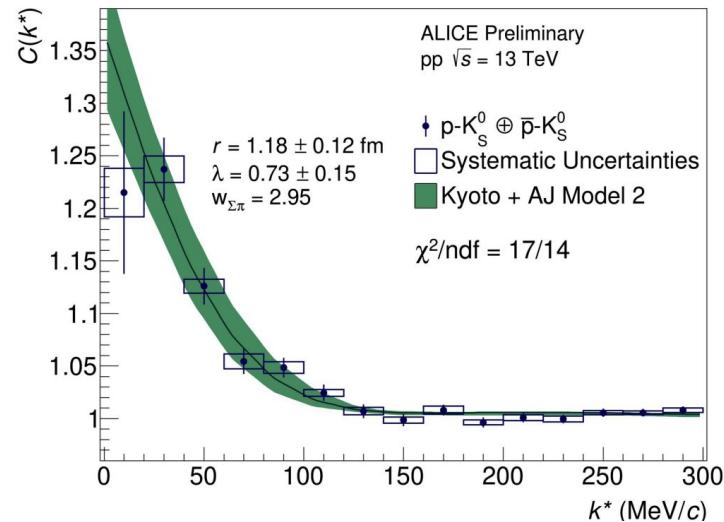
$KN, I = 0, 1$

$$C_{K_S^0 p} = [C_{\bar{K}^0 p} + C_{K^0 p}]/2$$



- $I = 1$ component only
 - Chiral amplitude
 - Ikeda, Hyodo, Weise, NPA881 (2012)
 - Effective potential
 - Miyahara, Hyodo, Weise, PRC 98 (2018)
- Well determined with scat. exp.
 - Chiral amplitude
 - K. Aoki and D. Jido, PTEP (2019)
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 - Constructed from chiral amp.

⇒ Access to $|I=1$ component of the KbarN interaction



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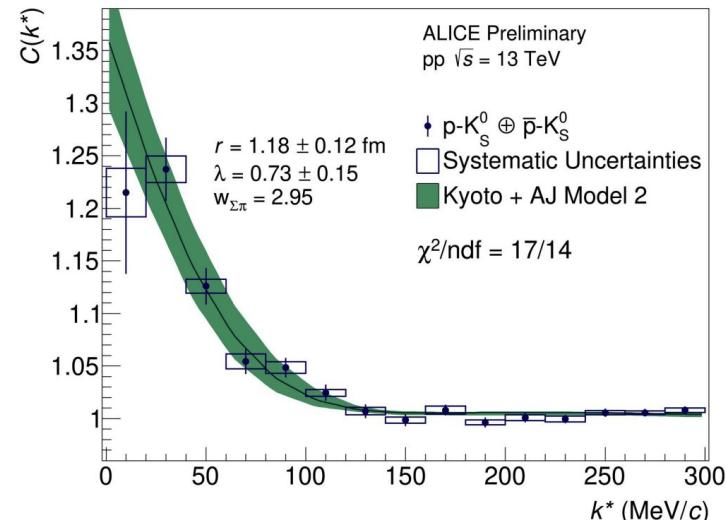
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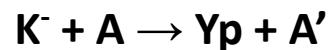
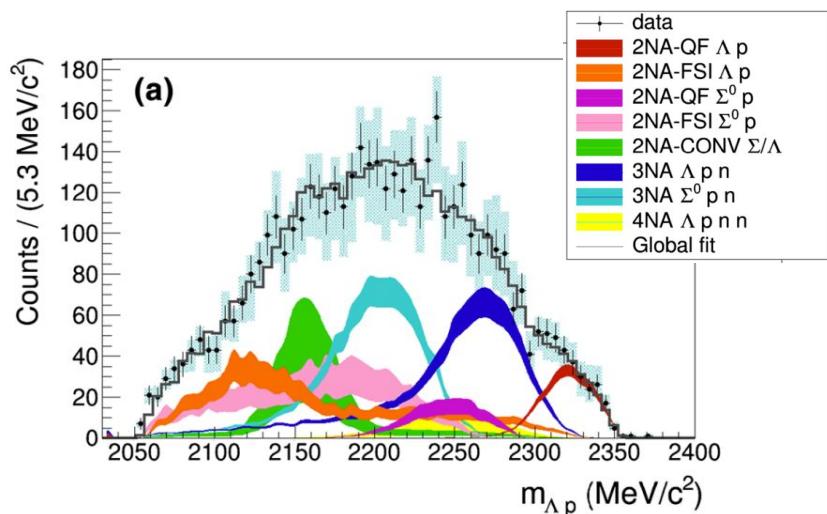


On the horizon by ALICE: K^+d , K^-d femtoscopy

For femtoscopy with deuterons (pd) by ALICE in small systems
see [B. Singh @ HYP2022](#) and by STAR in HIC: [H. Zbroszczyk @ HYP2022](#)

Kaonic nuclei

AMADEUS: K⁻ absorption in ⁴He and ¹²C



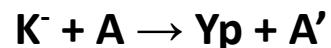
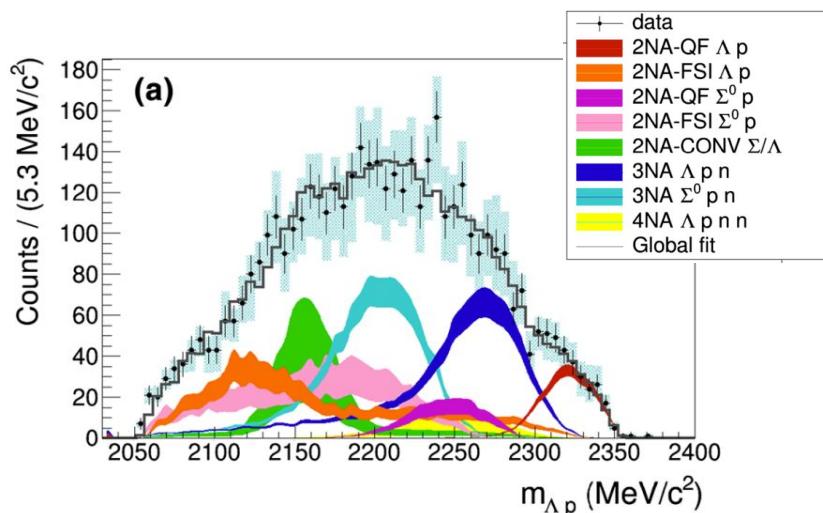
Multi-nucleon absorption processes dominate

AMADEUS Coll. PLB 758 (2016) 134

AMADEUS Coll. PLB 782 (2018) 339

AMADEUS Coll. FBS 62 (2021) 7.

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Multi-nucleon absorption processes dominate

AMADEUS Coll. PLB 758 (2016) 134

AMADEUS Coll. PLB 782 (2018) 339

AMADEUS Coll. FBS 62 (2021) 7.

Below threshold (-33 MeV)

K⁻n → π⁻Λ (l=1 non resonant)

K. Piscicchia @ HYP2022

$$|A_{K^-n \rightarrow \Lambda\pi^-}| = (0.334 \pm 0.018 \text{ stat})^{+0.034}_{-0.058} \text{ syst) fm.}$$

AMADEUS Coll., PLB 782 (2018) 339-345]

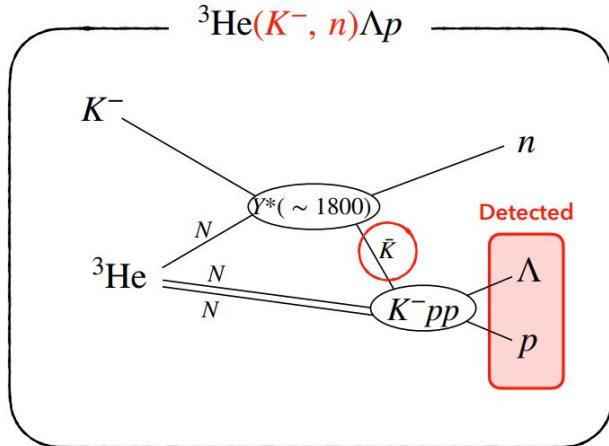
Above threshold

K⁻p → π⁰Λ, π⁰Σ⁰ cross sections (preliminary)

- $\sigma_{K^-p \rightarrow \Sigma^0\pi^0} = 42.8 \pm 1.5(\text{stat.})^{+2.4}_{-2.0}(\text{syst.}) \text{ mb}$
- $\sigma_{K^-p \rightarrow \Lambda\pi^0} = 31.0 \pm 0.5(\text{stat.})^{+1.2}_{-1.2}(\text{syst.}) \text{ mb,}$

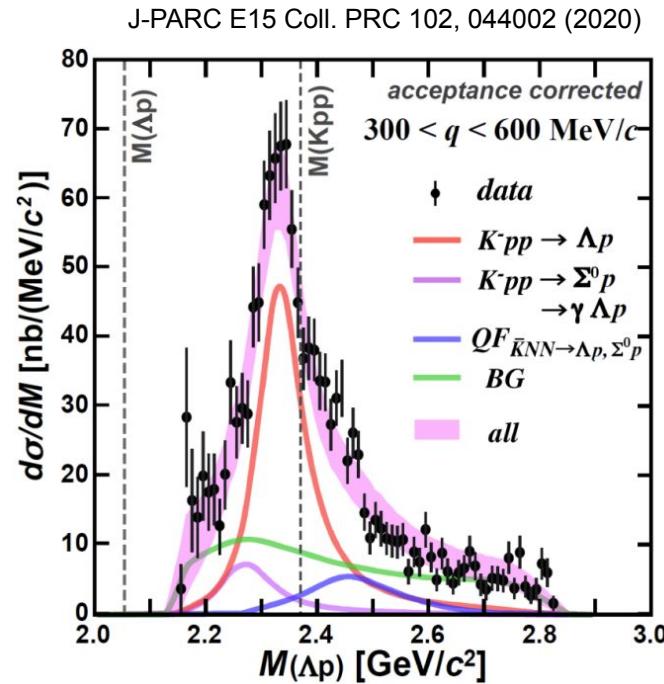
J-PARC E15

In-flight ${}^3\text{He}(K^-, n)\Lambda p$ reaction @ 1.0 GeV/c



$$B_K = 42 \pm 3(\text{stat.})^{+3}_{-4}(\text{syst.}) \text{ MeV}$$

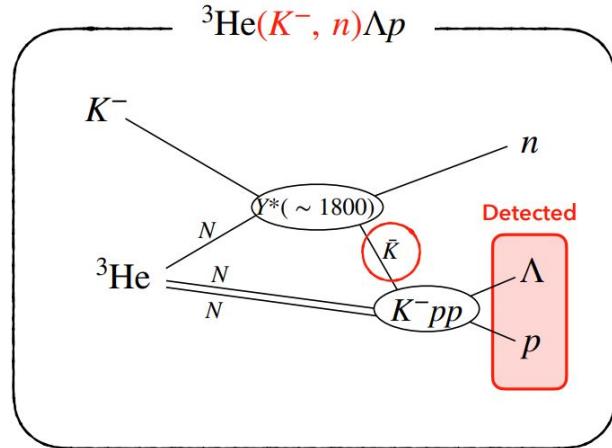
$$\Gamma_K = 100 \pm 7(\text{stat.})^{+19}_{-9}(\text{syst.}) \text{ MeV}$$



Similar interpretation by Sekihara, Oset, Ramos
JPS Conf. Proc. 26, 023009 (2019)
PTEP 2016, 12, 123D03

J-PARC E15

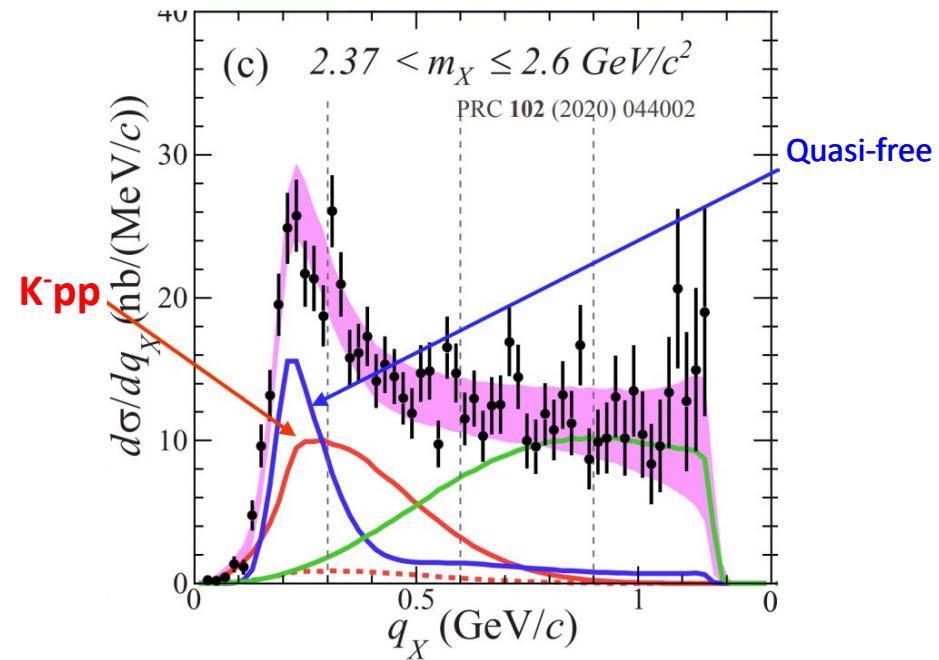
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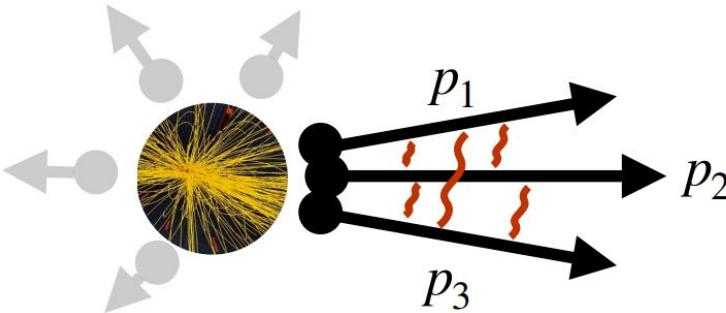
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J-PARC E15 Coll. PRC 102, 044002 (2020)



3-Body femtoscopy



L.Serksnyte@ HYP2022

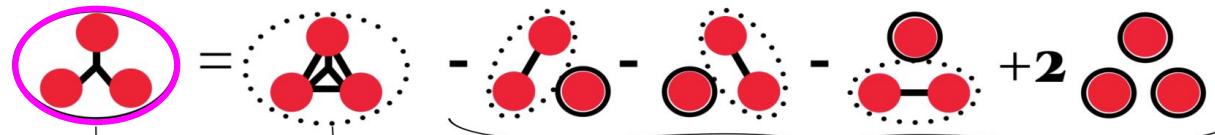
Three-particle correlation function:

$$C(\mathbf{p}_1, \mathbf{p}_2, \mathbf{p}_3) \equiv \frac{P(\mathbf{p}_1, \mathbf{p}_2, \mathbf{p}_3)}{P(\mathbf{p}_1) P(\mathbf{p}_2) P(\mathbf{p}_3)} = \mathcal{N} \frac{N_{\text{same}}(Q_3)}{N_{\text{mixed}}(Q_3)}$$

$$Q_3 = \sqrt{-q_{12}^2 - q_{23}^2 - q_{31}^2}$$

Isolation of the **three-body effects** in the correlation function:

$$c_3(Q_3) = C(Q_3) - C_{12}(Q_3) - C_{23}(Q_3) - C_{31}(Q_3) + 2 \quad \rightarrow \text{2-body correlations experimentally determined}$$



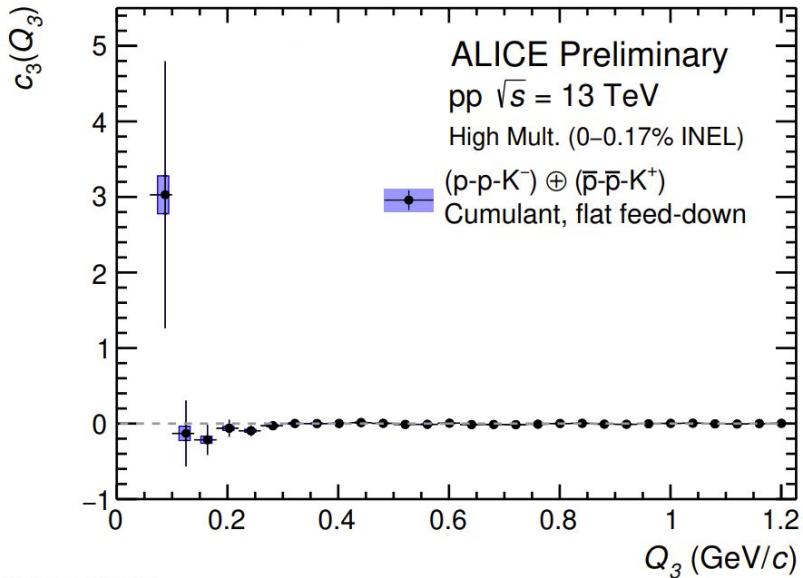
Genuine three-body correlations (cumulant)

Measured triplets

Lower-order correlations

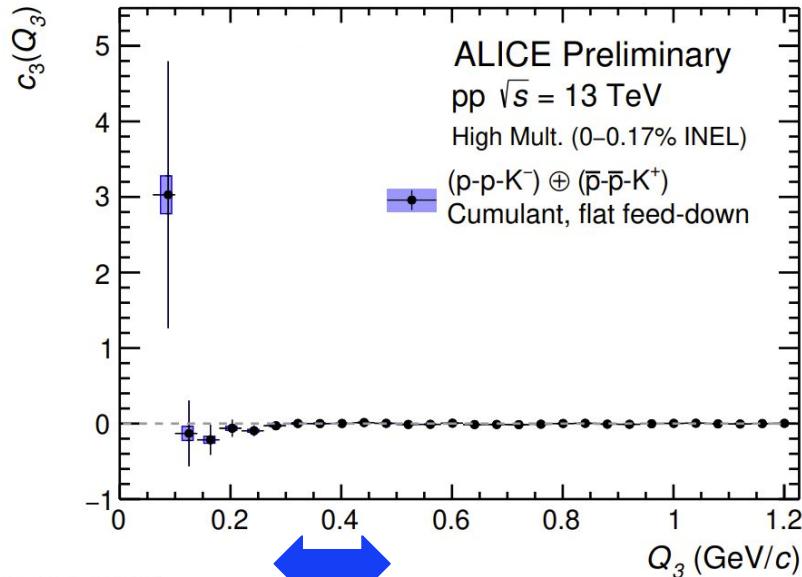
R. Kubo, J. Phys. Soc. Jpn. 17, 1100 (1962)

ppK⁻ cumulant



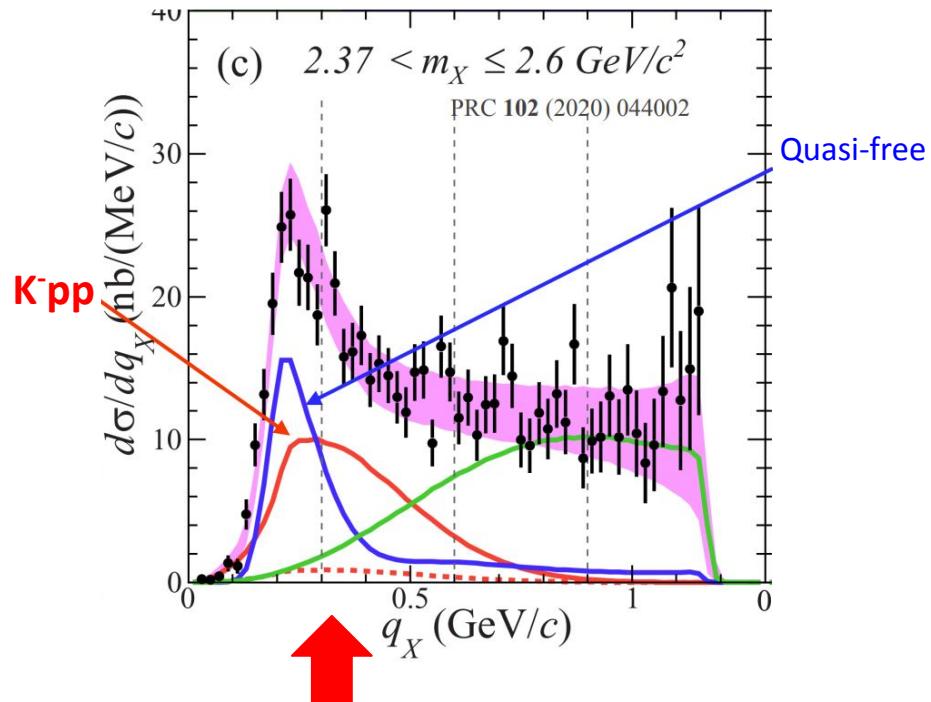
⇒ **ppK⁻ cumulant is compatible with zero.**

ppK⁻ cumulant

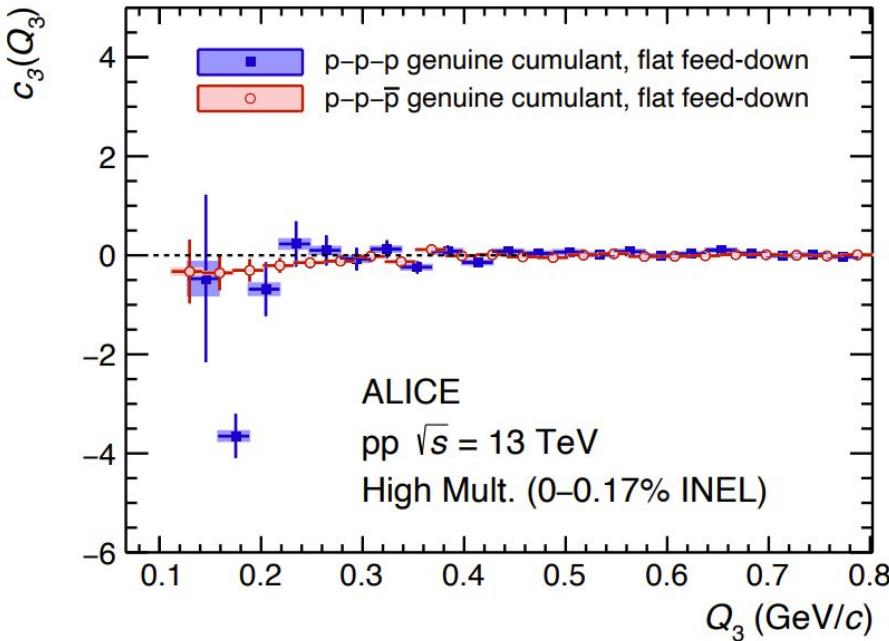


⇒ ppK⁻ cumulant is compatible with zero.

⇒ The measurement suggest that three-body effects are not relevant for the description of the K⁻pp system



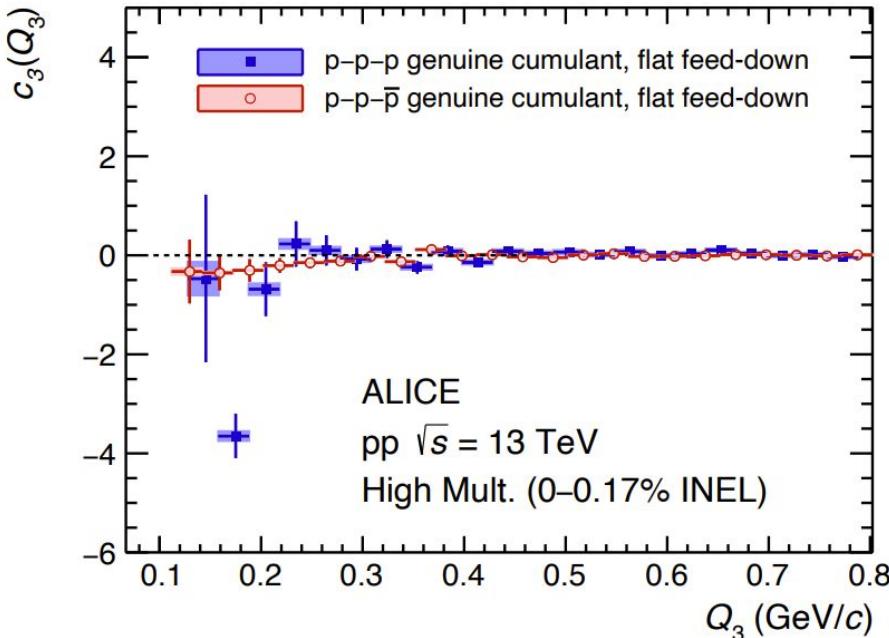
p-p-p cumulant



Statistical significance: $n_\sigma = 6.7$ for $Q_3 < 0.4$ GeV/c
 Origin of the 3-body effect?

- Pauli blocking at the three-particle level
- FSI

p-p-p cumulant



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[N. Shevchenko@ HYP2022](#)

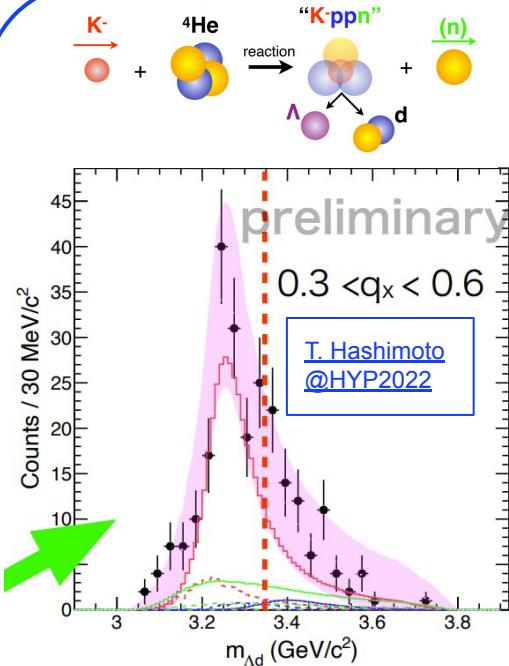
$K^- ppn - \bar{K}^0 nnp$ system:

$\bar{K}NNN$ with $I^{(4)} = 0, S^{(4)} = 1/2, L^{(4)} = 0$

	$V_{NN}^{\text{TSA-A}}$		$V_{NN}^{\text{TSA-B}}$		V_{NN}^{TSN}	
	B	Γ	B	Γ	B	Γ
$V_{\bar{K}N}^{1,\text{SIDD}}$	52.0	50.4	50.3	49.6	51.2	50.8
$V_{\bar{K}N}^{2,\text{SIDD}}$	47.0	39.6	46.4	38.2	46.4	39.9
$V_{\bar{K}N}^{\text{Chiral}}$	32.6	39.7	34.5	50.9	30.5	42.8

⇒ relevant for the description of the $K^- NNN$ system?

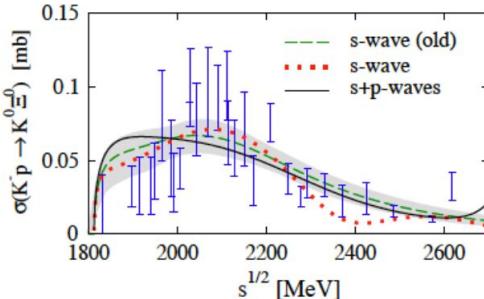
Data is coming



- Extended J-PARC program
for kaonic nuclei studies

[T. Yamaga@ HYP2022](#)

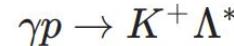
Calculations in p-wave with
fitting higher mass channels



[A. Feijoo @ HYP2022](#)

⇒ J-Lab proposal for the
secondary K_L beam

GLUEX:
 $\Lambda(1405)$ photoproduction



[N. Wickramaarachchi
@ HYP2022](#)

New approaches proposed:
fusion reaction $K^- d \rightarrow p\Sigma^-$
⇒ can deliver information on
the lower pole of $\Lambda(1405)$

[E. Oset @ HYP2022](#)

Outlook

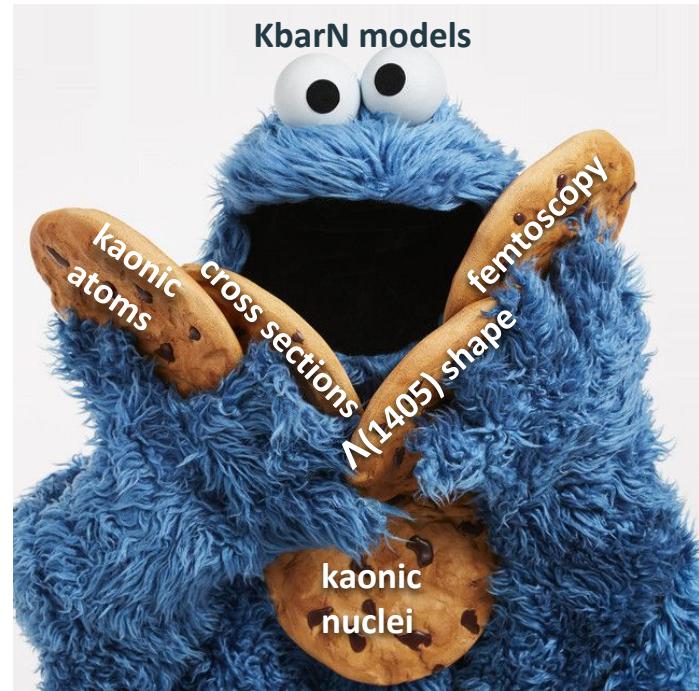
We are collecting new data (e.g. ALICE run 3 with $\geq x100$ stats and improved tracking, SIDDHARTA2 running)

In the near future:

- new experimental apparatus
- new facilities

In the KbarN field a boost similar to the precise measurement of kaonic hydrogen is expected

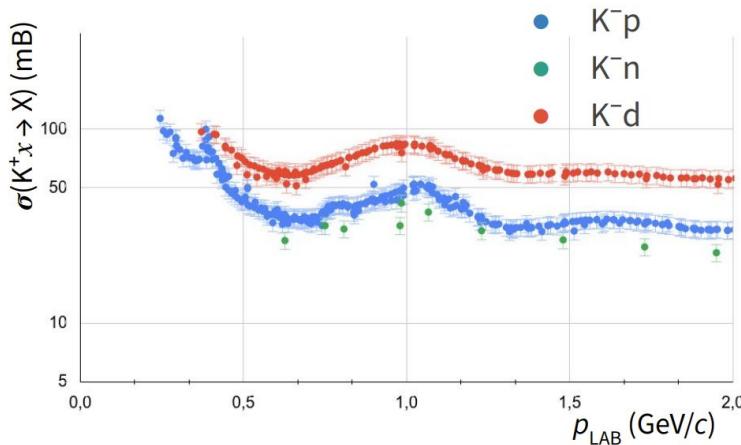
On the search for a description of the KbarN interaction that can accommodate all the data from above to below threshold



you really need to eat
ALL the cookies!

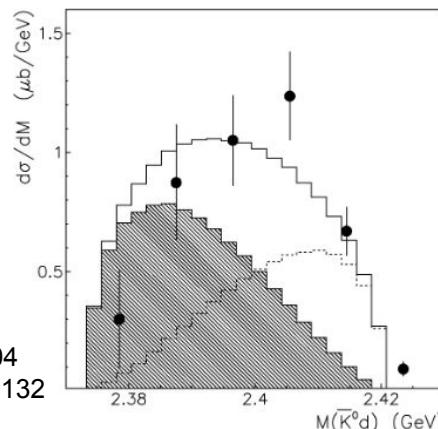
K^-n , K^-d experimental constraints

- K^-n , K^-d scattering data



- COSY at Jülich:
 $pp \rightarrow d\bar{K}^0 K^+$

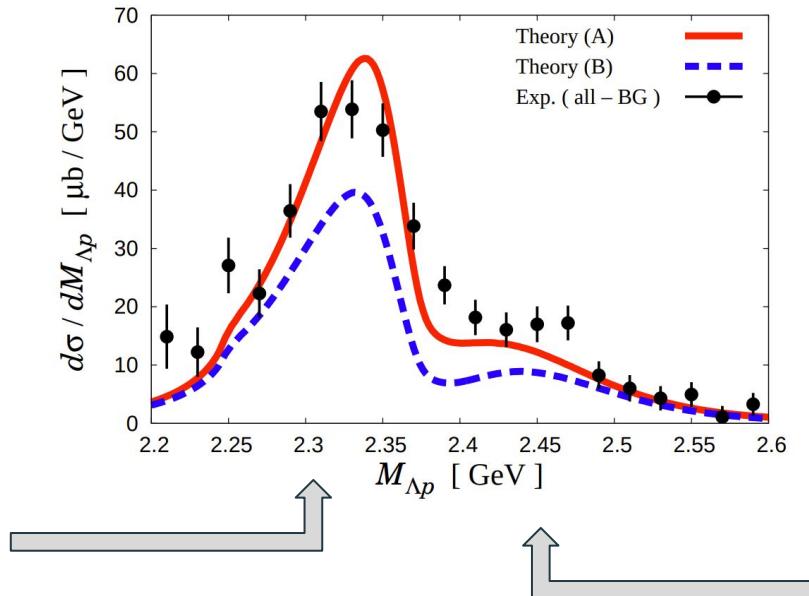
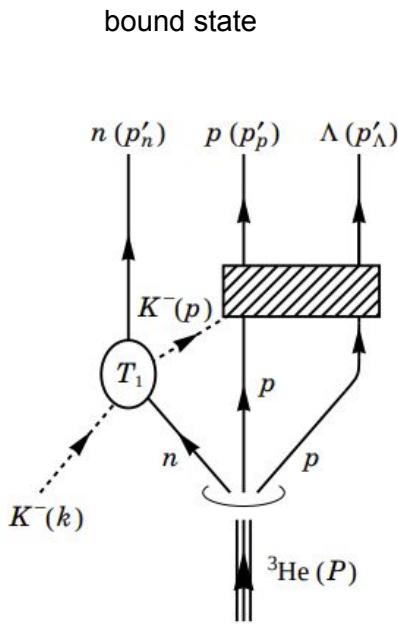
K^-d scattering length constraints:
 $|Re\,a| \leq 1.3$ fm, $Im\,a \leq 1.3$ fm



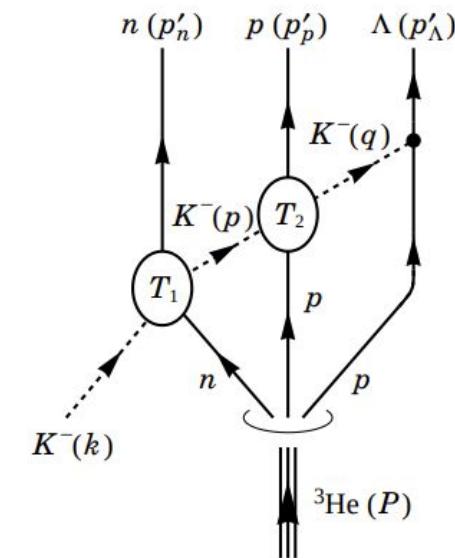
V. Kleber et al., PRL 91 (2003) 172304
A. Sibirtsev et al PLB 601 (2004), p. 132

E15

Sekihara, Oset, Ramos
 JPS Conf. Proc. 26, 023009 (2019)
 PTEP 2016, 12, 123D03



quasi-elastic scattering via $\Lambda(1405)$ formation



SIDDHARTA-2: antikaonic deuterium

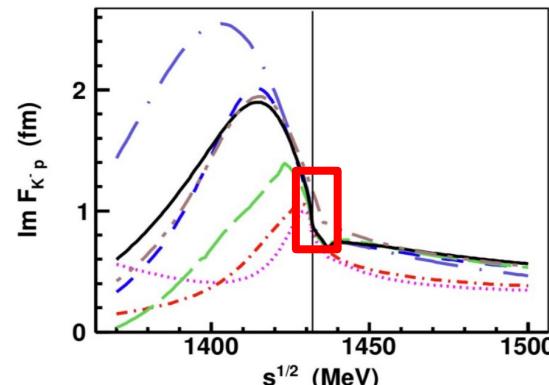
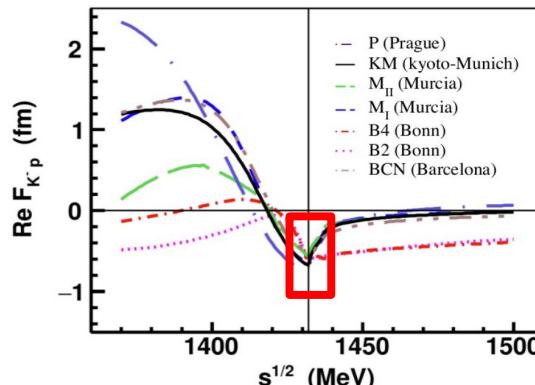
K-atom experiments @ J-PARC K1.8BR

	K-d (E57)	K- ^{3/4} He (E62)
Physics	$K^{\bar{N}}$ ($I=1$)	$K^{\bar{N}}$ -nucleus potential
X-ray transition	$2p \rightarrow 1s$	$3d \rightarrow 2p$
Energy	~ 8 keV	~ 6 keV
Width	~ 1000 eV	~ 2 eV
Yield (per stopped K-)	~ 0.1 % (0.04% of liquid D2 density)	~ 7 % (Liquid He)
X-ray detector	Silicon Drift Detector	TES microcalorimeter
FWHM resolution	~ 150 eV	~ 5 eV
Effective area	~ 200 cm ²	~ 0.2 cm ²

High sensitivity

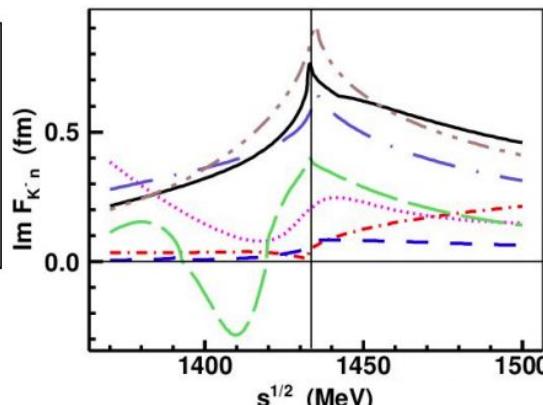
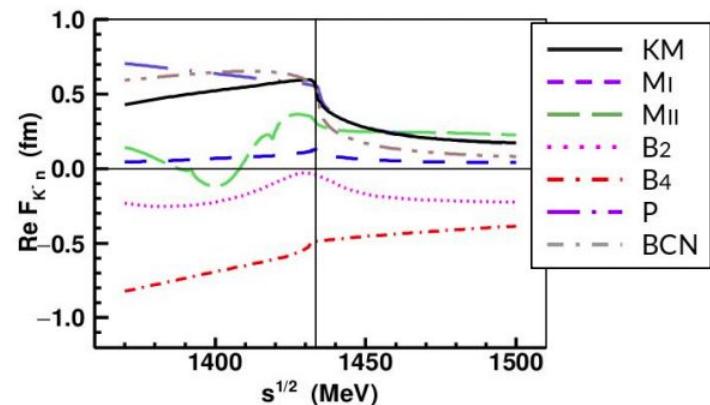
High resolution

antikaonic hydrogen: SIDDHARTA



⇒ Kaonic hydrogen
SIDDHARTA constraint

A. Cieply, J. Hrtánková, J. Mareš, E. Friedman, A. Gal and A. Ramos, AIP Conf. Proc. 2249, no. 1, 030014 (2020).



⇒ Lack of experimental
K-n constraint

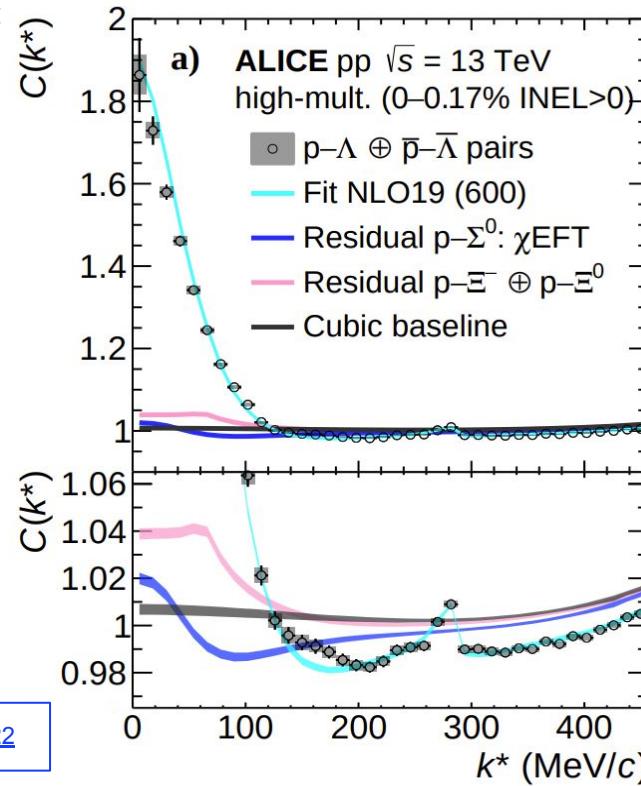
Femtoscopy

KbarN is also the ideal playground for femtoscopy studies

Connected to the main issues addressed using femtoscopy:

- Coupled channel dynamics
- Strangeness in NS
- Bound states

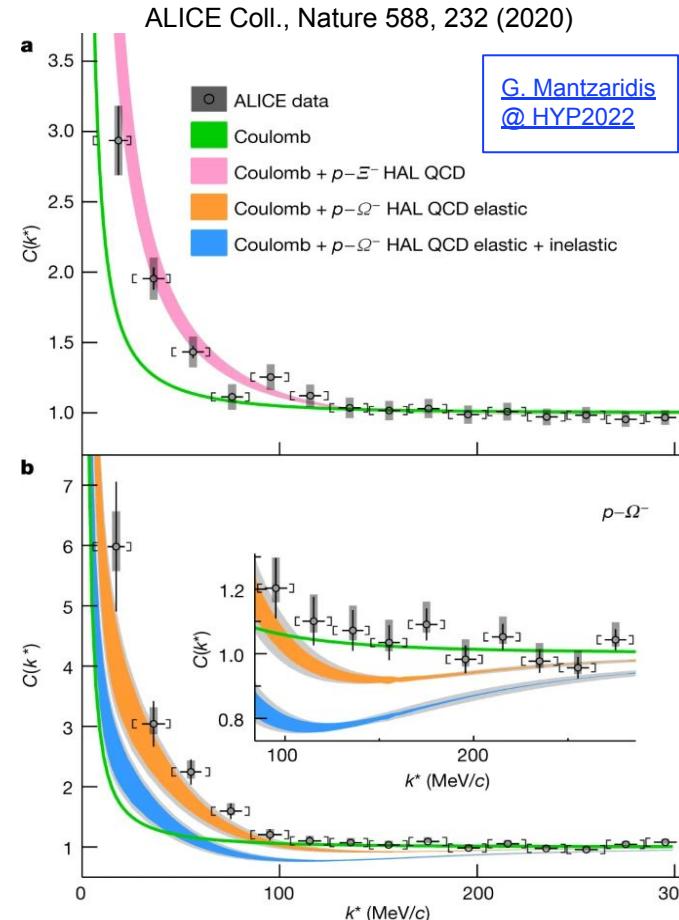
ALICE Coll., arXiv:2104.04427 [nucl-ex]. PLB in press.



Femtoscopy

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Single and Multi-Nucleonic absorptions

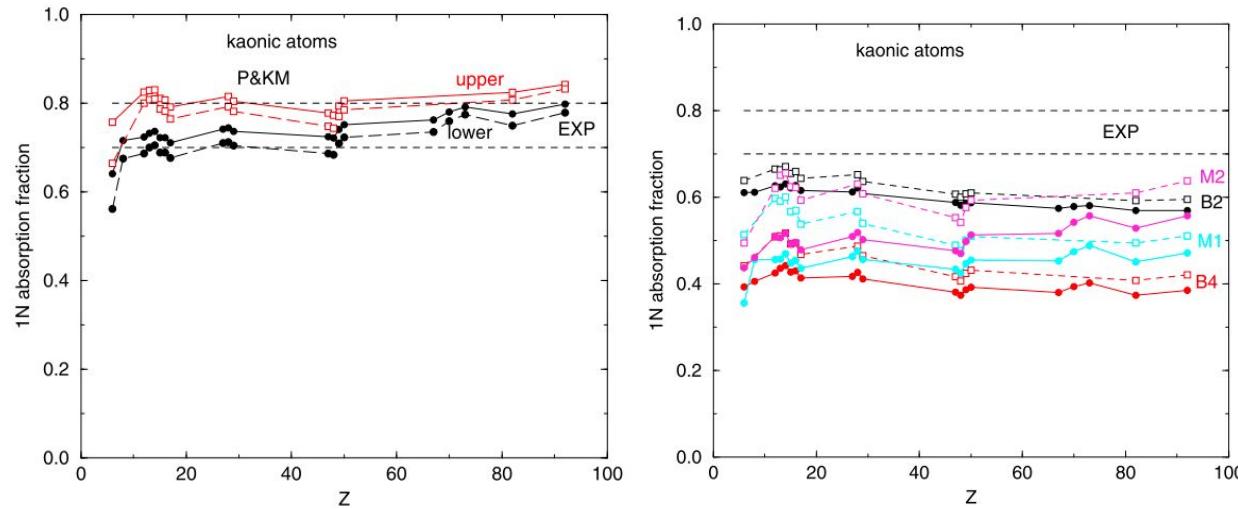
M multinucleonic absorption contributes between 20-30% constant in Z

Main background for kaonic nuclear states

KbarN potentials need to set constraints for such 1NA, 2NA, 3NA, etc

Chiral models tested in the framework of antikaonic atoms up to high Z:

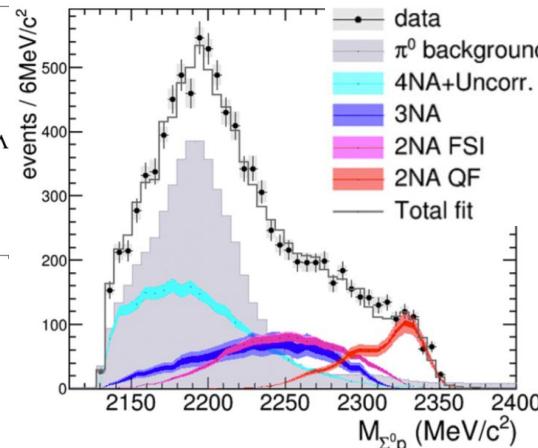
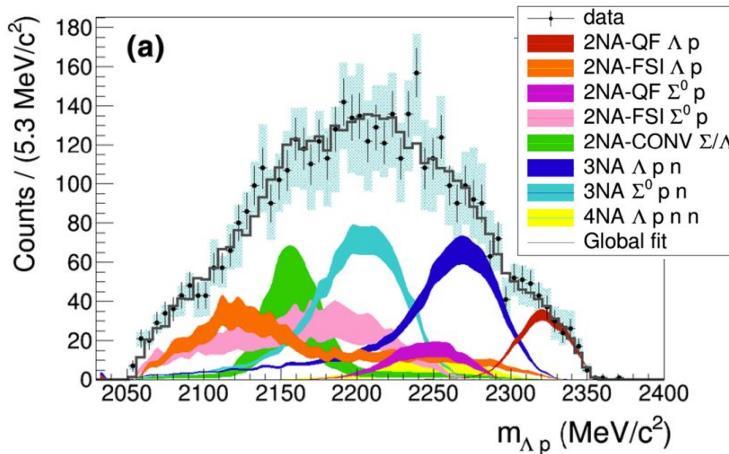
Nuclear Physics A 959 (2017) 66–82



⇒ to be incorporated to few-body calculations in order to predict BE vs Width for KNN, KNNN,

...

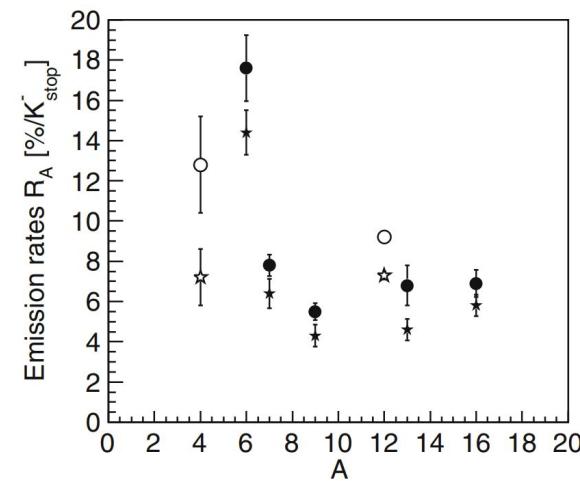
AMADEUS above and below threshold



AMADEUS Coll. PLB 758
(2016) 134, PLB 782 (2018)
339, FBS 62 (2021) 7.

$$(K^- 2NA \rightarrow YN) = (21.6 \pm 2.9(\text{stat.})^{+4.4}_{-5.6}(\text{syst.})) \%$$

FINUDA Coll. Hyperfine Interact (2012) 210:97–101
Sigma-p emission rates from A 6 to 16



emission rates of $K_{\text{stop}}^- + A \rightarrow \Sigma^+ + \pi^- + A'$, circles, and $K_{\text{stop}}^- + A \rightarrow \Sigma^- + \pi^+ + A'$,